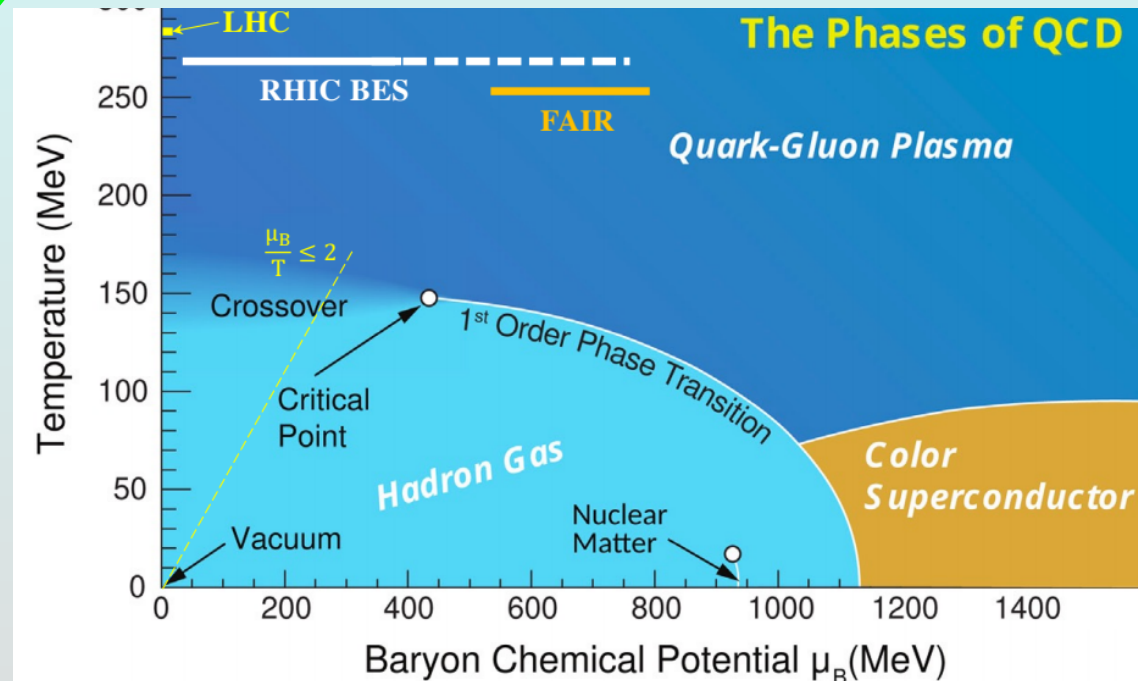
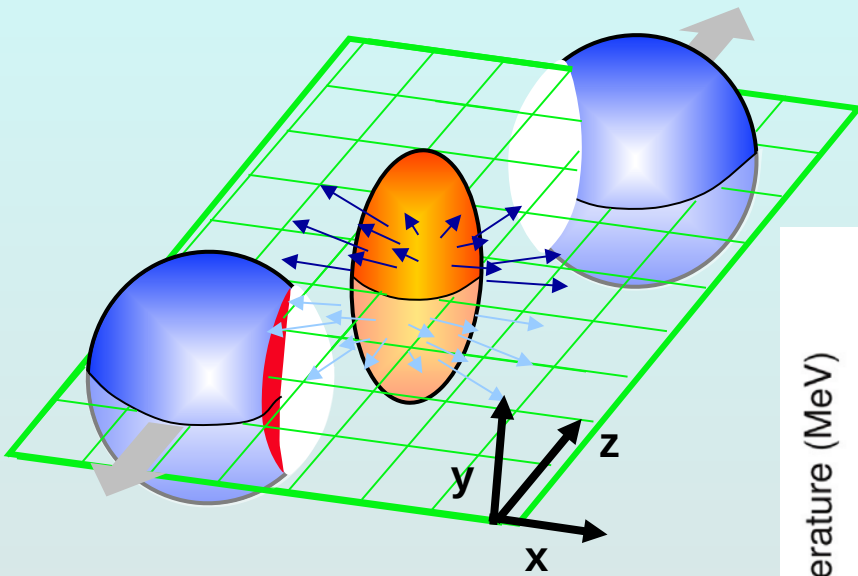


QCD Matter in Collisions of Heavy Ions

Lecture 1: Quark Gluon Plasma = many-body QCD



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June 10, 2023

Outline(s)

update

Lectures on QCD Matter in ion collisions

1. Intro, radiation, hydrodynamic flow. T , ε , η/s
2. Energy loss and opacity in the quark gluon plasma
3. Properties of strongly coupled plasmas and what the string theorists tell us
4. Electromagnetic probes, screening, and QGP outlook

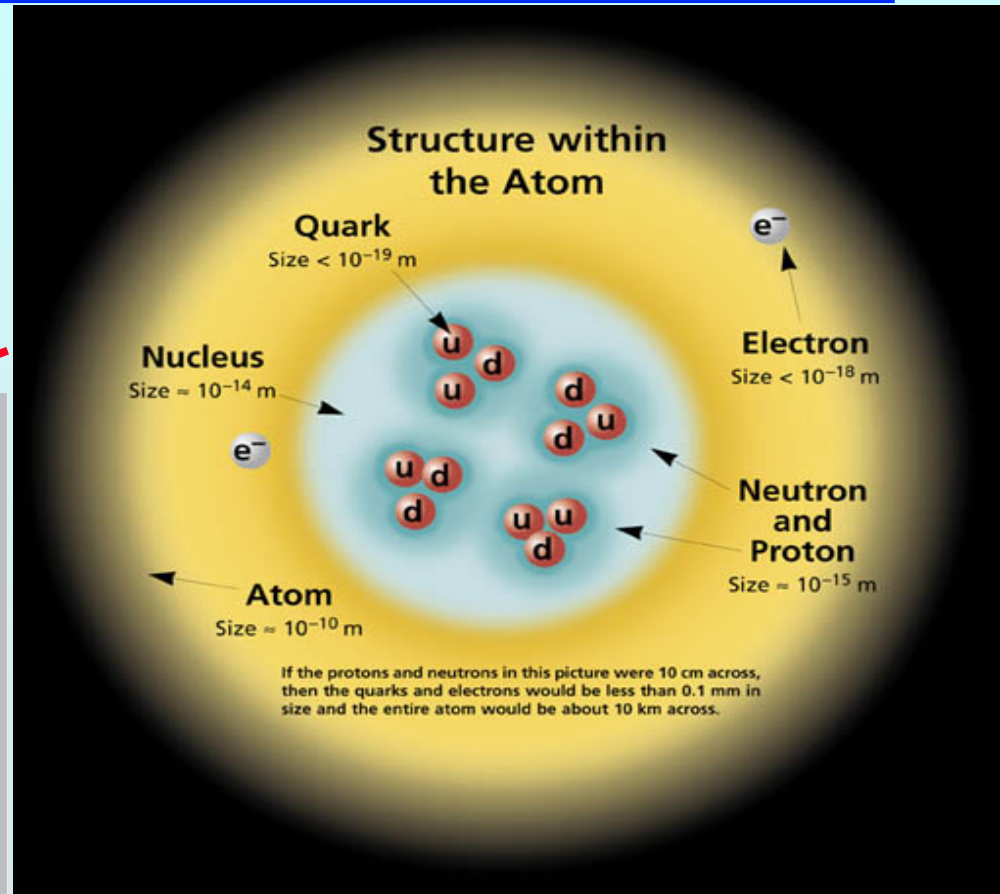
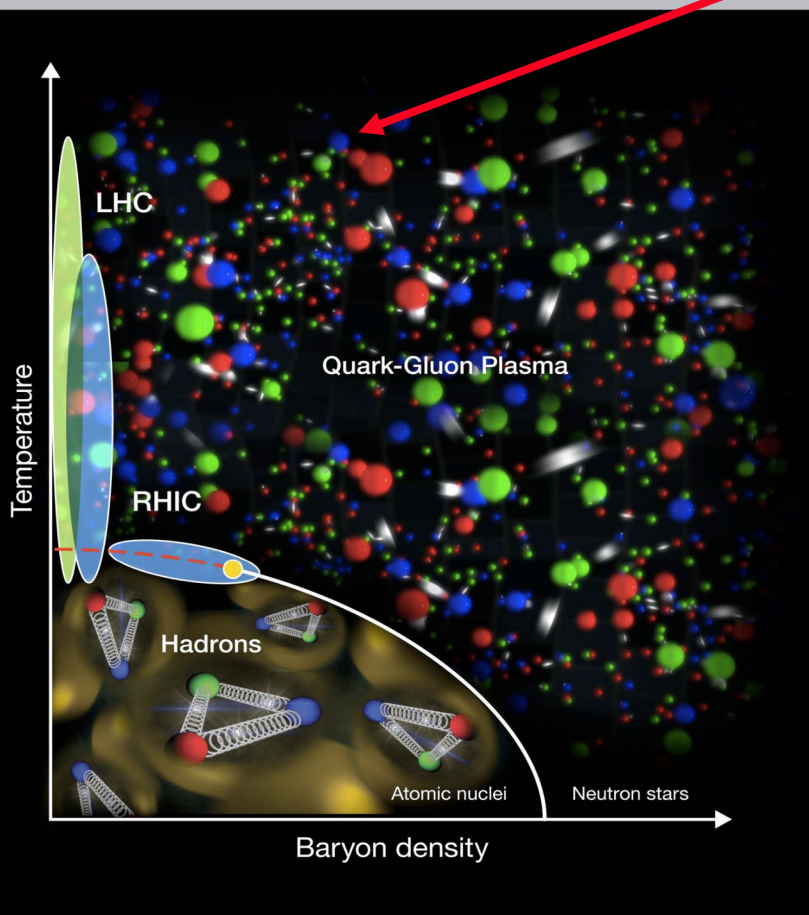
TODAY

- Introduction and motivation to study the QCD plasma
- Stages and timescales in heavy ion collisions
- Geometry of the collisions
- Collective flow, hydrodynamics, and “the perfect liquid”
- Effects of fluctuations

Create the hot, dense QCD matter in the lab

Heat to $T > 10^{12}$ K

last seen: $\sim 1 \mu$ second after the Big Bang!

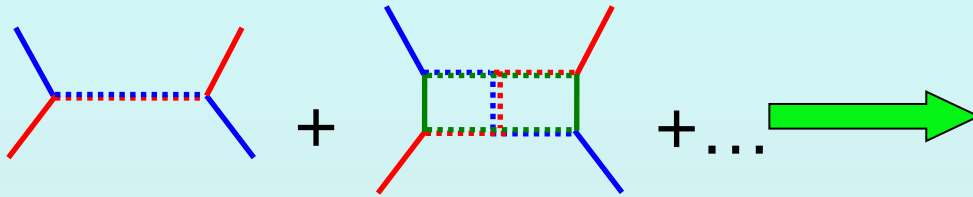


How does it work?
What are its properties?
Phase transitions? (KR)

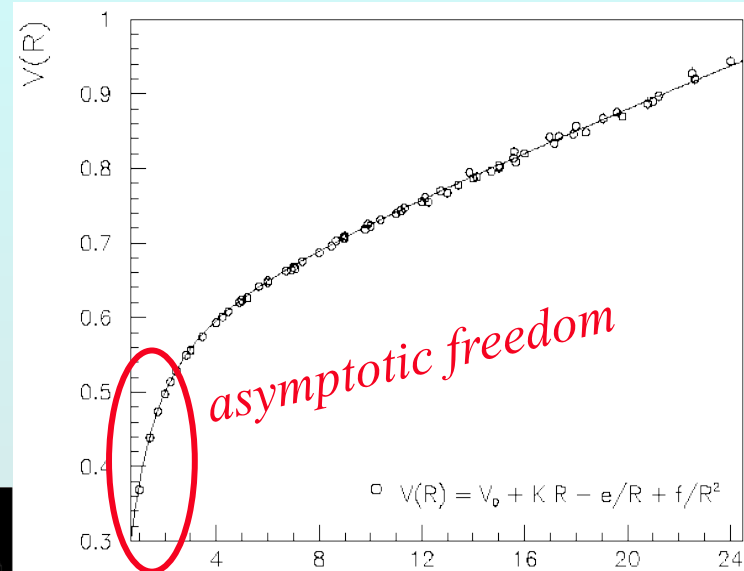
What to expect in hot, dense QCD matter?

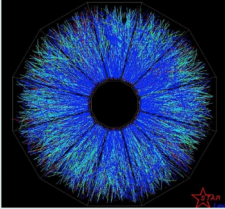
- Gluons carry color → interact among themselves

theory is non-abelian



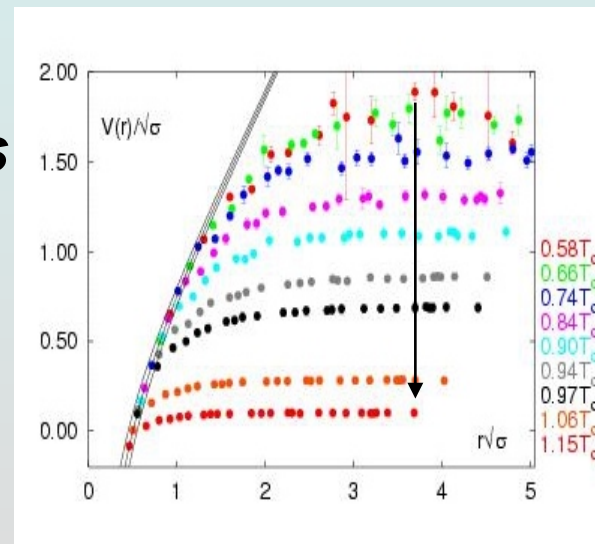
- Curious property at large distance: confinement of quarks in hadrons



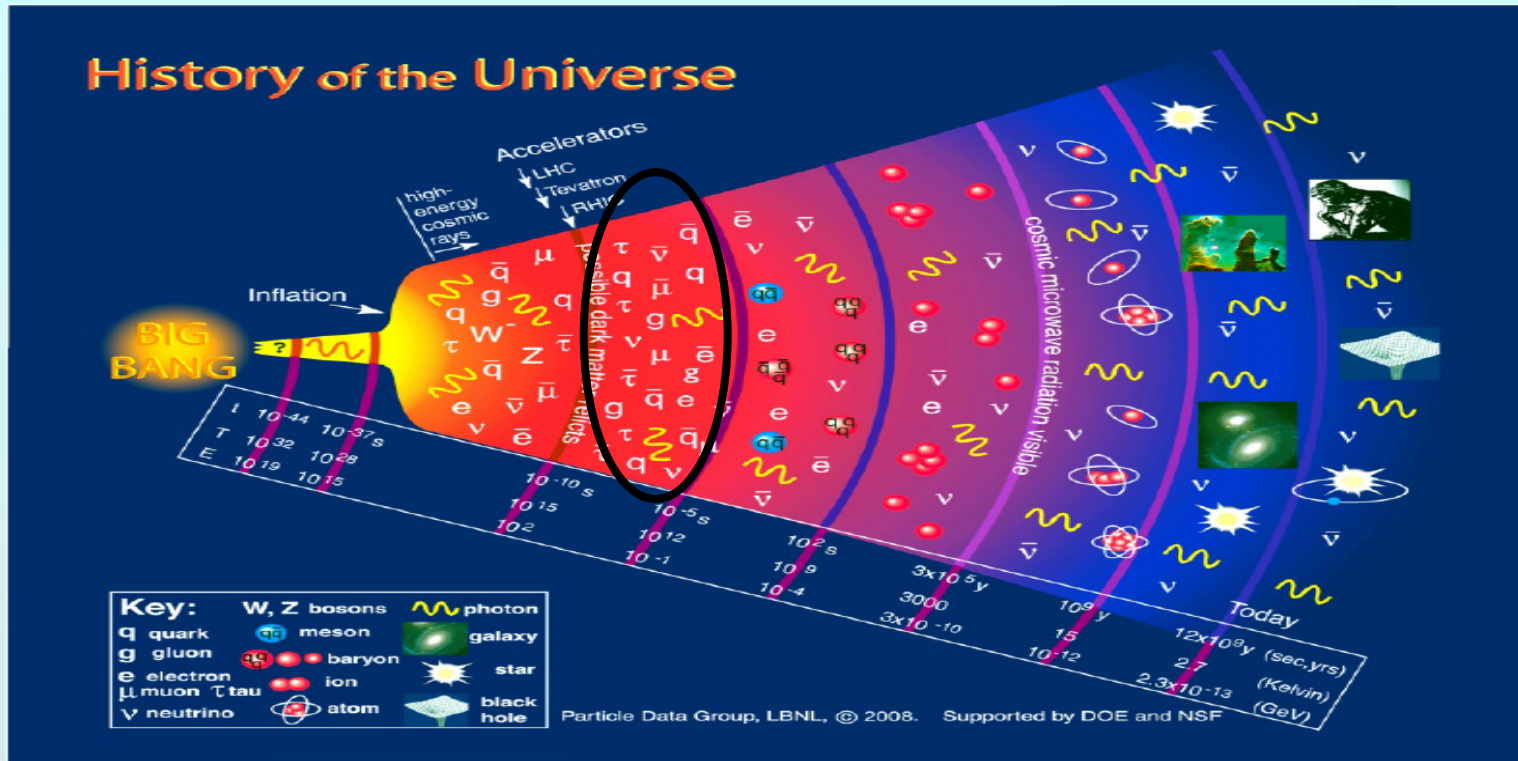
At high temperature/density  screening by produced colored particles

Expect phase transition to deconfined quark gluon plasma

Lattice QCD → $T_c \sim 150$ MeV



Quark gluon plasma in early universe

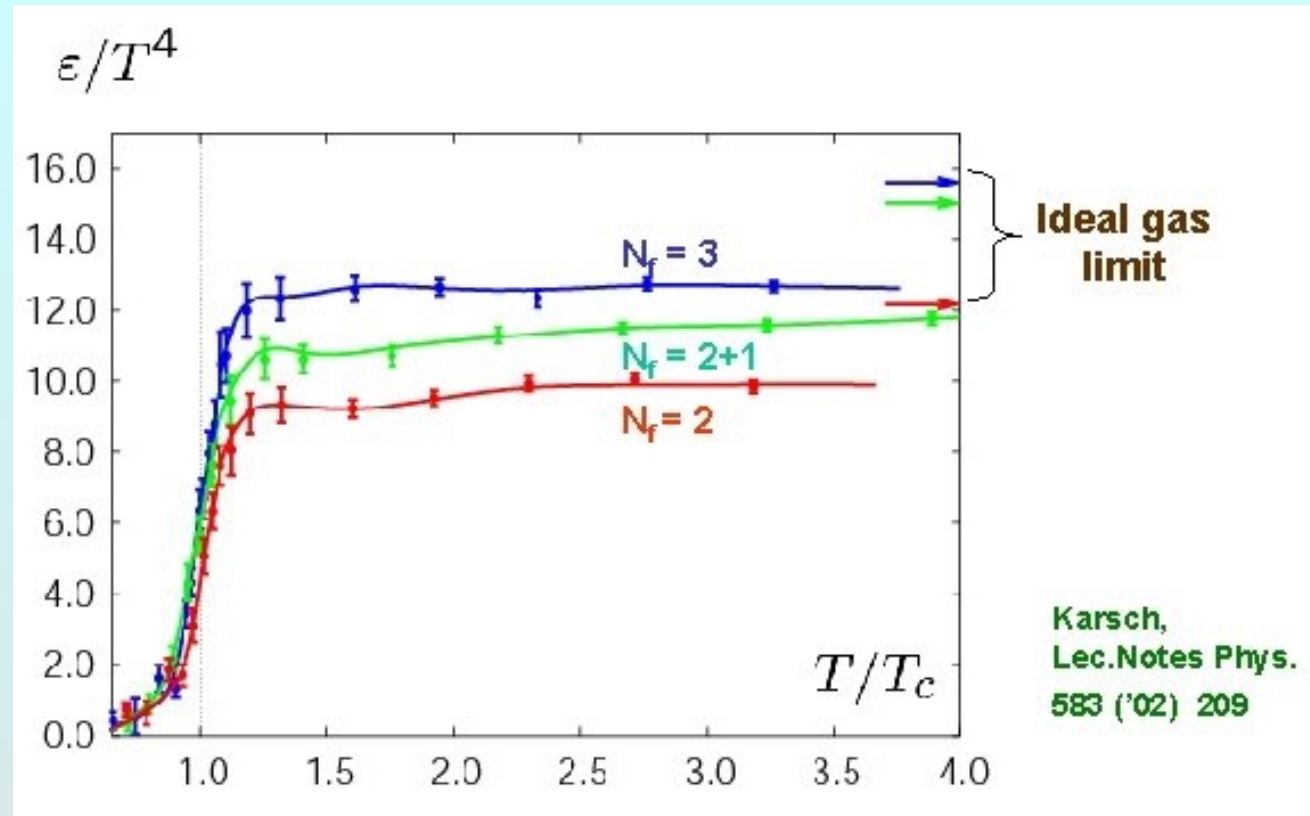


QGP likely also in core of massive neutron stars

- Electromagnetic interactions: photons are exchanged photons are electrically uncharged
- Interaction between quarks = strong interaction
Force is carried by exchanged gluons
Both quarks and gluons have a “color charge”

Why many-body interactions?

- Initial idea:
nearly an ideal
gas



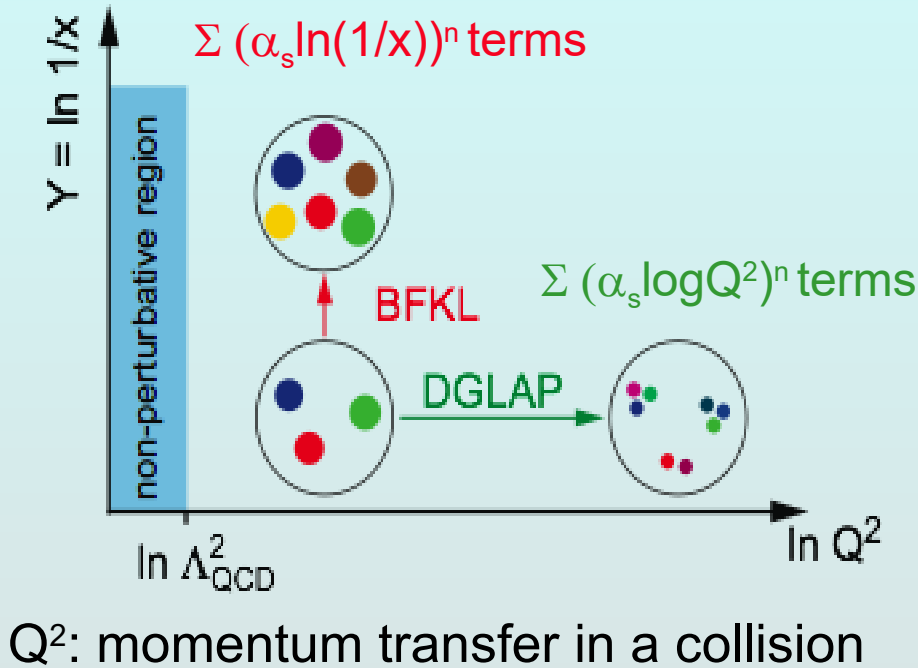
- Range of strong interaction: ~ 1 fm
- Number of gluons inside 1fm radius sphere: lots!
So, why an ideal gas? Should have many-body interactions

Starting point: inside a nucleon

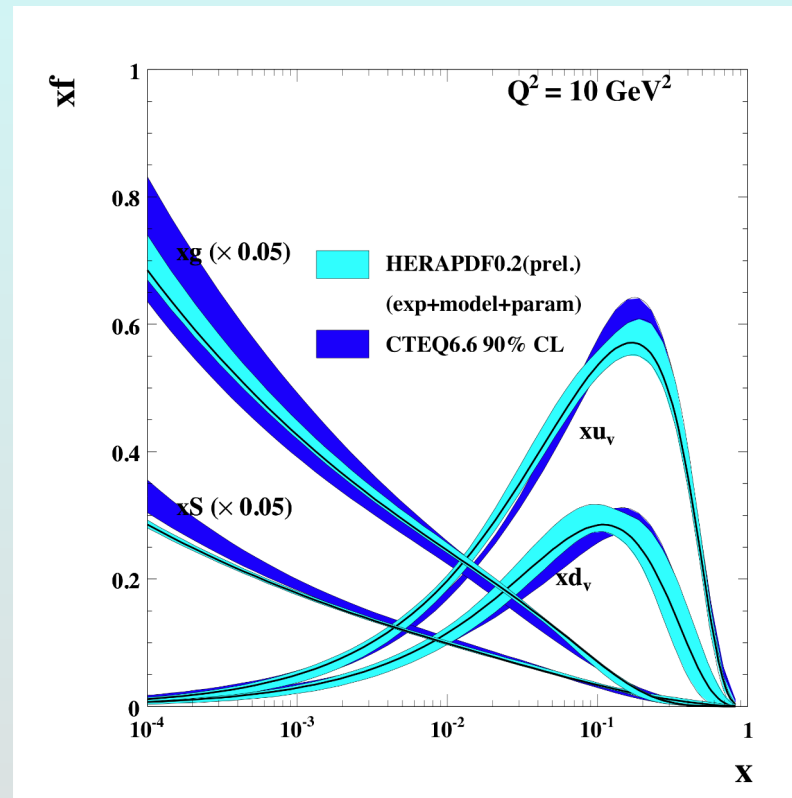
Theorists' view

Experimenter's view

x: momentum fraction carried by parton

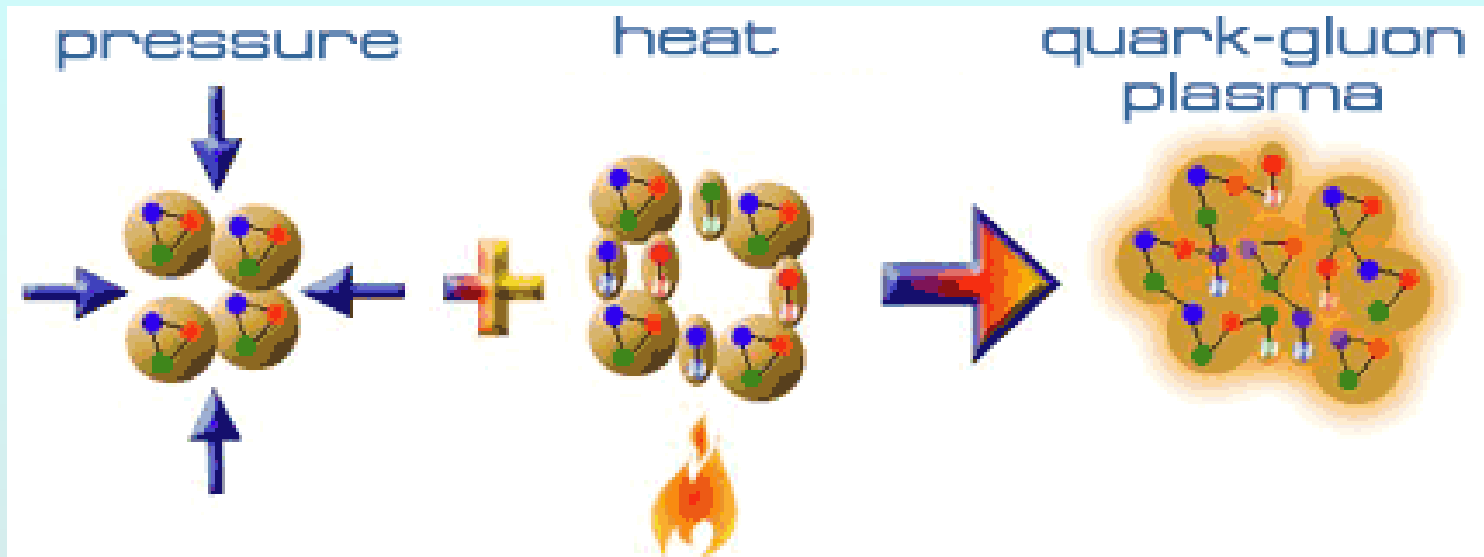


Scatter electrons off a p



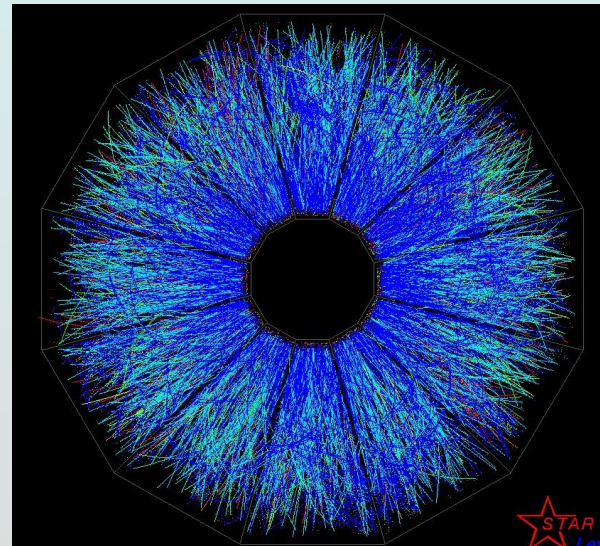
Gluon number is not fixed
Virtual $q\bar{q}$ pairs abundant

From the practical point of view



What we see in the lab:

All the particles which come out at the end of the collision



The heaters

Large Hadron Collider



CERN in Geneva

Pb+Pb @ 2.76 TeV/A

Relativistic Heavy Ion Collider

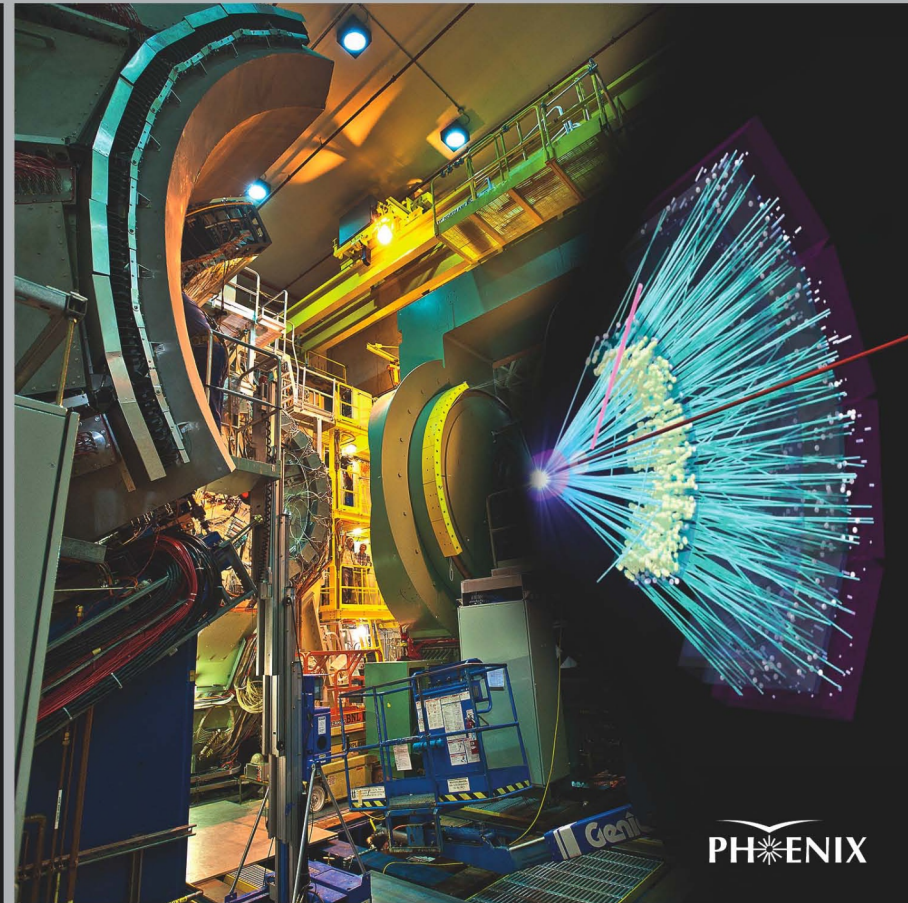


Brookhaven in New York

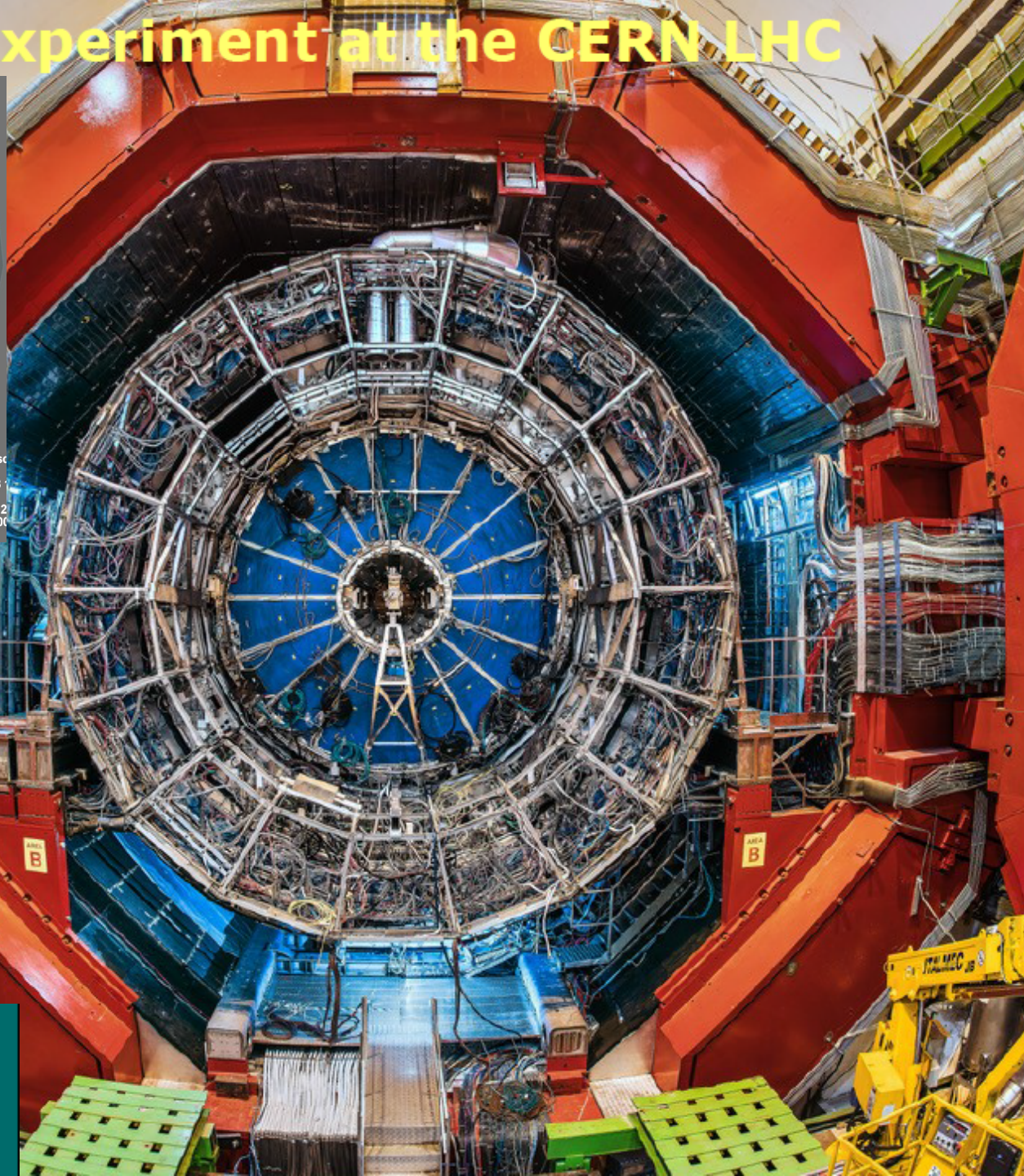
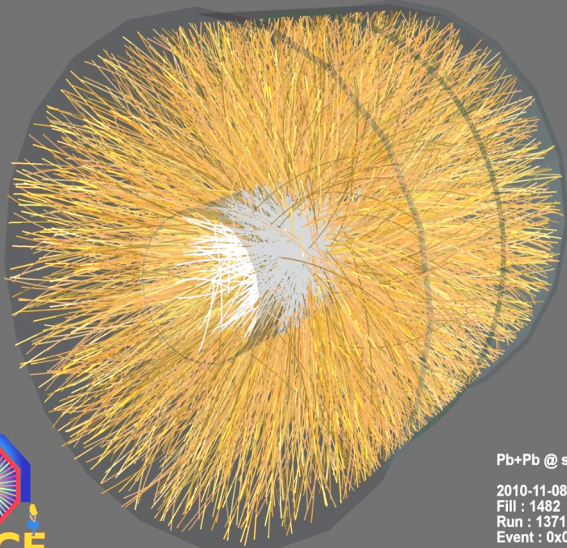
Au+Au @ 200 MeV/A

*Collide heavy nuclei for max temperature & volume
p+p and p+A for comparison*

Experiments at RHIC



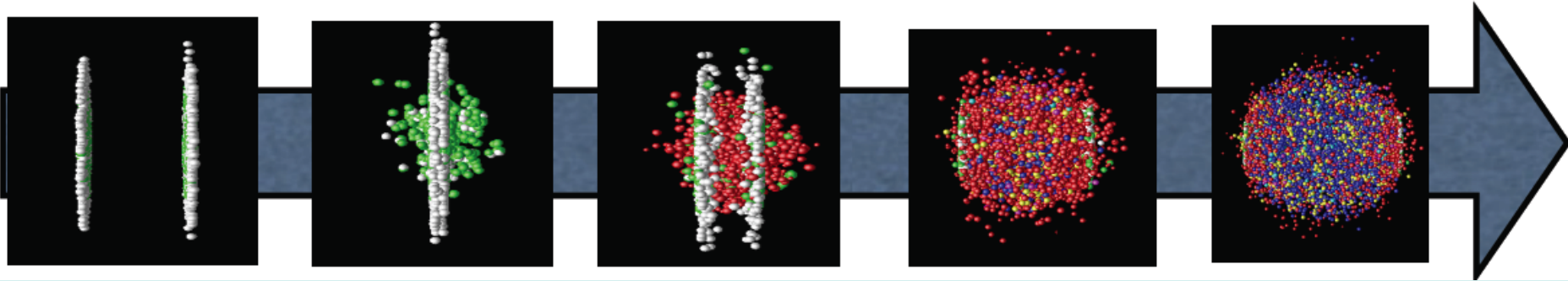
The ALICE experiment at the CERN LHC



**1000 scientists
in 172 institutes
from 40 countries**

Nuclear collision timeline

plasma lives $\sim 3 \times 10^{-23}$ seconds, $\sim 10^{-12}$ cm across



Lorentz contracted nuclei on their way in.

First scattering of q & g inside the nucleons. Some high momentum transfers.

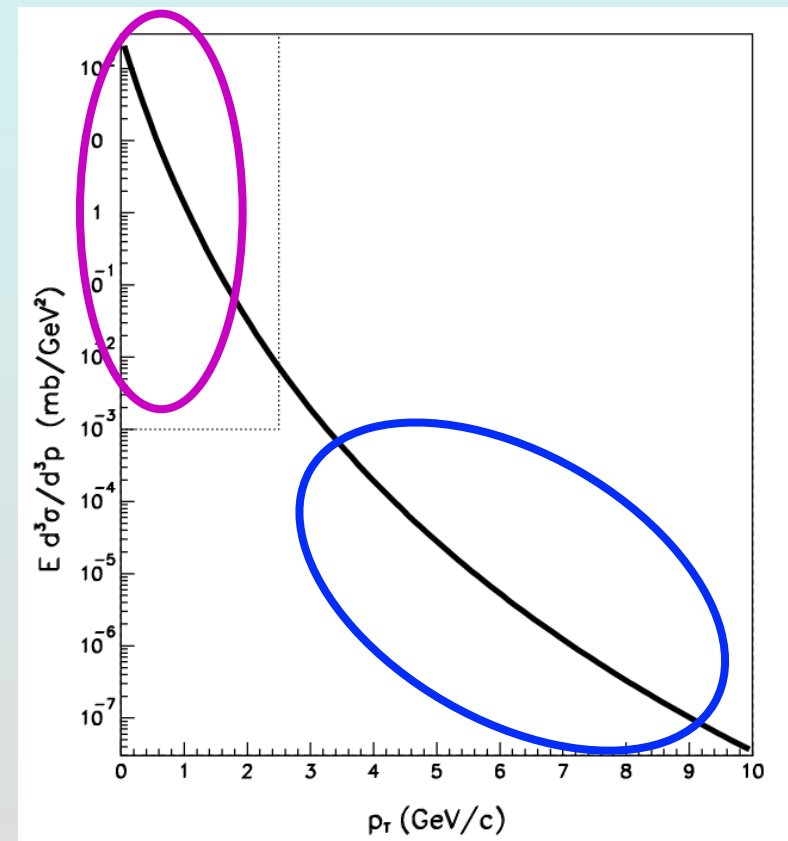
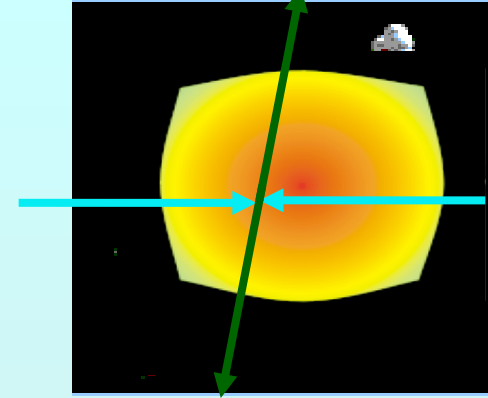
Secondary collisions, creating high density and temperature

Quark gluon plasma expands and cools, eventually condensing into hadrons.

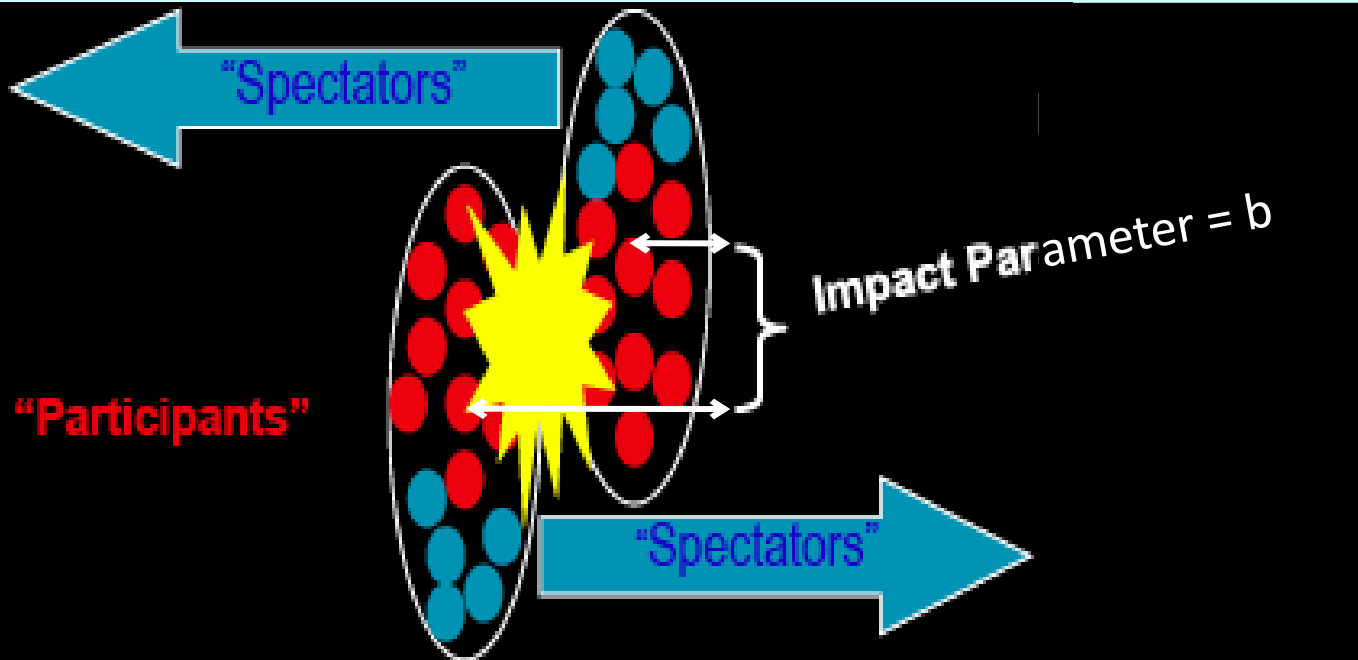
Hadron gas interacts, expands and cools further. Eventually collisions stop & hadrons stream freely.

study plasma with radiated & “probe” particles

- as a function of transverse momentum
90° is where the action is (max T, ρ)
 p_L between the two beams: midrapidity
- $p_T < 1.5$ GeV/c
“thermal” particles
radiated from bulk medium
“internal” plasma probes
- $p_T > 3$ GeV/c
large E_{tot} (high p_T or M)
set scale other than T(plasma)
autogenerated “external” probe
describe by perturbative QCD
- control probe: photons
EM, not strong interaction
produced in Au+Au by QCD
Compton scattering



Geometry matters



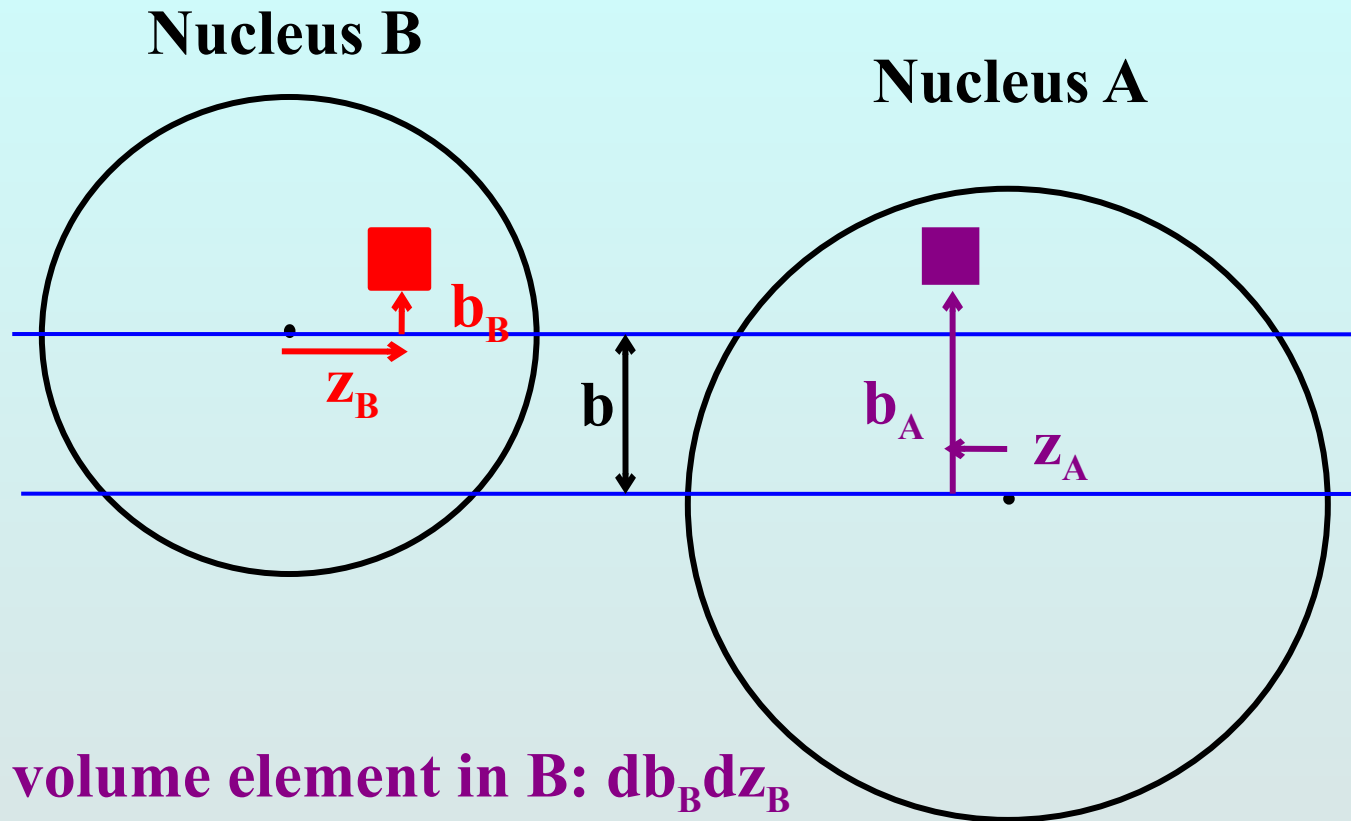
Use Glauber model of nucleons in the nucleus

calculate # of participant nucleons N_{part}

of binary NN collisions N_{coll}

Central collisions ($b \sim 0$) produce maximum volume of plasma

Glauber model: calculate probabilities



volume element in B: $db_B dz_B$

volume element in A: $db_A dz_A$

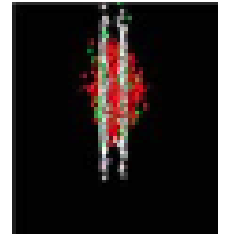
Probability of finding a nucleon in volume element B =

$$\rho_B(b_B, z_B) db_B dz_B \quad (\rho_B \text{ is nuclear density * nucleons in B})$$

OK, so what happens when two nuclei
collide?

NB: $b \neq 0$!

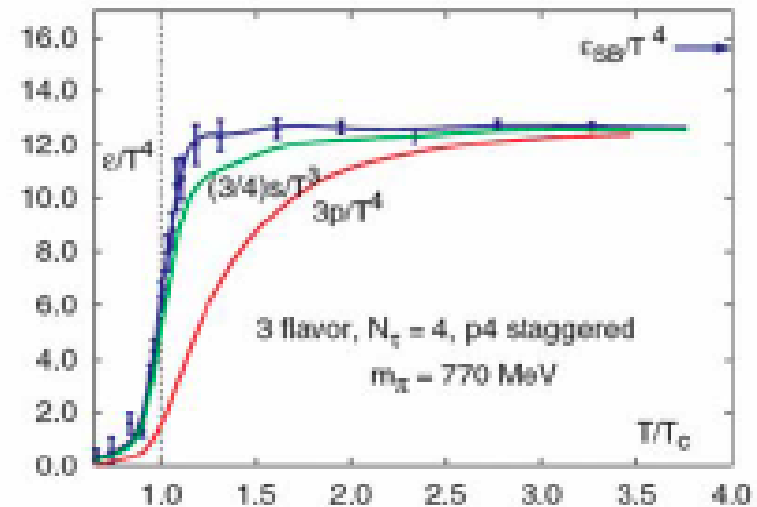
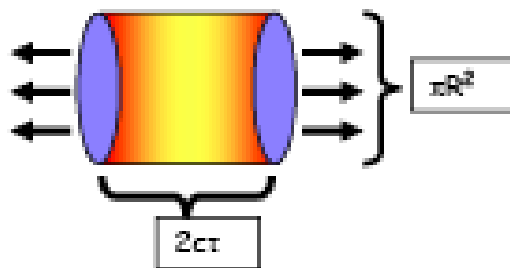
Energy Density



Energy density far above transition value predicted by lattice.

$$\varepsilon_{Bj} = \frac{1}{\pi R^2} \frac{1}{2c\tau} \left(2 \frac{dE_T}{dy} \right)$$

$R \sim 7 \text{ fm}$



Lattice: $T_c \sim 190 \text{ MeV}$ ($\varepsilon_c \sim 1 \text{ GeV/fm}^3$)

PHENIX: Central Au-Au yields

$$\left\langle \frac{dE_T}{d\eta} \right\rangle_{\eta=0} = 503 \pm 2 \text{ GeV}$$

$$\varepsilon \sim 15 \frac{\text{GeV}}{\text{fm}^3} @ \tau = 0.6 \text{ fm} / c \text{ (thermalization)}$$

Quark Gluon Plasma properties

- thermodynamic properties (equilibrium)

T, P, ρ

Equation Of State (relation btwn T, P, V, energy density)

v_{sound} , static screening length

- transport properties (non-equilibrium)*

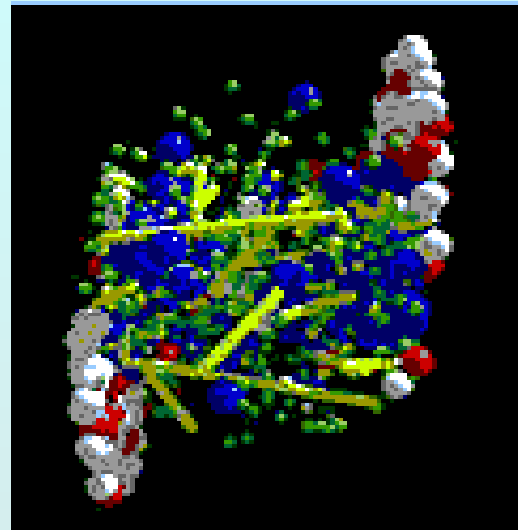
particle number, energy, momentum, charge

diffusion *sound* *viscosity* *conductivity*

*measuring these is new for nuclear/particle physics!

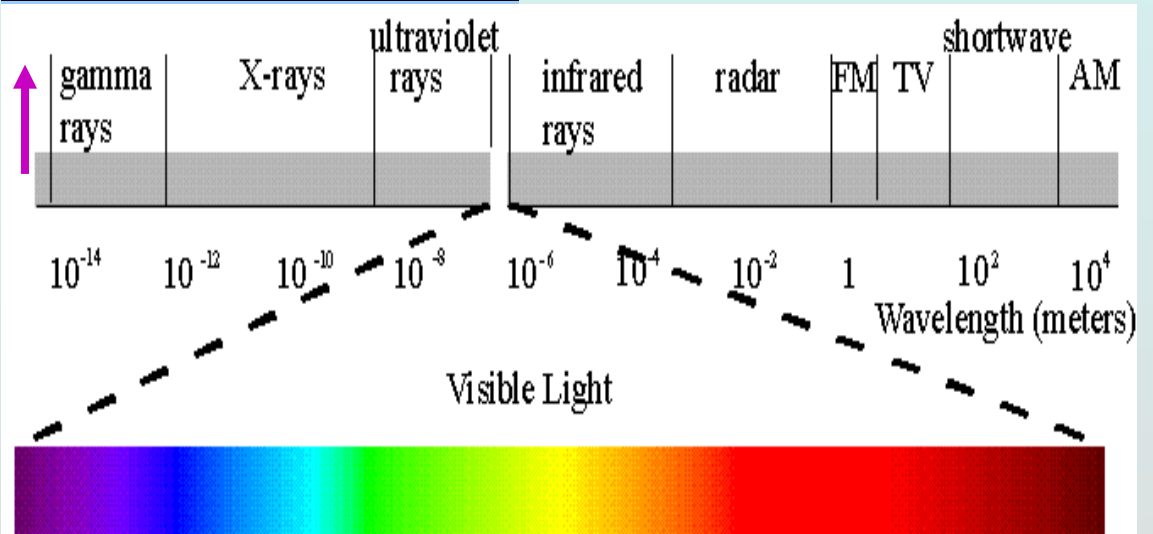
How hot is our QCD matter?

Hottest Science Experiment on the Planet*



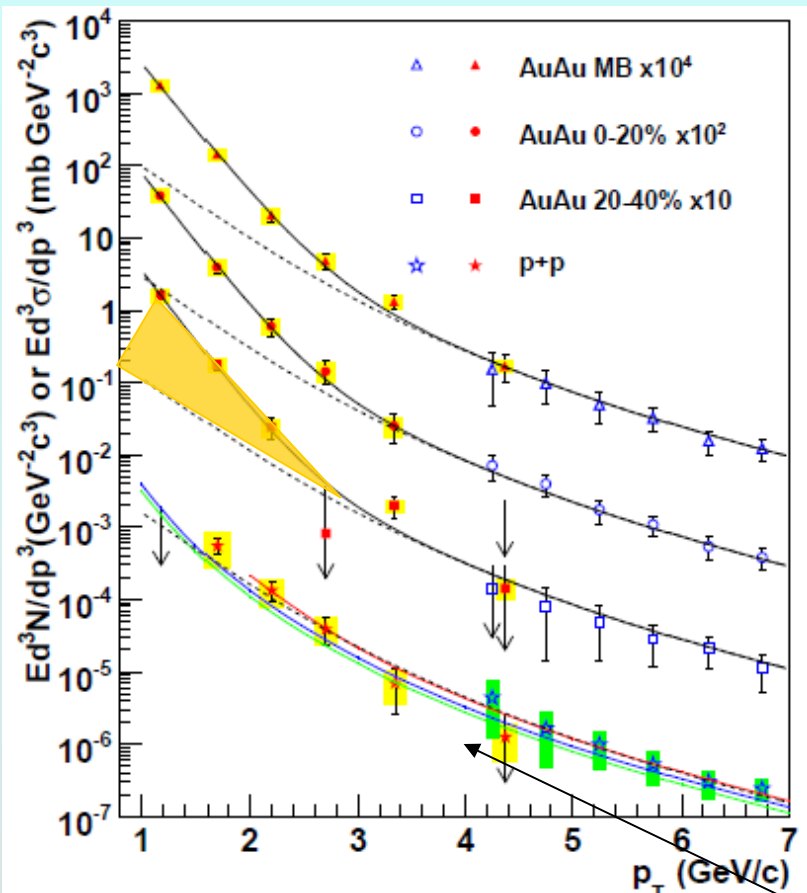
plasma lives for 3×10^{-23} s
droplet is 10^{-12} cm across

can't use a thermometer!
So we look at radiation

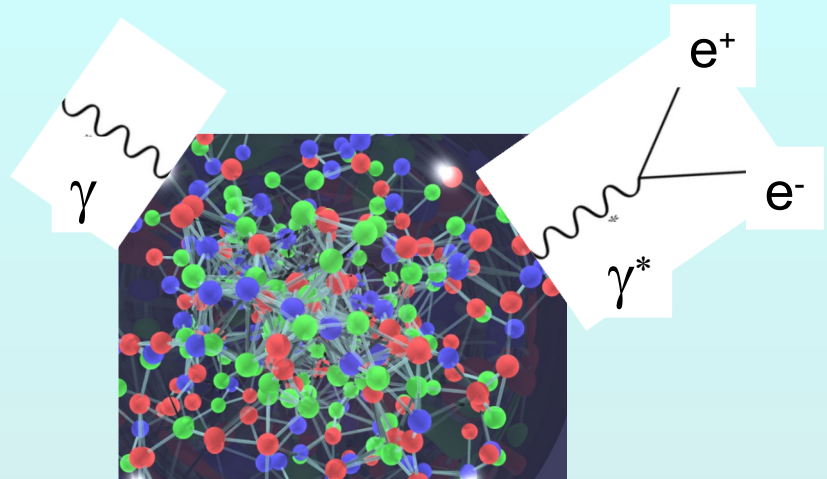


* According to Discover Magazine (2010)

Thermal radiation



PRL 104, 132301 (2010)



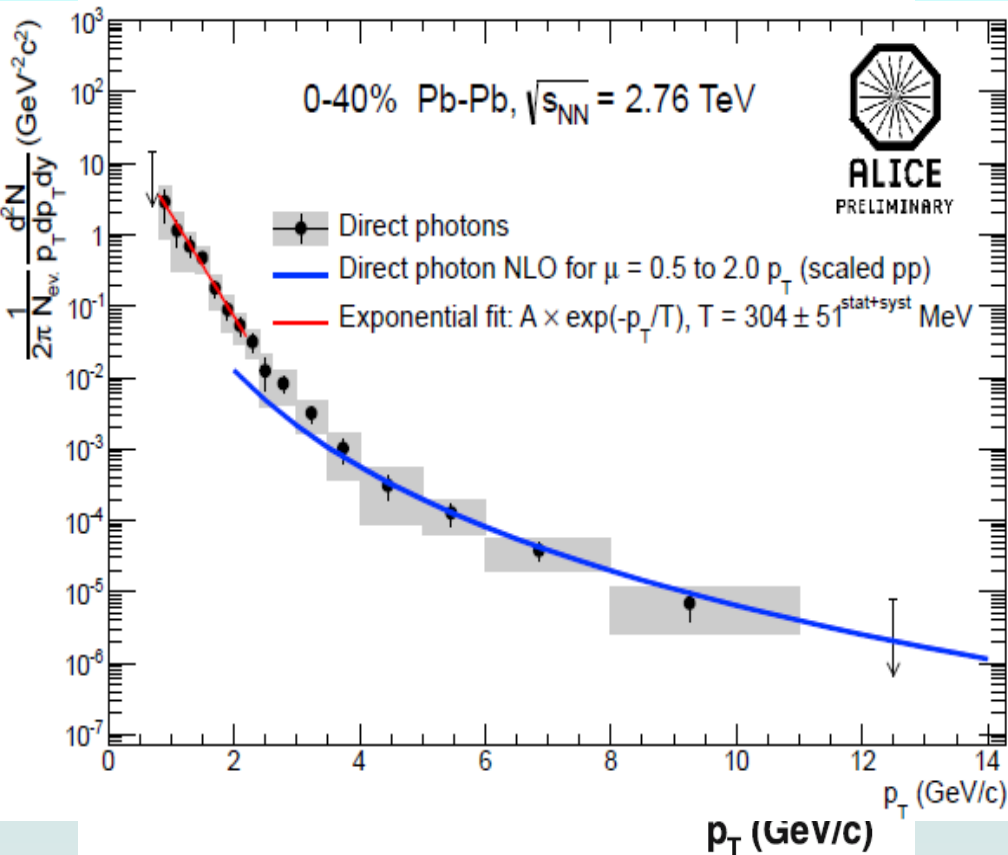
**Low mass, high p_T $e^+e^- \rightarrow$
nearly real photons
Large enhancement above
 $p+p$ in the thermal region**

pQCD γ spectrum
(Compton scattering @ NLO)
agrees with $p+p$ data

Analogy to the bronze is not quite right

- Similar to black body radiation, but...
- The photons are not bouncing around in equilibrium with QGP
 - Produced by interactions among partons in equilibrium
 - Exit the plasma with no further (strong) interactions
- The plasma is not static
 - It is expanding at $v=c$ longitudinally and $v\sim 0.5c$ radially
 - Photons arise from velocity boosted partons
- What to do about this?
- Try hydrodynamics to play the movie backwards!

direct photons: $T_{init} > T_c$!



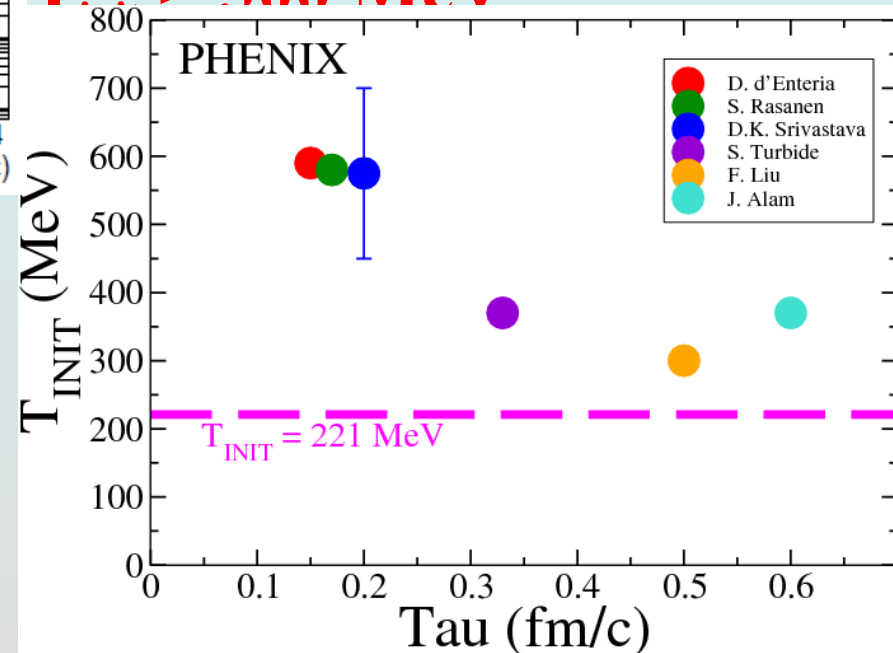
- Exponential fit in p_T :

$$T_{avg} = 221 \pm 23 \pm 18 \text{ MeV}$$

But, a hot system expands!

- Use hydrodynamics models to reproduce data

$$T_{init} > 300 \text{ MeV}$$

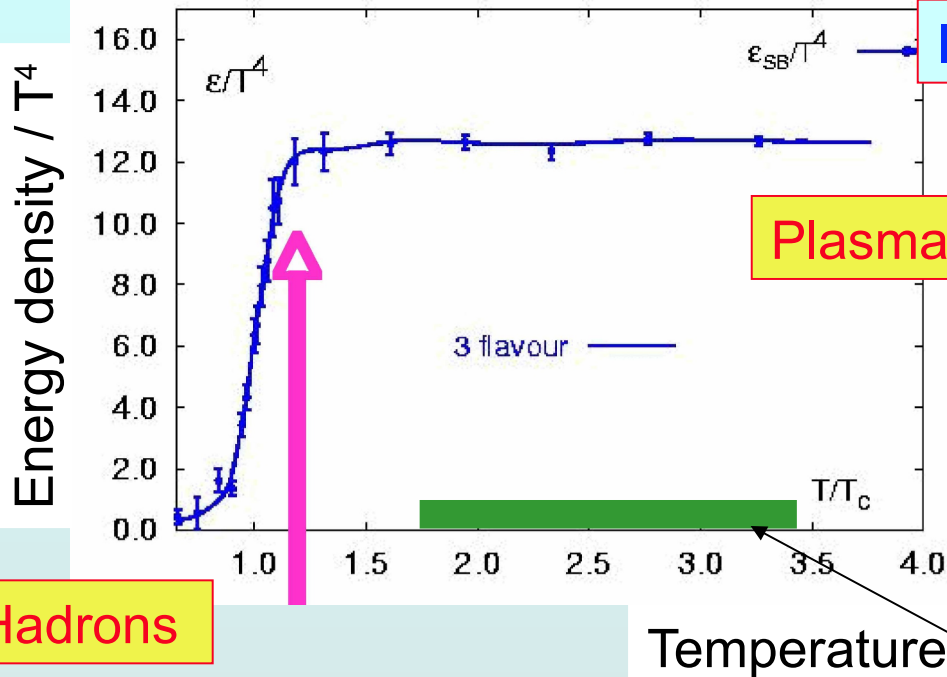


NB: $T_c \sim 150$ MeV

@ LHC $T_{avg} = 304 \pm 51$ MeV

$T \sim 30\%$ higher than at RHIC

Now we are on the map



$$\epsilon = g \frac{\pi^2}{30} T^4$$

Energy density $\propto T^4$
more degrees of freedom
in the plasma phase

$$T_c \sim 150 \pm 10 \text{ MeV}$$

$$\epsilon \sim 3 \text{ GeV/fm}^3$$

We are here
 $T \sim 250\text{--}500 \text{ MeV}$

Is QCD matter evolution well described
by hydrodynamics??

Hydrodynamics

a critical tool to extract information from data

Hydrodynamic Equations

$$\partial_\mu T^{\mu\nu} = 0, \quad \text{Energy-momentum conservation}$$

$$\partial_\mu n_i^\mu = 0 \quad \text{Charge conservations (baryon, strangeness, etc...)}$$

For perfect fluids (neglecting viscosity!),

$$T^{\mu\nu} = (e + P)u^\mu u^\nu - P g^{\mu\nu}$$

Energy density

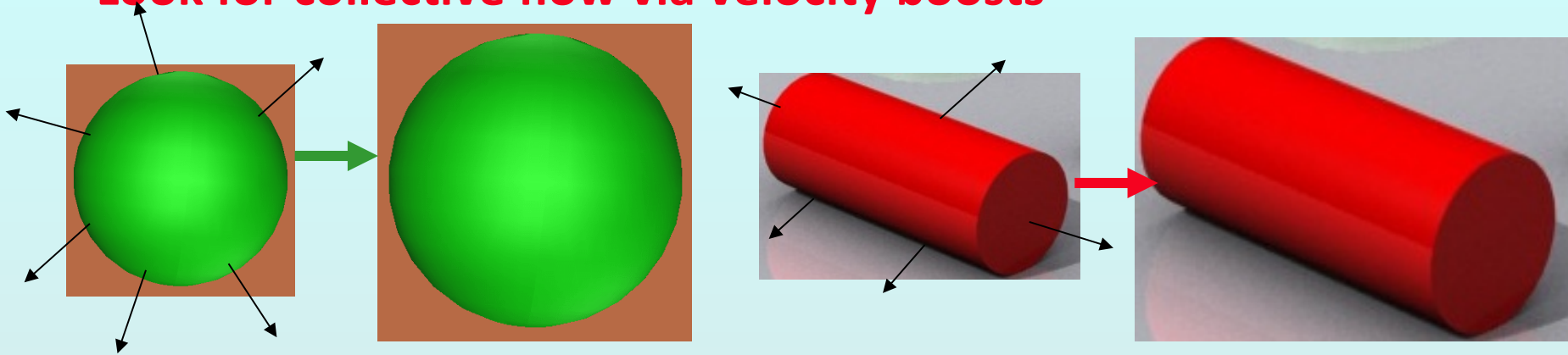
Pressure

4-velocity

Within ideal hydrodynamics, pressure gradient dP/dx is the driving force of collective flow.

Does the matter exhibit collectivity?

- Look for collective flow via velocity boosts



- Is the expansion hydrodynamical?

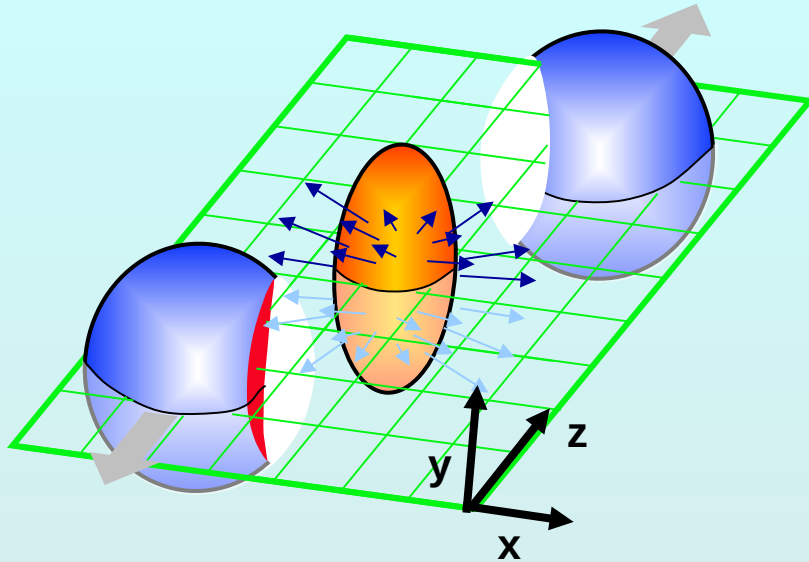
Model expansion of the system with fluid dynamics

$$\partial_t \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ e \end{pmatrix} + \partial_x \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ u(e + p) \end{pmatrix} + \partial_y \begin{pmatrix} \rho v \\ \rho uv \\ \rho v^2 + p \\ v(e + p) \end{pmatrix} -$$

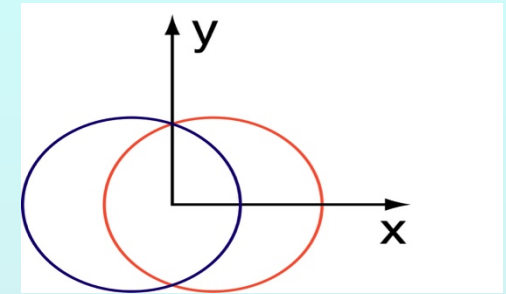
$$\partial_x \begin{pmatrix} 0 \\ \tau_{11} \\ \tau_{12} \\ \tau_{11}u + \tau_{12}v + k\partial_x\Theta \end{pmatrix} - \partial_x \begin{pmatrix} 0 \\ \tau_{21} \\ \tau_{22} \\ \tau_{21}u + \tau_{22}v + k\partial_y\Theta \end{pmatrix} = 0,$$

where u and v are the components of the velocity, ρ the density, p the pressure, e total energy density, τ_{ij} the components of the viscous part of the stress tensor, Θ the absolute temperature and k is the heat conductivity.

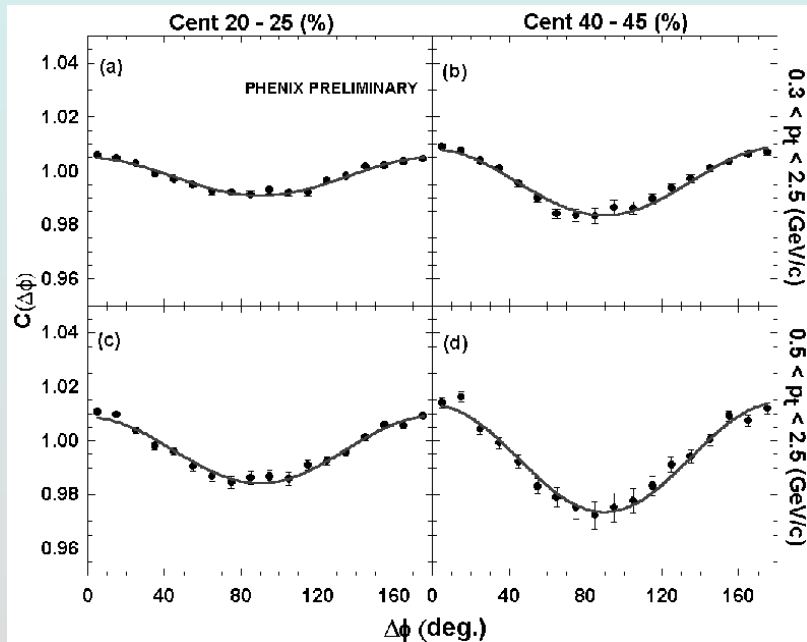
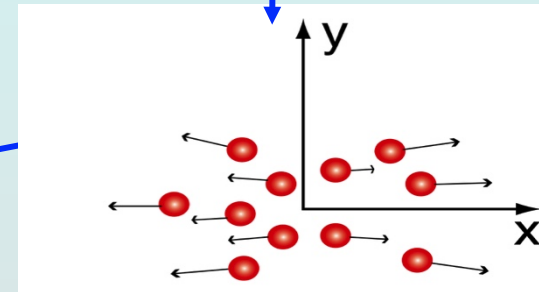
Measuring collective flow



Almond shape
overlap region
in **coordinate**
space

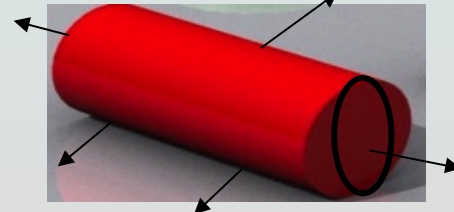


momentum
space

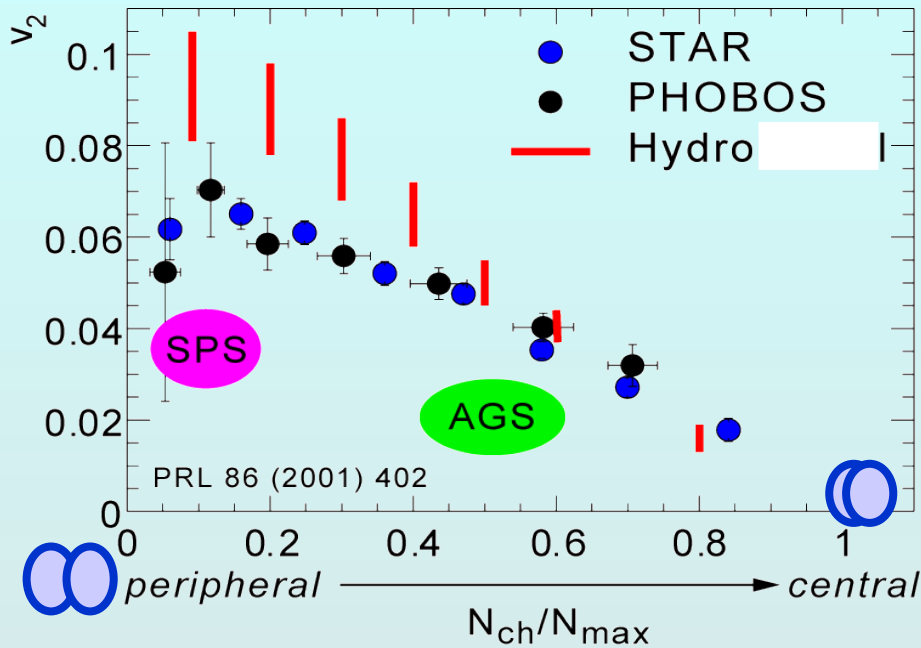


$$dN/d\phi \sim 1 + 2 v_2(p_T) \cos(2\phi) + \dots$$

“elliptic flow”

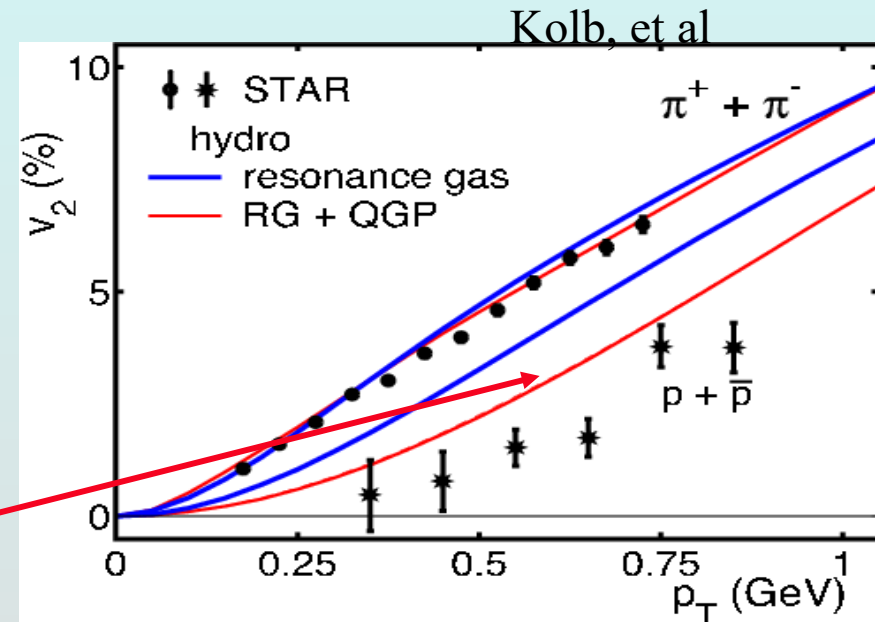


QGP flows hydrodynamically



- ✂ large anisotropy!
- ✂ huge pressure buildup
- ✂ build up quickly, else hydro misses data

Hydrodynamics reproduces
 elliptic flow of q - q and $3q$ states
 Mass dependence was first signal
QGP - NOT gas of hadrons

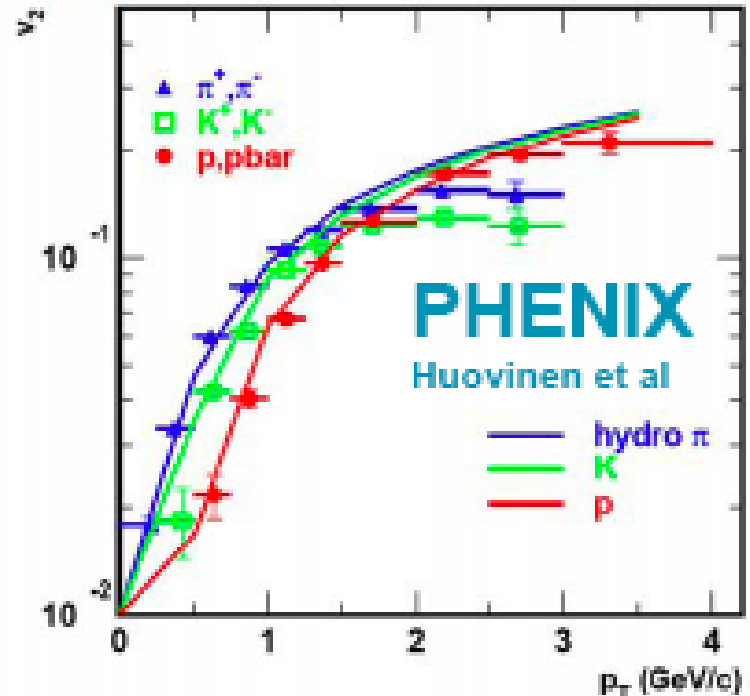
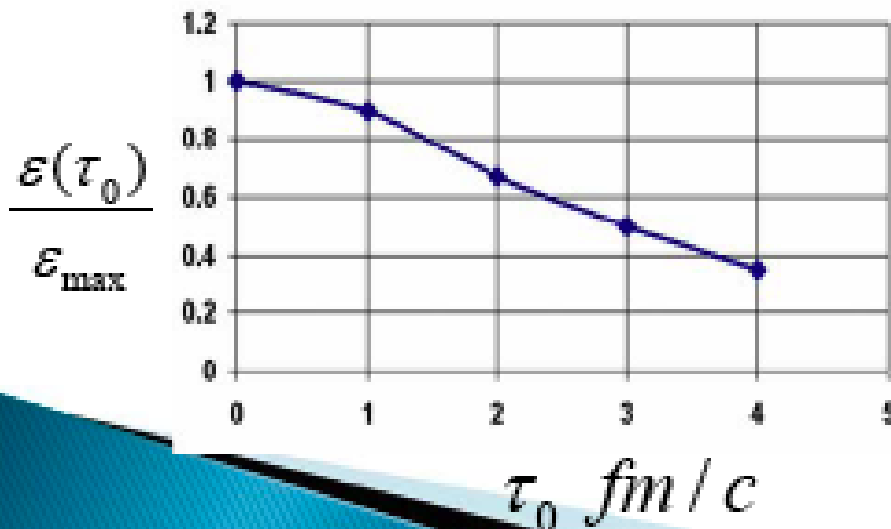


only works with very low viscosity/entropy
 “perfect” liquid (D. Teaney, PRC68, 2003)
 Many advances in relativistic viscous hydrodynamics in 20 years!

Strong flow Implies *early* thermalization

- ▶ If system free streams
 - spatial anisotropy is lost
 - v_2 is not developed

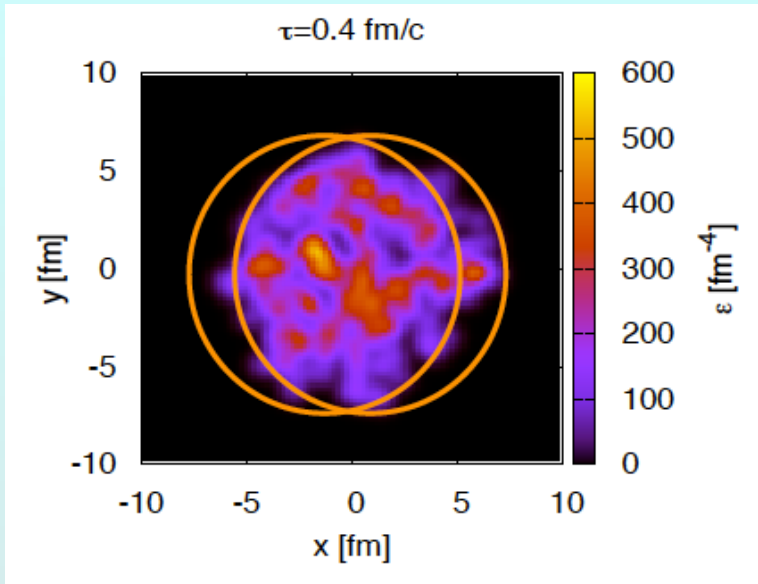
$$\text{Eccentricity: } \varepsilon = \frac{\langle y^2 \rangle - \langle x^2 \rangle}{\langle y^2 \rangle + \langle x^2 \rangle}$$



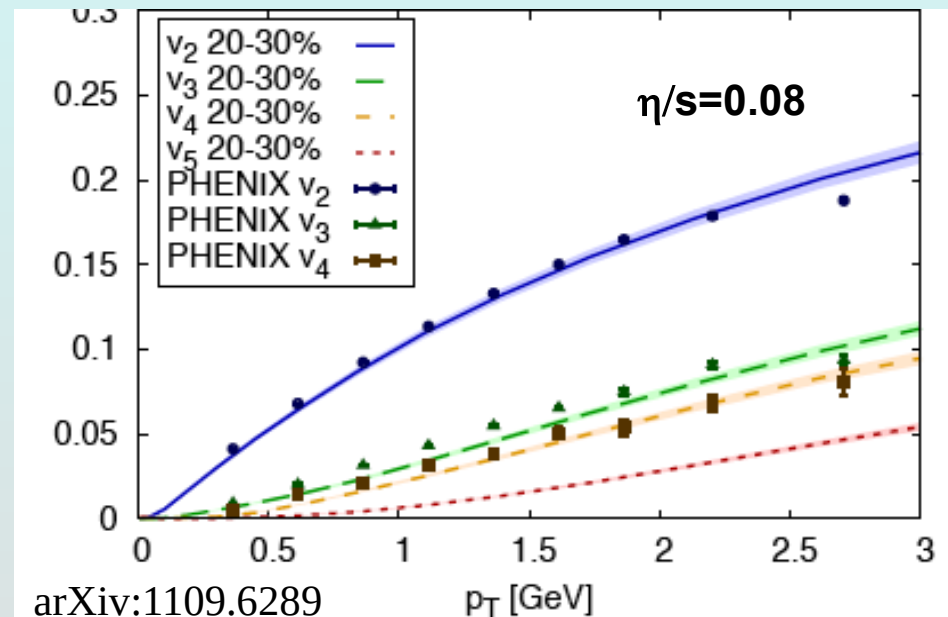
detailed hydro calculations
(QGP+mixed+RG, zero viscosity)

- $\tau_0 \sim 0.6 - 1.0$ fm/c
- $\varepsilon \sim 15 - 25$ GeV/fm³
- (ref: cold matter 0.16 GeV/fm³)

Fluctuations matter!



- Nucleons move around inside the nucleus
- > locations of NN scattering fluctuate
- > apparent symmetry effects yielding only even harmonics not realistic

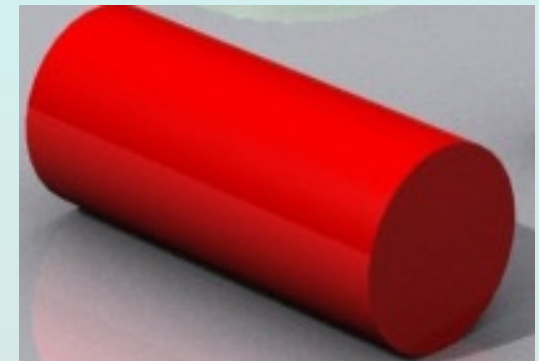
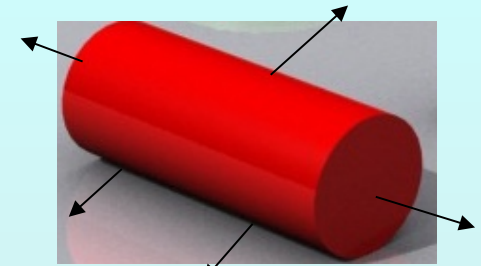
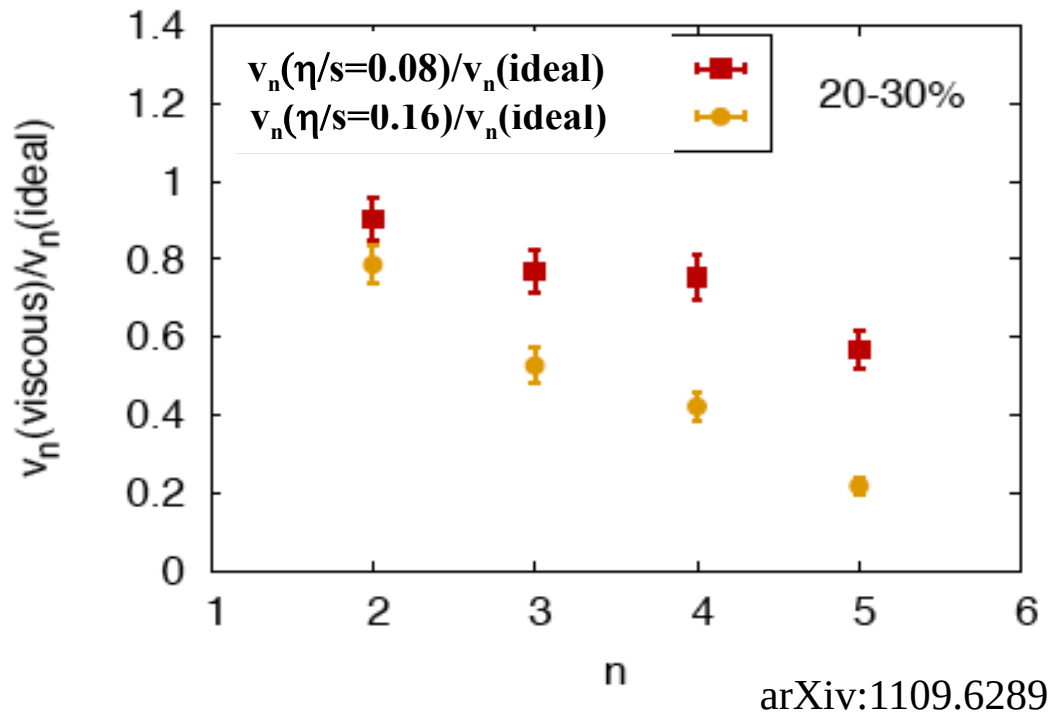


- Reproduce with hydro
- IF include fluctuating initial conditions
- Provides a tool to better pin down the viscosity/entropy ratio

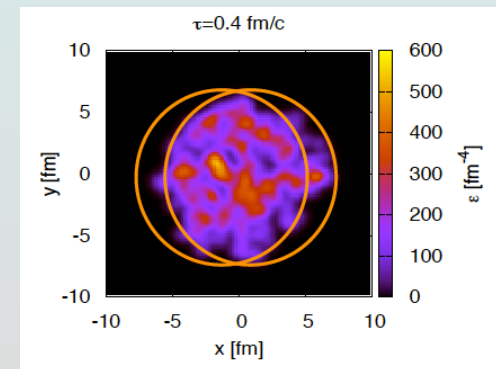
$$dN/d\phi \sim 1 + 2 v_2(p_T) \cos(2\phi) + \dots$$

“elliptic flow”

Higher moments more sensitive to viscosity



- Longitudinal expansion at $v \sim c$
- “freezes in” small shape perturbations
e.g. triangular fluctuations (v_3)
- Viscosity opposes dissipation!

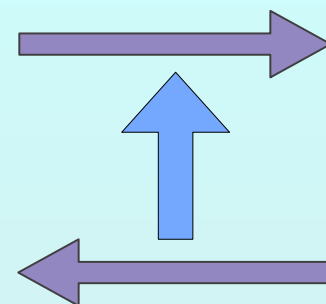


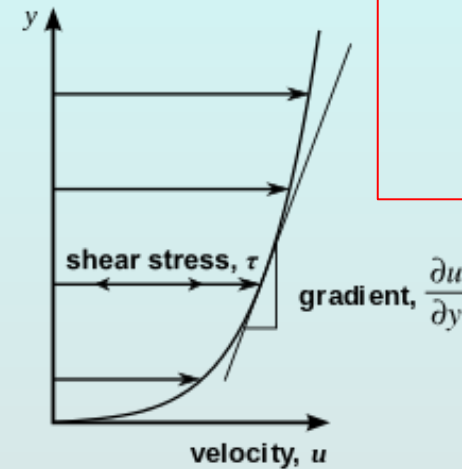
Small viscosity/entropy

Viscosity: *inability* to transport momentum & sustain a wave
internal friction damps waves

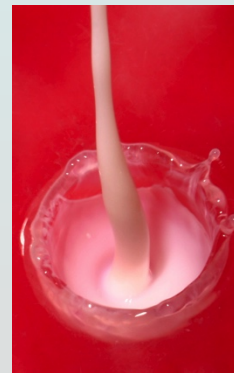
low $\eta \rightarrow$ large σ , transports momentum across fluid
normal QCD: σ not so large
large σ in QGP \rightarrow strongly coupled!
large $\sigma \rightarrow$ many-body interactions when the density is high

So dense QCD matter should also strongly affect quarks & gluons transiting it


$$\eta \approx \frac{1}{3} n \bar{p} \lambda_f$$
$$\bar{p} \sim T$$



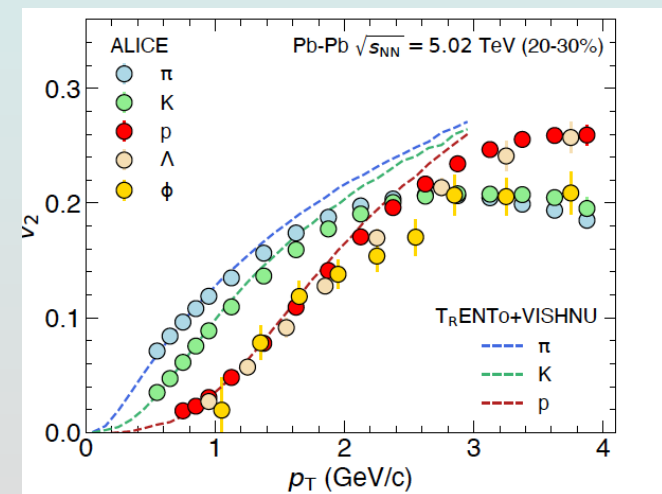
$$= \frac{\tau}{\frac{du}{dy}}$$



Example: milk.
Liquids with higher viscosities will not splash as high when poured at the same velocity.

Perfect liquid? Quantify viscosity of QGP!

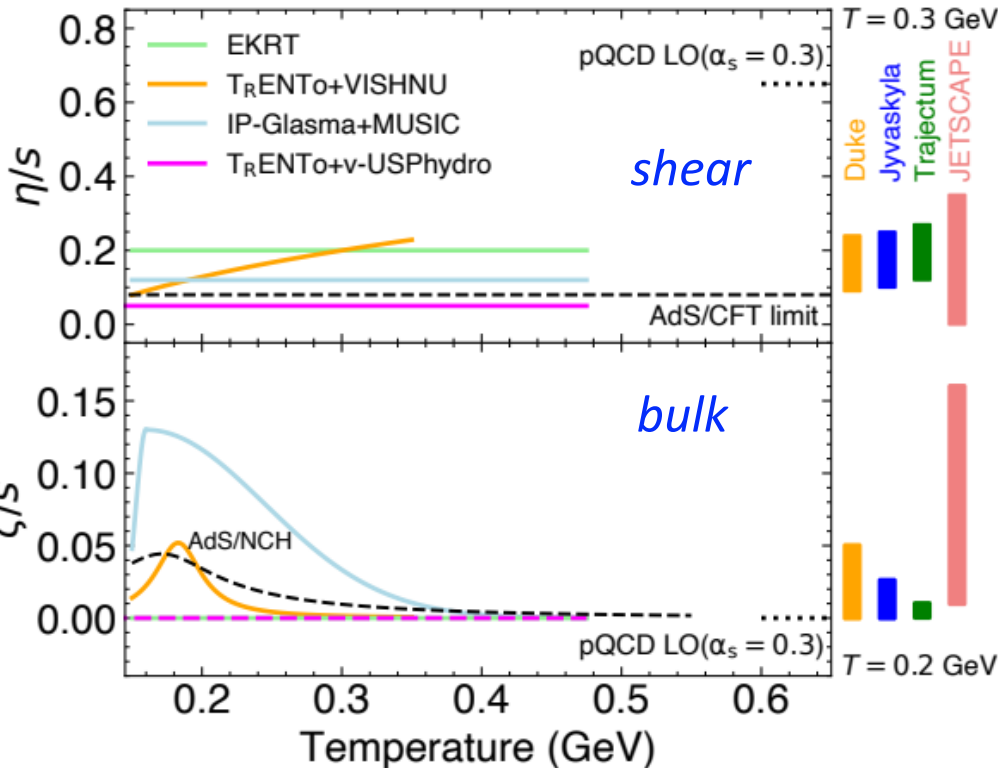
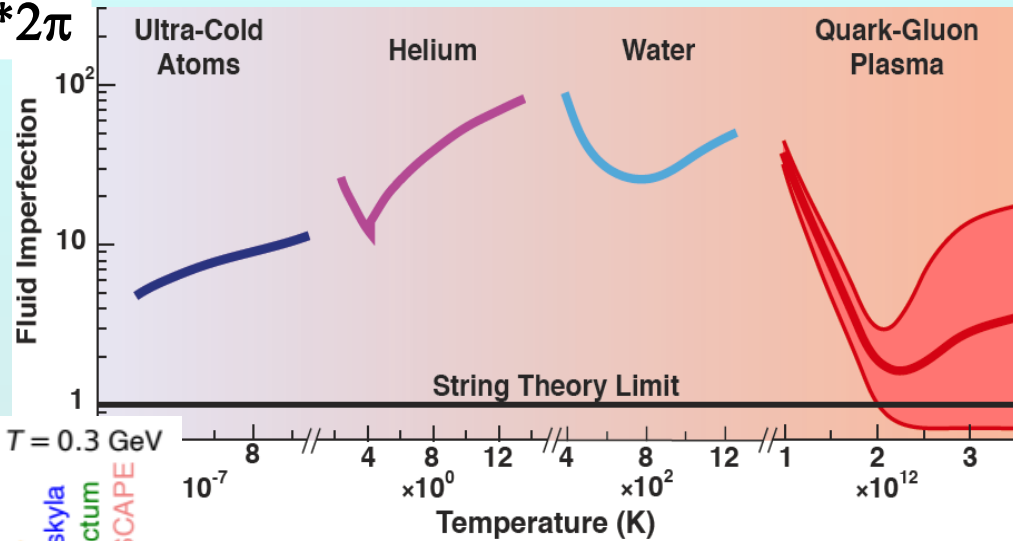
- Use hydro with lattice QCD-based EOS
- Set initial energy density to reproduce observed particle multiplicity
- Use various values of η/s
Quantum mechanical lower bound is $1/4\pi$
determined with help from AdS/CFT
- Constrain with data
Account for hadronic state viscous effects with a
hadron cascade afterburner
Precision data required!
and provided...



QGP property: viscosity per particle

hydrodynamic flow ->
 Nearly perfect liquid
 Strongly coupled QCD
 system!

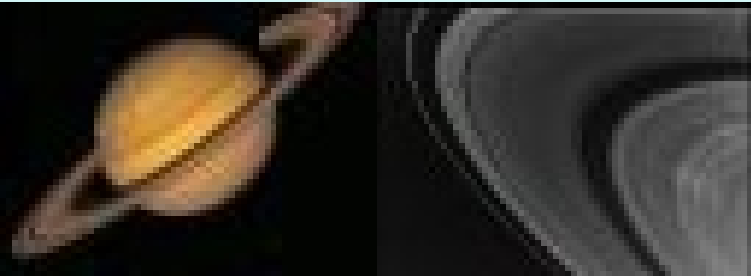
$$\eta/s * 2\pi$$



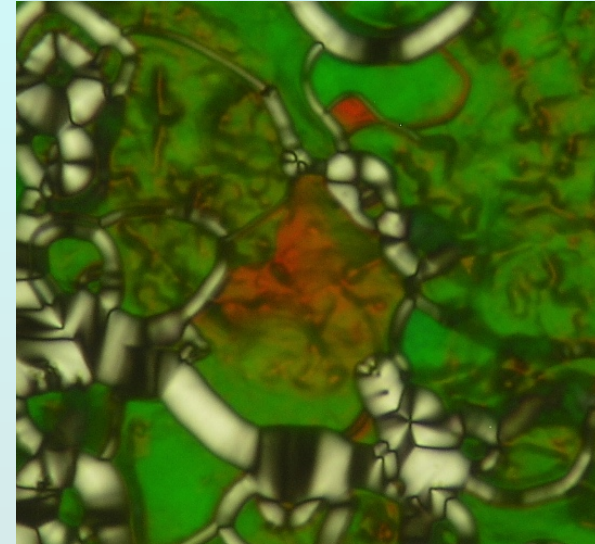
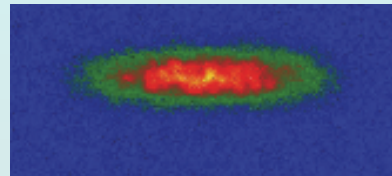
Bulk viscosity = resistance to volume growth

Many types of strongly coupled matter

Quark gluon plasma is like other systems with strong coupling - all flow and exhibit phase transitions



**Cold atoms:
coldest & hottest
matter on earth
are alike!**



Dusty plasmas &

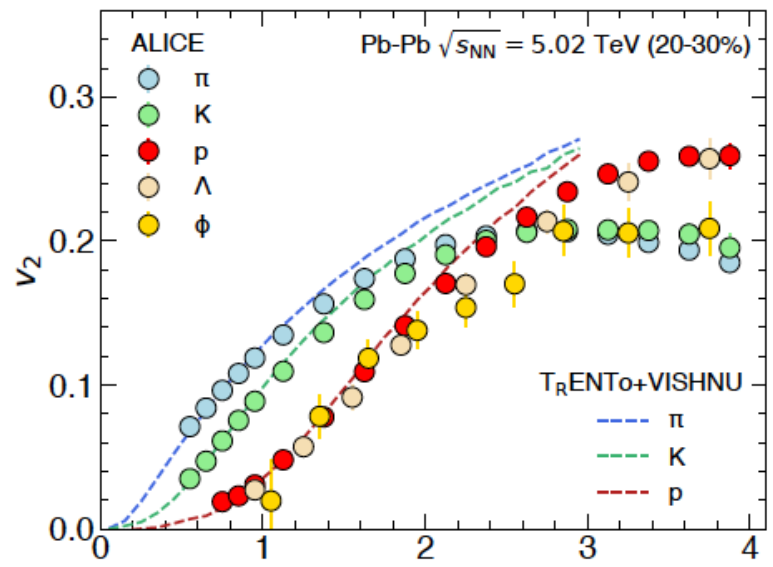
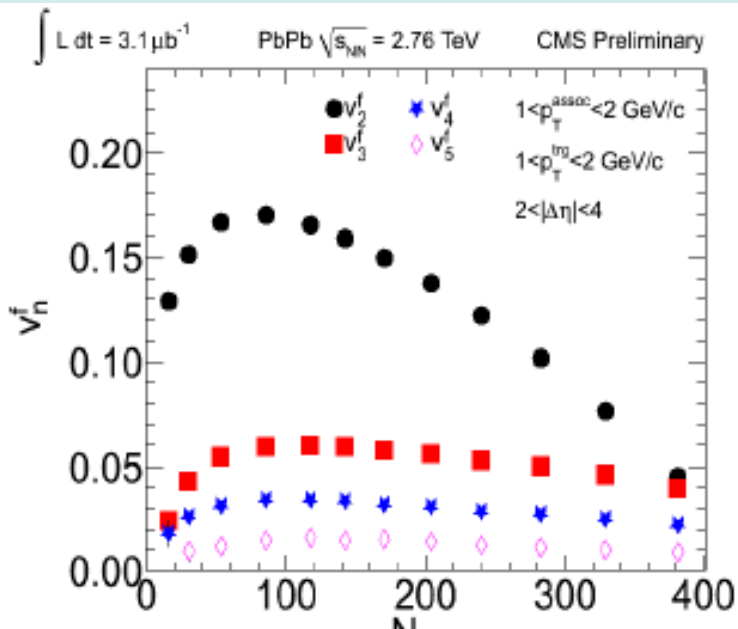
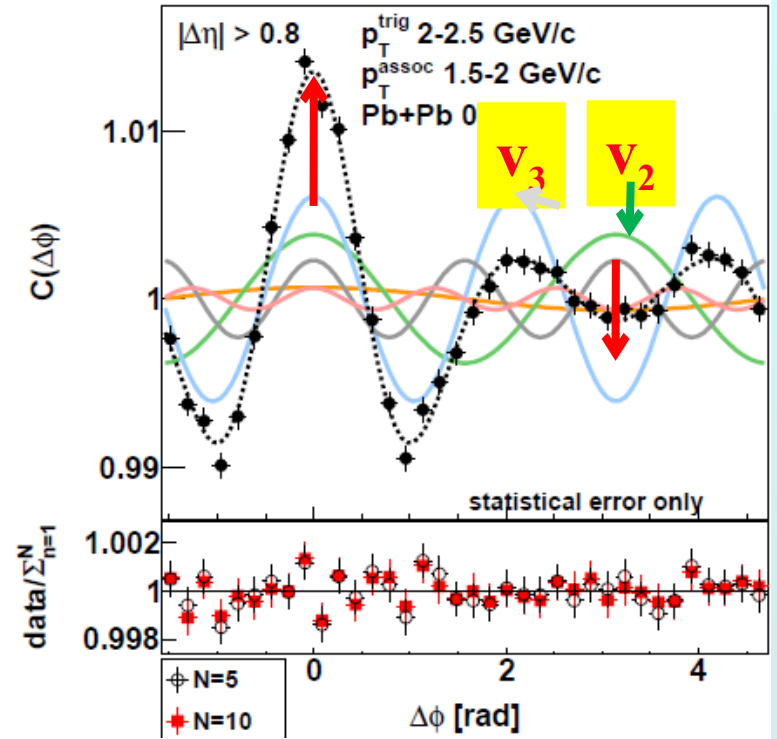
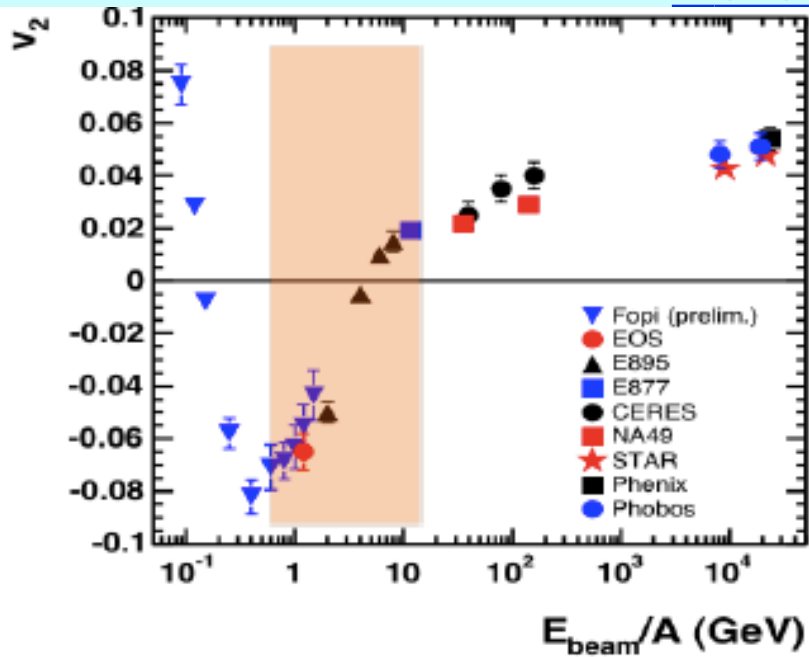
In all these cases have a competition:

Attractive forces \Leftrightarrow repulsive force or kinetic energy

High T_c superconductors: magnetic vs. potential energy

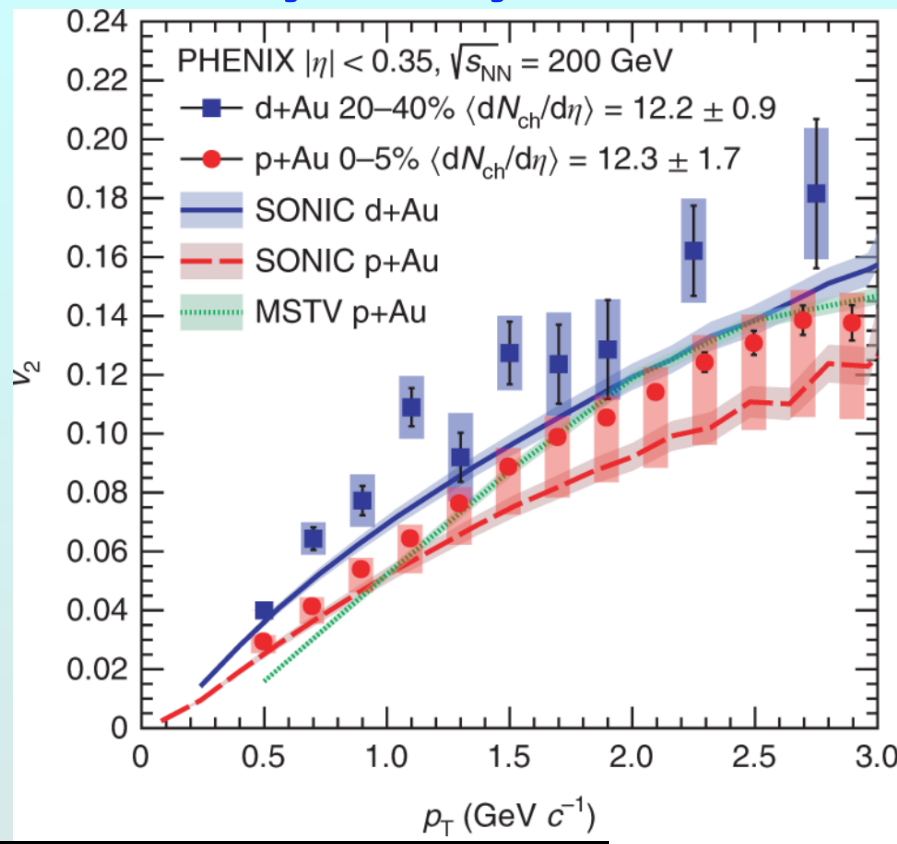
Result: many-body interactions, not pairwise!

v_s dependence



Eek! Hydrodynamics in small systems!

PHENIX
Collaboration
Nature Physics
(2018)

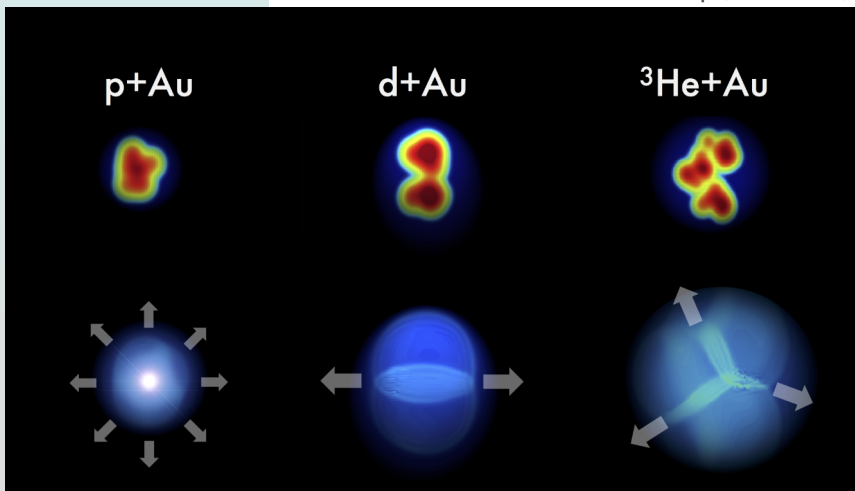


Not big & dense

But, we see collective flow!

Seeded by the initial geometry

A small droplet of QGP?!



Mechanism for fast thermalization?

- **Must be thermalized in $< 1 \text{ fm}/c$!**
Otherwise (viscous) hydro v_2 smaller than in data
- **Can this be achieved with gg, qg, and qq binary scatterings?**
NO!
Making this picture yield sufficient v_2 , requires boosting the pQCD parton-parton cross sections by a factor of ~ 50 !
- **Many-body interactions can do just that!**
- **But, how can hydro set in before equilibrium?**

Hydrodynamic attractor

Numerical solutions of viscous hydro

For conformal fluid

Lines = various initial conditions

Red & green = 1st and 2nd order hydrodynamics

Can consider hydro as a systematic gradient expansion in powers of w

τ = relaxation time

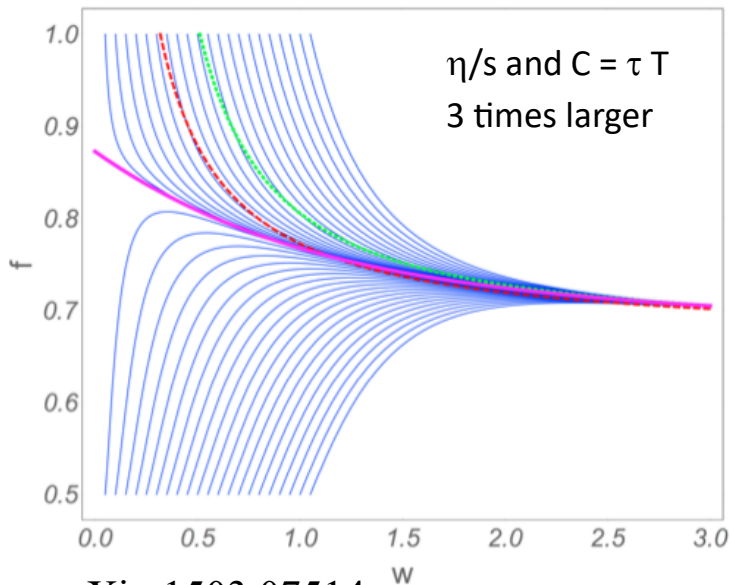
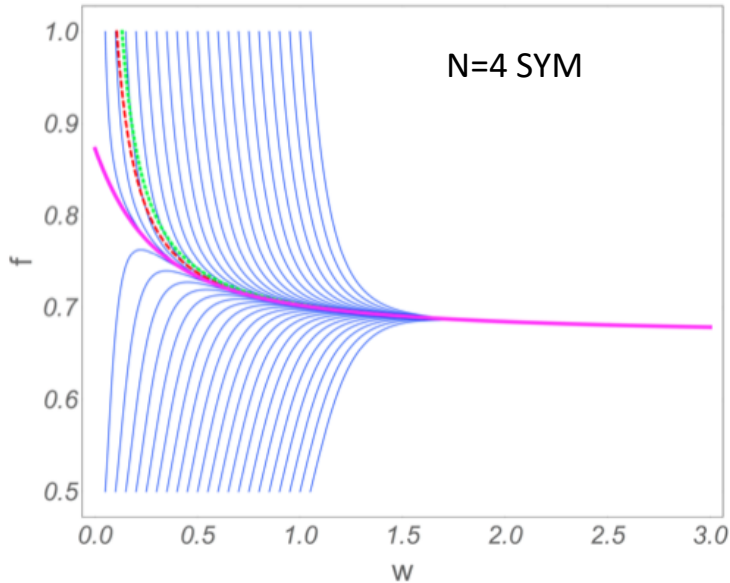
T = temperature

$w = T \times \text{time}$

Since non-hydro modes decay exponentially, system relaxes to an attractor

Driven by fast longitudinal expansion & competition of free streaming vs. dissipation

Attractor is present for larger η/s , but it takes longer for system to reach it

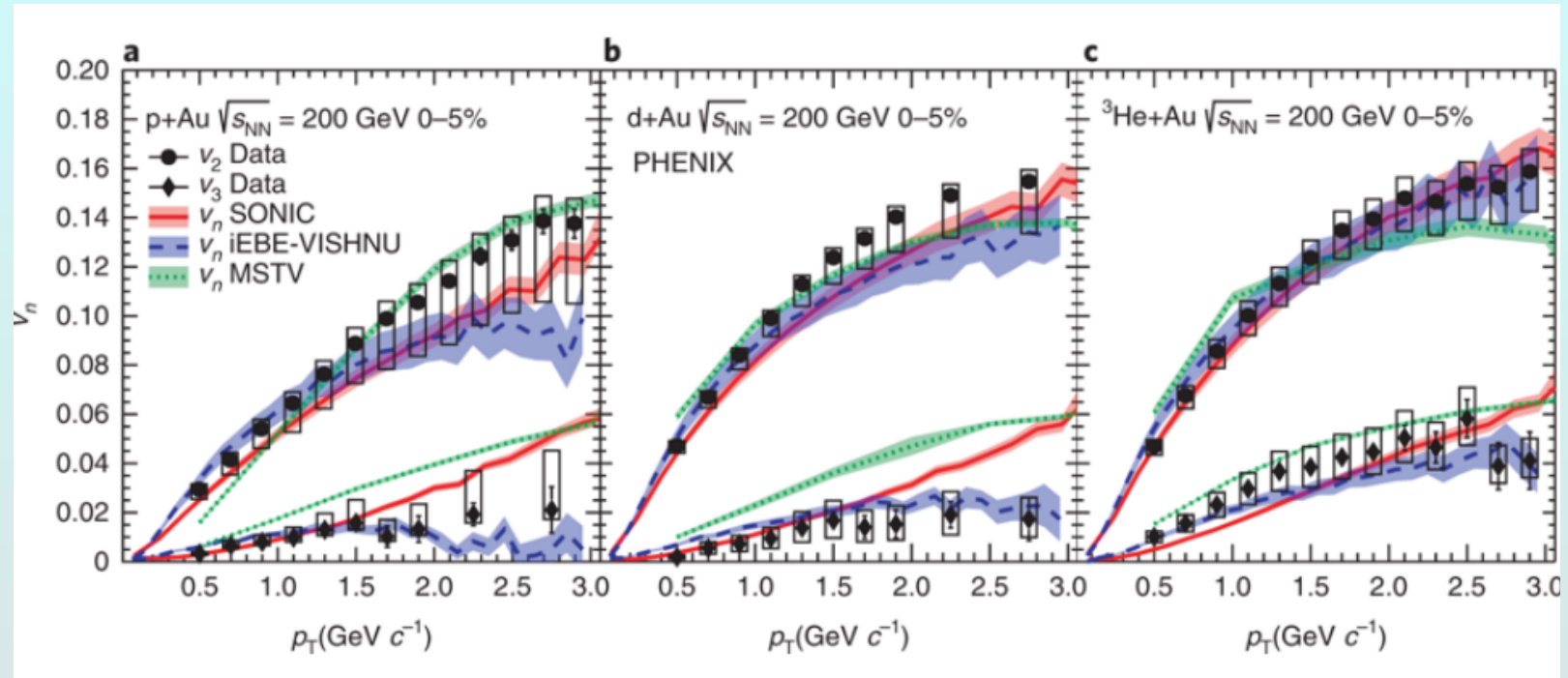


Take aways

- **Make hot QCD matter (quark gluon plasma) in heavy ion collisions**
- **Measure properties via radiated particles and photons**
 - $T_{\max} = 250 - 500 \text{ MeV}$, depending on collision energy
 - Then system expands and cools
- **Dynamics is well described by relativistic viscous hydrodynamics**
 - We measure particle correlations to constrain hydro**
 - Shear viscosity $\eta/s < 0.25$. Most “perfect” liquid
 - Bulk viscosity non-zero only near plasma phase transition
- **Even small systems show collective flows**
 - Likely due to fast longitudinal expansion driving relaxation to hydrodynamic attractor**
 - non-hydro modes decay exponentially so pre-equilibrium dynamics relax very quickly

- **backup slides**

Hydro for p, d, ^3He + Au



PHENIX Collaboration
Nature Physics (2018)

Lepton pair emission \leftrightarrow EM correlator

e.g. Rapp, Wambach Adv.Nucl.Phys 25 (2000)

Emission rate of dileptons per volume

$$\frac{dR_{ll}}{d^4q} = -\frac{\alpha^2}{3\pi^3} \frac{L(M)}{M^2} \text{Im}\Pi_{em,\mu}^\mu(M, q; T) f^B(q_0, T)$$

$$f^B(q_0, T) = 1/(e^{q_0/T} - 1)$$

$$L(M) = \sqrt{1 - \frac{4m_l^2}{M^2}} \left(1 + \frac{2m_l^2}{M^2}\right)$$

$\gamma^* \rightarrow ee$
decay

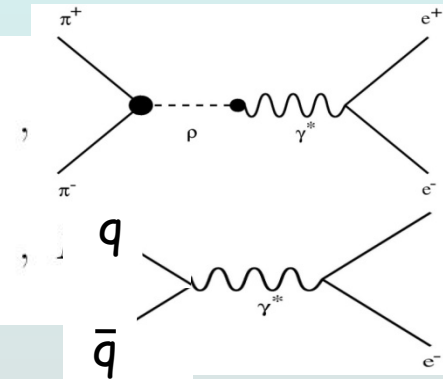
EM correlator
Medium property

Boltzmann factor
temperature

Hadronic contribution
Vector Meson Dominance

Medium modification of meson
Chiral restoration

$$\text{Im}\Pi_{em}^{\text{vac}}(M) = \left\{ \begin{array}{l} \sum_{V=\rho,\omega,\phi} \left(\frac{m_V^2}{g_V}\right)^2 \text{Im}D_V(M) \\ -\frac{M^2}{12\pi} \left(1 + \frac{\alpha_s(M)}{\pi} + \dots\right) N_c \sum_{q=u,d,s} (e_q)^2 \end{array} \right.$$

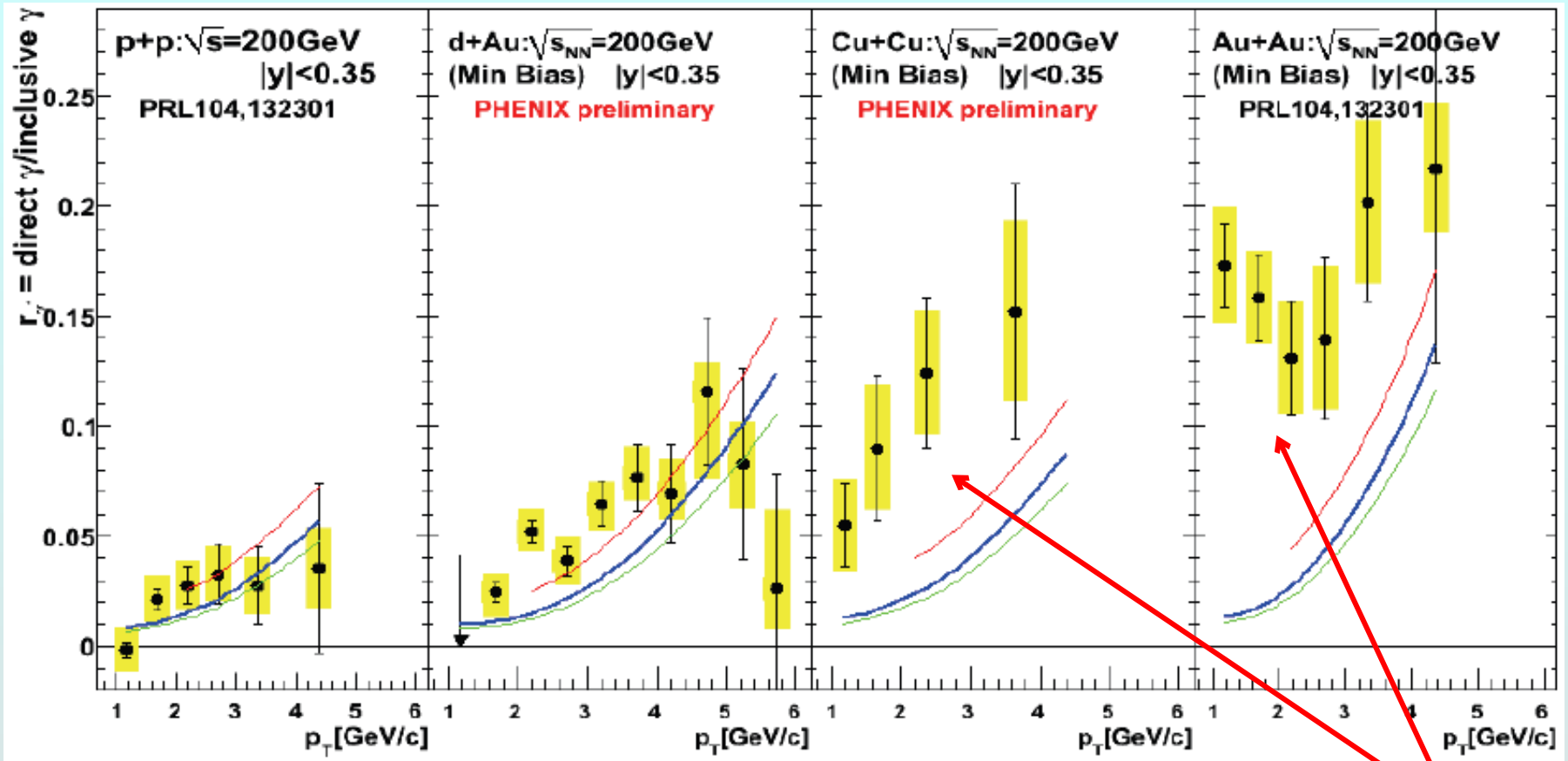


$\bar{q}q$ annihilation

Thermal radiation from
partonic phase (QGP)

From emission rate of dileptons, the medium effect on the EM correlator as well as temperature of the medium can be decoded.

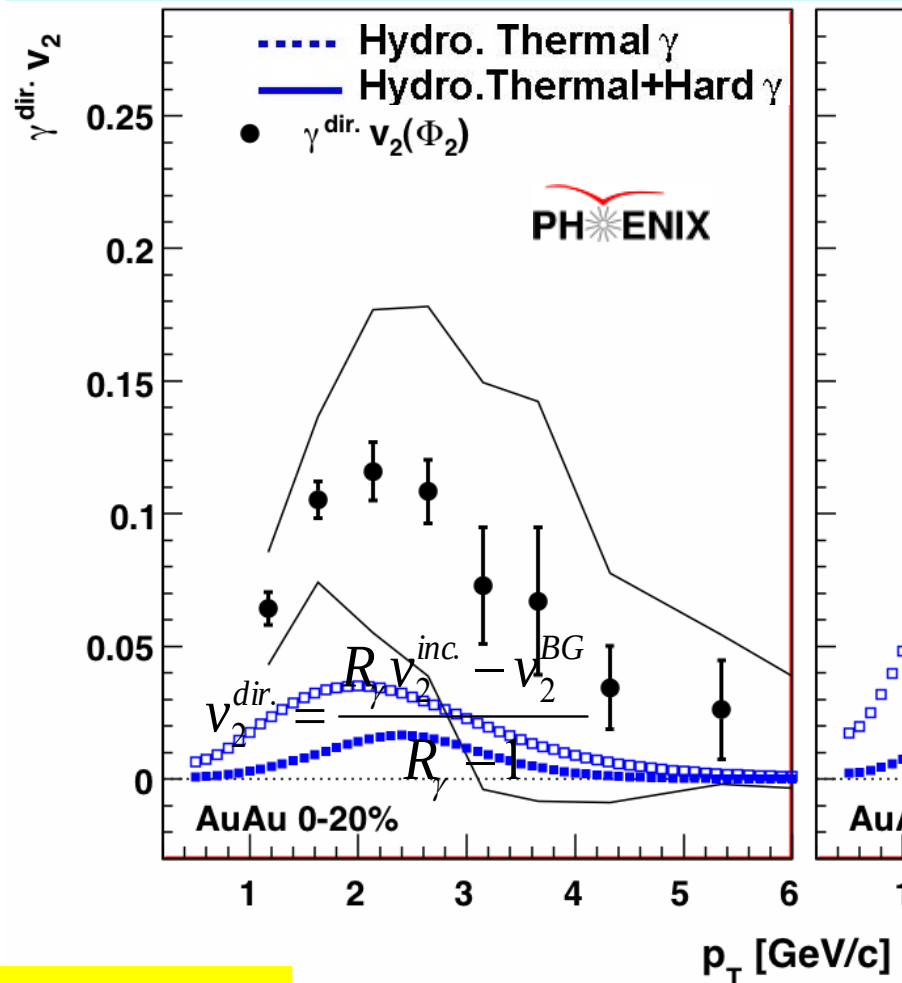
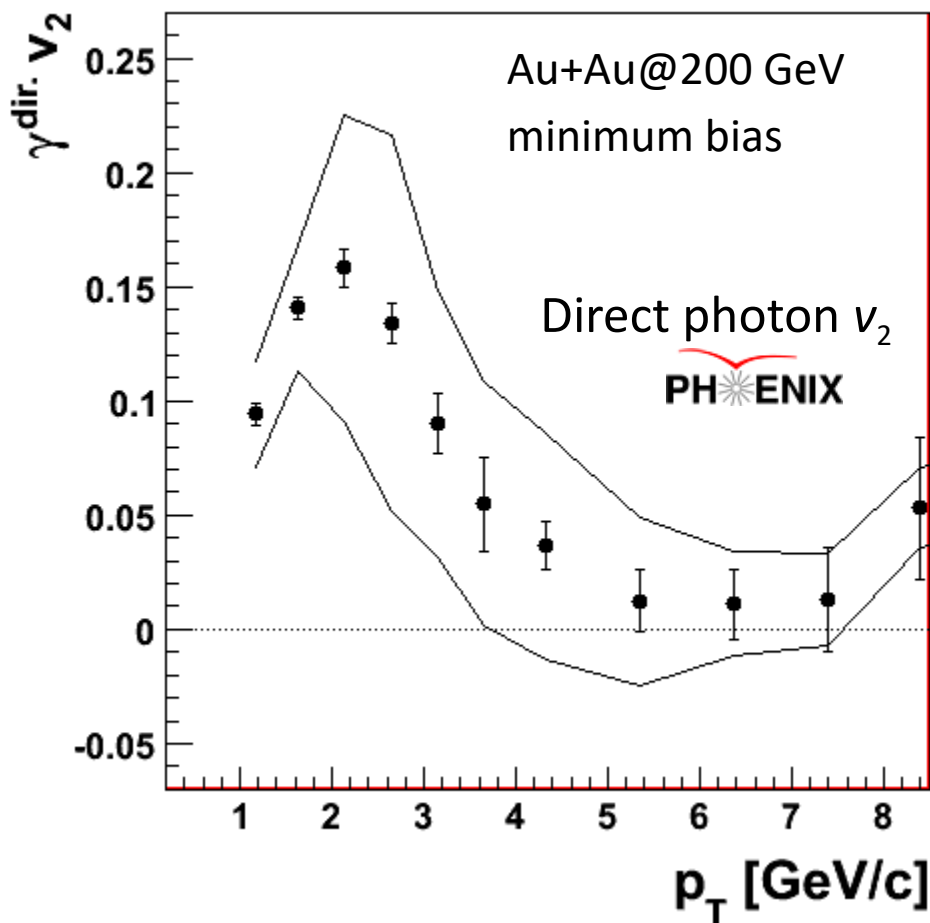
Thermal photons (virtual)



Observe excess photons beyond pQCD in AA collisions. In thermal p_T region

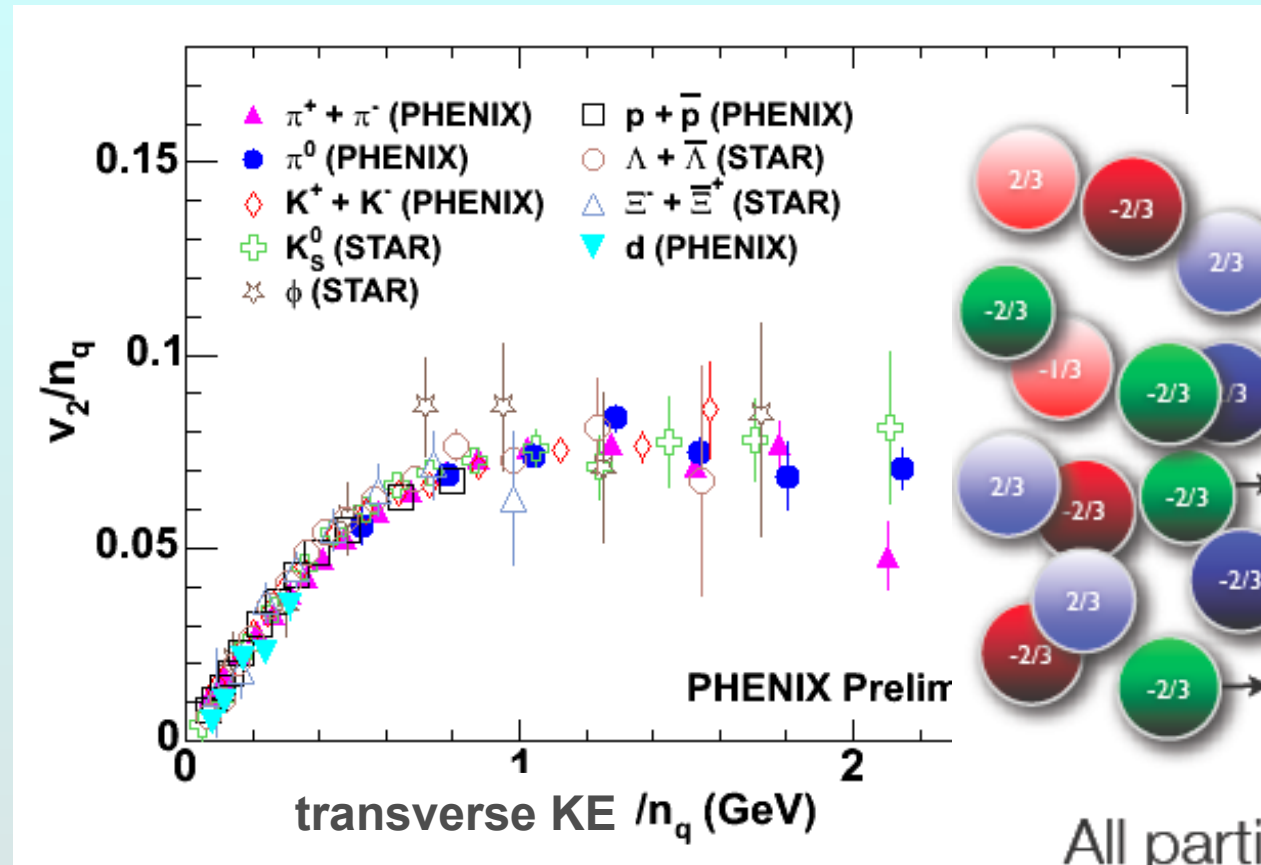
Thermal photons also flow!

arXiv:1105.4126

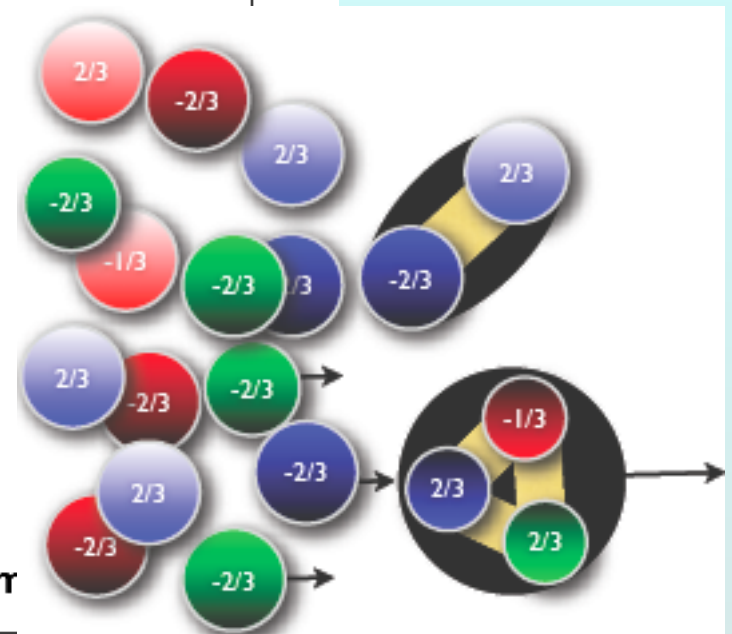


Large flow magnitude is very surprising!

Elliptic flow scales with number of quarks



implication: valence quarks, not hadron pressure builds early, dressed quarks are



All particles flow as if frozen out from a flowing soup of constituent quarks

Calculating transport in QGP

weak coupling limit

perturbative QCD

kinetic theory, cascades

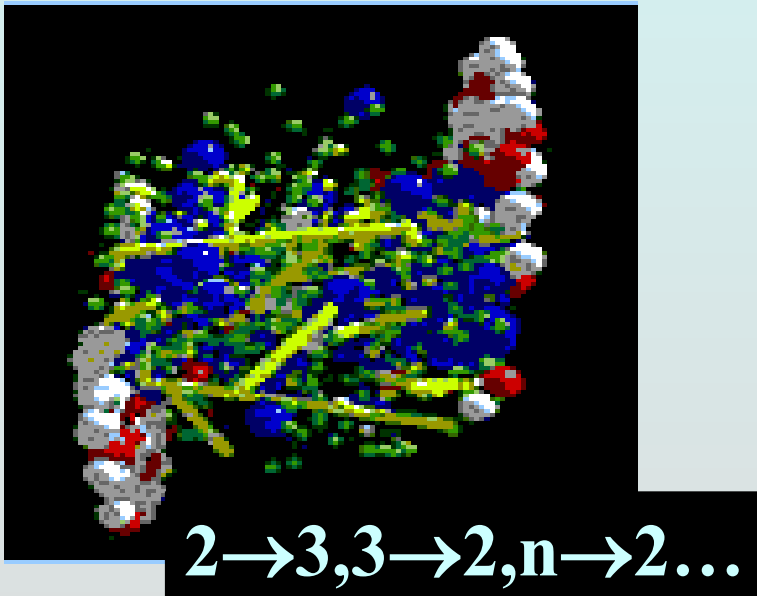
interaction of particles

∞ strong coupling limit

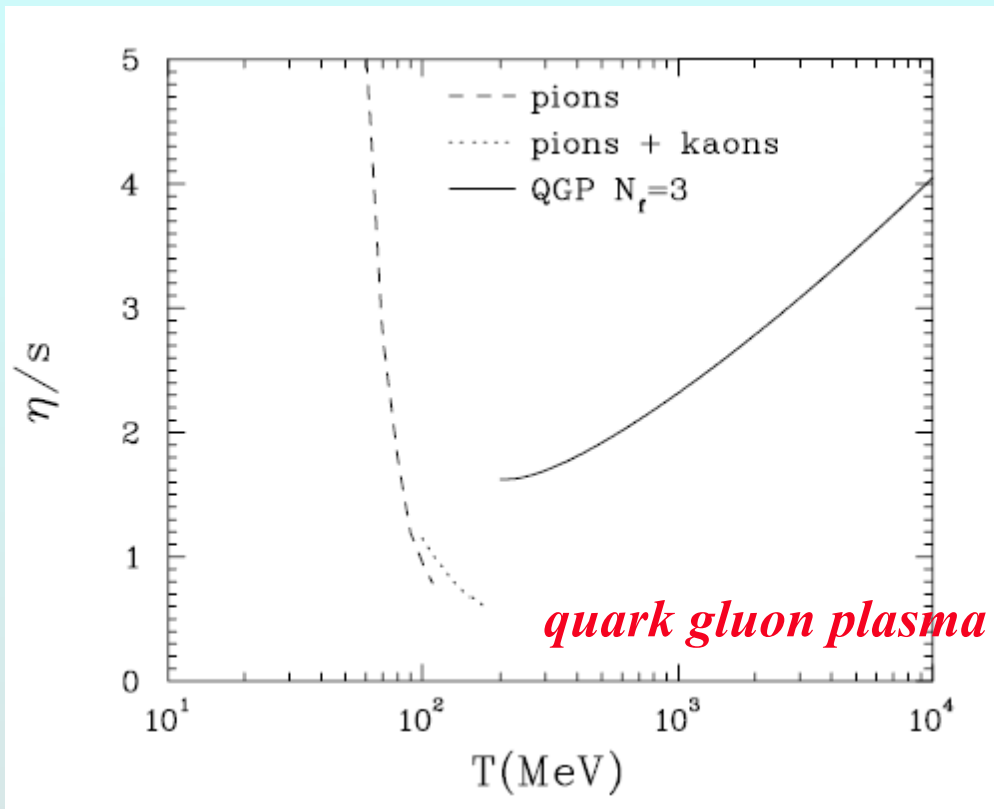
not easy! Try a pure field...

gravity \leftrightarrow supersym 4-d

(AdS/CFT)



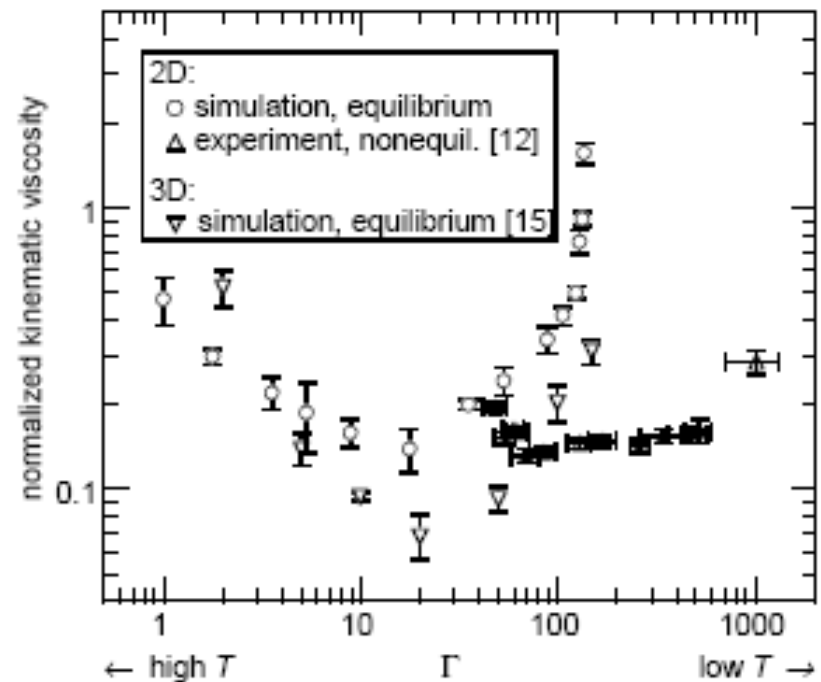
minimum η at phase boundary?



Csernai, Kapusta & McLerran
PRL97, 152303 (2006)

strongly coupled dusty plasma

B. Liu and J. Goree,

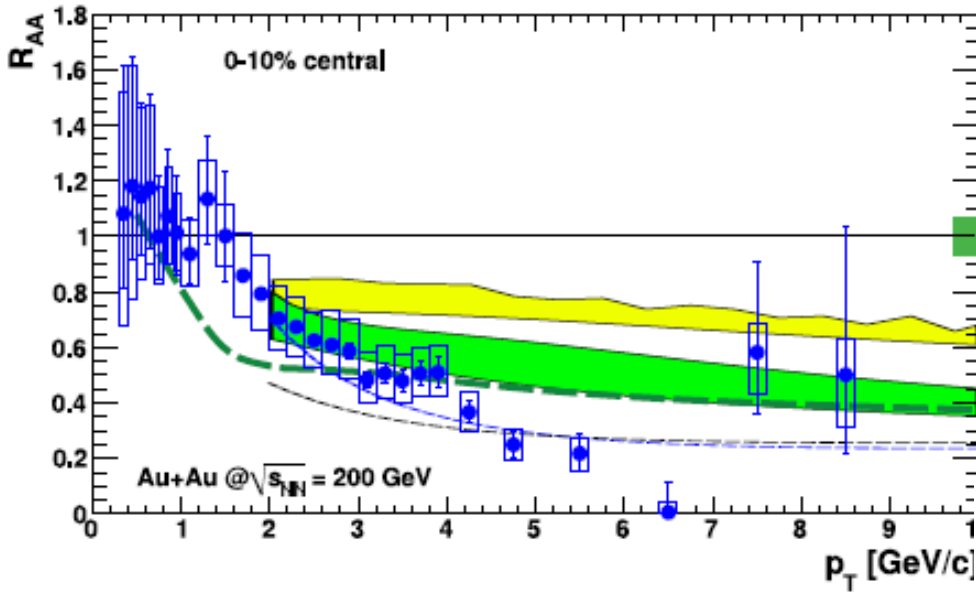


minimum observed in other strongly coupled systems –
kinetic part of η decreases with Γ while potential part increases

heavy quark suppression & flow?

PRL.98: 172301,2007

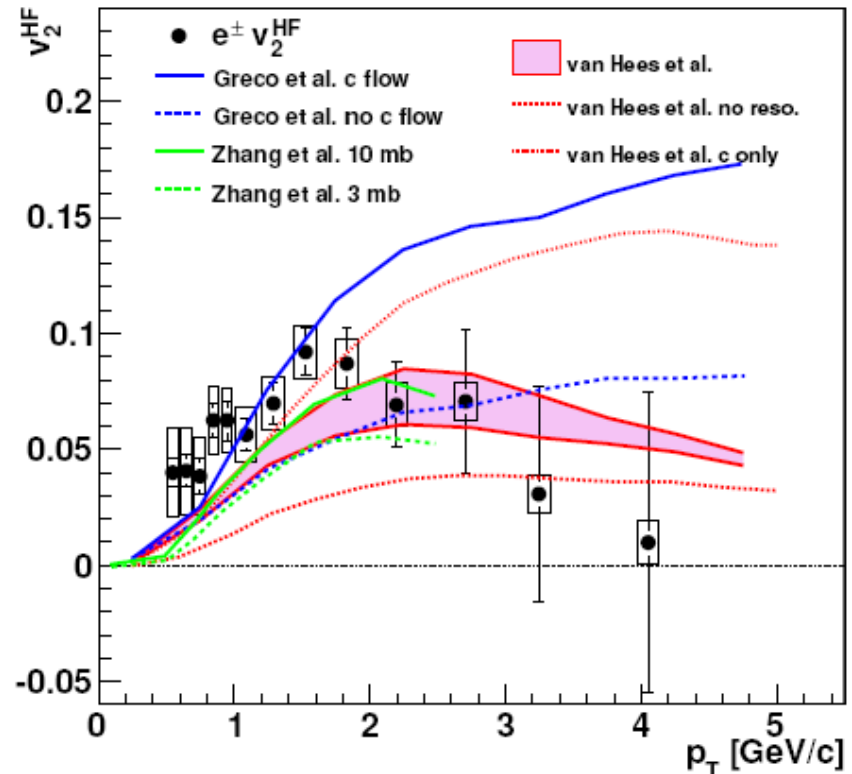
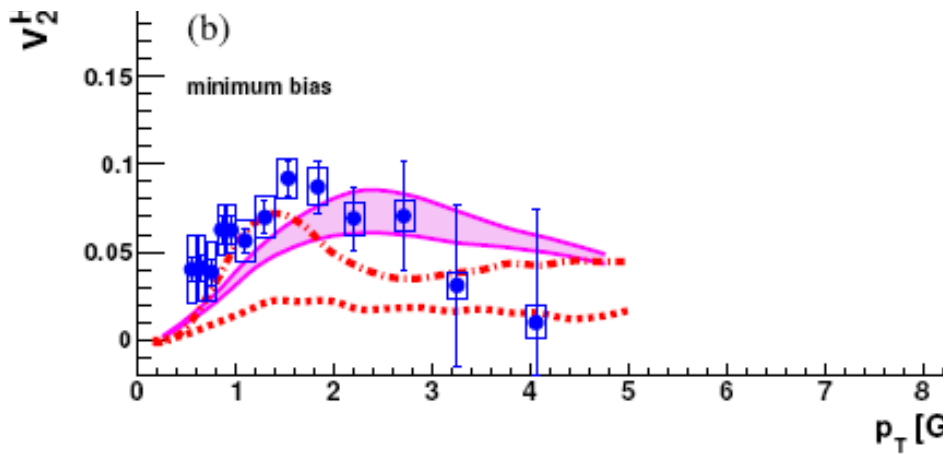
arXiv: 1005.1627



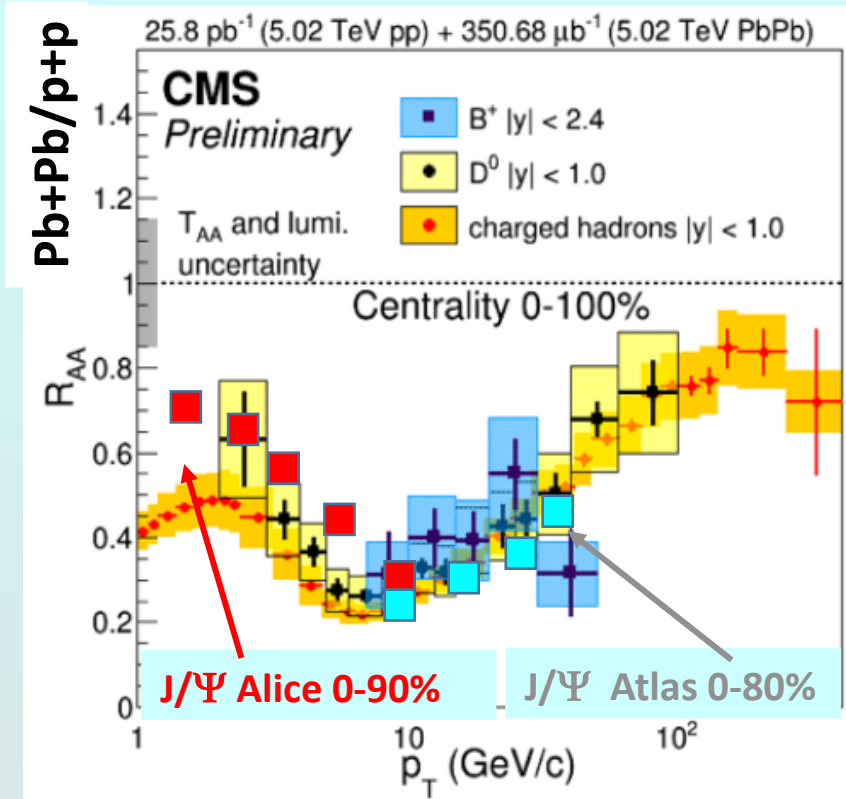
Collisional energy loss?

v_2 decrease with p_T ?

role of b quarks?



Heavy flavor R_{AA}

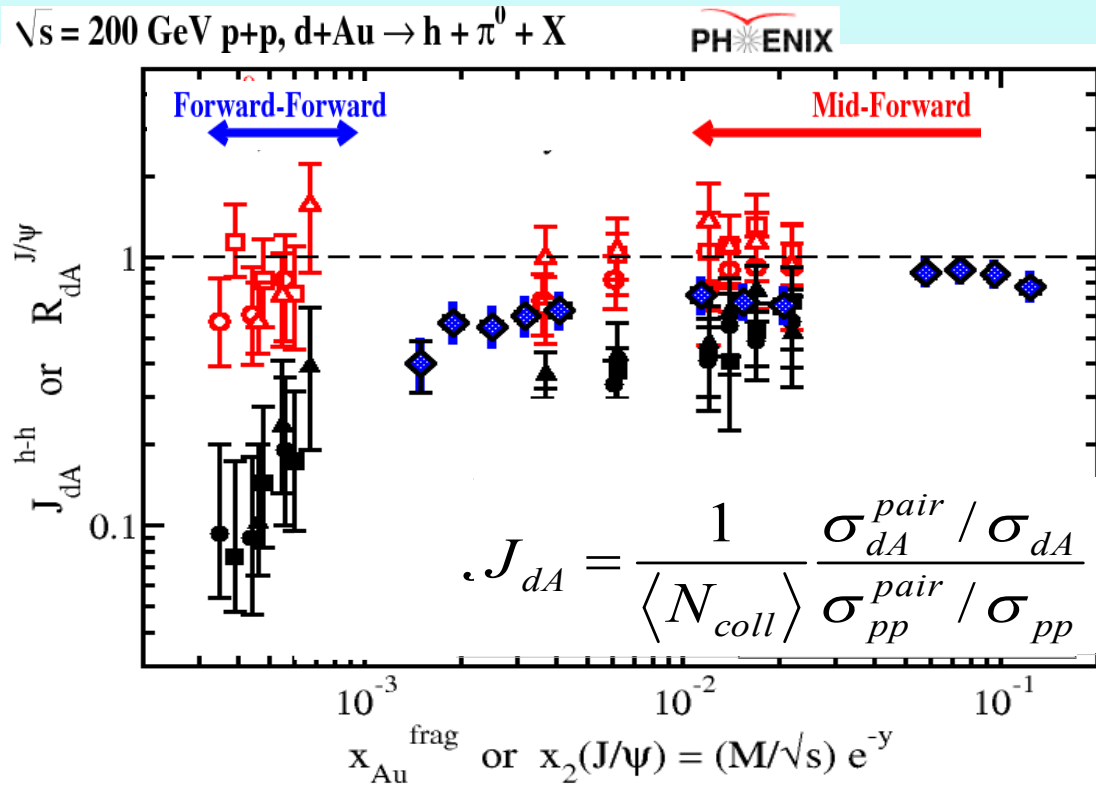
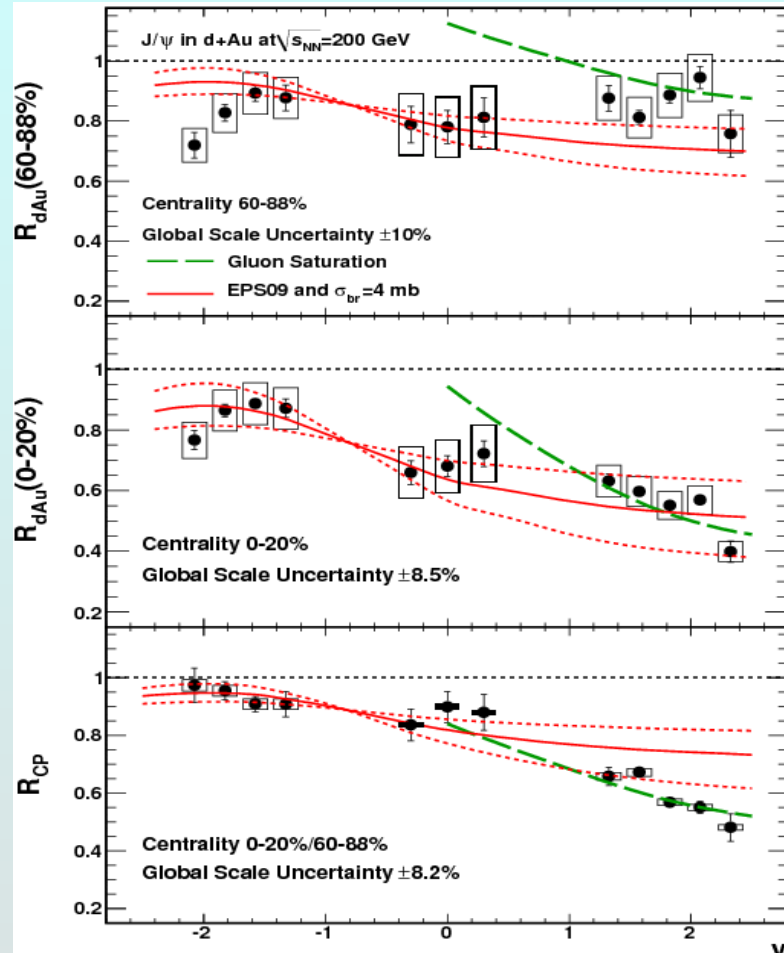


Dense gluonic matter (d+Au, forward γ):

large effects observed

arXiv:1010.1246

arXiv:1105.5112



Di-hadron suppression at low x
pocket formula (for $2 \rightarrow 2$):

$$x_{Au}^{frag} = \frac{\langle p_{T1} \rangle e^{-\langle \eta_1 \rangle} + \langle p_{T2} \rangle e^{-\langle \eta_2 \rangle}}{\sqrt{s}}$$

trend as, e.g. in CGC ...

Shadowing/absorption stronger than linear w/nuclear thickness

Viscosity/Entropy (natural units)

