HUNTING AXIONS WITH METAMATERIALS THE ALPHA HALOSCOPE

A. Gallo Rosso February 16, 2023

OVERVIEW

- 1. Axions
- 2. ALPHA Haloscope
- 3. Statistical inference

1. AXIONS What, how and why

arXiv:1801.08127, 2003.01100, 2012.05029 2104.07634, 2105.01406, 2109.07376

STRONG CP PROBLEM

Violating CP termNeutron electric dipole moment $\mathcal{L}_{\mathcal{LP}} = \theta \frac{\alpha_s}{8\pi} G^a_{\mu\nu} \widetilde{G}^{\mu\nu}_a$ $d_n = (2.4 \pm 1.0) \cdot 10^{-3} \, \mathrm{e\,fm} \times \theta$

The vacuum of the theory is determined by angle $\theta \in [0, 2\pi]$

STRONG CP PROBLEM

$$d_{
m n} < 1.8 \cdot 10^{-13} \, {
m e} \, {
m fm}$$

 $| heta| < 8 \cdot 10^{-11}$

¹C. Abel *et al.*, Phys. Rev. Lett. 124, 081803.

THE AXION

A clean solution

Dynamic variable $\theta(\vec{x}, t)$ Pseudo Nambu-Goldstone boson Peccei Quinn U(1) symmetry Spontaneously broken at f_A





²See e.g. arXiv: 1801.08127, 2003.01100, 2012.05029, 2104.07634, 2105.01406, 2109.07376.



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1 $E \gg f_A$

2 PQ transition

3 QCD transition



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1 $E \gg f_A$ 2 PQ transition 3 QCD transition θ_i ϕ_i ϕ

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A FREE MEAL

Strong CP problem Evolution from θ_i to CP-conserving value θ_{\min}

Dark matter

Non-thermal production from mismatch $\theta_i \neq \theta_{\min}$

Inflation

(Possible) explanation of dynamics of inflation field

²See e.g. arXiv: 1801.08127, 2003.01100, 2012.05029, 2104.07634, 2105.01406, 2109.07376.

PQ-AXION[©]

Models

- PQWW: $f_A \sim \text{EW-scale}$
- KSVZ: hadronic axions
- DFSZ: GUT axions

Massive boson

- $m_A \propto f_A$
- Light (< MeV)
- Non-relativistic

Coupling

- $g \propto f_A^{-1}$
- Hadrons & gluons
- Photons
- Electrons (DFSZ only)

²See e.g. arXiv: 1801.08127, 2003.01100, 2012.05029, 2104.07634, 2105.01406, 2109.07376.

neV μ eV meV eV keV MeV Axion mass $m_a c^2$

















Coherent oscillations Macroscopic wave behavior

$$\mathcal{L}_{a\gamma\gamma}=rac{g_{a\gamma}}{4}a\,F_{\mu
u} ilde{F}^{\mu
u}$$

$$\begin{aligned} \epsilon \nabla \cdot \mathbf{E} &= \rho - g_{a\gamma} \mathbf{B}_{e} \cdot \nabla a \\ \nabla \times \mathbf{H} - \dot{\mathbf{E}} &= \mathbf{j} + g_{a\gamma} (\mathbf{B}_{e} \dot{a} - \mathbf{E} \times \nabla a) \\ \ddot{a} - \nabla^{2} a + m_{a}^{2} a &= g_{a\gamma} \mathbf{E} \cdot \mathbf{B}_{e} \end{aligned}$$

Signature

EM radiation in vacuum in the presence of a magnetic field

$$\epsilon \nabla \cdot \mathbf{E} = \rho - g_{a\gamma} \mathbf{B}_{e} \cdot \nabla a$$
$$\nabla \times \mathbf{H} - \dot{\mathbf{E}} = \mathbf{j} + g_{a\gamma} (\mathbf{B}_{e} \dot{a} - \mathbf{E} \times \nabla a)$$
$$\ddot{a} - \nabla^{2} a + m_{a}^{2} a = g_{a\gamma} \mathbf{E} \cdot \mathbf{B}_{e}$$

$$|E| = \frac{g_{a\gamma}B_{e}a_{0}}{\epsilon}e^{-im_{a}t}$$
$$|E| = 1.3 \cdot 10^{-12} |C_{a\gamma}| \left(\frac{B_{e}}{10 \text{ T}}\right) \left(\frac{0.3 \text{ GeV cm}^{-3}}{\rho_{a}}\right)^{1/2} \frac{\text{V}}{\text{m}}$$

$$|\mathbf{E}| = g_{a\gamma}B_{e}a_{0}\left(1 - \frac{\omega_{p}}{\omega_{a}^{2} - i\omega_{a}^{2}\Gamma_{p}}\right)^{-1}$$

³P. Sikivie, Phys. Rev. Lett. 51, 1415 (1983).







CURRENT LIMITS



⁴O' Hare, cajohare/AxionLimits:AxionLimits.

MATCHING WAVELENGTHS

Desiderata

- Cryogenic temperature
- Tunability
- Large volume
- "Low" plasma frequency



WIRE METAMATERIALS

Metamaterials

Composite materials with different properties than their single parts



⁵P.A. Belov *et al.*, J. Electromagn. Waves. Appl. 16, 8 (2002).

WIRE METAMATERIALS



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WIRE METAMATERIALS

TM₁₁₀ mode structure

Behavior as an effective medium Field distortion near wires might complicate readout



⁶A. Millar *et al.*, arXiv:2210:00017.
WIRE METAMATERIALS

PROs

- Cryogenic
- Solenoidal magnet
- Much larger volume than cavities

CONs

• Behind on R&D

⁶A. Millar *et al.,* arXiv:2210:00017.



2. ALPHA

Axion Longitudinal Plasma HAloscope

A. Millar et al., arXiv:2210.00017, accepted by PRD

AXION LONGITUDINAL PLASMA HALOSCOPE

ALPHA CONSORTIUM

Fermilab IIT Chicago IIT Kanpur **ITMO University** MIT Cambridge Niels Bohr Institute **Oak Ridge National Labs** Stockholm University & OKC UC Berkeley UC Davis UCL London University of Maryland University of Oxford Uppsala University

POWER IN THE DETECTOR

Quality factor *Q* System dampening

$$Q = \frac{\omega \mathcal{U}}{P_{\text{loss}}} = \frac{\omega}{c \, \Gamma_p}$$

$$P_s = \left(\frac{\rho_a}{m_a}g_{a\gamma}^2\right)\kappa B_e^2 Q \mathcal{G} V$$

Geometry factor *G* Normalization of stored energy

$$\mathcal{G} = \frac{1}{a_0^2 g_{a\gamma}^2 B_e^2 V Q^2} \int |\mathbf{E}|^2 \,\mathrm{d}V$$

Volume V

QUALITY FACTOR

- Tightly bound to design

 → (material, shape, volume...)
- Theory ✓ Simulation ✓ Experiment
- $Q \sim \text{some 1000s}$ $\hookrightarrow \times 10 \text{ than expected}$
- Wire losses > surface \hookrightarrow Asymptotically $P_s \propto V$

SINGLE CAVITY MODE BREAKS DOWN



⁷R. Balafendie *et al.*, PRB 106, 075106 (2022).

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GEOMETRY FACTOR

For single-mode systems

$$\mathcal{G} = \mathcal{C} = rac{1}{\overline{B}_e^2 V^2} \left(\int B_e \mathcal{E}_i \, \mathrm{d}V
ight)^2$$

For multi-mode systems

MORE THAN ONE ANTENNA NEEDED





TUNING



⁶A. Millar *et al.*, arXiv:2210:00017.

3. STATISTICAL INFERENCE

Searching strategies and sensitivity projections

A. Gallo Rosso et al., arXiv:2210.16095

TUNING IN





TUNING IN





SCANNING STRATEGY





SCANNING STRATEGY





SCANNING STRATEGY







ONE FREQUENCY, ONE SCAN





GEOMETRIC TEST



GEOMETRIC TEST

LIKELIHOOD BASED



LIKELIHOOD BASED



GEOMETRIC TEST

 $[P(x_{\ell} > x_{\text{thr}} | H_0)]^{k_{\alpha} - 1} \le \alpha$ Rescan probability













$$\begin{aligned} \frac{\alpha}{k_{\alpha}} &= P(D_{\ell} | H_{0}) = \\ P(s_{\ell} \geq c_{\ell}, b_{\ell-1} \leq s_{\ell-1} < c_{\ell-1}, \dots, b_{1} \leq s_{1} < c_{1}) = \\ \int_{c_{\ell}}^{\infty} ds_{\ell} \int_{b_{\ell-1}}^{c_{\ell-1}} ds_{\ell-1} \cdots \int_{b_{1}}^{c_{1}} ds_{1} \mathcal{N}_{\ell}(s | \mathbf{0}, \mathbf{\Sigma}_{s}) \\ \hline P\left(C_{\ell}^{c} \cap B_{\ell} \Big| \bigcap_{h=1}^{\ell-1} \{C_{h}^{c} \cap B_{h}\}, H_{0}\right) = \\ P\left(\bigcap_{h=1}^{\ell} \{C_{h}^{c} \cap B_{h}\} \Big| H_{0}\right) \Big/ P\left(\bigcap_{h=1}^{\ell-1} \{C_{h}^{c} \cap B_{h}\} \Big| H_{0}\right) = \\ \frac{\int_{b_{\ell}}^{c_{\ell}} ds_{\ell} \cdots \int_{b_{1}}^{c_{1}} ds_{1} \mathcal{N}_{\ell}(s | \boldsymbol{\mu}_{s}, \mathbf{\Sigma}_{s})}{\int_{b_{\ell-1}}^{c_{\ell-1}} ds_{\ell-1} \cdots \int_{b_{1}}^{c_{1}} ds_{1} \mathcal{N}_{\ell-1}(s | \boldsymbol{\mu}_{s}, \mathbf{\Sigma}_{s})} = q_{\ell} \end{aligned}$$

$$\begin{array}{c} \checkmark \quad c_1 \\ \checkmark \quad b_1 \\ \checkmark \quad c_2 \\ \checkmark \quad b_2 \\ \checkmark \quad \dots \end{array}$$

ONE FREQUENCY, MANY SCANS



⁸A. Gallo Rosso *et al.*, arXiv:2210.16095.



MANY FREQUENCIES, MANY SCANS



$$\max_{t=1,...,T} \left\{ s_L^t - c_L^t \right\} = \mathcal{S}$$

MANY FREQUENCIES, MANY SCANS





$$\max_{t=1,...,T}\left\{s_{L}^{t}-c_{L}^{t}\right\}=\mathcal{S}$$

$$\tilde{c}_L^t = c_L^t + \mathcal{S}_{1-lpha}$$



IT WORKS



⁸A. Gallo Rosso *et al.*, arXiv:2210.16095.



LOOK ELSEWHERE CORRECTIONS



$$P(D_{t\ell}|H_0) = \frac{\alpha}{R_T^{\alpha}k_{\alpha}}$$



⁸A. Gallo Rosso *et al.*, arXiv:2210.16095.

A SIMPLE EXAMPLE



⁸A. Gallo Rosso *et al.*, arXiv:2210.16095.

A SIMPLE EXAMPLE



⁸A. Gallo Rosso *et al.*, arXiv:2210.16095.
DISCOVERY REACH

ALPHA PHASE I

- 2 years run
- $(5 \div 40)$ GHz
- HEMT amplifiers
- Single scan mode

ALPHA PHASE II

- 2 years run
- (5÷45) GHz
- Quantum noise
- Single scan mode



⁷A. Millar *et al.*, arXiv:2210:00017.

CONCLUSIONS

ALPHA

- Overall design validation
 - Theory/simulation/experiment
 - ✓ $Q \times 10$ than previously expected
 - ✓ Feasible tuning
- KSVZ and DFSZ at reach

DAQ & ANALYSIS

- Framework for inference on sequential tests
 - ✓ One frequency: Closed form for power
 - ✓ Multiple frequencies: Monte Carlo for LEE
- Protocol and computational optimizations

IMPROVE SYNERGY FOR SIMULATION & ANALYSIS

CONCLUSIONS



IMPROVE SYNERGY FOR SIMULATION & ANALYSIS