

Holographic duality from Howe duality: Chern-Simons gravity as an ensemble of code CFTs

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Error-correcting codes

Narain CFT

Holography

Chern–Simons TQFT

1-form symmetry

Modular forms

Quantum information

Representation theory

Finite fields

Howe duality

Based on [\[2504.08724\]](#) with Brian McPeak and Anatoly Dymarsky

Recently, the factorisation puzzle in holography has motivated the idea of **ensemble averaging** [Saad, Shenker, Stanford 2019]. It is important to exhibit explicit instances of this proposal, and understand better why they work

Narain ensemble

Code CFT ensemble

[Afkhani-Jeddi et al 2020; Maloney, Witten 2020]

$$\int Z(\mu, \Omega) d\mu = \frac{\mathcal{E}_{\frac{n}{2}, \frac{n}{2}}(\Omega)}{\Phi(\Omega)^n}$$

Continuous Narain moduli space, Zamolodchikov metric

$$\frac{\mathcal{E}_{\frac{n}{2}, \frac{n}{2}}(\Omega)}{\Phi(\Omega)^n} = \sum_{\gamma \in \Gamma \backslash Sp(2g, \mathbb{Z})} \chi_{\text{vac}}^{U(1)}(\gamma\Omega)$$

Equality follows from Siegel–Weil formula

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Narain ensemble	Code CFT ensemble
[Afkhami-Jeddi et al 2020; Maloney, Witten 2020]	
$\int Z(\mu, \Omega) d\mu = \frac{\mathcal{E}_{\frac{n}{2}, \frac{n}{2}}(\Omega)}{\Phi(\Omega)^n}$	$\frac{1}{N} \sum_i Z_{\text{CFT}_i}(\Omega) = \sum_{\gamma \in \Gamma \backslash Sp(2g)} \psi_{\bar{0}}(\gamma\Omega)$
Continuous Narain moduli space, Zamolodchikov metric	Discrete ensemble, equal weights
Infinite sum over modular images $\frac{\mathcal{E}_{\frac{n}{2}, \frac{n}{2}}(\Omega)}{\Phi(\Omega)^n} = \sum_{\gamma \in \Gamma \backslash Sp(2g, \mathbb{Z})} \chi_{\text{vac}}^{U(1)}(\gamma\Omega)$	Finite sum over modular images [JH, McPeak 2022; Aharony, Dymarsky, Shapere 2023]
Equality follows from Siegel–Weil formula	Equality follows from Howe duality [Dymarsky, JH, McPeak 2025]

Code = collection of (binary/ p -ary) vectors

E.g. $\{(00000000), (00001111), (00110011), (00111100),$
 $(01010101), (01011010), (01100110), (01101001),$
 $(10010110), (10011001), (10100101), (10101010),$
 $(11000011), (11001100), (11110000), (11111111)\}$

Studied since the 1950's (Voyager mission,
classification, good codes $\leftrightarrow$ large distance/"gap,"
Siegel modular forms, ...)

Two important aspects:

- Codes can be used to **construct CFTs**
[Dolan, Goddard, Montague 1990]
Code $\rightarrow$ lattice $\rightarrow$ CFT
- Codes come in natural **ensembles** $\Rightarrow$
Averaging over codes [Pless, Sloane 1975]

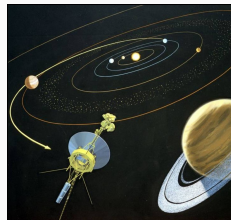


Image: NASA

CLASSIFICATION AND ENUMERATION OF SELF-DUAL CODES 319

4. THE SUM OF ALL WEIGHT ENUMERATORS

From the results of Section 3 it follows easily that (a) the sum of the weight enumerators of all self-dual codes of length n is (for n even)

$$\prod_{j=1}^{(n/2)-2} (2^j + 1) \cdot \left[2^{(n/2)-1}(1 + x^n) + \sum_{\substack{i|n \\ i \neq 1}} \binom{n}{i} x^i \right];$$

and (b) the corresponding sum when the weights are divisible by 4 is (for n divisible by 8)

$$\prod_{i=0}^{(n/2)-3} (2^i + 1) \cdot \left[2^{(n/2)-2}(1 + x^n) + \sum_{\substack{i|n \\ i \neq 1, 2}} \binom{n}{i} x^i \right].$$

Today consider codes over $\mathbb{F}_p$

Definition: $\mathcal{C} = \{v | v \in \mathbb{F}_p^{2n}\}$

Code defines a (Lorentzian) lattice by embedding $\mathbb{F}_p^{2n}$ in $\mathbb{Z}^{2n} \Rightarrow$ Code CFTs are discrete points in Narain moduli space

Enumerator polynomial

$$\Psi_{\mathcal{C}} = \sum_{v \in \mathcal{C}} \psi_v \in \mathcal{H} = \text{span}\{\psi_v | v \in \{0, \dots, p-1\}^{2n}\}, \quad \dim \mathcal{H} = p^{2n}$$

becomes CFT partition function under the substitution

$$\psi_v \mapsto \Psi_v(\tau, \xi, \bar{\xi}) = \frac{1}{|\eta(\tau)|^2} \sum_{n,m} e^{i\pi(\tau p_L^2 - \bar{\tau} p_R^2) + 2\pi i(p_L \xi - p_R \bar{\xi}) + \frac{\pi}{2\tau_2}(\xi^2 + \bar{\xi}^2)}$$

Ψ_v non-holomorphic conformal blocks of abelian CS theory [\[Gukov et al 2004\]](#)

$$p_{L,R} = \sqrt{\frac{p}{2}} \left(\frac{1}{r}(n + \alpha/p) \pm r(m + \beta/p) \right), \quad v = (\alpha, \beta)$$

Number of even self-dual codes grows with n and p

n	# codes ($p = 3$)	# codes ($p = 5$)
1	2	2
2	8	12
3	80	312
4	2240	39312
5	183680	24609312
6	44817920	76928709312
7	32717081600	1202088011709312
8	71584974540800	93914328002801709312
9	469740602936729600	36685378290422420501709312
10	9246374028206585446400	71651166158859580464821501709312

At small n , p , possible to list all even self-dual codes and compute **average** explicitly

Example: $n = 1$, $p = 3$:

$$\Psi_{C_1} = \psi_{00} + \psi_{01} + \psi_{02}, \quad \Psi_{C_2} = \psi_{00} + \psi_{10} + \psi_{20}$$

$$\text{Average: } \overline{\Psi} = \psi_{00} + \frac{1}{2}(\psi_{01} + \psi_{02} + \psi_{10} + \psi_{20})$$

Next: recover this by a “holographic” computation

Averaging over codes

Code average formulas can be rewritten as a **sum of modular images** of $\psi_{\vec{0}}$

[JH, McPeak 2022; Aharony, Dymarsky, Shapere 2023]

Modular invariance [MacWilliams 1962]: Ψ_C invariant under

$$S: \psi_v \mapsto \frac{1}{p^n} \sum_{v'} e^{-\frac{2\pi i}{p} v \eta v'} \psi_{v'} \quad T: \psi_v \mapsto e^{\frac{i\pi}{p} v \eta v} \psi_v \quad \eta = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Example $n = 1, p = 3$. Seed ψ_{00}

$$\begin{aligned} \bar{\Psi} &= \frac{1}{2} (\psi_{00} + S\psi_{00} + TS\psi_{00} + T^2S\psi_{00}) \\ &= \frac{1}{2} (\psi_{00} + \frac{1}{3} [\psi_{00} + \psi_{01} + \psi_{02} + \psi_{10} + \psi_{11} + \psi_{12} + \psi_{20} + \psi_{21} + \psi_{22}] \\ &\quad + \frac{1}{3} [\psi_{00} + \psi_{01} + \psi_{02} + \psi_{10} + e^{2\pi i/3} \psi_{11} + e^{4\pi i/3} \psi_{12} + \psi_{20} + e^{4\pi i/3} \psi_{21} + e^{2\pi i/3} \psi_{22}] \\ &\quad + \frac{1}{3} [\psi_{00} + \psi_{01} + \psi_{02} + \psi_{10} + e^{4\pi i/3} \psi_{11} + e^{2\pi i/3} \psi_{12} + \psi_{20} + e^{2\pi i/3} \psi_{21} + e^{4\pi i/3} \psi_{22}]) \\ &= \psi_{00} + \frac{1}{2} (\psi_{01} + \psi_{02} + \psi_{10} + \psi_{20}) = \frac{1}{2} (\Psi_1 + \Psi_2) \end{aligned}$$

At genus g : $Sp(2g)$ images of $\psi_{\vec{0}, \dots, \vec{0}}$ (40 terms at $g = 2$)

Average at any length and genus

By explicit computation, we can check

Ensemble average = sum over modular images

case by case in n , p and g . But can we **prove** it in general?

- Everything takes place inside finite-dimensional vector space

$$\mathcal{H}^{(g)} = \text{span} \left\{ \psi_{v_1 \dots v_g} \mid v_i \in \{0, \dots, p-1\}^{2n} \right\}, \quad \dim \mathcal{H}^{(g)} = p^{2gn}$$

- Two finite groups act on $\mathcal{H}^{(g)}$
 - $Sp(2g, \mathbb{F}_p)$ modular transformations
 - $O(n, n, \mathbb{F}_p)$: $\psi_{v_1, \dots, v_g} \mapsto \psi_{hv_1, \dots, hv_g}$. Acts transitively on the set of codes

Rewrite equality in a more symmetric form

$$\frac{1}{N'} \sum_{h \in O(n, n)} U_h \Psi_{C_0} \stackrel{?}{=} \sum_{\gamma \in Sp(2g)} U_\gamma \psi_{\vec{0}}$$

*If we can show that there is a **unique one-dimensional subspace** of $\mathcal{H}^{(g)}$ invariant under $Sp \times O$, then it follows that the two sides are proportional*

Howe duality/theta correspondence [Howe 1973, ...] explains certain structure of paired representations. **Ingredients:**

- Reductive dual pair $G \times H \subset \mathbf{G}$ mutually centralising. For us $Sp(2g) \times O(2n) \subset Sp(4gn)$
- ω oscillator representation: a certain “minimal” representation of $Sp(4gn)$:
For us $\omega \cong \mathcal{H}^{(g)}$

Central statement: *Irreps parametrised by joint label*

$$\omega|_{G \times H} = \sum_i (\rho_i)_G \otimes (\pi_i)_H$$

Similar to Schur–Weyl duality

$$(\mathbb{C}^n)^{\otimes N} = \sum_{\lambda} (r_{\lambda})_{S_N} \otimes (s_{\lambda})_{GL_n}$$

$$\text{Example: } \mathbb{C}^n \otimes \mathbb{C}^n = S^2 + \Lambda^2, \quad M_{\mu\nu} = M_{(\mu\nu)} + M_{[\mu\nu]}$$

We would like to prove

$$\frac{1}{N'} \sum_{h \in O(n,n)} U_h \Psi_{C_0} \stackrel{?}{=} \sum_{\gamma \in Sp(2g)} U_\gamma \psi_{\bar{0}} \quad (1)$$

Let us consider Howe duality with $H = O(n, n, \mathbb{F}_p)$, $G = Sp(2g, \mathbb{F}_p)$ and $\omega \cong \mathcal{H}^{(g)}$

$$\omega|_{Sp \times O} = \sum_i (\rho_i)_{Sp} \otimes (\pi_i)_O = \mathbf{1}_{Sp} \otimes \mathbf{1}_O + (\text{non-invariant terms}) \quad (2)$$

$\Rightarrow$ One-dimensional vector space invariant under $Sp \times O$, must contain both sides of (1)

Precise statements require **Howe duality over finite fields**. Structure of (2) follows from results conjectured in [Aubert, Michel, Rouquier 1996] and proven in [Pan 2024; Ma, Qiu, Zou 2024].

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CORRESPONDANCE DE HOWE POUR LES GROUPES RÉDUCTIFS SUR LES CORPS FINIS

ANNE-MARIE AUBERT, JEAN MICHEL ET RAPHAËL ROUQUIER

Soit (G, G') une paire de sous-groupes d'un groupe symplectique $\mathbf{Sp}_{2n}(\mathbb{F}_q)$ (où $\mathbb{F}_q$ est un corps fini de caractéristique p impaire) dont chacun est le centralisateur de l'autre dans $\mathbf{Sp}_{2n}(\mathbb{F}_q)$. Suivant R. Howe (voir [H1] et [H2]), il y a une correspondance entre représentations de G et de G' donnée par la représentation de Weil ω du groupe symplectique $\mathbf{Sp}_{2n}(q) := \mathbf{Sp}_{2n}(\mathbb{F}_q)$.

Plus précisément, si (G_m, G_m') est une paire réductrice duale irréductible dans $\mathbf{Sp}_{2n}(q)$, i.e., l'une des paires $(\mathbf{Sp}_{2n}(q), \mathbf{O}_{2n'}^\pm(q))$ (avec $n = 2mm'$), $(\mathbf{Sp}_{2n}(q), \mathbf{O}_{2m'+1}(q))$ (avec $n = m(2m' + 1)$), $(\mathbf{U}_m(q), \mathbf{U}_{m'}(q))$, $(\mathbf{GL}_m(q), \mathbf{GL}_{m'}(q))$ (avec $n = mm'$), la restriction de ω à $G_m \cdot G_m'$ définit une application entre groupes de Grothendieck de représentations complexes $\mathcal{R}(G_m) \rightarrow \mathcal{R}(G_m')$ que nous appellerons correspondance de Howe.

- We have investigated “holography by averaging” following the Narain ensemble blueprint:

ensemble average = sum over modular images

- We exhibited explicit instances using the **ensemble of error-correcting codes**
- We proved equality using Howe duality over finite fields

$$\mathcal{H}^{(g)} = \mathbf{1}_{Sp} \otimes \mathbf{1}_O + \dots$$

$\Rightarrow$ both sides of equality must lie in one-dimensional vector space invariant under $Sp \times O$

Next: What does this have to do with “Chern–Simons gravity”?

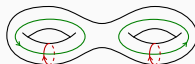
How to construct Chern–Simons gravity?

- Consider level- p abelian CS theory on B ,
 $\partial B = \Sigma_g$

- Hilbert space $\mathcal{H}^{(g)} = \text{span } \Psi_{v_1, \dots, v_g}$,

v_i charges of Wilson lines on

non-contractible cycles



- This TQFT has a global $\mathbb{Z}_p \times \mathbb{Z}_p$ 1-form symmetry carried by the Wilson lines
- “No global symmetry:” gauge away symmetry (maximal non-anomalous subgroup) [Benini, Copetti, Di Pietro 2022]. $\Psi = \sum_v \Psi_v$

Maximal non-anomalous subgroup $\leftrightarrow$ even self-dual code $\mathcal{C}$

- No unique choice: Average over all gaugings/codes with equal weight

$$\frac{1}{N} \sum_{\mathcal{C}} \Psi_{\mathcal{C}} \quad [\text{Barbar, Dymarsky, Shapere 2023}]$$

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Alternatively, one could follow the logic

- Gravity implies a sum over topologies
- Sum over all B with $\partial B = \Sigma_g$ (handlebodies)
- Sum reorganises to sum over modular images of fixed topology (c.f.

[Maloney, Witten 2007])

$$\sum_{\gamma} Z_{B_0}(\gamma\Omega)$$

$$Z_{B_0}(\Omega) = \Psi_{\vec{0}, \dots, \vec{0}}(\Omega)$$

Summary and outlook

Summary

- We have seen an explicit instance of holography by averaging
$$\text{ensemble average} = \text{sum over modular images} \quad (1)$$

- The ensemble arises naturally by gauging maximal subgroups of $\mathbb{Z}_p \times \mathbb{Z}_p$ 1-form symmetry of abelian CS theory

sum over gaugings = sum over topologies

Thus (1) represents **equality of two dual pictures** of “Chern–Simons gravity”

- We proved (1) using Howe duality over finite fields

$$\mathcal{H}^{(g)} = \mathbf{1}_{Sp} \otimes \mathbf{1}_O + \dots$$

Outlook

- Different manifestations of Howe duality: Siegel–Weil formulas, quantum information (ω =qudit Hilbert space), spherical potential $\psi_{nlm} = R_{nl}(r)Y_{lm}(\theta, \phi)$
[Ashwinkumar et al 2021], [Gross, Nezami, Walter 2017], [Rowe et al 2012; Basile et al 2020]
- Study other observables in ensemble averaging

$$\langle \mathcal{O}_i \dots \rangle_{\text{theories containing } \mathcal{O}_i} \stackrel{?}{=} \frac{\delta}{\delta \xi_i} \dots \sum_{\gamma} U_{\gamma}(\Psi_{vac} + \Psi_{\mathcal{O}_i})$$

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Thank you for your attention!

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