

EXPLORING BLACK HOLE HILBERT SPACES VIA THE GRAVITY PATH INTEGRAL

arXiv: 2410.00091

arXiv: 2412.06884

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The Question

- Bekenstein and Hawking: black holes are thermodynamic objects with an entropy

$$S_{\text{BH}} = \frac{\mathcal{A}}{4G_N}$$

- If this entropy counts microstates, then

$$\dim(\mathcal{H}_{\text{BH}}) = e^{S_{\text{BH}}}$$

- **Exponentially large number of black hole microstates?**

Paths to Resolution

- String theory approach [Strominger and Vafa 1996]
- "Bag of gold" geometries [Wheeler 1964]

Black hole microstates



A Breakthrough

Can we give a universal microscopic explanation for the entropy of black holes?

Microscopic Origin of the Entropy of Black Holes in General Relativity

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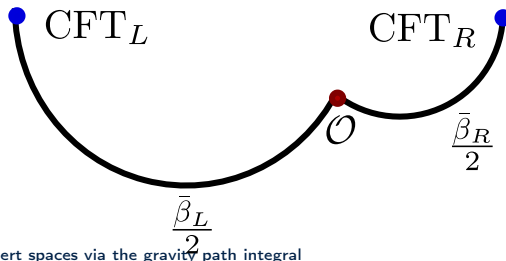
Family of States

Family of States in Holographic CFT

Consider states on the tensor product of two identical copies of a d -dimensional CFT defined on $\mathbb{S}^{d-1} \times \mathbb{R}$

$$|\Psi\rangle = \sum_{n,m} \Psi_{mn} |E_m\rangle_L \otimes |E_n\rangle_R$$

$$\Psi_{mn} = \frac{e^{-\frac{1}{2}\bar{\beta}_L E_m - \frac{1}{2}\bar{\beta}_R E_n}}{\sqrt{Z_\Psi}} \langle m | \mathcal{O} | n \rangle \text{ and } Z_\Psi = \text{Tr} \left[\mathcal{O}^\dagger e^{-\bar{\beta}_L H} \mathcal{O} e^{-\bar{\beta}_R H} \right]$$

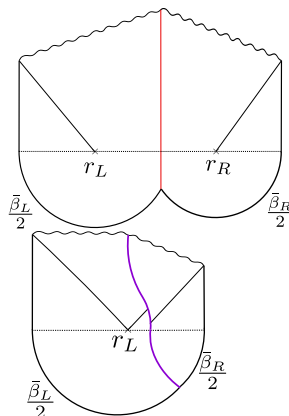
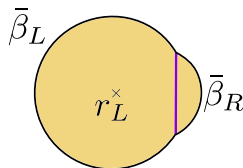
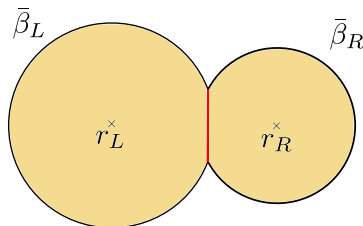


Family of States

Family of States in Gravity

Operators $\mathcal{O}$ in CFT

$\leftrightarrow$ Shell of matter in gravity



Microcanonical Projection

The Projection Operator

- Assume $\mathcal{H} = \bigoplus_E \mathcal{H}^E$
- Define the microcanonical projector $\Pi_E : \bigoplus_{E'} \mathcal{H}^{E'} \rightarrow \mathcal{H}^E$
- Black hole microstates

$$|\psi^E\rangle = \frac{1}{\sqrt{Z_\Psi}} \sum_{|E_m\rangle_L \otimes |E_n\rangle_R \in \mathcal{H}^E} \mathcal{O}_{mn} |E_m\rangle_L \otimes |E_n\rangle_R$$

- In practice, use inverse Laplace transform

$$f(E) \propto \int d\beta e^{\beta E} F(\beta)$$

Counting states

Strategy

- Introduce states $|\Psi_k^E\rangle$, where $\mathcal{O}^{(k)}$ are a product of n scalar operators $\mathcal{O}_{\Delta_k}$ that create shells with mass $m_k = km_{\Delta_1}$

$$\mathcal{H}_{\text{bulk}}^E(\Omega) \equiv \text{Span}\{|\Psi_k^E\rangle, \quad k = 1, \dots, \Omega\}$$

- Define the Gram matrix $G_{ij} = \langle \Psi_i^E | \Psi_j^E \rangle$

$$\dim \left(\mathcal{H}_{\text{bulk}}^E(\Omega) \right) = \text{rank}(G) = \Omega - \dim(\text{Ker}(G))$$

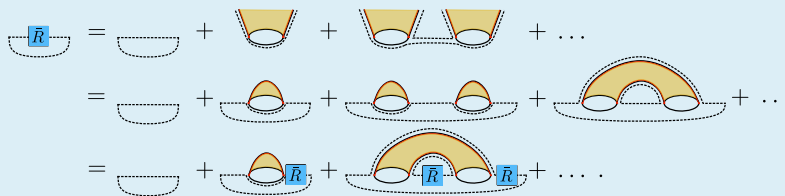
- Find $\text{Ker}(G)$ using the resolvent matrix $R_{ij}(\lambda) = \left(\frac{1}{\lambda \mathbb{1} - G} \right)_{ij}$

$$R_{ij}(\lambda) = \frac{1}{\lambda} \delta_{ij} + \sum_{n=1}^{\infty} \frac{1}{\lambda^{n+1}} (G^n)_{ij}$$

Counting States

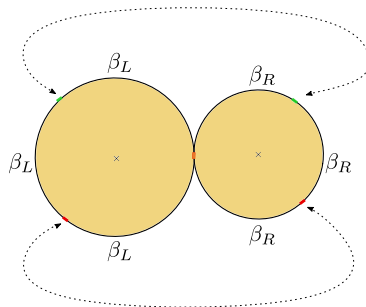
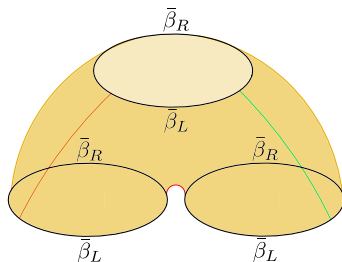
$$R(\lambda) = \frac{\Omega}{\lambda} + \frac{1}{\lambda} \sum_{n=1}^{\infty} \sum_{i=1}^{\Omega} \frac{1}{\lambda^n} (G^n)_{ii}$$

Semiclassical Approximation for Trace of the Resolvent



$$\bar{R}(\lambda) = \frac{\Omega}{\lambda} + \frac{1}{\lambda} \sum_{n=1}^{\infty} \sum_{i=1}^{\Omega} \frac{\bar{R}^n}{\Omega^n} (\overline{G^n})_{ii} \Big|_{\text{conn.}}$$

Counting States



$$\overline{\langle \Psi_{k_1} | \Psi_{k_2} \rangle \langle \Psi_{k_2} | \Psi_{k_3} \rangle \dots \langle \Psi_{k_n} | \Psi_{k_1} \rangle} \Big|_{\text{conn.}} = \frac{Z(n\beta_L)Z(n\beta_R)e^{-nI_{\text{univ.}}}}{(Z(\beta_L)Z(\beta_R)e^{-I_{\text{univ.}}})^n}$$

$$\overline{\langle \Psi_{k_1}^E | \Psi_{k_2}^E \rangle \langle \Psi_{k_2}^E | \Psi_{k_3}^E \rangle \dots \langle \Psi_{k_n}^E | \Psi_{k_1}^E \rangle} \Big|_{\text{conn.}} = e^{(1-n)(S_{\text{BH}}(E_L) + S_{\text{BH}}(E_R))}$$

Counting States

Solving the Schwinger-Dyson Equation

$$\lambda \overline{R} = \Omega + \frac{\overline{R} e^{S_L + S_R}}{e^{S_L + S_R} - \overline{R}}$$

Solution:

$$\overline{R}_+ = \frac{\Omega - e^{S_L + S_R}}{\lambda} \Theta \left(\Omega - e^{S_L + S_R} \right) + R_0 + O(\lambda)$$

$$\Rightarrow \overline{\dim \left(\mathcal{H}_{\text{bulk}}^E(\Omega) \right)} = \begin{cases} \Omega, & \Omega < e^{S_L + S_R} \\ e^{S_L + S_R}, & \Omega \geq e^{S_L + S_R} \end{cases}$$

Factorization of Microcanonical Bulk Traces

Microcanonical Bulk Trace

$$\text{Tr}_{\mathcal{H}_{\text{bulk}}^E} (f(H_L)g(H_R)) = f(E_L)g(E_R)\text{dim}(\mathcal{H}_{\text{bulk}}^E)$$

- coarse-grained

$$\overline{\text{dim}(\mathcal{H}_{\text{bulk}}^E(\Omega))} = \begin{cases} \Omega, & \Omega < e^{S_L+S_R} \\ e^{S_L+S_R}, & \Omega \geq e^{S_L+S_R} \end{cases}$$

- fine-grained (*)

$$\overline{\left(\text{Tr}_{\mathcal{H}_{\text{bulk}}^E} (f(H_L)g(H_R)) - \overline{\text{Tr}_{\mathcal{H}_{\text{bulk}}^E} (f(H_L)g(H_R))} \right)^2} = 0$$

Factorization of Microcanonical Bulk Traces

Fine-grained Factorization

- Boruch et al. 2024 **How the Hilbert space of two-sided black holes factorises** [arXiv:2406.04396]
- Balasubramanian et al. 2025 **Factorization of the Hilbert space of eternal black holes in general relativity** [arXiv:2410.00091]
- Balasubramanian and Yildirim 2025a **A Nonperturbative Toolkit for Quantum Gravity** [arXiv:2504.16986]
- Balasubramanian and Yildirim 2025b **The Nonperturbative Hilbert Space of Quantum Gravity With One Boundary** [arXiv:2506.04319]

Further Investigation and Outlook

Further Investigation

- Climent et al. 2024 **Universal construction of black hole microstates** [arXiv:2401.08775]
- Balasubramanian et al. 2024 **Counting microstates of out-of-equilibrium black hole fluctuations** [arXiv:2412.06884]
- Abdalla et al. 2025 **The gravitational path integral from an observer's point of view** [arXiv:2501.02632]
- ...

Outlook

- Hilbert space of black holes formed from collapse
- Quantum corrections
- de Sitter microstates
- ...

References



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