

Jackson R. Fliss

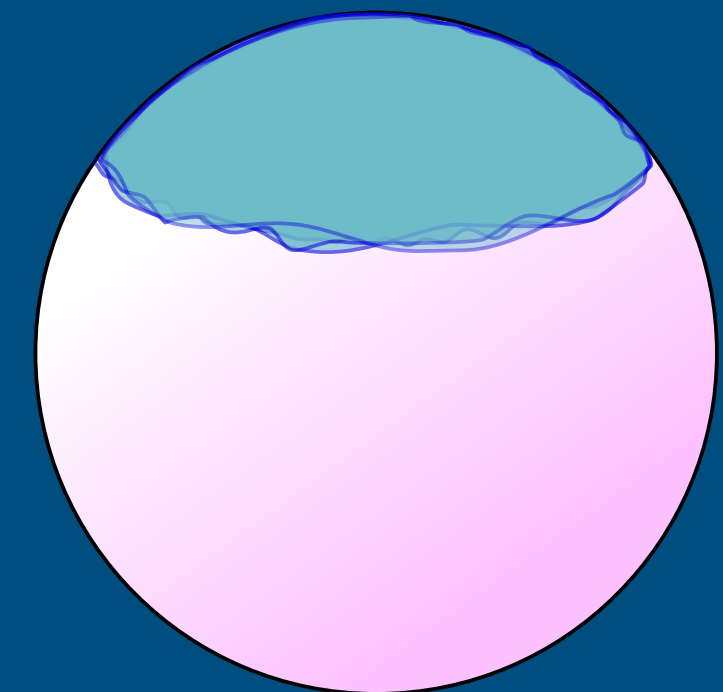
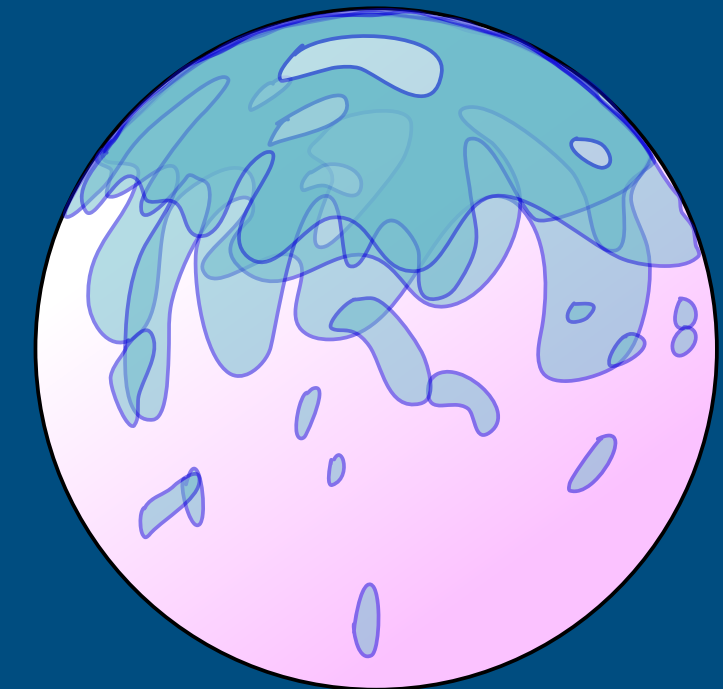
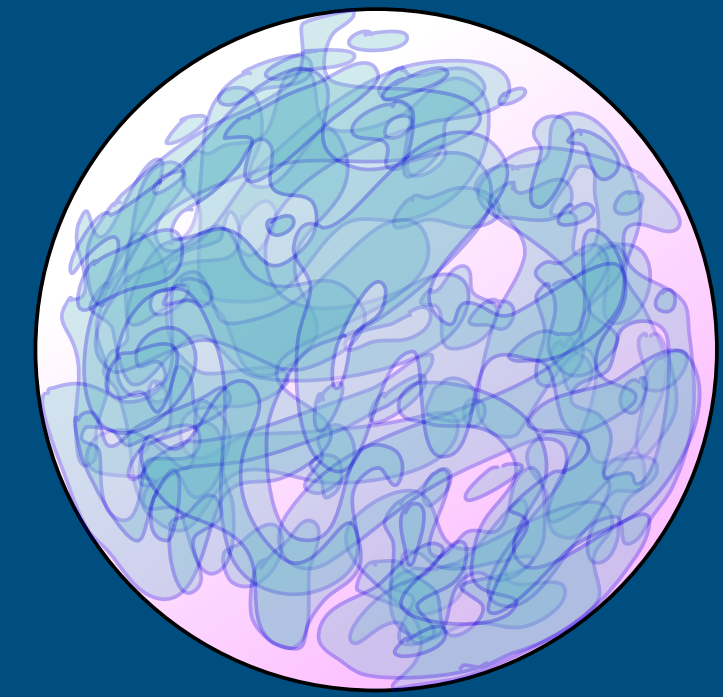


UNIVERSITY OF
CAMBRIDGE

Minimal areas from entangled matrices

2408.05274

w/ Alex Frenkel, Sean Hartnoll, and Ronak Soni



Eurostrings 2025

How does geometry emerge from information?

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 - Black holes are thermodynamic with an entropy $\sim$ horizon area.

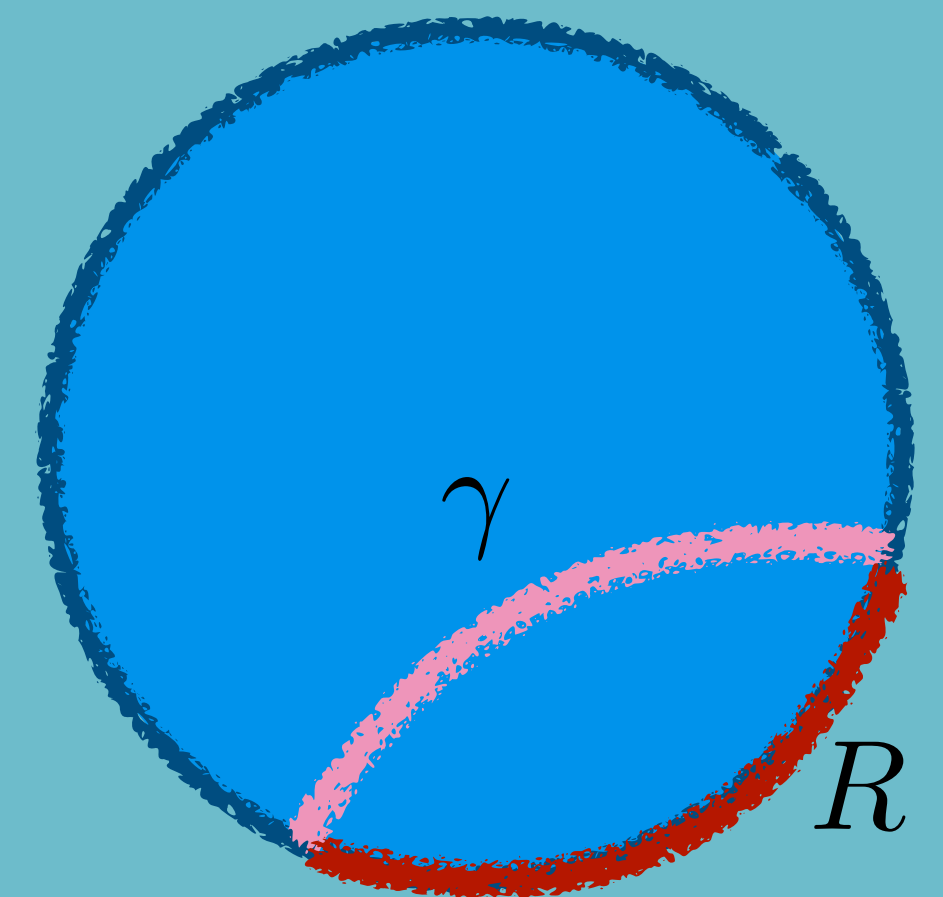
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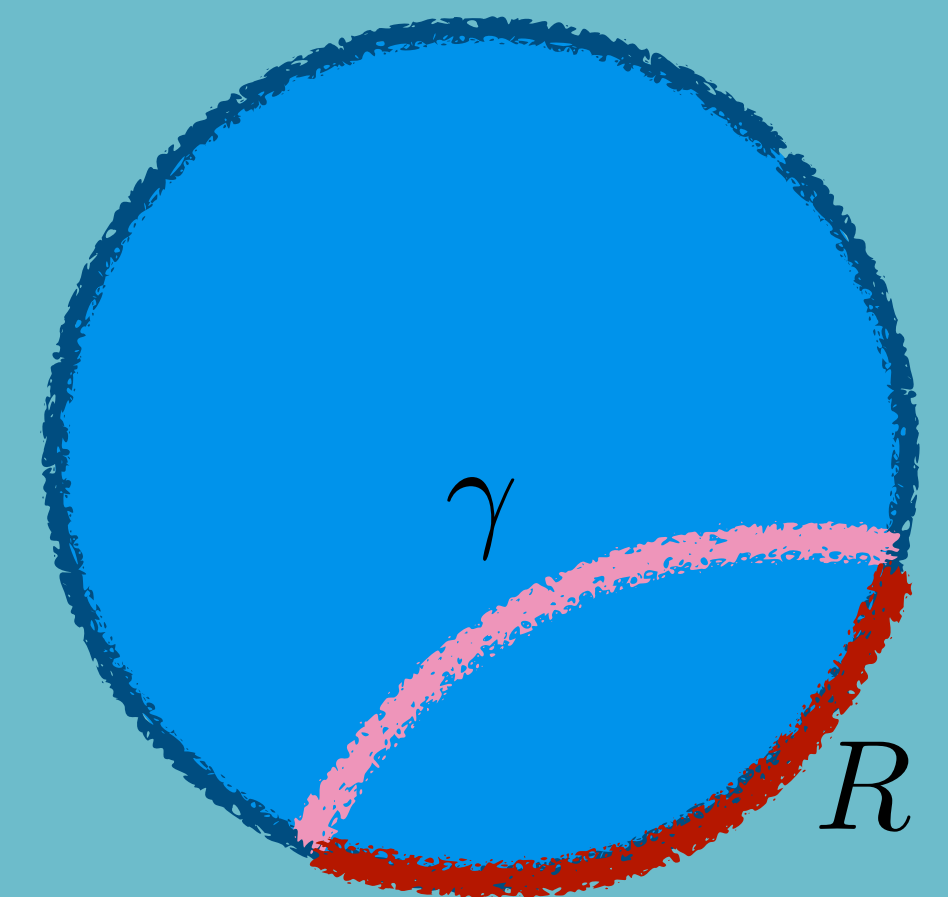
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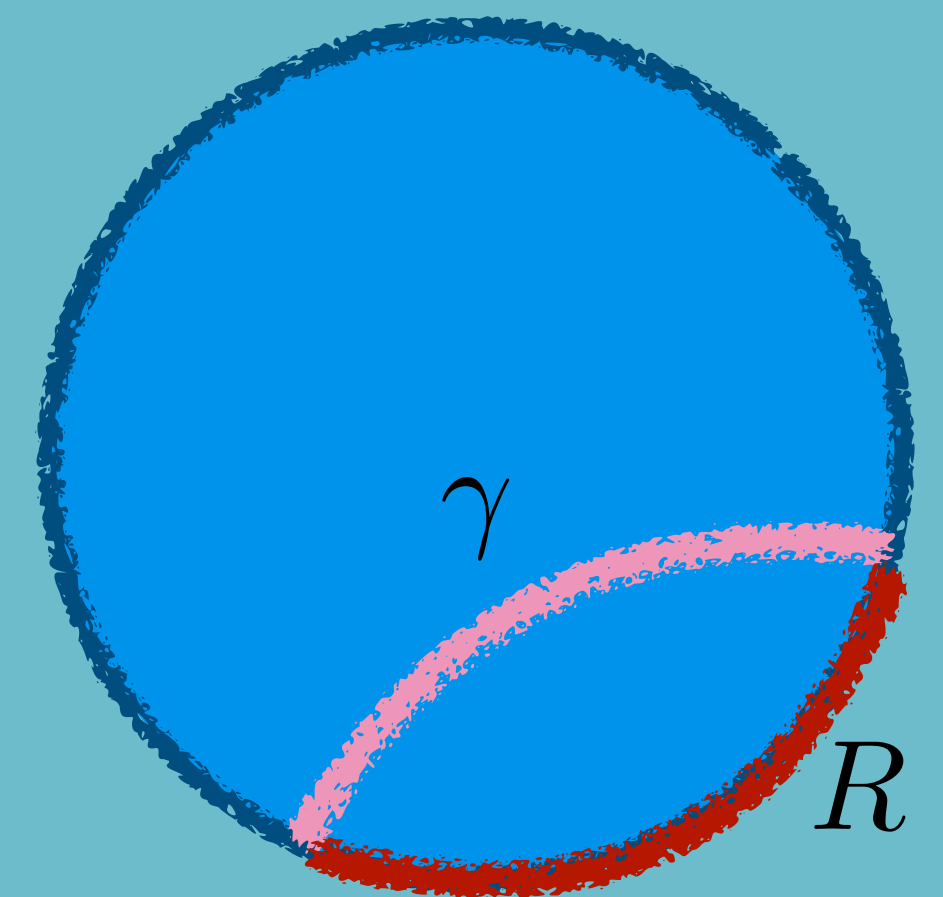
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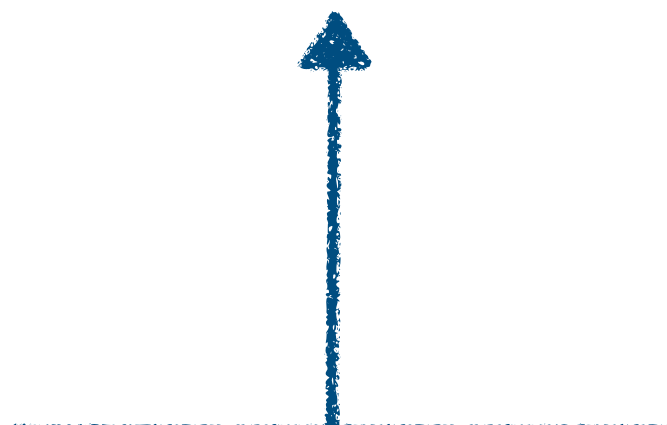


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Look “under the hood” of this mechanism...

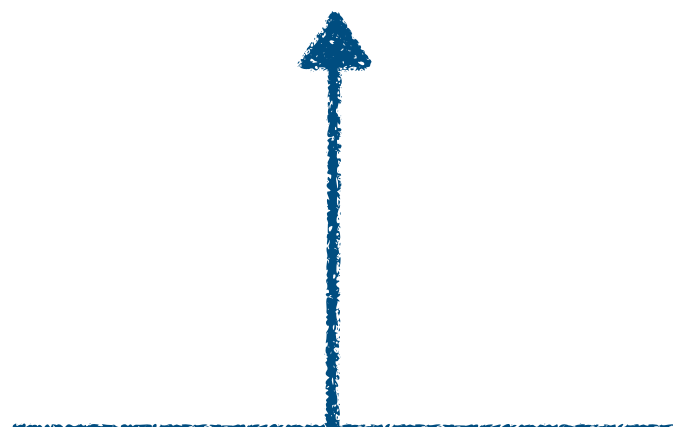
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Area as state counting

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A blue rectangular box highlights the term $\text{Area}(\gamma)$ in the numerator of the fraction. A blue arrow points vertically upwards from the top center of this box to the title *Area as state counting*.

Area as state counting

a la horizon
microstate counting

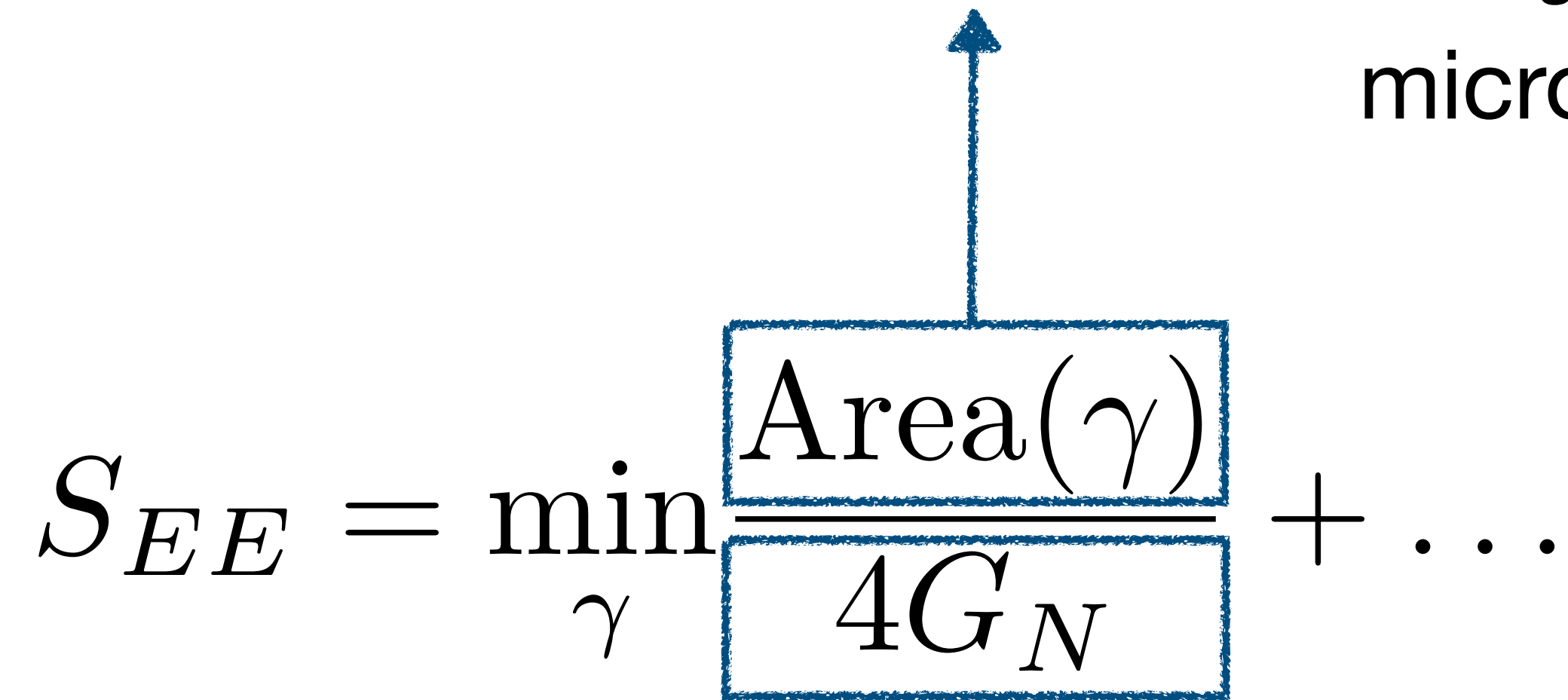


The diagram consists of a blue rectangular box containing the text $\text{Area}(\gamma)$. A blue arrow points vertically upwards from the top center of this box towards the word "Area" in the title "Area as state counting" located above the equation.

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Inverse coupling (instead of UV cutoff)

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Appearance of low-energy
effective field theory

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Consequence of diffeomorphism
invariance, subject to some
gauge-invariant condition

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($c=1$ matrix model)*

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$$\mathcal{L} = \text{Tr} \left[\sum_{a=1}^d \left(\dot{X}^a \right)^2 + \sum_{a < b} [X^a, X^b]^2 + \dots \right] \quad X^a: \mathbf{N} \times \mathbf{N} \text{ matrices}$$

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Not (necessarily) quantum gravity or holography.

Regardless, interesting, tractable, models displaying similar key features

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- Non-commutative sphere of radius $\frac{\nu N}{2}$

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- At large N acts on membrane theory as ***volume preserving diffeomorphisms***.
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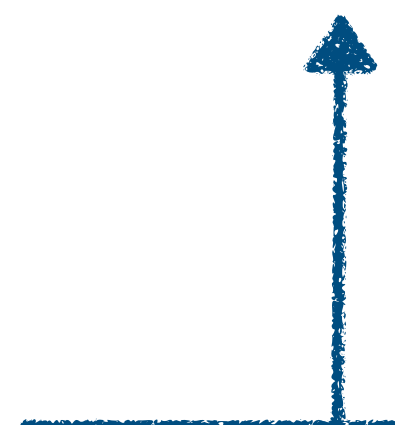
The goal: VPD-invariant entanglement entropies

The result:

$$S_{EE} = \min_{\substack{\Sigma, \\ |\Sigma| \text{ fixed}}} \frac{|\partial\Sigma|}{g_{\text{eff}}} \log \left(g_{\text{eff}} N \frac{|\Sigma|}{|\partial\Sigma|} \right)$$

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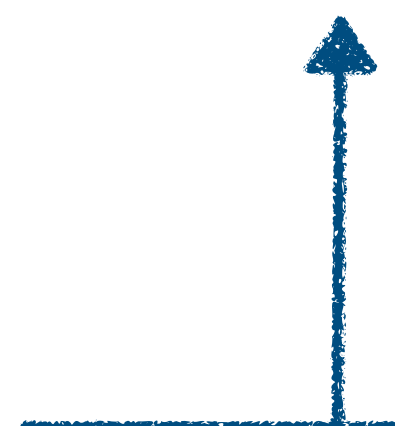
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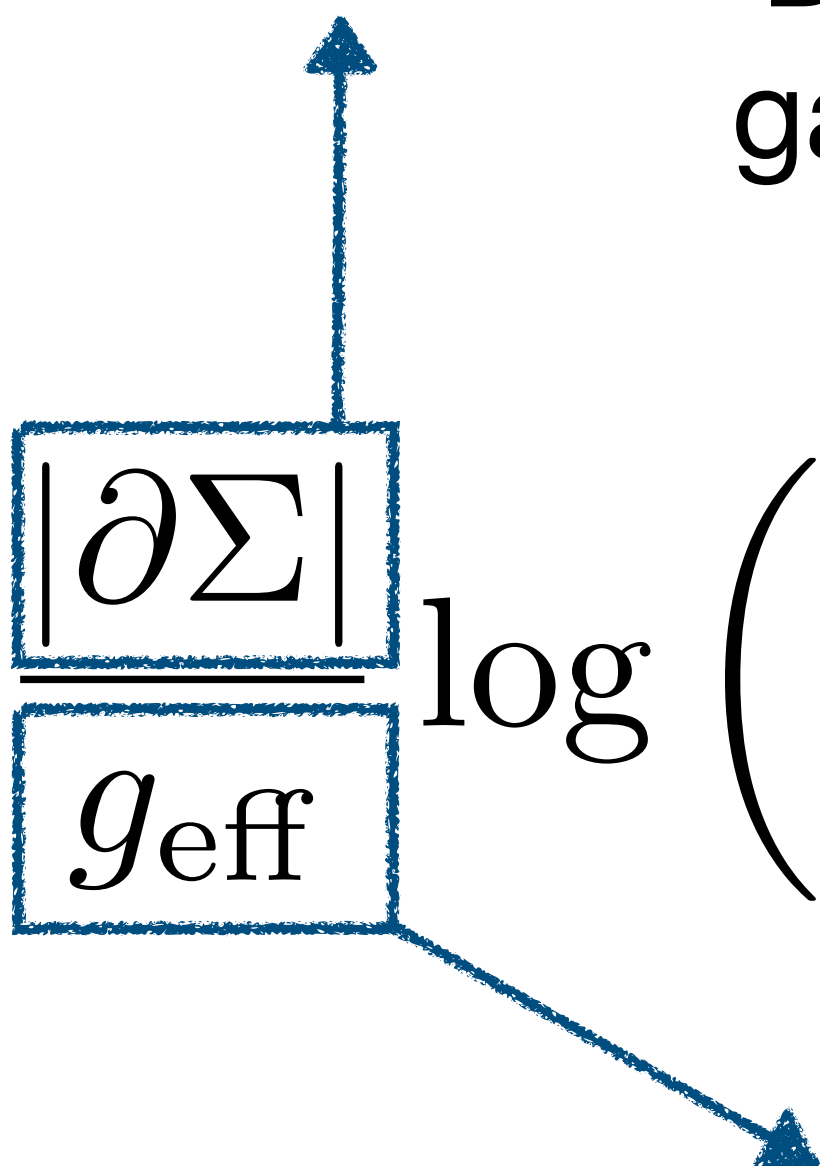
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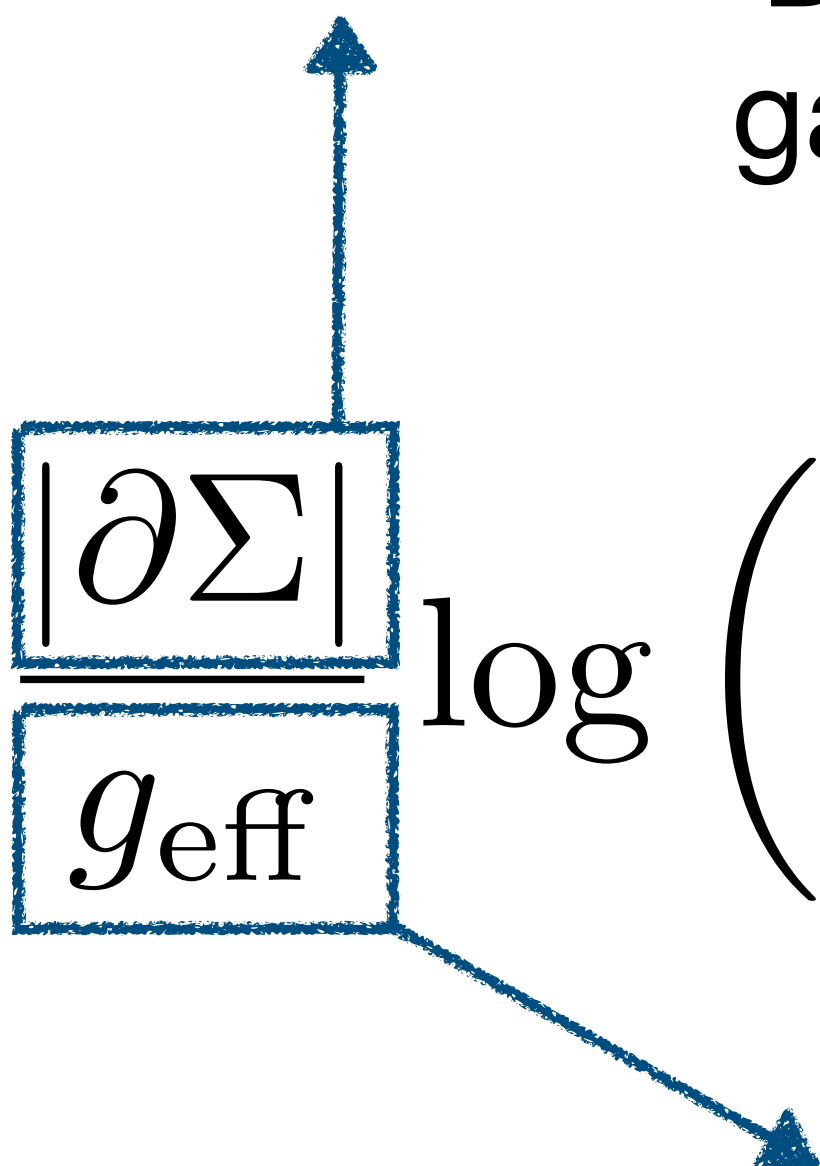
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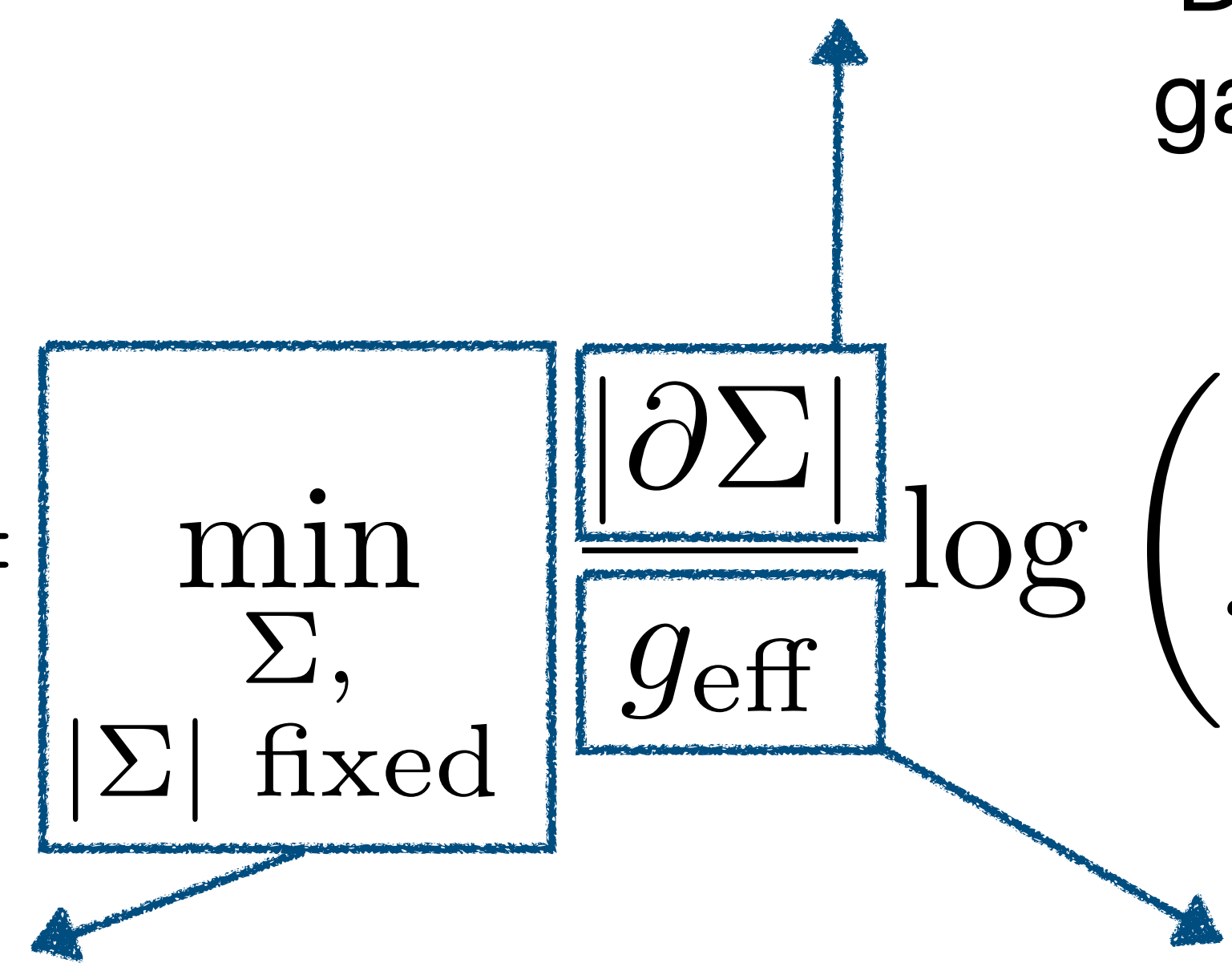
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Saddle-point in average over VPDs that change the subsystem

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Subsystems and gauge (in)variance

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[Buividovich,
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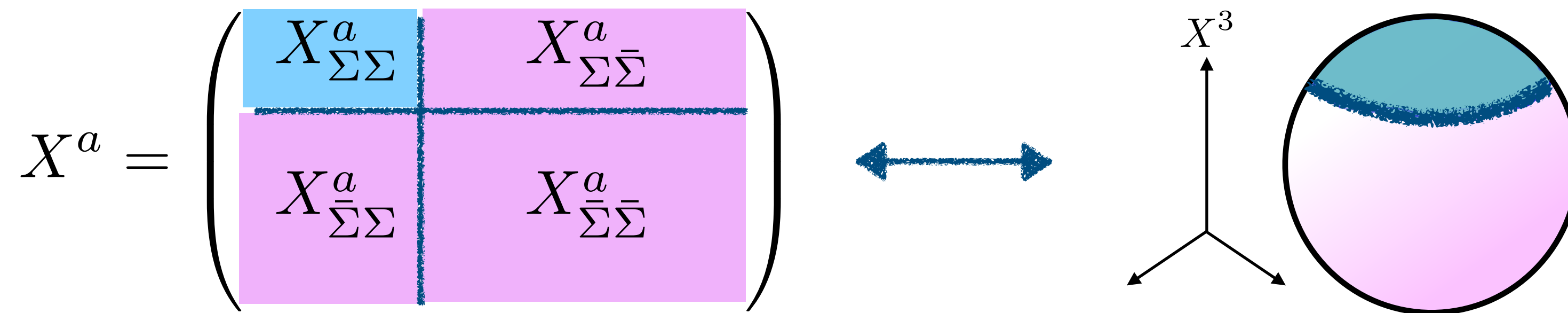
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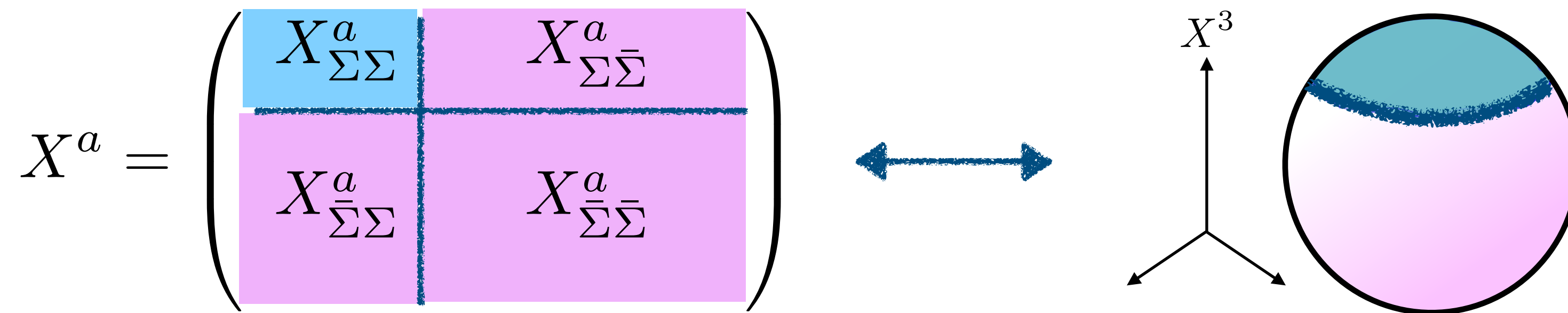
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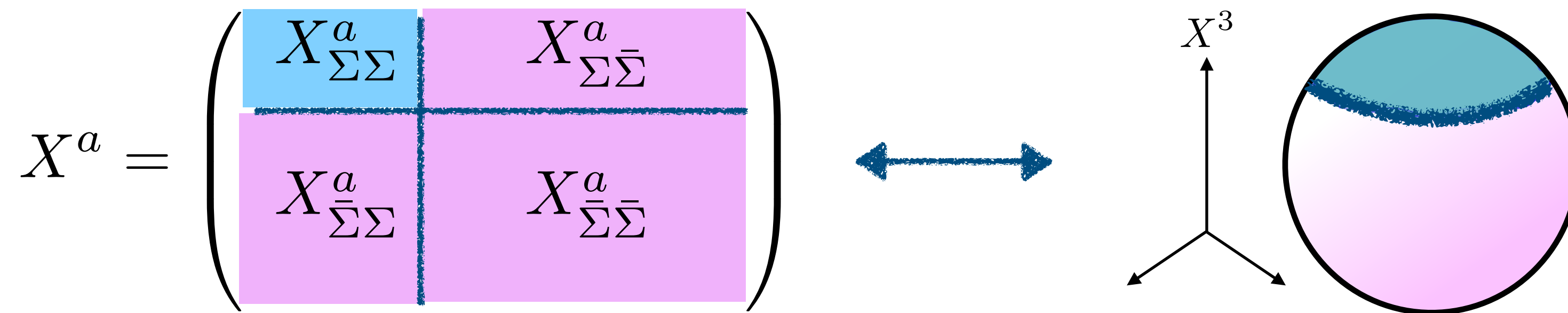
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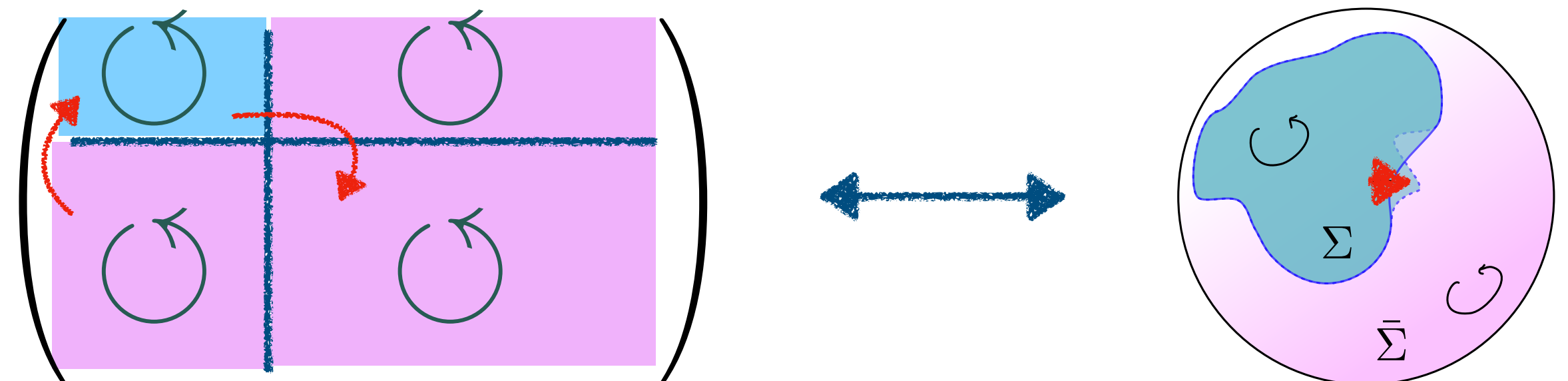
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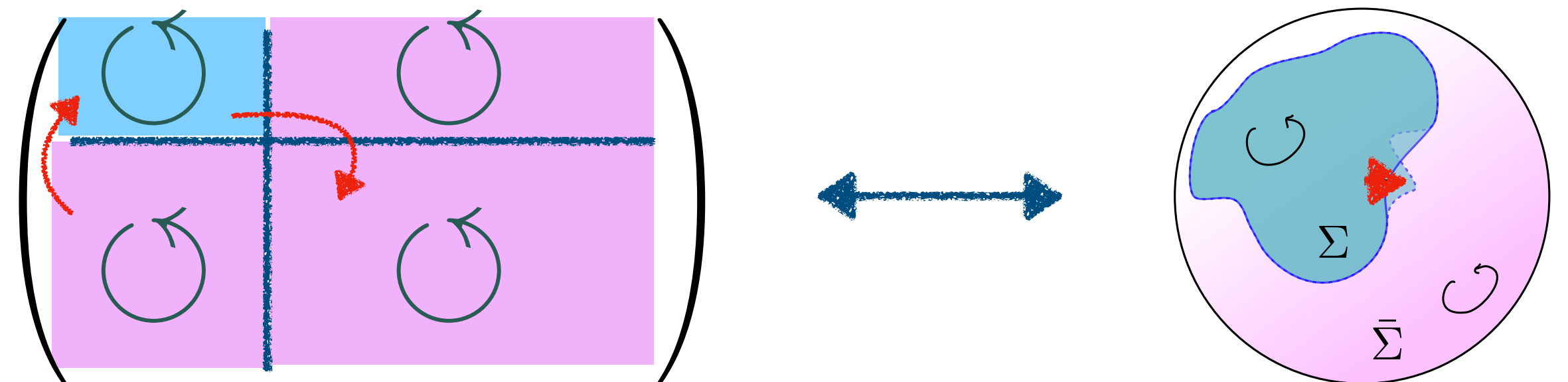
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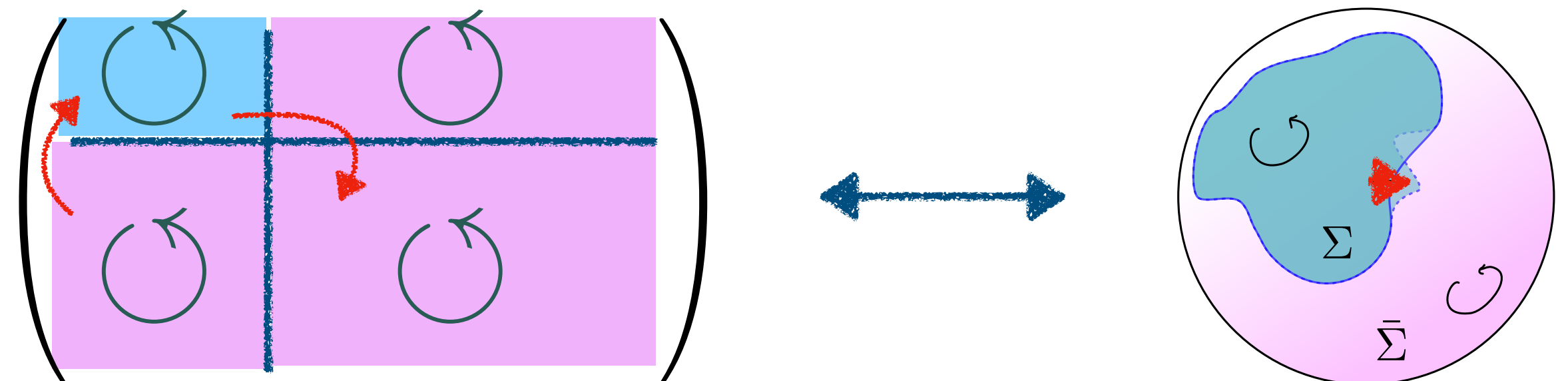
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$$U(N) / (U(M) \times U(N - M)) \equiv F$$



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what is the entropy of *any*
 $M \times M$ matrix block?



what is the entropy of *any*
region of fixed volume?

Gauge invariant entanglement

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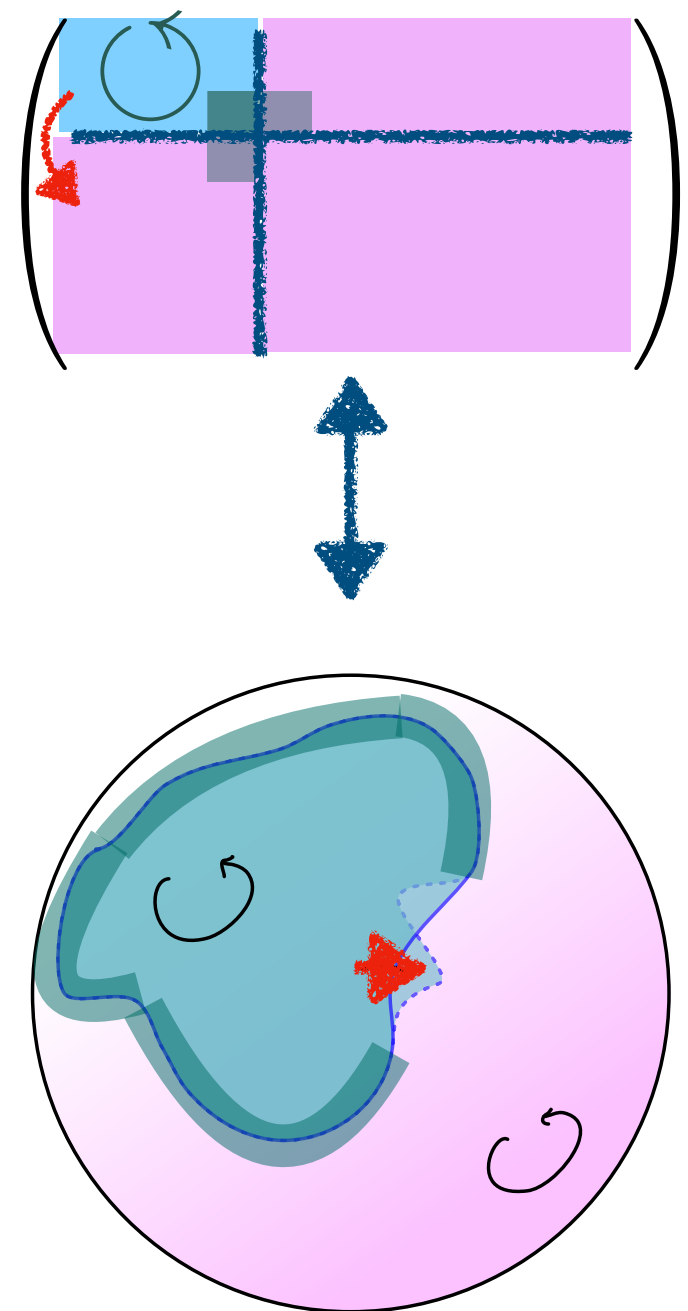
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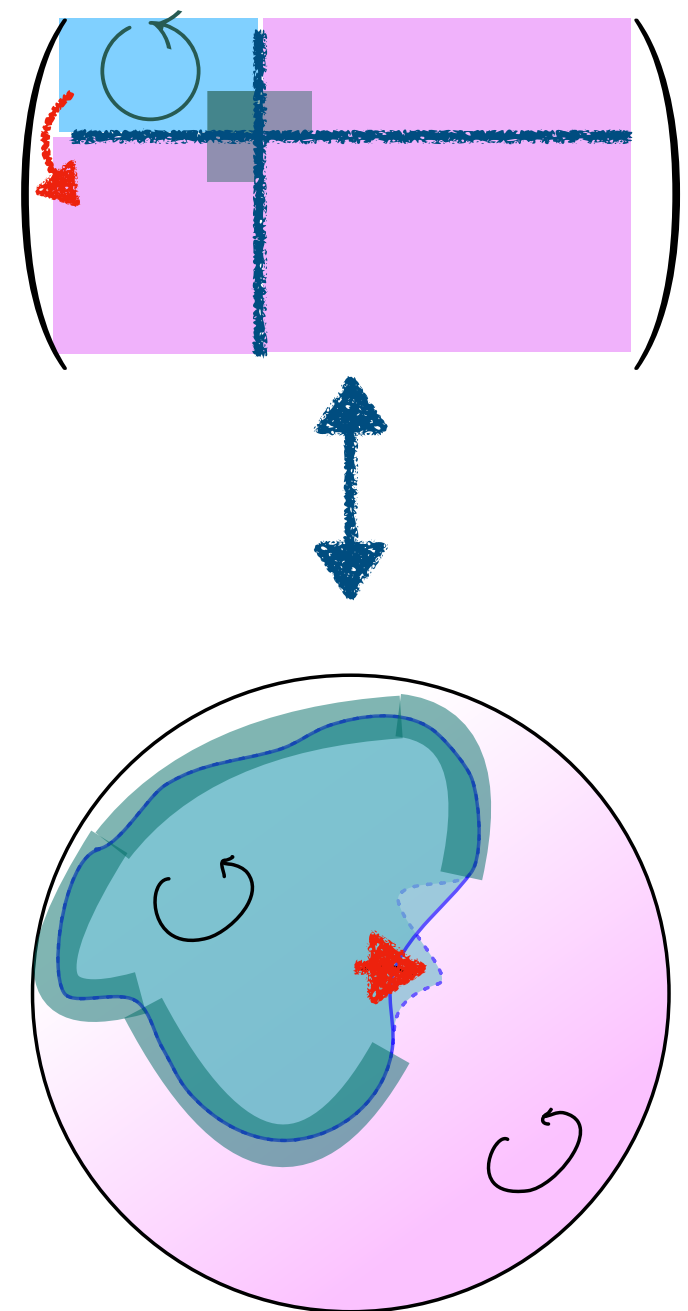
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- $U(M) \times U(N - M)$: similar to story in lattice gauge theory
 - Maximally entangled edge modes in irreps μ_Σ of $U(M)$.
- F : no analogue in local gauge theory
 - $V \in F$ leads to different subsystems being traced out.

Area from the $U(M)$ charge

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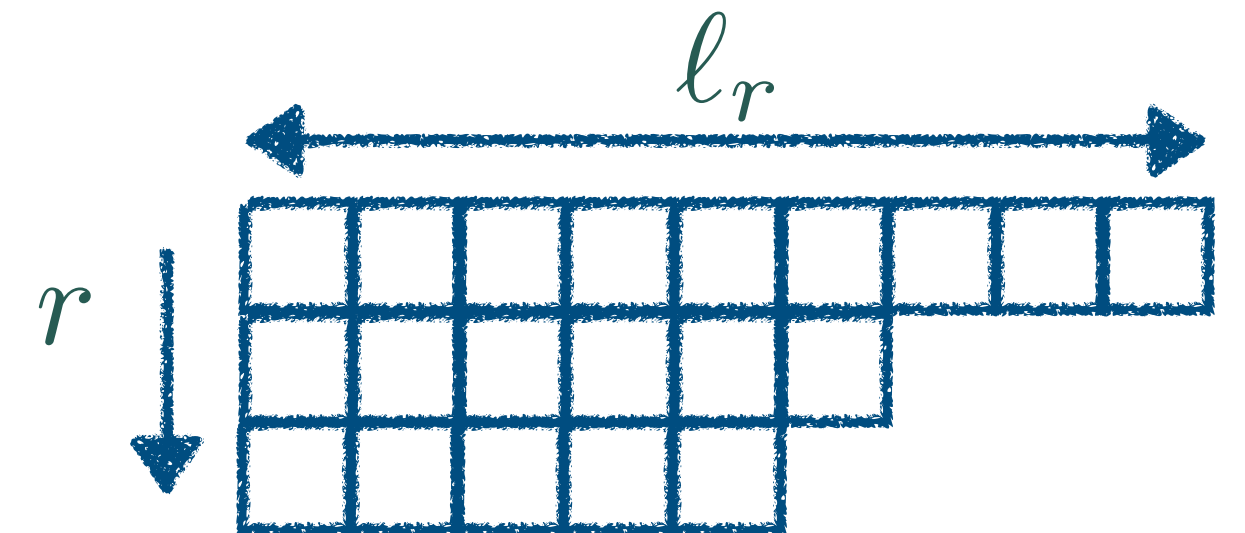
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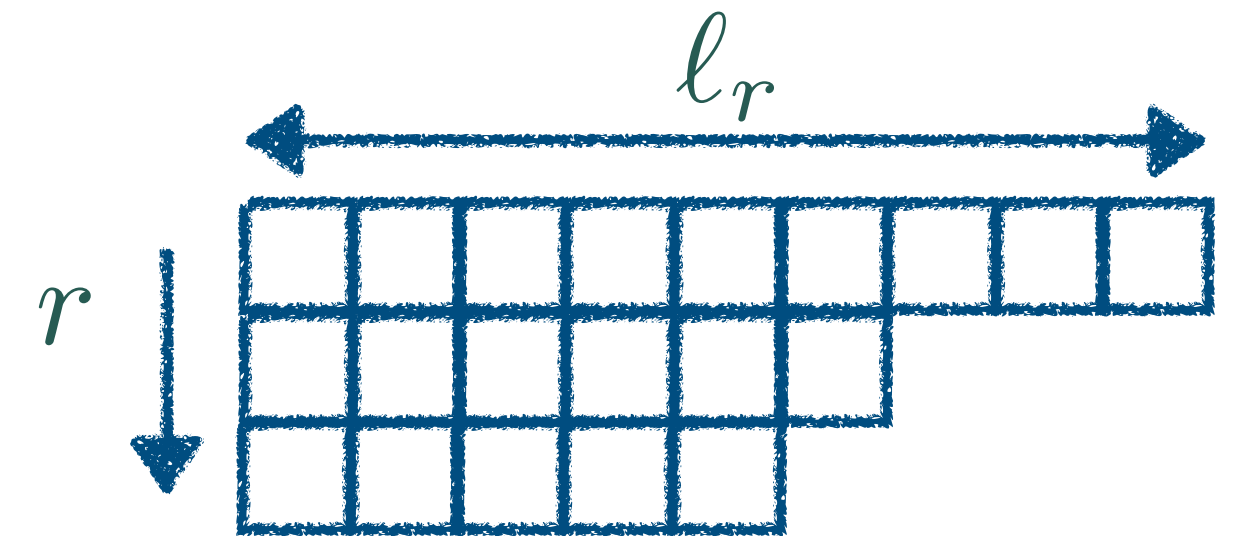
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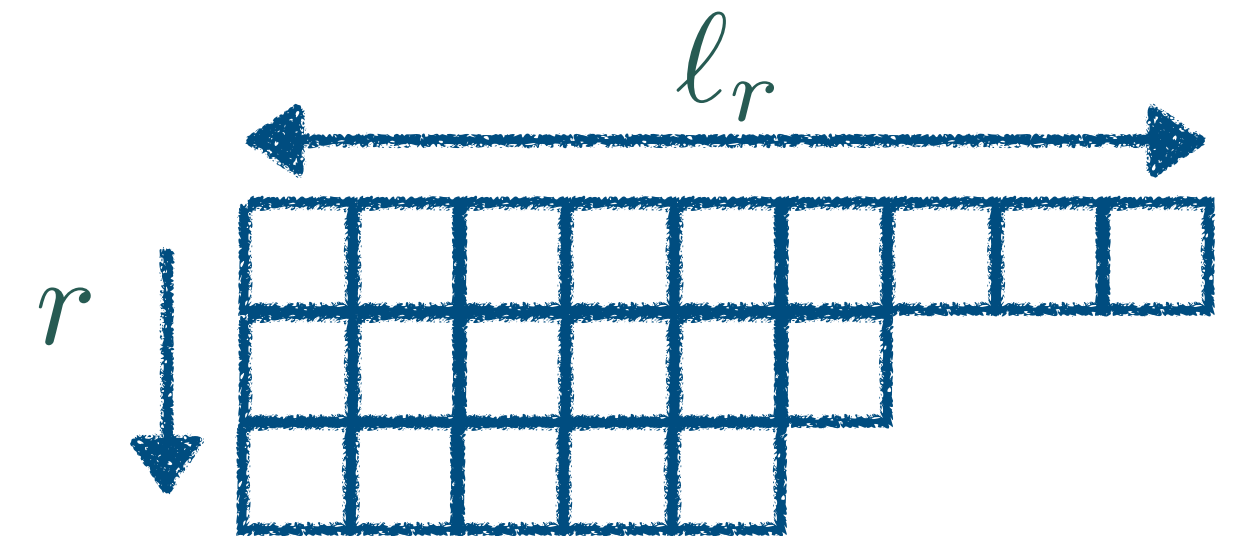
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[Frenkel; '23]

$$\log d_{\mu^*} \approx \sum_r \ell_r \log \frac{M}{\ell_r} = \sum_r |\partial \Sigma_{(r)}| \log \left(\frac{N |\Sigma|}{|\partial \Sigma_{(r)}|} \right)$$

what is the entropy of *any* $M \times M$ matrix block?

Incorporating F

- ***Assumption:*** reduced states of different subregions are *distinguishable* (to leading order in N)

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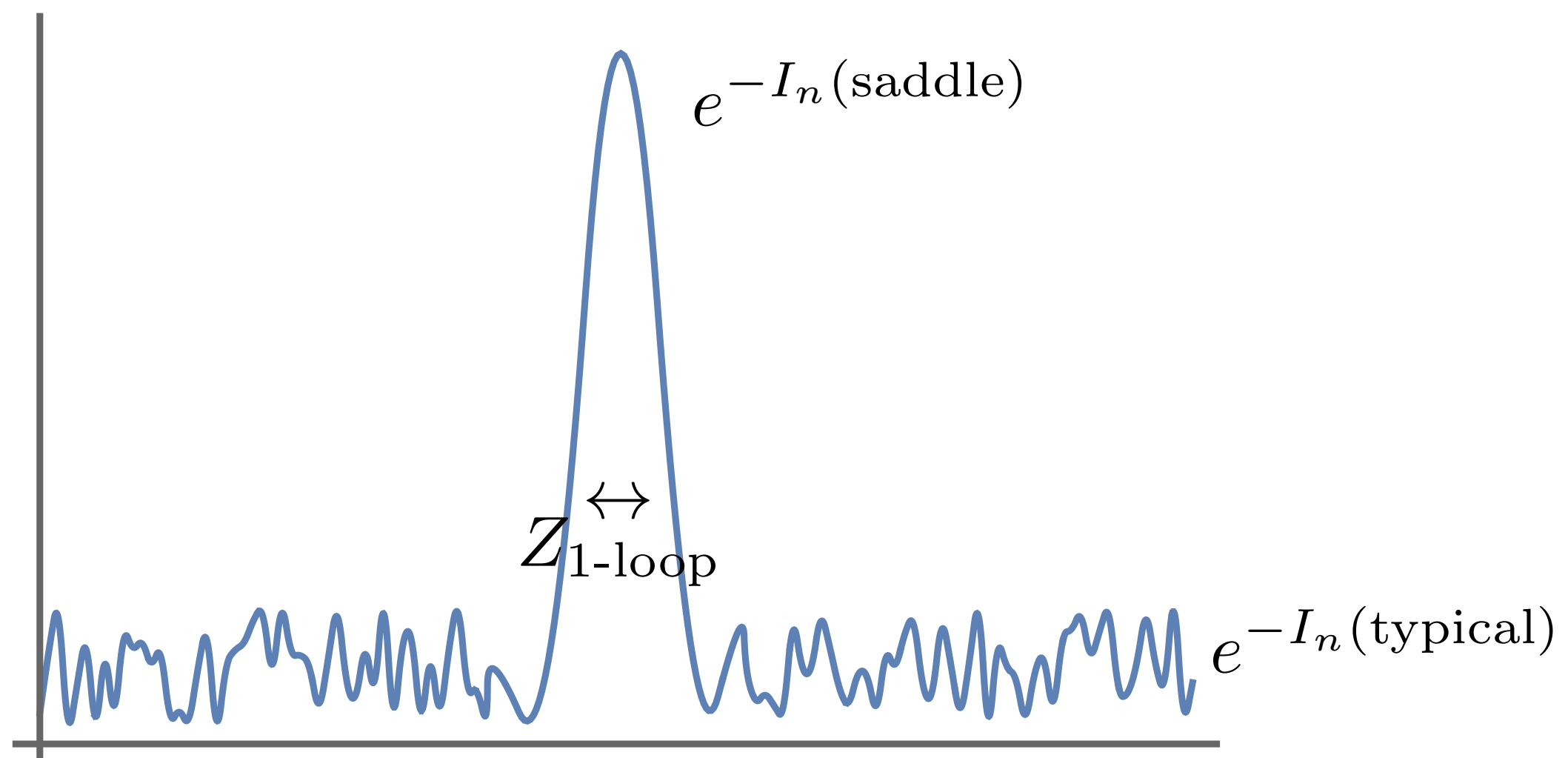
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$$\approx Z_{1\text{-loop}} e^{-I_n(\text{saddle})} + e^{-I_n(\text{typical})}$$

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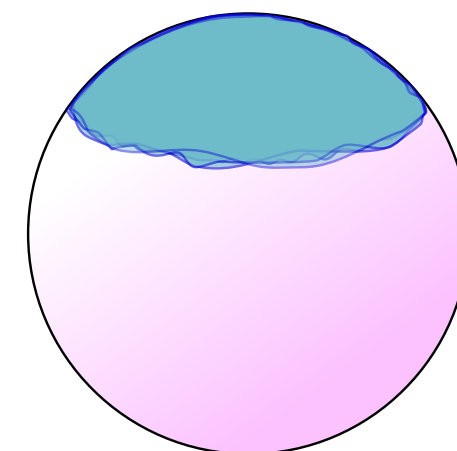
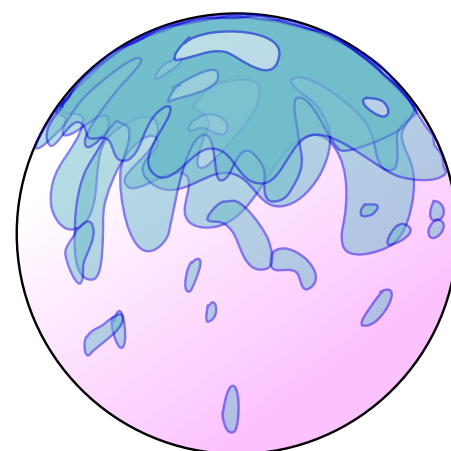
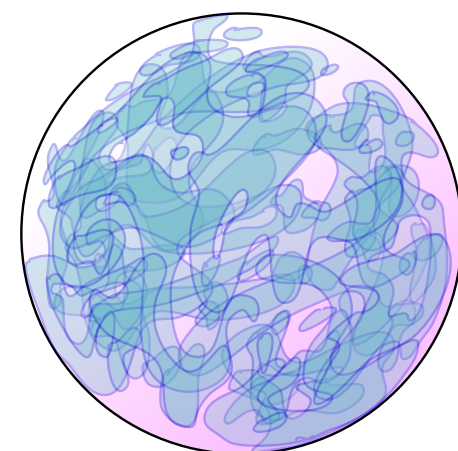
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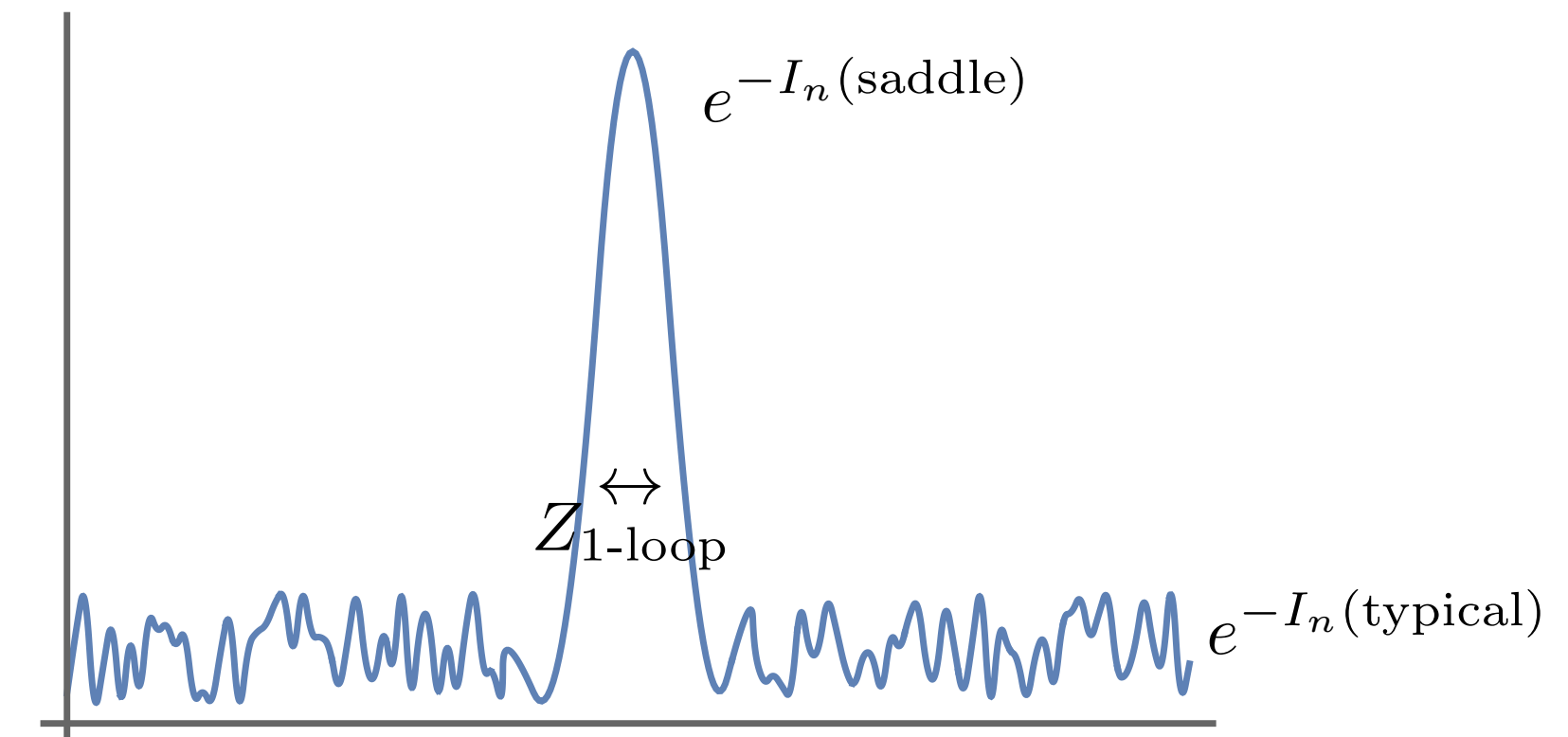
- *Revisit replica symmetry: accurate for $n \sim 1$, but RS breaking important for Rényis*
- *Incorporating Lorentz / SUSY: MQMs relevant for flat-space supergravity, e.g. BFSS, BMN*
- *Towards an RT formula outside of holography: gauge fixing to more abstract features than a boundary. QRFs.*

Thank you for your attention.

Extra slides

Fine print on the saddle point

$$\mathrm{Tr} \rho_{\mathrm{in}}^n \approx Z_{1\text{-loop}} e^{-I_n(\text{saddle})} + e^{-I_n(\text{typical})}$$



- *Establishing the existence + dominance of a saddle point.*

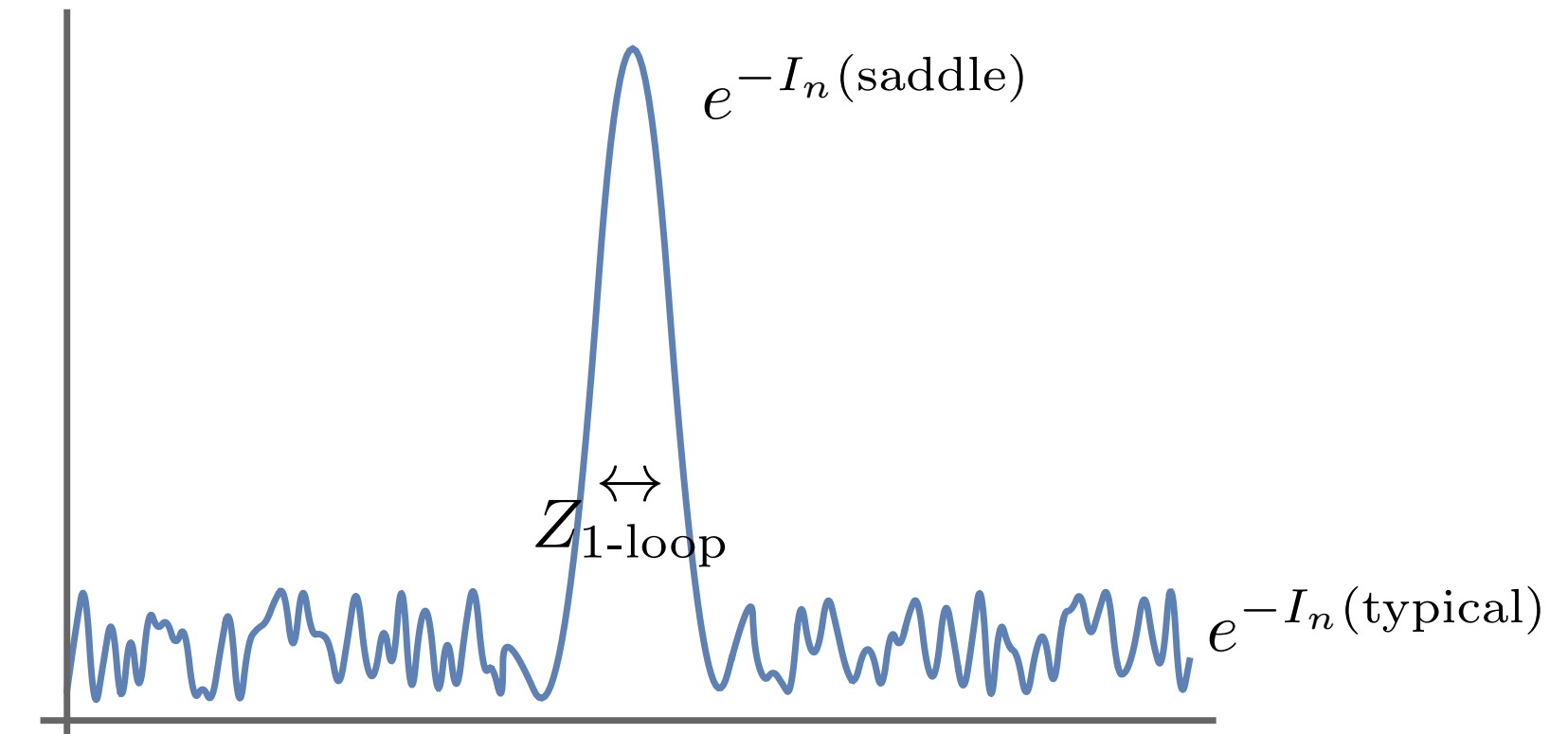
- the value of $I_n(\text{saddle})$ itself and that it is a minimum
- $Z_{1\text{-loop}}$: the order of fluctuations about the saddle point
- $I_n(\text{typical})$: value of $\log d_{\mu_V^*}$ for a typical configuration

*Haar averaged
integrals*

Can be checked, e.g. in the fuzzy sphere state

more serious fine print...

$$\mathrm{Tr} \rho_{\mathrm{in}}^n \approx Z_{1\text{-loop}} e^{-I_n(\text{saddle})} + e^{-I_n(\text{typical})}$$



- $U(N)$ is much larger than smooth continuum VPDs.
 - F contains many elements that act non-geometrically (send Σ to fractal and/or fragmented Planck-sized regions)
 - *These non-geometric maps proliferate and wash out the saddle-point*

We find it necessary to “coarse-grain” the integral over F to elements that act geometrically on Σ

Does not seem special to matrix character of the problem... may be generic

On coarse-grained diffeos

Return to the factorization map

$$|\psi\rangle = \int_F dV \int_{U(M)} dU \int_{U(N-M)} d\tilde{U} \hat{\pi}(U\tilde{U}V)|\psi_{\text{gf}}\rangle$$

- Integrals over $U(M)$, $U(N-M)$ $\longrightarrow$ reduced state at any V is *indistinguishable* against VPDs preserving a subregion.
- Integral over F $\longrightarrow$ different reduced states that are *distinguishable* for different V 's.
- The reduced state of a low-energy observer in any sector should not be orthogonal to one related by a Planck-sized diffeo $\longrightarrow$ remove from integral.

In practicality: integrate over a quotient of F

- Still large enough to account for geometric VPDs as $N \rightarrow \infty$, but suppressed enough that the geometric area dominates as a saddle.

E.g. in the fuzzy sphere state:

- A “low-tech” solution...

$$F = U(N') / (U(M') \otimes U(N' - M')) \qquad N' = N/p \quad M' = M/p$$

- By comparing $I_n(\text{saddle})$ to $Z_{1\text{-loop}}$ and to $I_n(\text{typical})$ we find a range of p such that the minimal area dominates.

$$N^{3/4} \ll p \ll N$$