

Better Alone

On the String Theory of a Single NS5-brane

Andrea Dei — University of Chicago

In collaboration with Emil Martinec — arXiv: 2506.22300

Decoupling Limit

Bulk

Boundary

N **D3**

$AdS_5 \times S^5$
String Theory

4d $\mathcal{N} = 4$ SYM

n_5 **NS5**

Asymptotically
Linear **D**ilaton
String Theory

6d **L**ittle **S**tring **T**heory

[Giveon, Kutasov '99]

Decoupling the NS5 branes

$$ds^2 = dx_{1,5}^2 + \left(1 + \frac{n_5 \alpha'}{r^2}\right) (dr^2 + r^2 d\Omega_3^2)$$

$$e^{2\Phi} = g_s^2 \left(1 + \frac{n_5 \alpha'}{r^2}\right)$$

Decouple the dynamics on the NS5 brane from the bulk by $g_s \rightarrow 0 \quad r \rightarrow 0$

$$ds^2 = dx_{5,1}^2 + n_5 \alpha' (d\rho^2 + d\Omega_3^2)$$

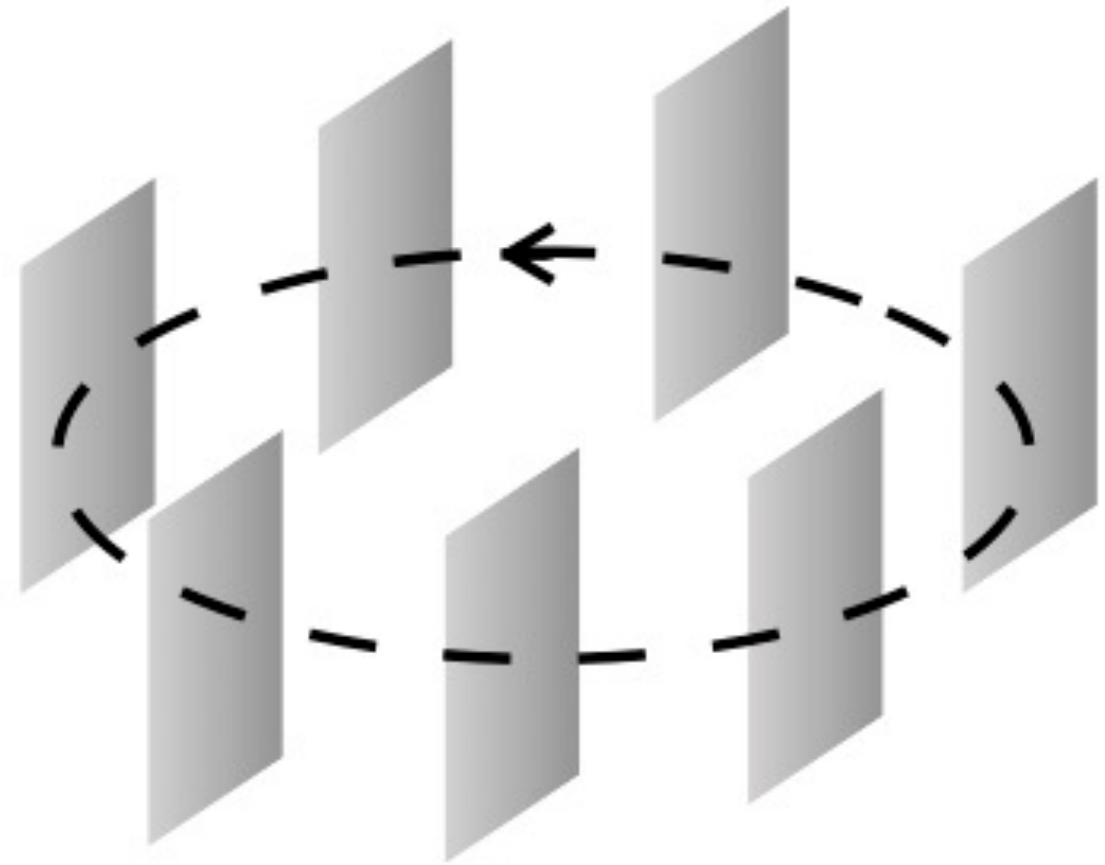
$$\Phi = -\rho$$

Worldsheet CFT

$$\mathbb{R}^{5,1} \times \mathbb{R}_\rho^{(1)} \times SU(2)_{n_5}^{(1)}$$

[Callan, Harvey, Strominger '91]

Separating the NS5 branes



$$g_s \rightarrow 0 \quad R \rightarrow 0 \quad \text{with} \quad \alpha' \text{ and } \frac{R}{g_s} \text{ fixed}$$

$$ds^2 = dx_{1,5}^2 + n_5 \alpha' (d\rho^2 + \tanh^2 \rho d\theta^2 + d\sigma^2 + \tan^2 \sigma d\phi^2)$$

$$\Phi = \Phi_0 - \log \cosh \rho$$



$$\uparrow R_\theta = \sqrt{n_5} \ell_s$$

[Giveon, Kutasov '99]

Worksheet CFT

The resulting worldsheet CFT is

$$\left(\frac{SL(2, \mathbb{R})_{n_5}^{(1)}}{U(1)} \times \frac{SU(2)_{n_5}^{(1)}}{U(1)} \right) / \mathbb{Z}_{n_5} \times \mathbb{R}^{5,1}$$

[Giveon, Kutasov '99]

An equivalent formulation is the null gauged WZW model

$$\frac{SL(2, \mathbb{R})_{n_5}^{(1)} \times SU(2)_{n_5}^{(1)}}{U(1) \times U(1)} \times \mathbb{R}^{5,1}$$

[Israel, Kounnas, Pakman, Troost '04; Martinec, Massai '17; Martinec, Massai, Turton '18-'25; AD, Martinec '24]

Little String Theory

- **NOT** a local Quantum Field Theory
- **Strongly Coupled**: No small parameter in which to organize a perturbative expansion of the boundary theory. String tension and brane separation are the only scales left.
- **Non-Abelian**: multiple NS5 branes

The Virtue of Solitude

For a single NS5 brane, $n_5 = 1$, one expects Little String Theory:

- **No longer strongly-coupled**: the size of the linear dilaton throat is comparable to the string length and the string cannot probe the strongly coupled region.
- **Abelian**: we only have one brane

The Goal

- Define the **worldsheet** string theory of a **single NS5 brane**
- Identify its **boundary dual**
- Solve an **exact** incarnation of Little String Theory **holography**

The Goal



Define the **worldsheet** string theory of a **single NS5 brane**



Identify its **boundary dual**



Solve an **exact** incarnation of Little String Theory **holography**

Hard to Be Alone

Naive Guess

$$\frac{SL(2, \mathbb{R})_{n_5}^{(1)} \times SU(2)_{n_5}^{(1)}}{U(1) \times U(1)} \times \mathbb{R}^{5,1}$$

Hard to Be Alone

Naive Guess

$$\frac{SL(2, \mathbb{R})_{n_5=1}^{(1)} \times SU(2)_{n_5=1}^{(1)}}{U(1) \times U(1)} \times \mathbb{R}^{5,1}$$

Problem

$$SU(2)_{n_5=1}^{(1)} = SU(2)_{n_5=-1} \oplus 3 \text{ free fermions} \quad \Rightarrow \quad \text{Not unitary!}$$

A similar puzzle for tensionless $AdS_3 \times S^3 \times \mathbb{T}^4$

At $n_5 \geq 2$

$$SL(2, \mathbb{R})_{n_5}^{(1)} \times SU(2)_{n_5}^{(1)} \longleftrightarrow PSU(1, 1 | 2)_{n_5}$$

Field content

- Bosonic $SL(2, \mathbb{R})_{n_5+2} \times SU(2)_{n_5-2}$
- 6 free fermions
- Bosonic $SL(2, \mathbb{R})_{n_5+2} \times SU(2)_{n_5-2}$
- 4 bc systems of dim (1,0)

A similar puzzle for tensionless $AdS_3 \times S^3 \times \mathbb{T}^4$

At $n_5 = 1$ it is then natural to consider the worldsheet theory

$$PSU(1,1 | 2)_1 \times \mathbb{T}^4$$

[Eberhardt, Gaberdiel, Gopakumar '18]

$PSU(1,1 | 2)_1$ admits a free field realization in terms of

- 1 $\beta\gamma$ system with dim (1,0)
- 2 bc systems with dim (1,0)

[Beem, Ferrari '23; AD, Knighton, Naderi '23]

Like a Bad Scrabble Rack

Refined Guess

$$\frac{PSU(1,1|2)_1}{U(1) \times U(1)} \times \mathbb{R}^{5,1}$$

The current we are going to gauge cannot be written in terms of the free fields generating $PSU(1,1|2)_1$

Our alphabet is too small

$n_5 \geq 2$ and $n_5 = 1$ comparison

$PSU(1,1|2)_{n_5}$ is generated by

$$n_5 \geq 2$$

- Bosonic $SL(2, \mathbb{R})_{n_5+2} \times SU(2)_{n_5-2}$
- 4 bc systems of dim (1,0)

$$n_5 = 1$$

- 1 $\beta\gamma$ system with dim (1,0)
- 2 bc systems with dim (1,0)

Our strategy

- We insist on **extending the map** between hybrid and RNS and find a free field realization for bosonic $SL(2, \mathbb{R})_3 \times SU(2)_{-1}$ in terms of 3 $\beta\gamma$ systems
- We obtain a new free field realization of $PSU(1,1 | 2)_1$ in terms of 3 $\beta\gamma$ **systems** and 4 bc **systems**
- **Not Unitary** and **Not equivalent** to the tensionless string
- **Extra constraint** relating the 4 bc systems restores unitarity and equivalence with the tensionless string

A larger alphabet

We overcome the obstacle in three steps:

1. **Alternative free field realization** of $PSU(1,1 | 2)_1$ and construct a map to RNS variables
2. **Additional constraint** restores unitarity
3. **Complete equivalence** with the original formulation of the tensionless string

The Single Brane String Theory Unlocked

We provide a **worldsheet formulation** of the single NS5 brane string theory in terms of a **free field worldsheet CFT** subject to **three** mutually compatible **BRST constraints**:

- Familiar String BRST operator
- U(1) gauging BRST operator
- Unitarity BRST operator unique to $n_5 = 1$

Future directions

- **Dual 5+1d abelian gauge theory**: how does the holographic map work in detail?
- **Solve the BRST cohomology**: what is the spacetime symmetry algebra?
- **Correlation functions**: Localizations? Free theory?
- **String theory dual of $T\bar{T}$ deformation**: a variant of the gauging embedding describes the string dual of the $T\bar{T}$ -deformed symmetric product orbifold