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## Hydrodynamic simulations as a test bed for turbulent convection models

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*aim: comparison of 1D turbulent convection model to 3D hydrodynamic* 

*simulations*

## Turbulent convection models I

three equations for three convective variables (Kuhfuß 1986,1987, Flaskamp 2003)

$$
\omega = \frac{1}{2} \overline{u'^2}, \ \Pi = \overline{u's'}, \ \Phi = \frac{1}{2} \overline{s'^2} \qquad \qquad T = \text{ temperature}
$$
\n
$$
d_t \omega = \frac{\nabla_{\text{ad}} T}{H_P} \Pi \qquad - \qquad \varepsilon \qquad - \qquad \frac{1}{\rho} \text{div}(-D_{\omega} \nabla \omega) \qquad \qquad \frac{H_P}{\rho} = \text{ pressure scale height}
$$
\nbuoyant driving dissipation non-local flux

\nO Kolmogorov cascade

\n
$$
\varepsilon = C_D \frac{\omega^{3/2}}{\Lambda} \qquad \text{in overshoothing zone:} \quad \Lambda \propto \frac{1}{\text{buoyancy freq.}} \propto \frac{1}{\nabla - \nabla_{\text{ad}}}
$$

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Kupka, Ahlborn and Weiss (2022), Ahlborn et al. (2022).

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## Turbulent convection models II

different flavours

1-equation model: mixing length approximation for convective flux

$$
\Pi~= \alpha_s \Lambda \sqrt{\omega} \frac{c_p}{H_p} (\nabla - \nabla_{\rm ad}) \propto - \Lambda \sqrt{\omega} \frac{\partial s}{\partial r}
$$

3-equation model: two additional equations

$$
\mathrm{d}_t\Pi=\frac{2\nabla_\mathrm{ad}T}{H_P}\Phi+\frac{2}{3}\frac{c_p}{H_p}(\nabla-\nabla_\mathrm{ad})\omega-\frac{1}{\rho}\mathrm{div}(-D_\Pi\nabla\Pi)-\frac{\Pi}{\tau_\mathrm{rad}}\\ \mathrm{d}_t\Phi=\frac{c_p}{H_p}(\nabla-\nabla_\mathrm{ad})\Pi-\frac{1}{\rho}\mathrm{div}(-D_\Phi\nabla\Phi)-\frac{2\Phi}{\tau_\mathrm{rad}}
$$

implemented into the Garching stellar evolution code (Weiss and Schlattl 2008)



### Hydrodynamic simulations

- Seven Leagues Hydro (SLH) (e.g. Miczek 2015, Edelmann et al. 2021)
- three-dimensional wedge geometry  $O$  384 x 96 x 96
- nominal luminosity
- $\sim$  10000 h simulation time
- **Reynolds Averaged Navier Stokes** (RANS) analysis

Ahlborn and Higl in prep. <sup>7</sup>

## Initial stellar models

- $13\mathrm{M}_\odot$  stellar model as initial model  $\circ$  beginning of main-sequence two 1D stellar models for comparison ○ MLT ○ Kuhfuss 3-equation different final states
	- $\circlearrowright$  thermal timescale too long
	- which one is correct?  $\circlearrowright$  probably none



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## Turbulent kinetic energy

- compare turbulent convection model and MLT with RANS data
- good agreement between TCM and RANS



## Flow anisotropy

- ratio of radial to total kinetic energy
- contradicts assumption of isotropic flow
- 4th dynamic equation? (e.g. Xiong or Canuto)



## Non-local term

1D non-local term:  $\frac{1}{2} - \frac{1}{2} \text{div}(-D_{\omega} \nabla \omega)$ 

agreement in terms of shape



#### Pressure fluctuation term

#### ● non-zero

- neglected in the Kuhfuss model
- reminiscent of the non-local term
- solutions from literature (e.g. Rotta 1951, Canuto 1992, Canuto 1993, Sander 1998 and references therein)

#### Ahlborn and Higl in prep.



## Deardorff layer

- infer thermal structure from entropy profile  $-\frac{\partial \overline{s}}{\partial r}=\frac{c_p}{H_n}(\nabla-\nabla_{\rm ad})$
- subadiabatic region with positive convective flux
- convection driven by non-local effects

## Applications

- convective core on the main sequence
	- Ahlborn, Kupka, Weiss and Flaskamp (2022)
	- $\circ$  Kupka, Ahlborn and Weiss (2022)
- standard solar model
	- $\circ$  Braun, Ahlborn and Weiss (2024) (T. Braun will be here in the 3rd and 4th week)
- Cepheid mass discrepancy problem
	- $\bigcirc$  Deka, Ahlborn, Braun and Weiss in prep.

## Conclusions

- self-consistent convective boundary layers using a turbulent convection model
	- Kuhfuss model (Kuhfuss 1986, 1987)
	- $\circ$  implemented in the Garching Stellar evolution code (GARSTEC)
- hydrodynamic simulations of a  $3M_{\odot}$  star
	- nominal luminosity
	- compute 1D averages

#### confirming

- $\circlearrowright$  turbulent kinetic energy equation
- Deardorff layer

## disagreeing final stratification pressure fluctuation terms

○ isotropy