

Heidelberg Institute for Theoretical Studies

Think beyond the limits!







Hydrodynamic simulations as a test bed for turbulent convection models

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aim: comparison of 1D turbulent convection model to 3D hydrodynamic

simulations

Turbulent convection models I

three equations for three convective variables (Kuhfuß 1986,1987, Flaskamp 2003)

$$\begin{split} \omega &= \frac{1}{2} \overline{u'^2}, \ \Pi = \overline{u's'}, \ \Phi = \frac{1}{2} \overline{s'^2} & T = \text{temperature} \\ \Pi &= \text{convective flux} \\ \text{d}_t \omega &= \frac{\nabla_{\text{ad}} T}{H_P} \Pi &- \varepsilon &- \frac{1}{\rho} \text{div} (-D_\omega \nabla \omega) & \rho = \text{density} \\ \text{buoyant driving} & \text{dissipation} & \text{non-local flux} \\ \end{split}$$

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dis

Kupka, Ahlborn and Weiss (2022), Ahlborn et al. (2022).

5

Turbulent convection models II

different flavours

1-equation model: mixing length approximation for convective flux

$$\Pi ~= lpha_s \Lambda \sqrt{\omega} rac{c_p}{H_p} (
abla -
abla_{
m ad}) \propto -\Lambda \sqrt{\omega} rac{\partial s}{\partial r}$$

3-equation model: two additional equations

$$egin{aligned} &\mathrm{d}_t \Pi = rac{2
abla_{\mathrm{ad}} T}{H_P} \Phi + rac{2}{3} rac{c_p}{H_p} (
abla -
abla_{\mathrm{ad}}) \omega - rac{1}{
ho} \mathrm{div}(-D_\Pi
abla \Pi) - rac{\Pi}{ au_{\mathrm{rad}}} \ &\mathrm{d}_t \Phi \ &= & rac{c_p}{H_p} (
abla -
abla_{\mathrm{ad}}) \Pi - rac{1}{
ho} \mathrm{div}(-D_\Phi
abla \Phi) - rac{2 \Phi}{ au_{\mathrm{rad}}} \end{aligned}$$

implemented into the Garching stellar evolution code (Weiss and Schlattl 2008)



Hydrodynamic simulations

- Seven Leagues Hydro (SLH) (e.g. Miczek 2015, Edelmann et al. 2021)
- three-dimensional wedge geometry
 384 x 96 x 96
- nominal luminosity
- ~10000 h simulation time
- Reynolds Averaged Navier Stokes (RANS) analysis

Initial stellar models

^{3M}_☉ stellar model as initial model
 ○ beginning of main-sequence
 • two 1D stellar models for

comparison

O MLT

- Kuhfuss 3-equation
- different final states

 thermal timescale too long
 which one is correct?
 probably none



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Turbulent kinetic energy

- compare turbulent convection model and MLT with RANS data
- good agreement between TCM and RANS



Flow anisotropy

- ratio of radial to total kinetic energy
- contradicts assumption of isotropic flow
- 4th dynamic equation? (e.g. Xiong or Canuto)



Non-local term

• 1D non-local term: $-\frac{1}{
ho} \operatorname{div}(-D_{\omega} \nabla \omega)$

agreement in terms of shape



Pressure fluctuation term

non-zero

- neglected in the Kuhfuss model
- reminiscent of the non-local term
- Solutions from literature (e.g. Rotta 1951, Canuto 1992, Canuto 1993, Sander 1998 and references therein)

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Deardorff layer

• infer thermal structure from entropy profile $\partial \overline{s} = c_p (\nabla \nabla \nabla \nabla)$

$$-rac{\partial s}{\partial r}=rac{c_p}{H_p}(
abla-
abla_{
m ad})$$

- subadiabatic region with positive convective flux
- convection driven by non-local effects

Applications

- convective core on the main sequence
 - Ahlborn, Kupka, Weiss and Flaskamp (2022)
 - Kupka, Ahlborn and Weiss (2022)
- standard solar model
 - Braun, Ahlborn and Weiss (2024) (T. Braun will be here in the 3rd and 4th week)
- Cepheid mass discrepancy problem
 - Deka, Ahlborn, Braun and Weiss in prep.

Conclusions

- self-consistent convective boundary layers using a turbulent convection model
 - O Kuhfuss model (Kuhfuss 1986, 1987)
 - implemented in the Garching Stellar evolution code (GARSTEC)
- hydrodynamic simulations of a $3M_{\odot}$ star
 - nominal luminosity
 - compute 1D averages

confirming

- turbulent kinetic energy equation
- Deardorff layer

disagreeing

- final stratification
- pressure fluctuation terms
- isotropy