



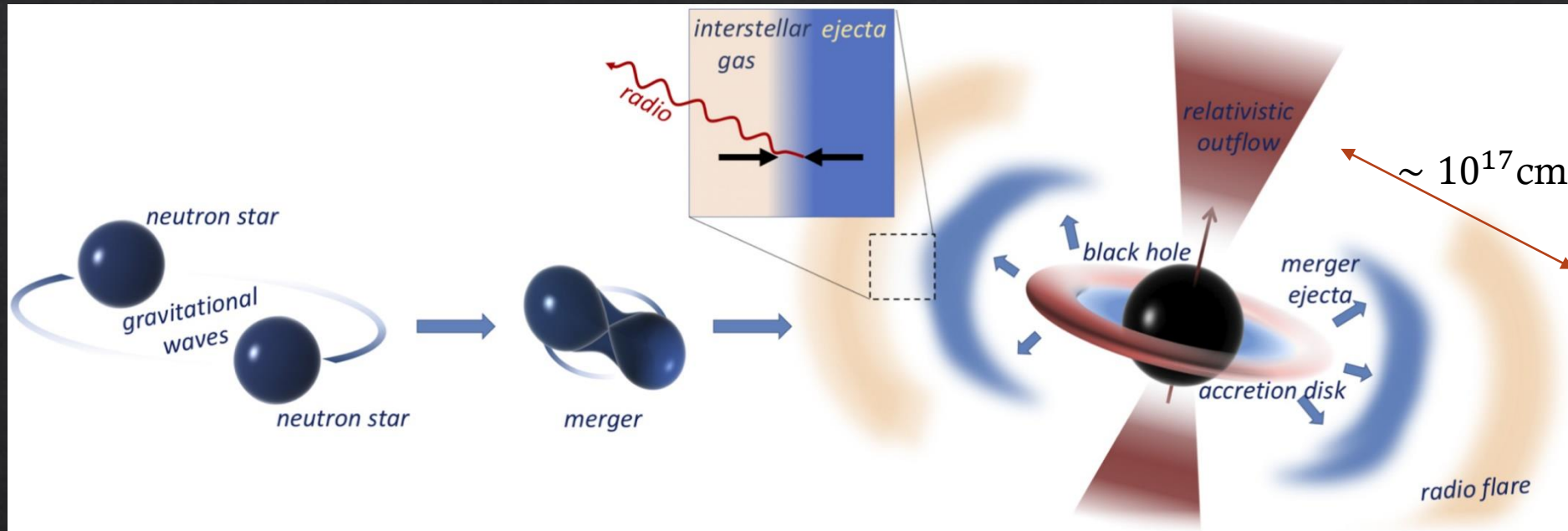
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Non-thermal emission from mildly relativistic dynamical ejecta of BNS merger

Overview




Credit: AAS Nova

- ◇ BNS mergers are expected to eject mass dynamically at mildly relativistic velocities, $\beta > 0.6$
- ◇ A collisionless shock driven by this ejecta will accelerate electrons and produce synchrotron radiation
- ◇ Previous models are based on extrapolation of results valid for $\gamma\beta \gg 1$ & $\gamma\beta \ll 1$ to $\gamma\beta \sim 1$

Why Kilonova ejecta?

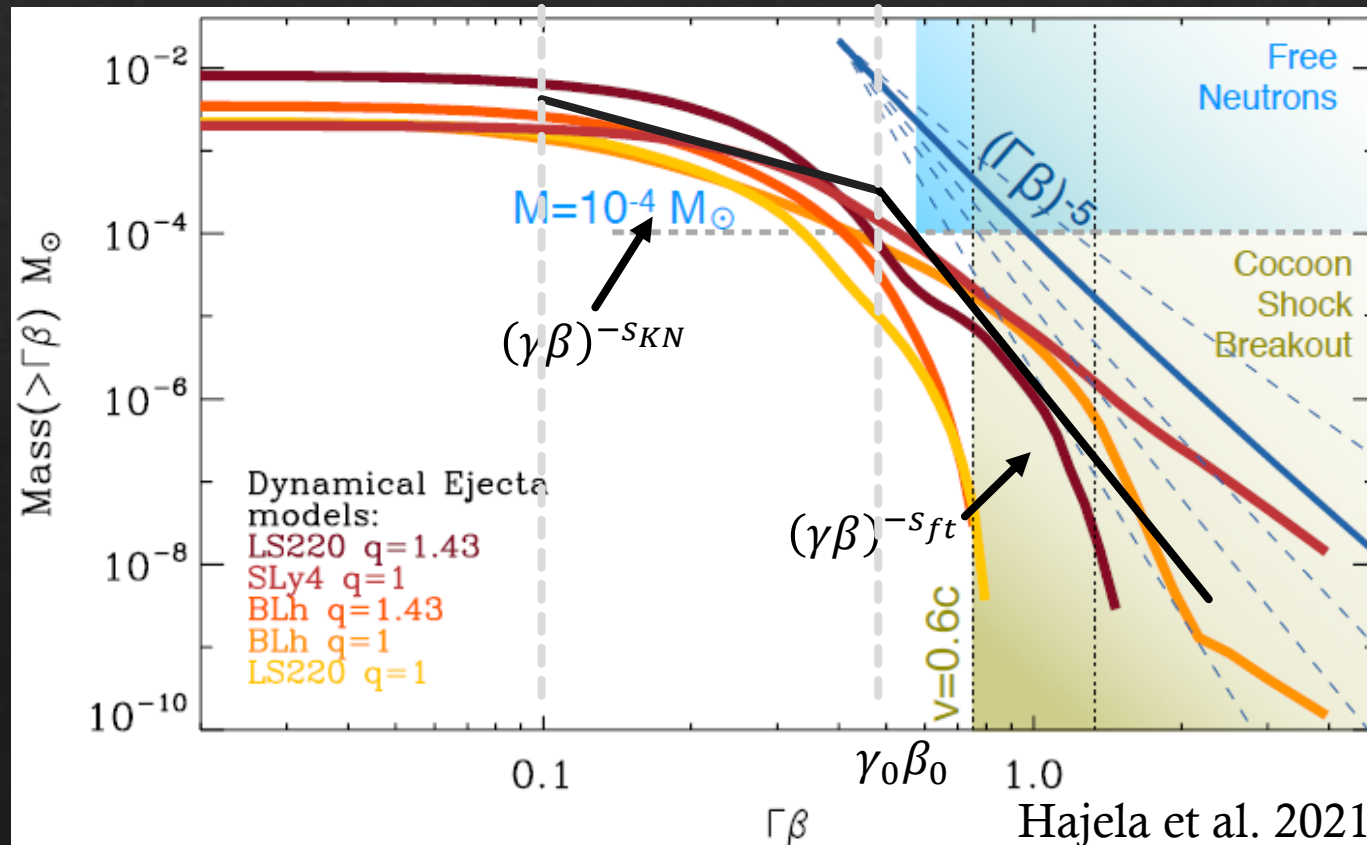
| Properties | Kilonova ejecta | Relativistic jet |
|------------|-------------------------------|-----------------------------------|
| Mass | $10^{-2} - 10^{-3} M_{\odot}$ | $\sim 10^{-8} M_{\odot}$ |
| Velocity | $0.1 < \gamma\beta < 4$ | $\gamma \gg 10$ |
| Geometry | Quasi-spherical | Cone, opening angle of ~ 0.1 |

- ◇ Large numerical uncertainty, especially for $\gamma\beta > 1$
- ◇ Dominate emission at late time
- ◇ Constrain the properties of Kilonova ejecta  binary system parameters + EoS + merger dynamics

Published work: spherical ejecta

$$\diamond M(> \gamma\beta) = M_0 \begin{cases} \left(\frac{\gamma\beta}{\gamma_0\beta_0}\right)^{-s_{ft}}, & \gamma_0\beta_0 < \gamma\beta \\ \left(\frac{\gamma\beta}{\gamma_0\beta_0}\right)^{-s_{KN}}, & 0.1 < \gamma\beta < \gamma_0\beta_0 \end{cases} \quad \begin{matrix} s_{ft} > 5, \\ 1 < s_{KN} < 3, \\ 0.3 < \gamma_0\beta_0 < 1 \end{matrix}$$

Ejected mass: BNS merger simulation



Non-thermal emission

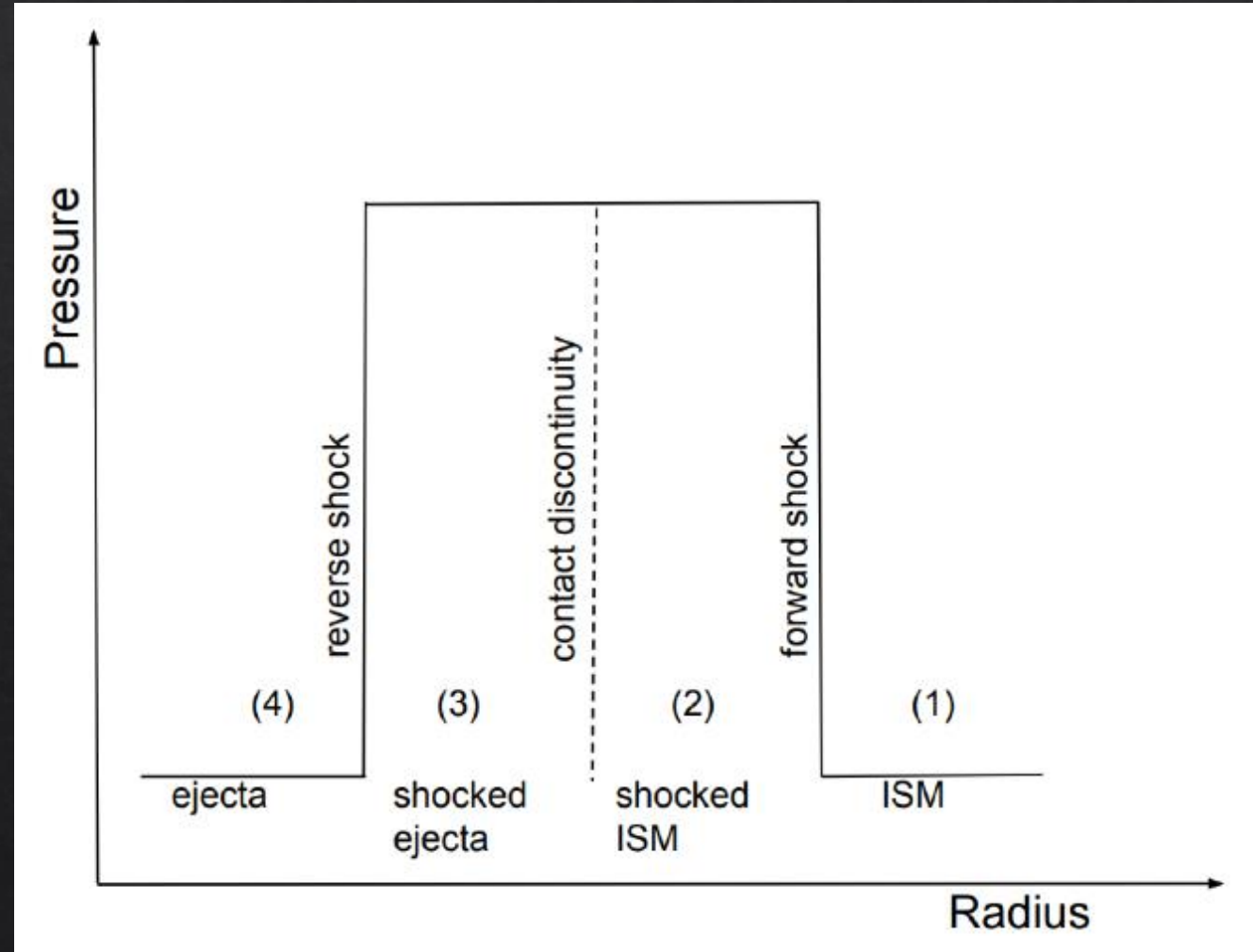
◆ Assumptions:

- Fractions ε_e and ε_B of the post shock internal energy carried by non-thermal electrons and magnetic fields.

- Electrons follow energy power-law distribution $\frac{dn_e}{d\gamma_e} \propto \gamma_e^{-p}$

◆ Two competing effects:

- Increase of the shocked mass
- Deceleration



Solution methods

- ◇ Analytic, based on an approximate solution for the dynamics
 - Including emission from forward shock only
- ◇ Numeric, relativistic hydrodynamic calculation of the dynamics
 - Integrating over the emissivity, properly accounting for relativistic effects and radiation arrival time
 - Including both forward & reverse shock emission

The analytic solution provides an excellent approximation to the numeric results

Key analytic results

Peak flux obtained when the reverse shock crosses the fast tail ($s_{ft} > 5$), reaching the “break” in the ejecta profile at β_0

$$t < t_{peak}$$

Reverse shock in the steep ejecta ($s_{ft} > 5, \beta > \beta_0$)
Rising flux $F_\nu \propto t^{q_{ft}}$

$$t_{peak} < t < t_{ST}$$

Reverse shock in the moderate ejecta
($1 < s_{KN} < 3, \beta < \beta_0$)
Deceleration is faster for a shallow mass distribution
Declining flux $F_\nu \propto t^{q_{KN}}$

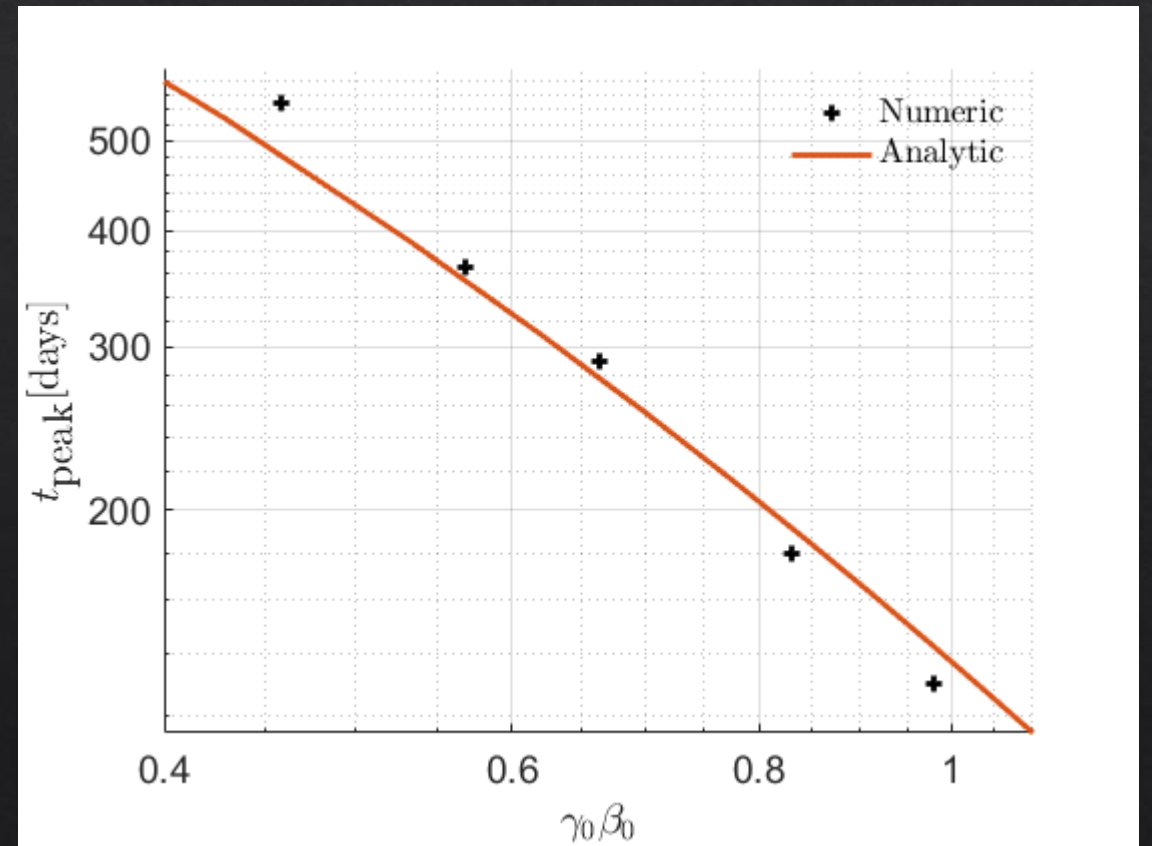
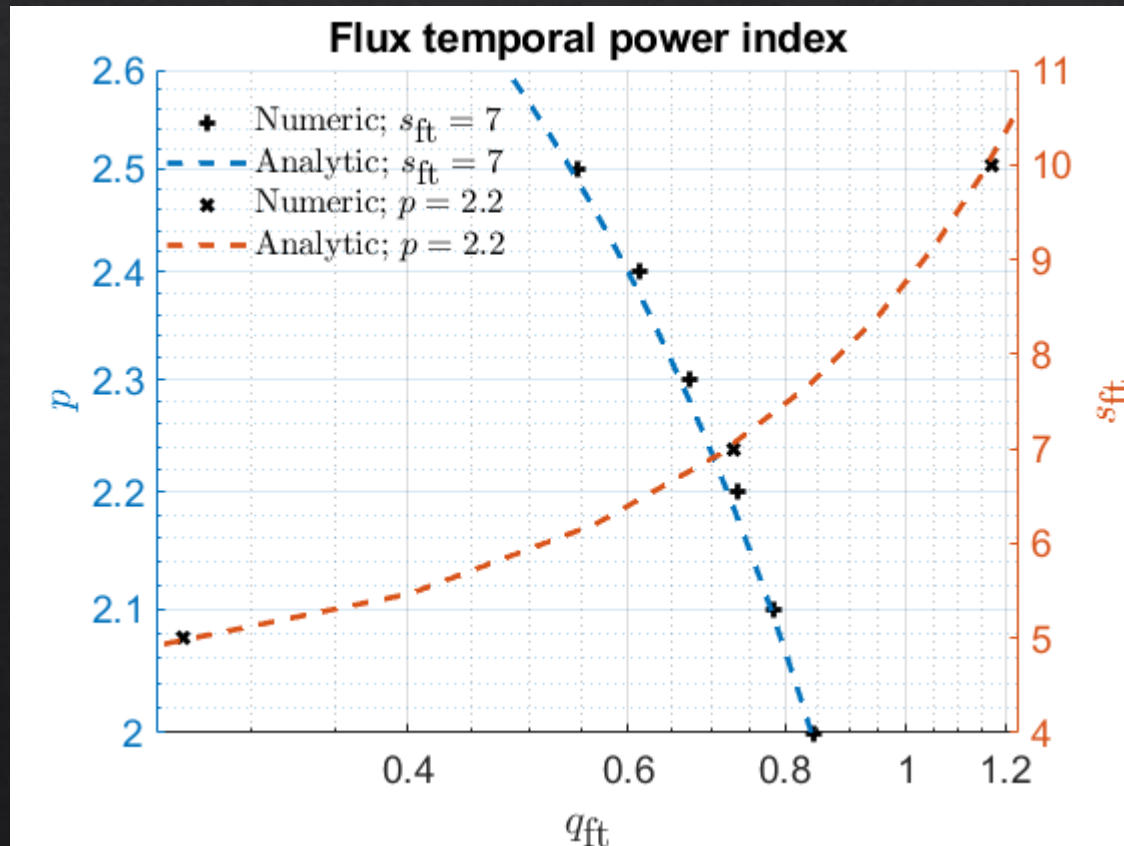
$$t_{ST} < t$$

Non-relativistic
Sedov-Taylor flow, $F_\nu \propto t^{q_{ST}}$

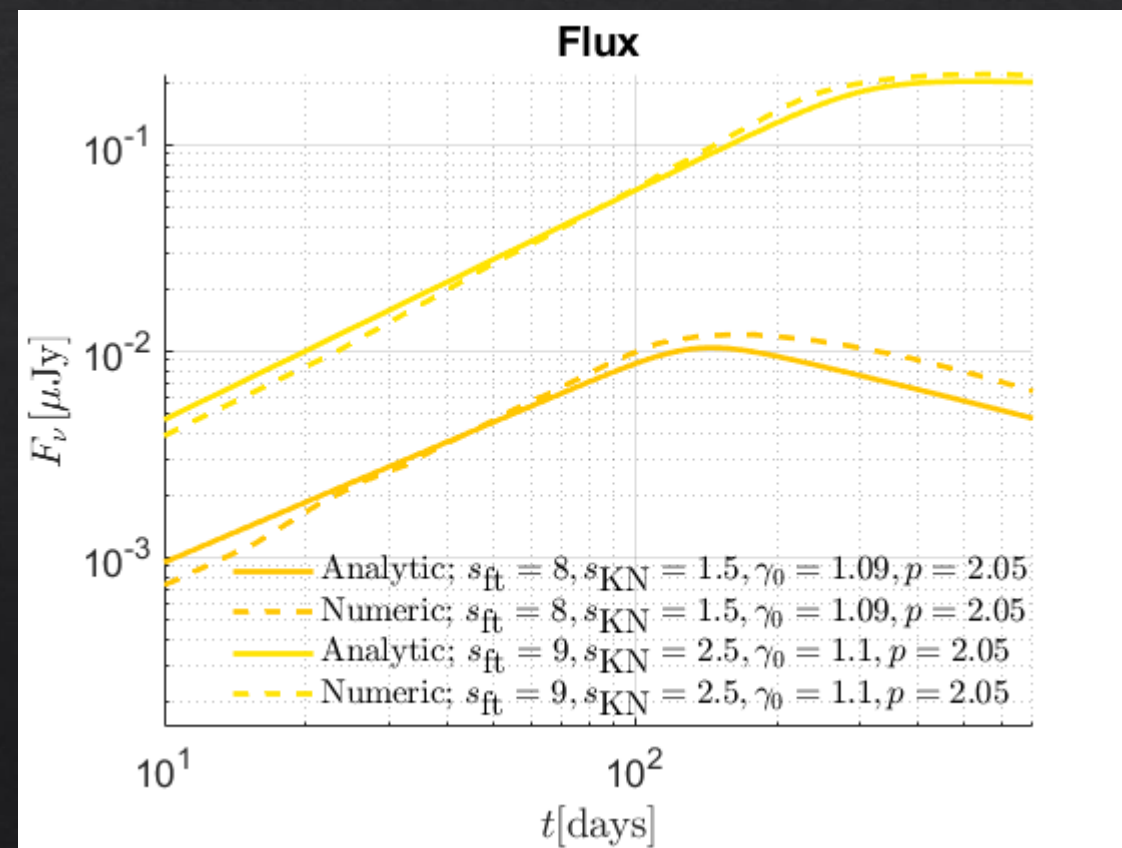
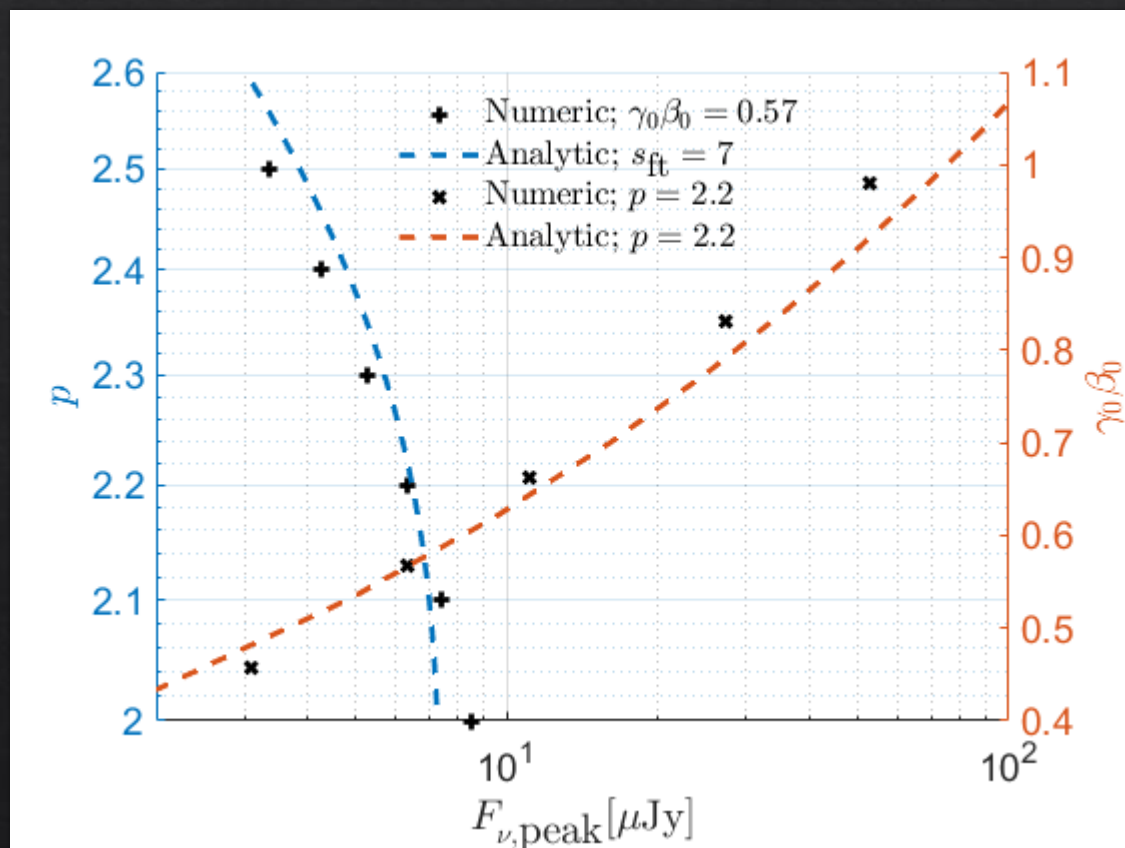
$$\diamond F_{\nu, peak} \approx 12 D_{26.5}^{-2} \varepsilon_{e,-1}^{p-1} \varepsilon_{B,-2}^{\frac{p+1}{4}} n_{-2}^{\frac{p+1}{4}} M_{0,-4} \nu_{9.5}^{\frac{1-p}{2}} (\gamma_0 \beta_0)^{2.2p-0.5} \mu\text{Jy} \quad (\nu < \nu_c)$$

$$\diamond t_{peak} = 550 \left(\frac{M_{0,-4}}{n_{-2}} \right)^{\frac{1}{3}} g(\beta_0) \text{days}, \quad g(\beta_0) = \frac{1.5 - \sqrt{0.25 + 2\beta_0^2}}{\gamma_0^{\frac{1}{3}} \beta_0}$$

Excellent agreement with full numeric results



Excellent agreement with full numeric results

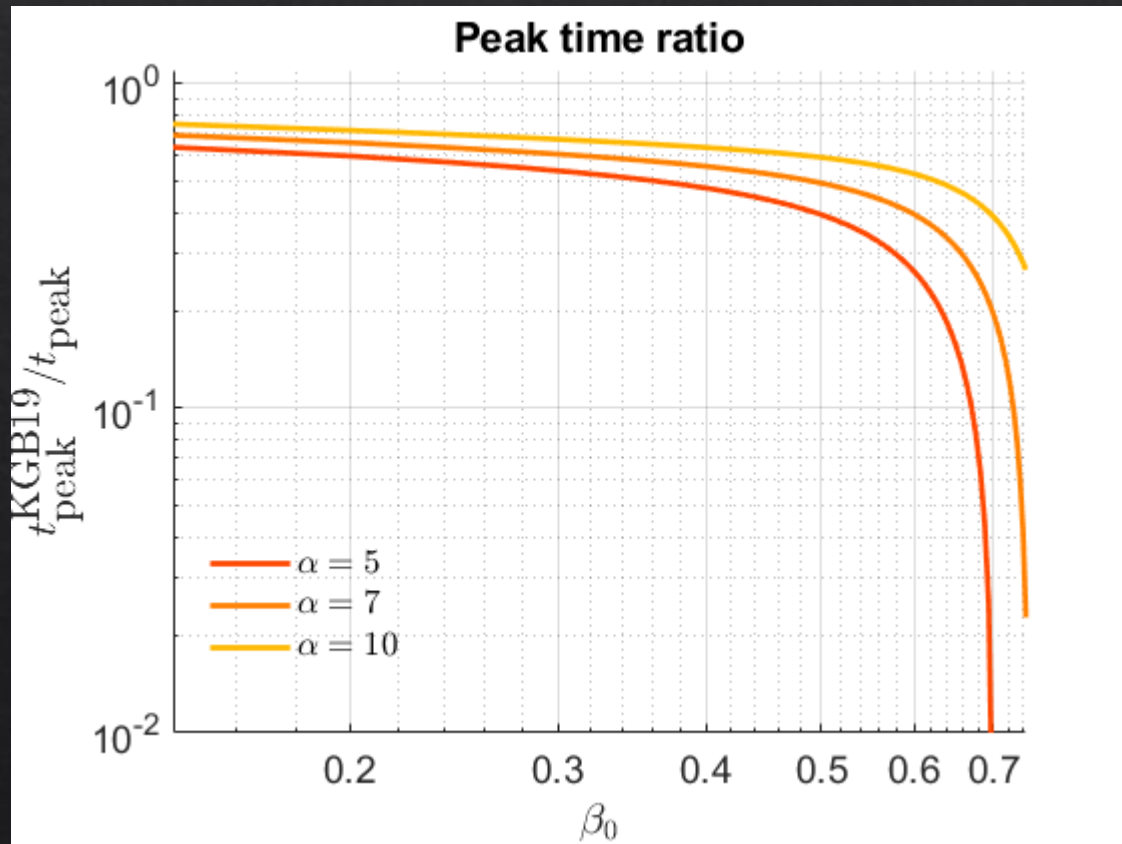




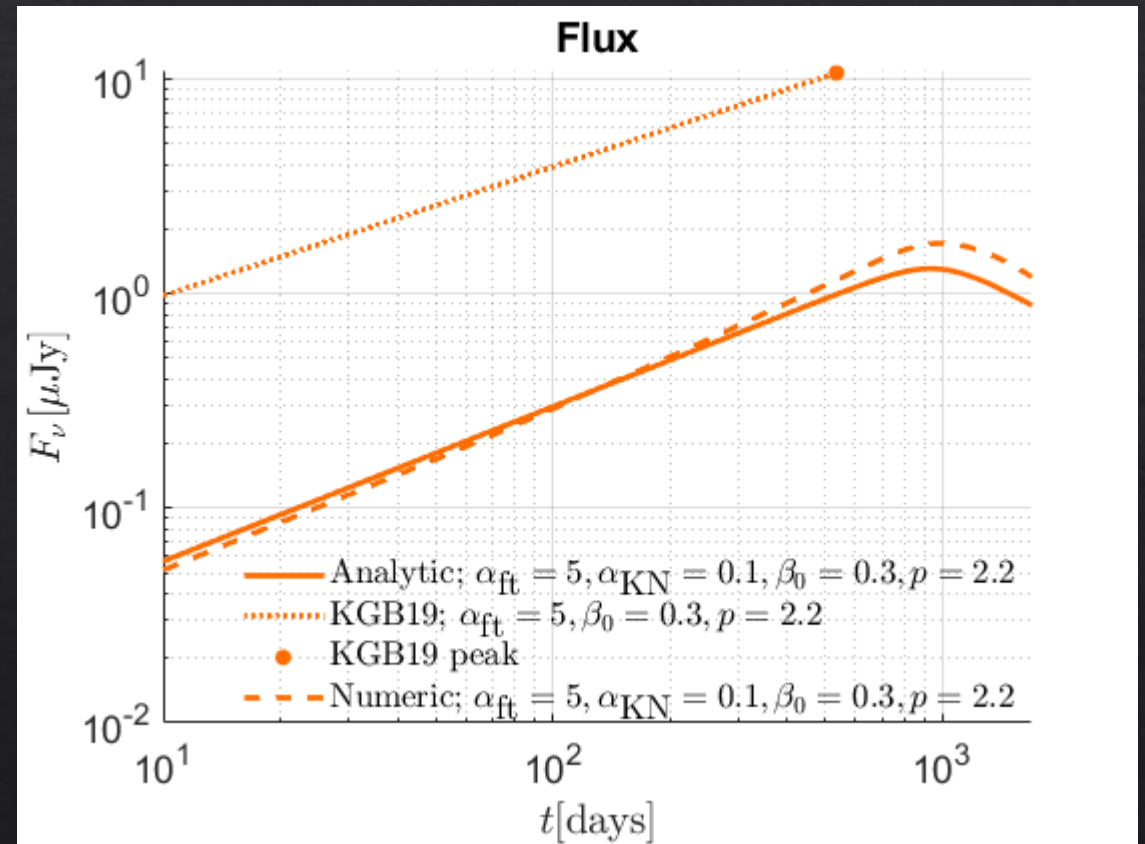
Comparison with earlier work

- ◇ Earlier analytic work by Kathigamaraju et al (2019), based on Nakar and Piran (2011)
 - A single power-law for $t < t_{peak}$, not including the deceleration to $\beta_0 \ll 1$
 - Valid for $\beta_0 < 0.5$
- ◇ The peak flux obtained in their work is typically an order of magnitude larger than our result, $\frac{F_{\nu,peak}}{F_{\nu,peak}^{our}} \approx 6\gamma_0^{2.5-2.2p} \beta_0^{0.3p-1}$

Comparison with earlier work



As expected, large deviations for $\beta_0 > 0.5$



$\beta_0 = 0.3$

GW170817

- ◇ The early ($t < 1000$ days) non-thermal emission following GW170817 is consistent with a jet
- ◇ A late ($t > 1000$ days) non-thermal emission may indicate a new component, due to a “fast tail” of the KN ejecta
- ◇ Using GW170817 ejecta parameters inferred from the radioactive emission for $\beta < 0.3$,

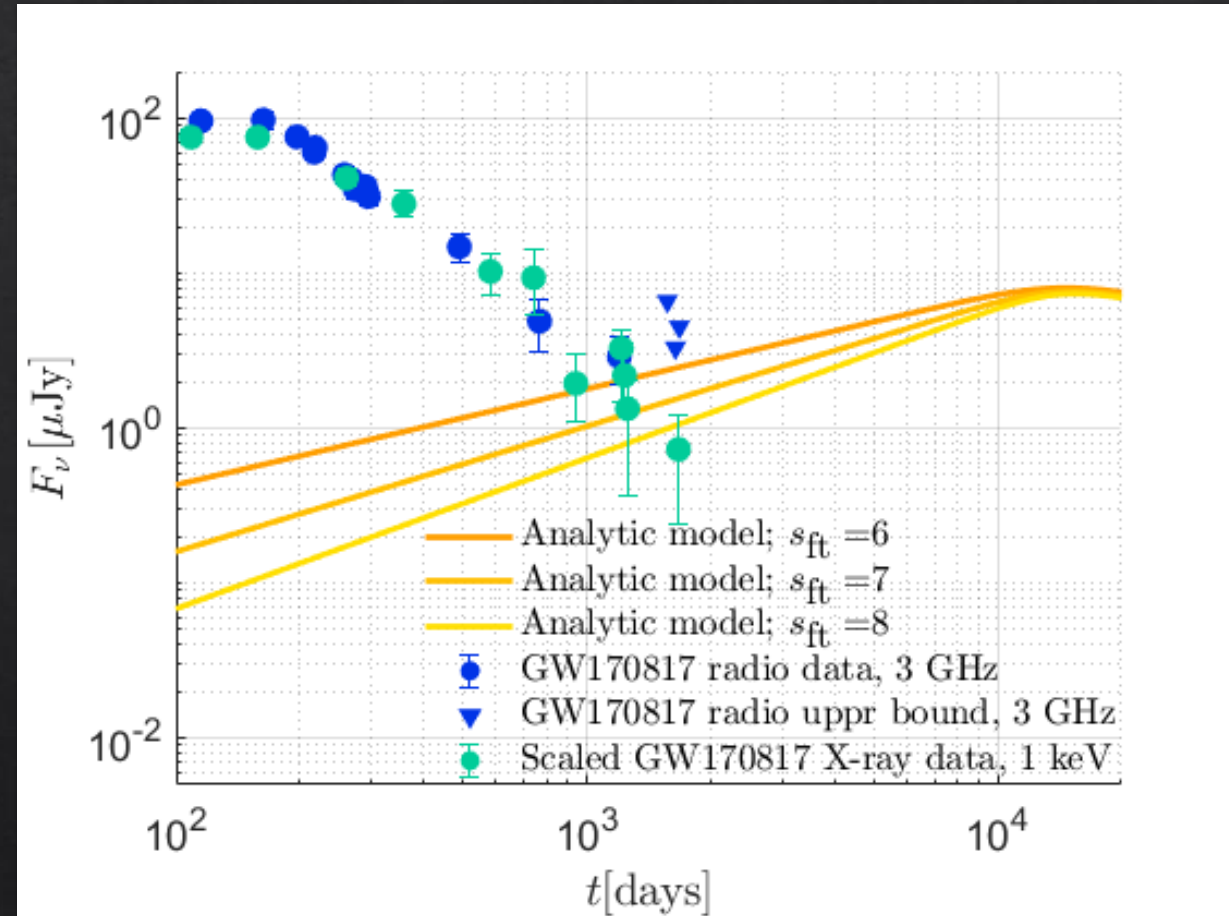
$$M = 0.05M_{\odot}, \beta_0 = 0.3, s_{KN} = 1.6$$

[Waxman et al. 2018)]

We infer from the late time observations

$$\varepsilon_{B,-2}^{-1.5} n_{-2}^{-1.17} > 6.7,$$

$$\varepsilon_{e,-1}^{1.05} \varepsilon_{B,-2}^{0.76} n_{-2}^{1.03} < 0.12$$



Summary

- ◇ We derived an analytic model for the non-thermal emission from a mildly relativistic ejecta with a broken power-law mass distribution
- ◇ The model provides an excellent approximation to the results of numeric calculations for a wide range of ejecta parameter values, that represent the mass profiles produced in merger simulations
- ◇ The analytic model
 - Expands the validity range of existing models, $\gamma\beta \gg 1, \gamma\beta \ll 1$, to $\gamma\beta \sim 1$
 - Includes a description of the phase of deceleration to sub-relativistic expansion
 - Is accurate to 10's of percent, a significant improvement over previous work, which overestimate the flux by \sim an order of magnitude for relevant parameter values
- ◇ An analytic description is highly valuable for the “inverse problem”, inferring model parameters from observations: It does not require resource consuming numeric calculations for the derivation of model predictions for each point in the parameter space.



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**Thank
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