

# Primordial black holes as dark matter: the role of stochastic effects

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based on 2012.06551, 2111.07437, 2312.XXXXX with  
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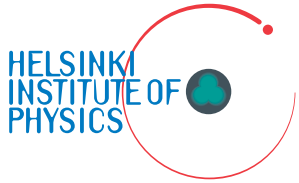
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# Dark matter without dark matter

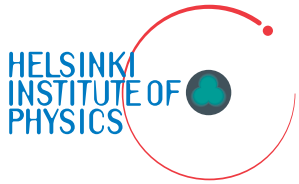
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- Primordial black holes (PBH) are a candidate for dark matter. (Hawking 1971)
- They could also seed supermassive black holes in centers of galaxies.
- Let's consider asteroid-mass PBHs as dark matter.



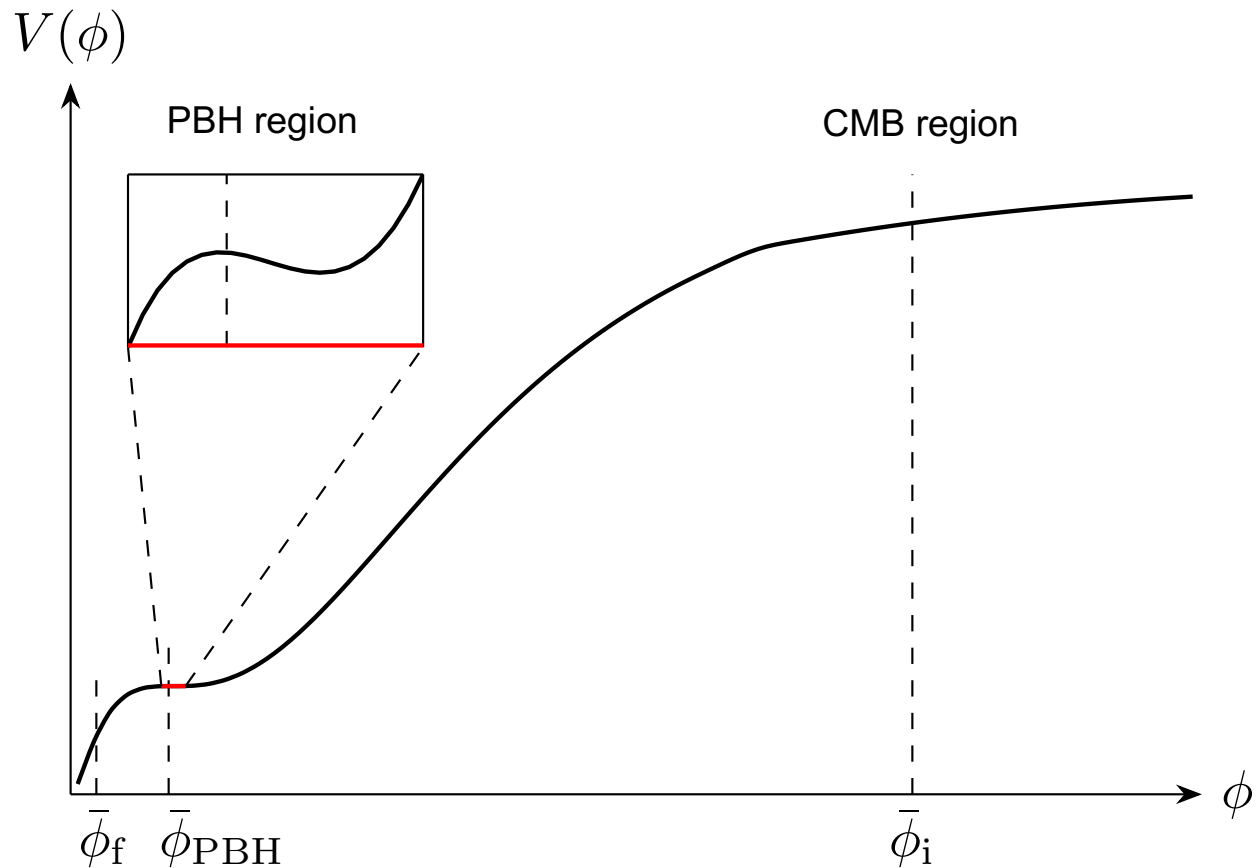
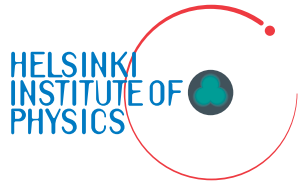
# Facilitating collapse



- PBHs require large ( $\sim 1$ ) fluctuations on small scales.
- They could be generated by the same process as the small ( $\sim 10^{-5}$ ) fluctuations on large scales.
- Most successful scenario is inflation.
- The curvature perturbation (in single field slow-roll) is
$$\zeta = -H\delta\phi/\dot{\phi} \sim H^2/\dot{\phi}$$
.
- The slower the field, the larger the perturbations.



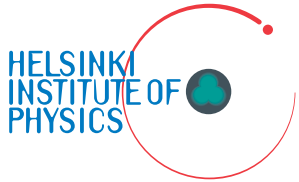
# Potential for asteroid-mass PBHs





# Making black holes

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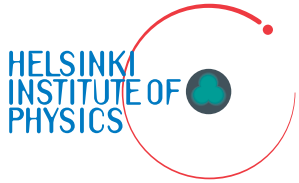


- During inflation,  $k$ -modes stretch to super-Hubble scales and freeze.
- After inflation, they cross back inside the Hubble radius and start evolving.
- If a Hubble patch is overmassive enough, it collapses into a PBH.
  - PBH mass is close to the mass inside the Hubble patch.



# Slow therefore stochastic

$$\ddot{\phi} + 3H\dot{\phi} = -V' + \xi$$

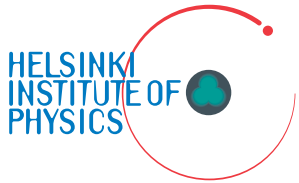


- Because the field moves slowly, stochastic effects are important. (Large kicks, small force.)
- Stochastic effects increase PBH production:
  1. Patches with large fluctuations are more likely: distribution tail is exponential, not Gaussian.
  2. Individual patches are choppy.



# All inflation is stochastic inflation

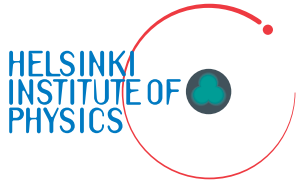
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- Inflaton evolution is stochastic when the coupling of small- and long-wavelength modes is taken into account. (Starobinsky 1986)
- As modes become super-Hubble and classicalise, they change the background in which shorter modes evolve.
- Amplitude of every  $\mathbf{k}$ -mode is independently drawn from a Gaussian distribution, so the background is subject to Gaussian white noise.



# One step at a time

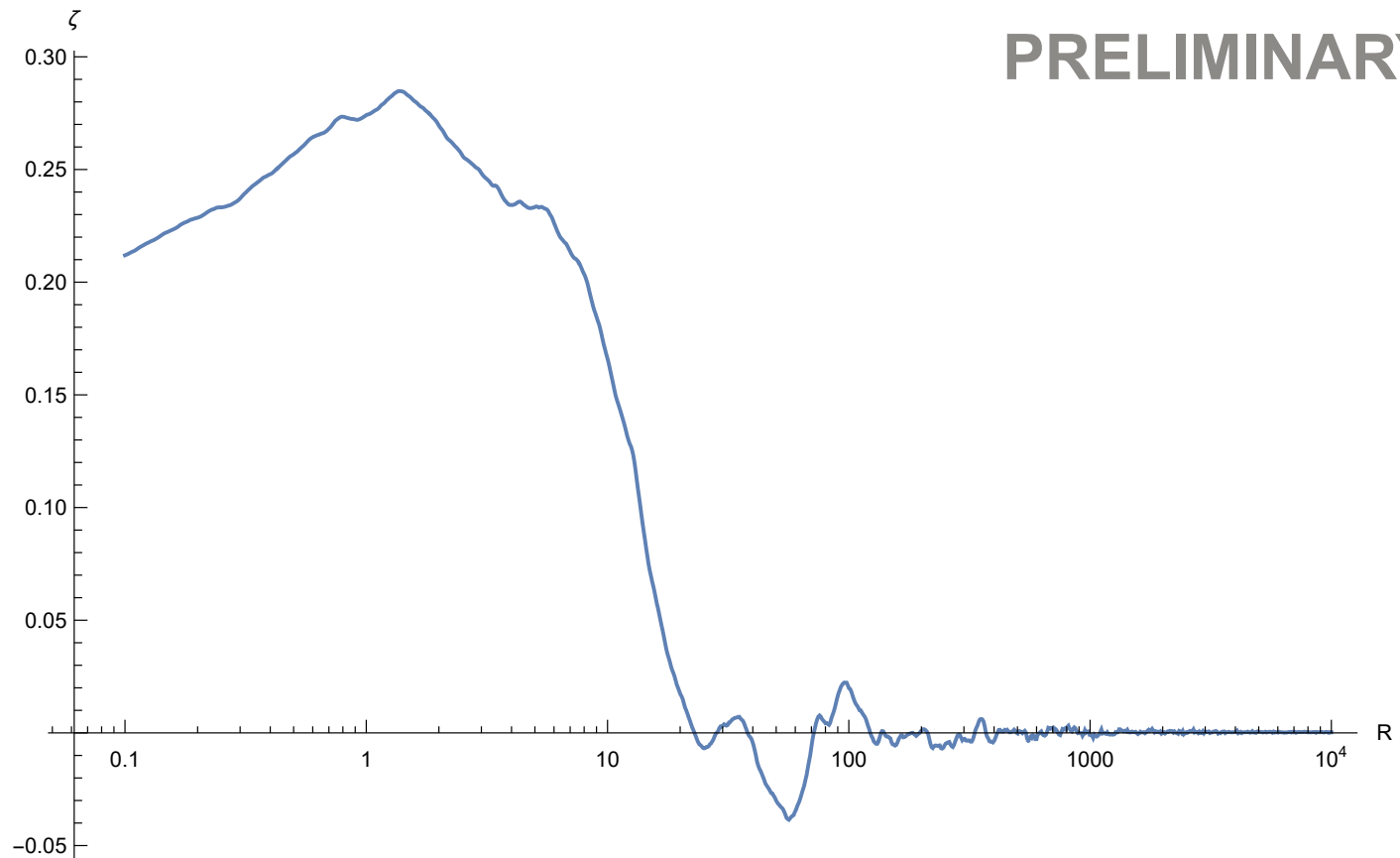
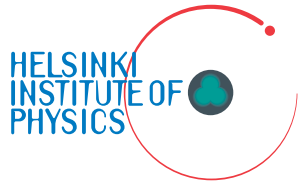


- We solve the full coupled background + perturbation evolution numerically:
  - Coarse-grain a mode, recalculate the background and perturbations at every timestep.
  - Continue until inflation ends, record the total number of e-folds  $N$ .
  - Repeat  $10^{11}$  times to gather statistics.
  - Find distribution  $P(\Delta N)$ , where  $\Delta N = N - \bar{N} = \zeta$ .
- We also reconstruct the profile  $\zeta(r)$  inside each patch, using a simplified treatment. (Tomberg: 2210.17441, 2304.10903)



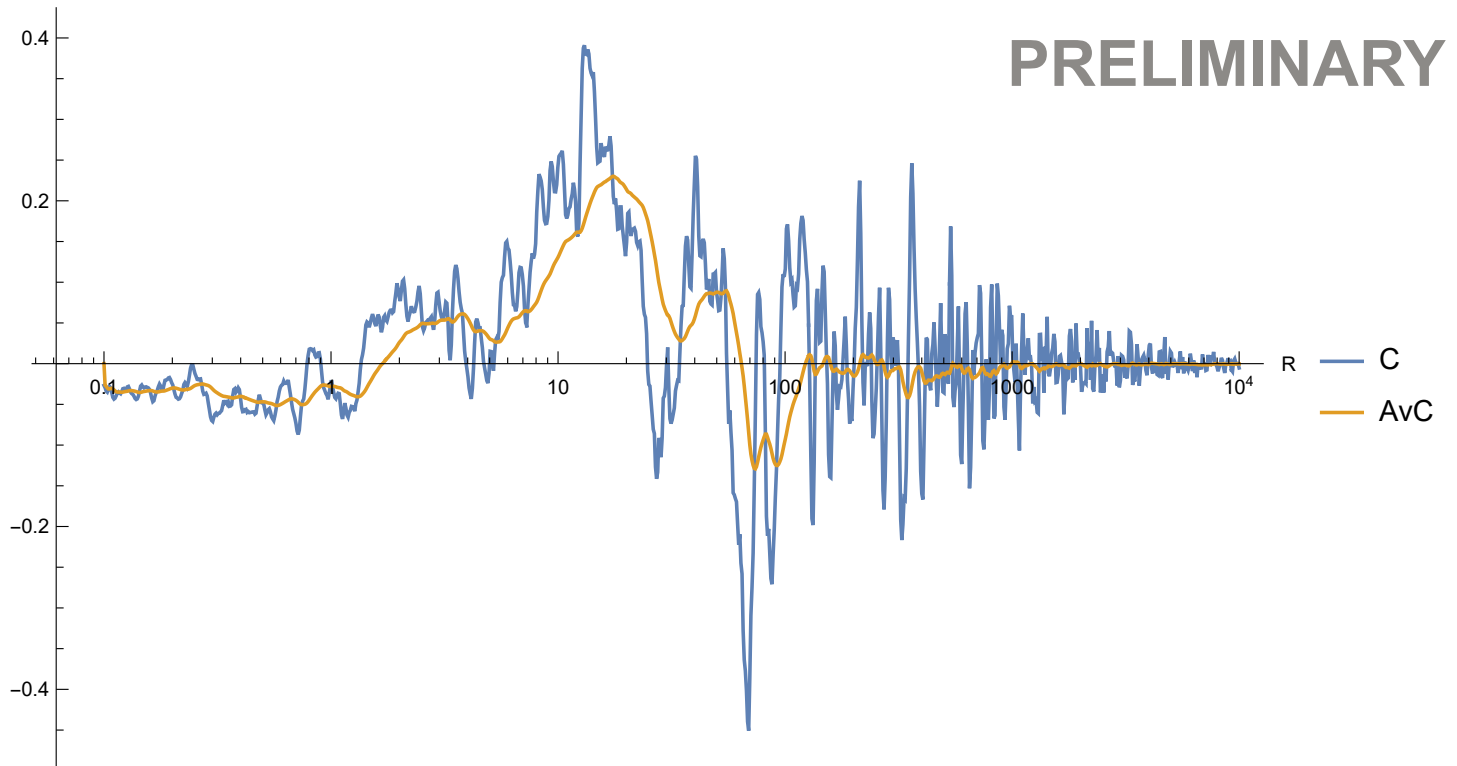
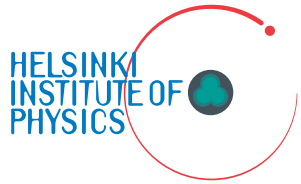


# One realisation of the curvature perturbation $\zeta(r)$





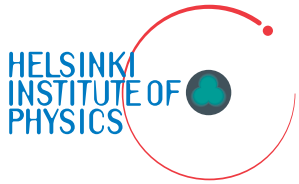
# One realisation of the compaction function $\mathcal{C}(r)=2G_N\Delta M(r)$





# Tail and spikes

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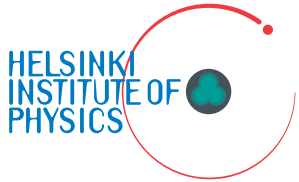


- Taking into account the stochastic tail for the patch distribution enhances PBH abundance by  $\sim 10^5$ .
- Taking into account the stochastic nature of individual patches naively enhances PBH abundance by an extra factor of  $\sim 10^2$ - $10^8$ . (Numbers preliminary.)
  - Caveat: once inside Hubble radius, spikes lead to large pressure gradients that smoothen the profile.
  - PBH formation simulations have to be redone.



# Spiking the conclusions

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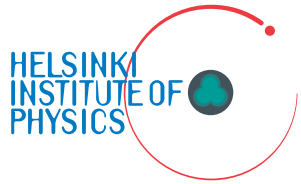


- PBHs can constitute all dark matter.
- Generating PBHs requires slowing down the inflaton, in single-field models.
- Stochastic effects can change PBH abundance by many orders of magnitude.



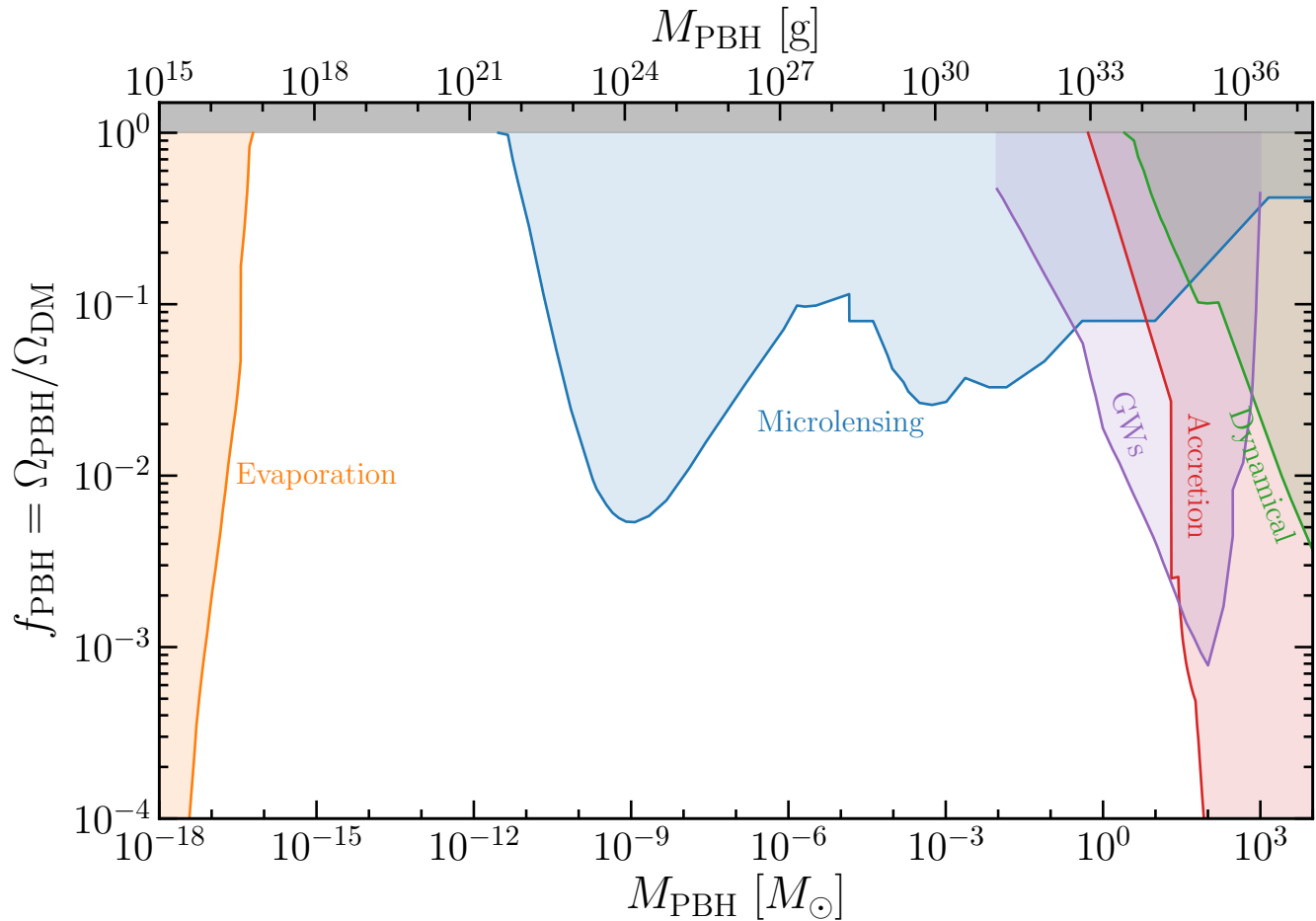
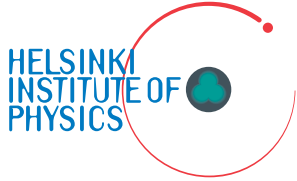
# Backup slides

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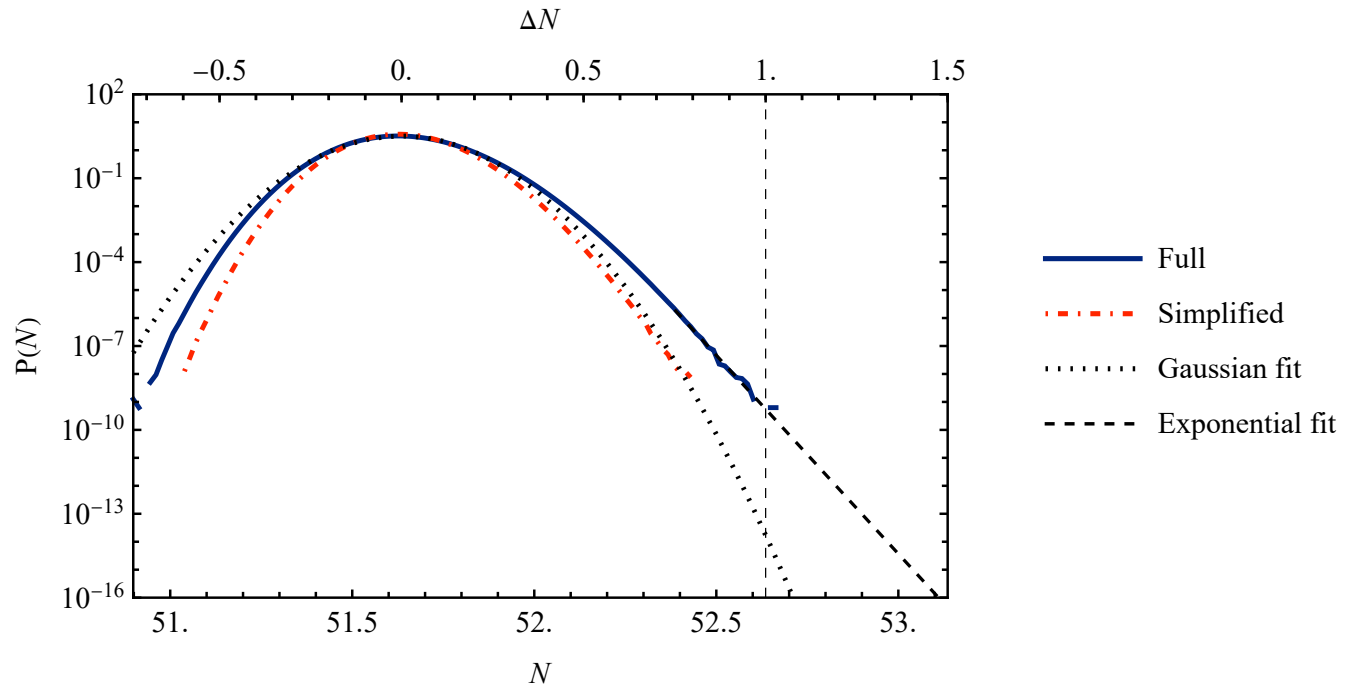
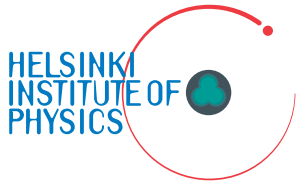


# Asteroid-mass or Planck scale





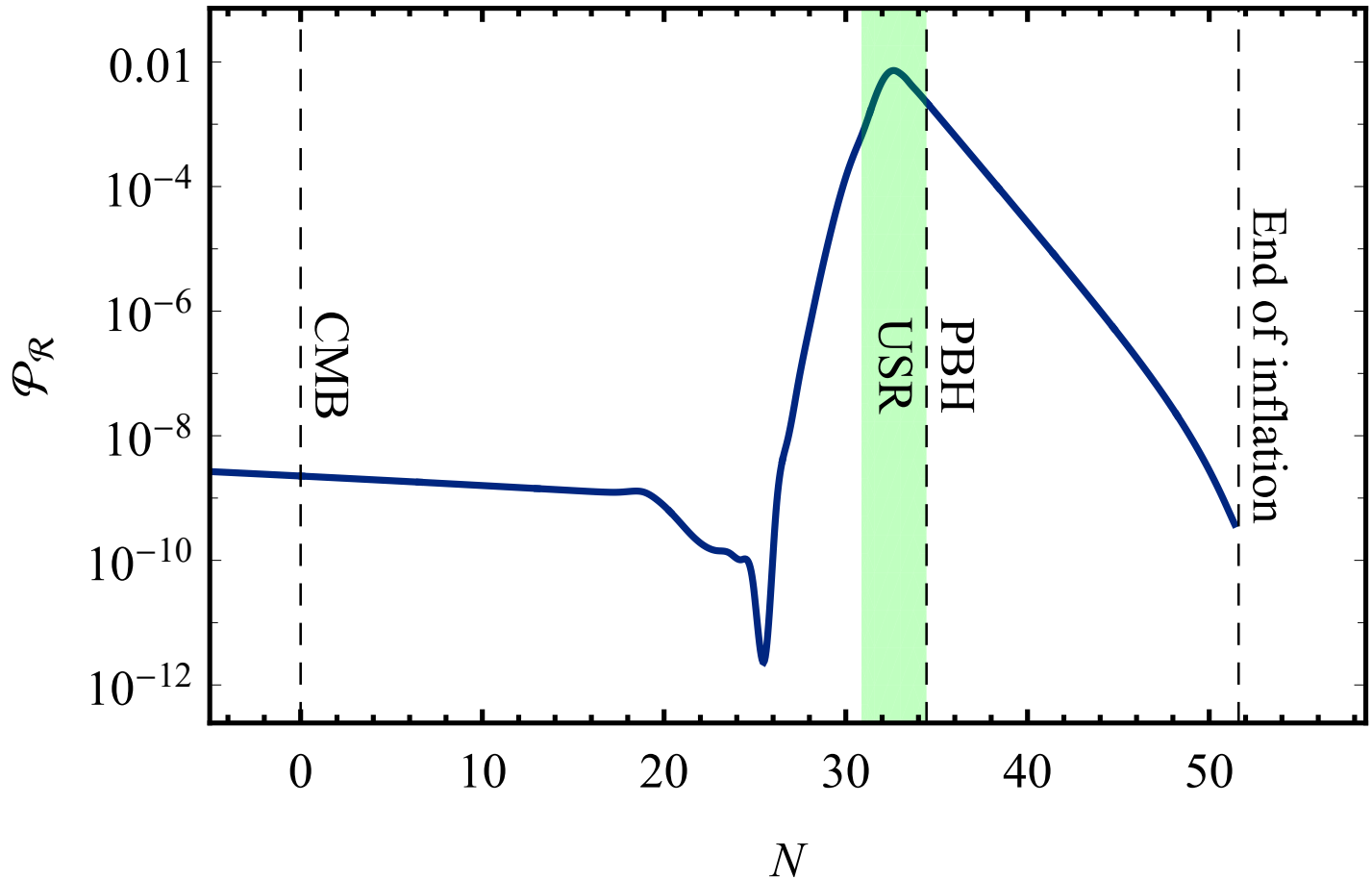
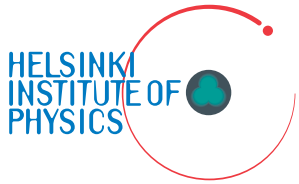
# $P(\Delta N = \zeta)$



- Stochastic kicks generate exponential tail.  
(Pattison, Vennin, Assadullahi, Wands: 1707.00537; Ezquiaga, Garcia-Bellido, Vennin: 1912.05399)
- This enhances PBH abundance by  $\sim 10^5$ .



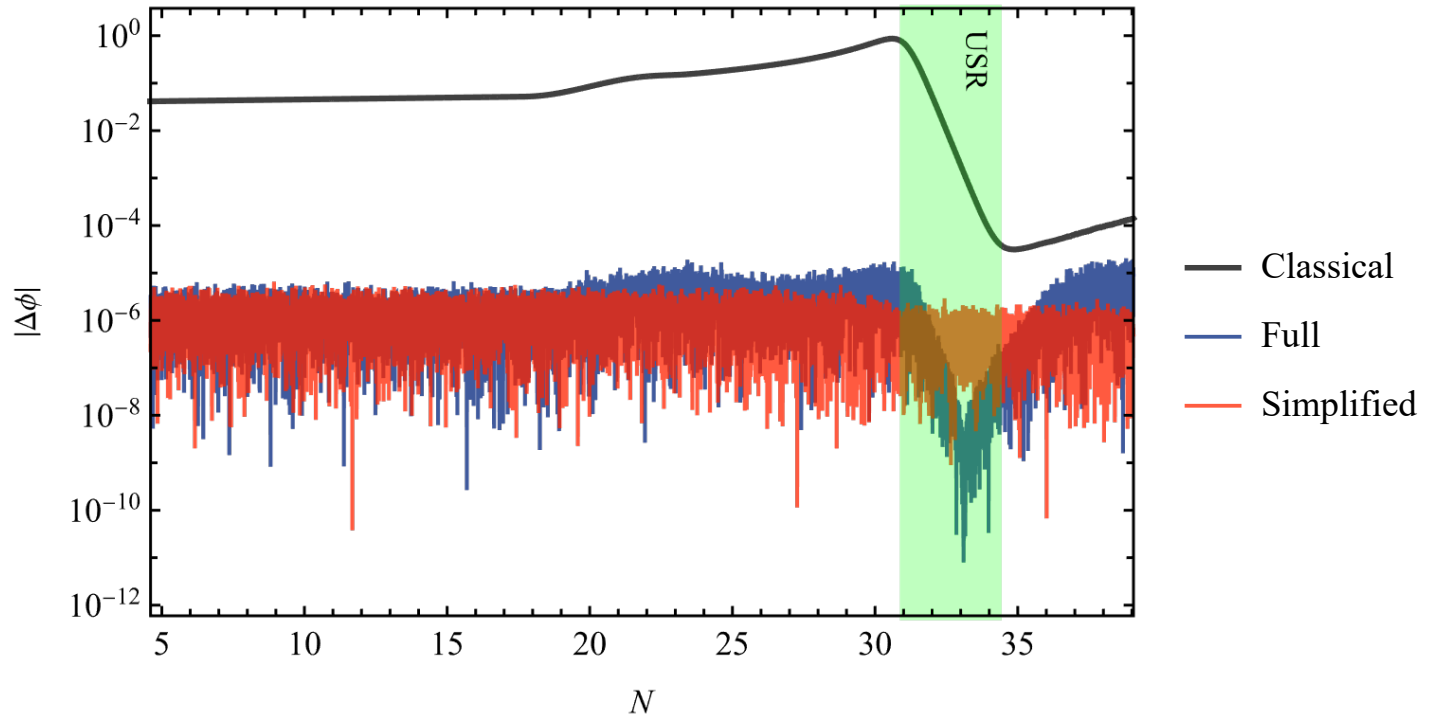
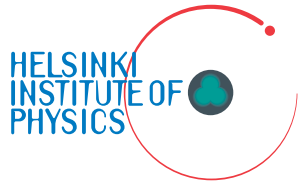
# Curvature power spectrum







# Classical force vs quantum kicks





# Correlation of $\max(\zeta)$ and $\max(\bar{c})$

