

REAL SCALAR PHASE TRANSITIONS: BUBBLE NUCLEATION, NONPERTURBATIVELY

with Oliver Gould & David J. Weir

Anna Kormu

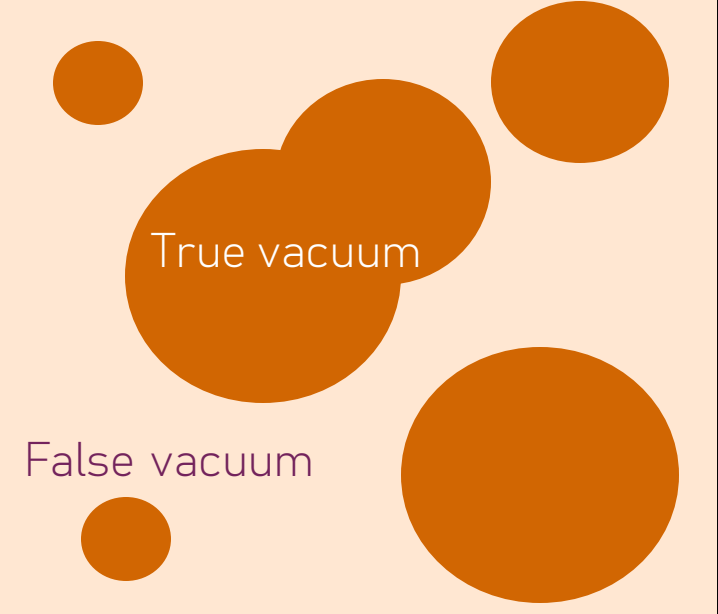
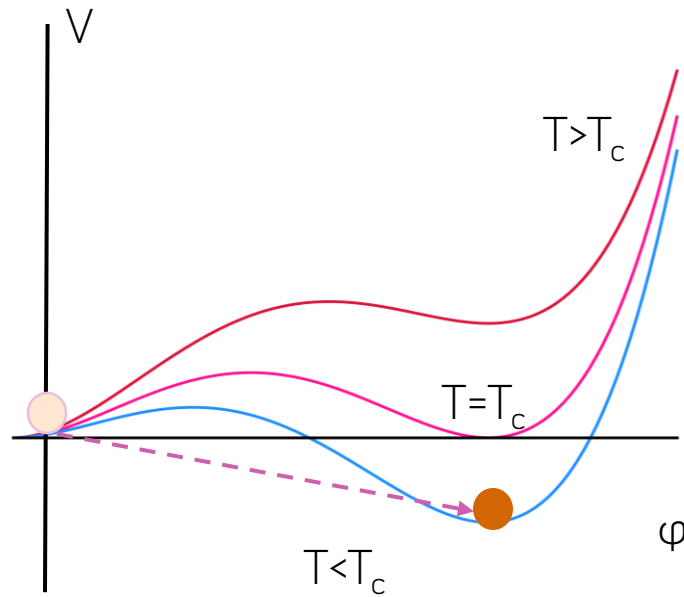
(she/her)

First Nordic Cosmology Meeting

24.10.2023

University of Helsinki & Helsinki Institute of Physics

Universe at 100GeV



1ST ORDER PHASE TRANSITIONS IN THE EARLY UNIVERSE

FATE OF THE FALSE VACUUM

- Relativistic field theory generalisation Callan & Coleman ([Phys. Rev. D 16, 1762 \(1977\)](#))
- Finite temperature approach introduced later by Affleck & Linde ([Phys. Rev. Lett. 46, 388 \(1981\)](#)), [Phys. Lett. B 100, 37 \(1981\)](#))

Volume
averaged
nucleation rate

$$\frac{\Gamma}{V} = A e^{-S_{3,b}(T)/T}$$

3D action

Prefactor

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Difficult!!

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- Moore, Rummukainen & Tranberg introduce a simulation method ([hep-lat/0103036](#), [hep-ph/0009132](#))

LISA PIPELINE

This work: from particle physics model to PT params



@AnnaKormu

1

Particle physics model

2

Phase transition parameters

Initial conditions
Critical temperature, bubble wall speed

3

GW power spectrum

Numerical simulations

4

LISA sensitivity

Configuration + noise level

5

Signal-to-noise ratio (SNR)

SINGLET SCALAR THEORY

Gould, [arXiv:2101.05528](https://arxiv.org/abs/2101.05528)

- Toy model possessing key features of BSM models
 - Potential has a tree-level barrier
 - Strong phase transition
- Dimensional reduction 4D cont \rightarrow 3D cont \rightarrow 3D lattice (imaginary time, high temp)

$$S_{\text{lat}} = \sum_x a^3 \left[-\frac{1}{2} Z_\phi \phi_x (\nabla_{\text{lat}}^2 \phi)_x + \sigma_{\text{lat}} \phi_x + \frac{1}{2} Z_\phi Z_m m_{\text{lat}}^2 \phi_x^2 + \frac{1}{3!} g_{\text{lat}} \phi_x^3 + \frac{1}{4!} Z_\phi^2 \lambda_{\text{lat}} \phi_x^4 \right]$$

NONPERTURBATIVELY

1

Pick an order parameter that behaves differently in the two phases

2

Simulate the probability of being in the critical bubble configuration

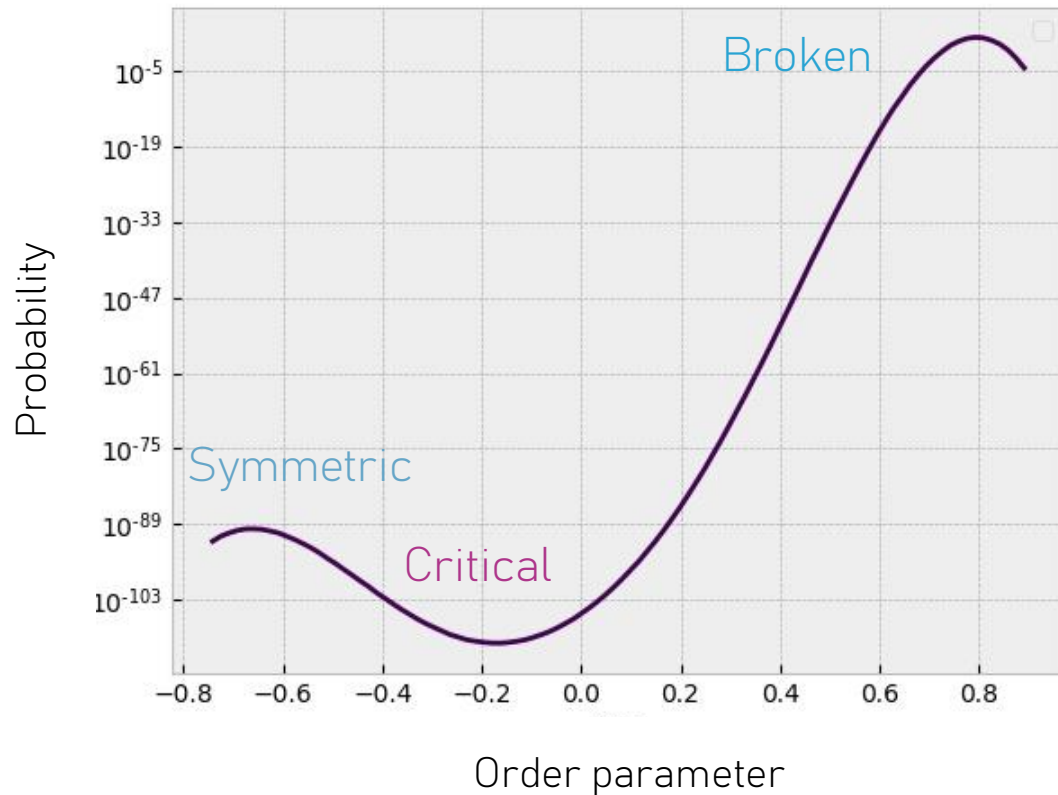
3

Perform real time evolution to determine whether the critical bubble tunnels or not

4

Calculate the total nucleation rate, dynamical prefactor \times probability info

NONPERTURBATIVELY



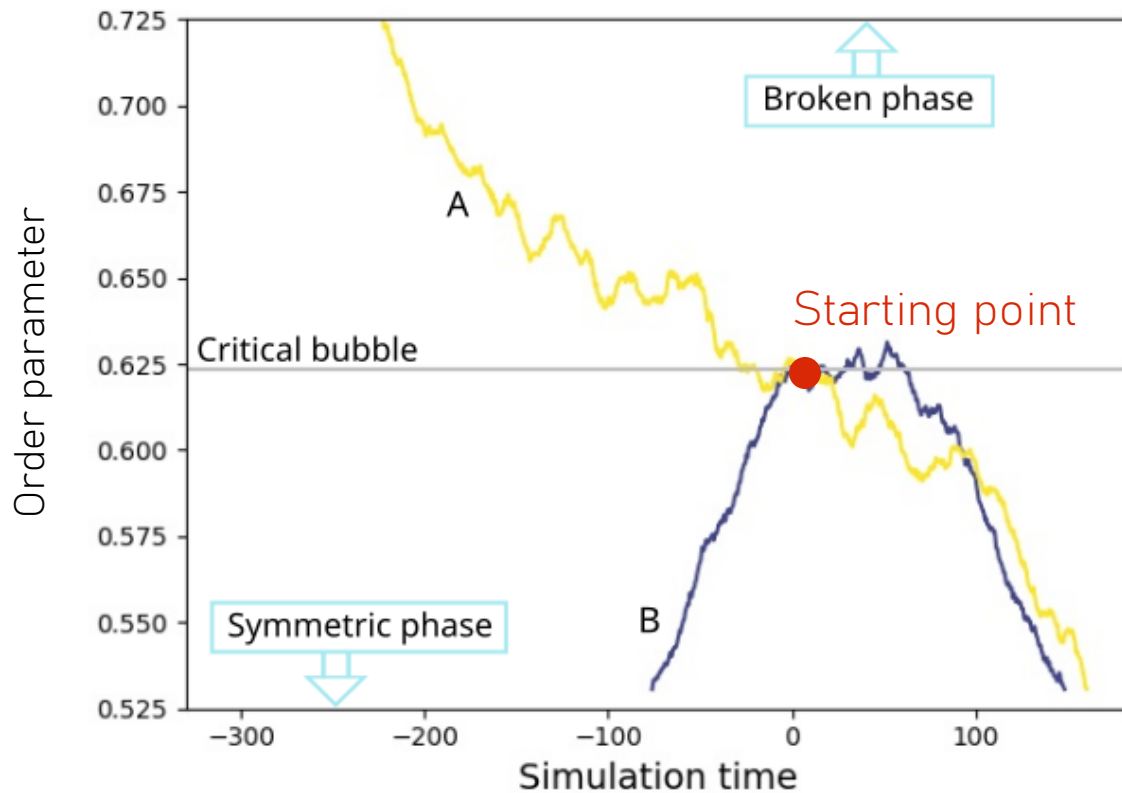
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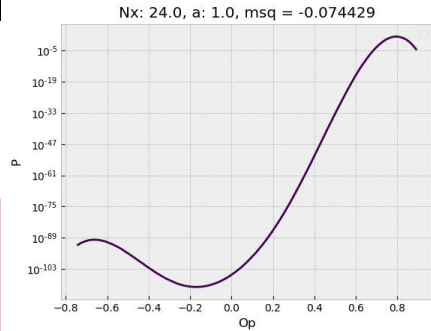
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$$\Gamma V = \frac{1}{2} P_C^\epsilon \left\langle \left| \frac{\Delta\theta(\alpha)}{\Delta t} \right| \times \mathbf{d}^\alpha \right\rangle$$

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Calculate the total nucleation rate, dynamical prefactor \times probability info

NONPERTURBATIVELY



Flux, rate of change of op

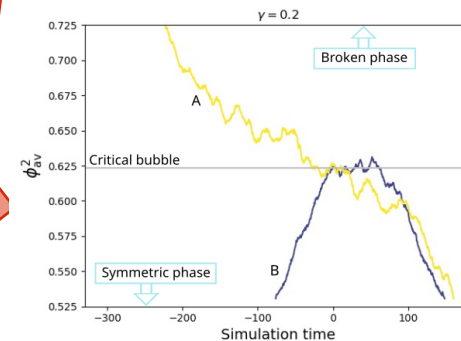
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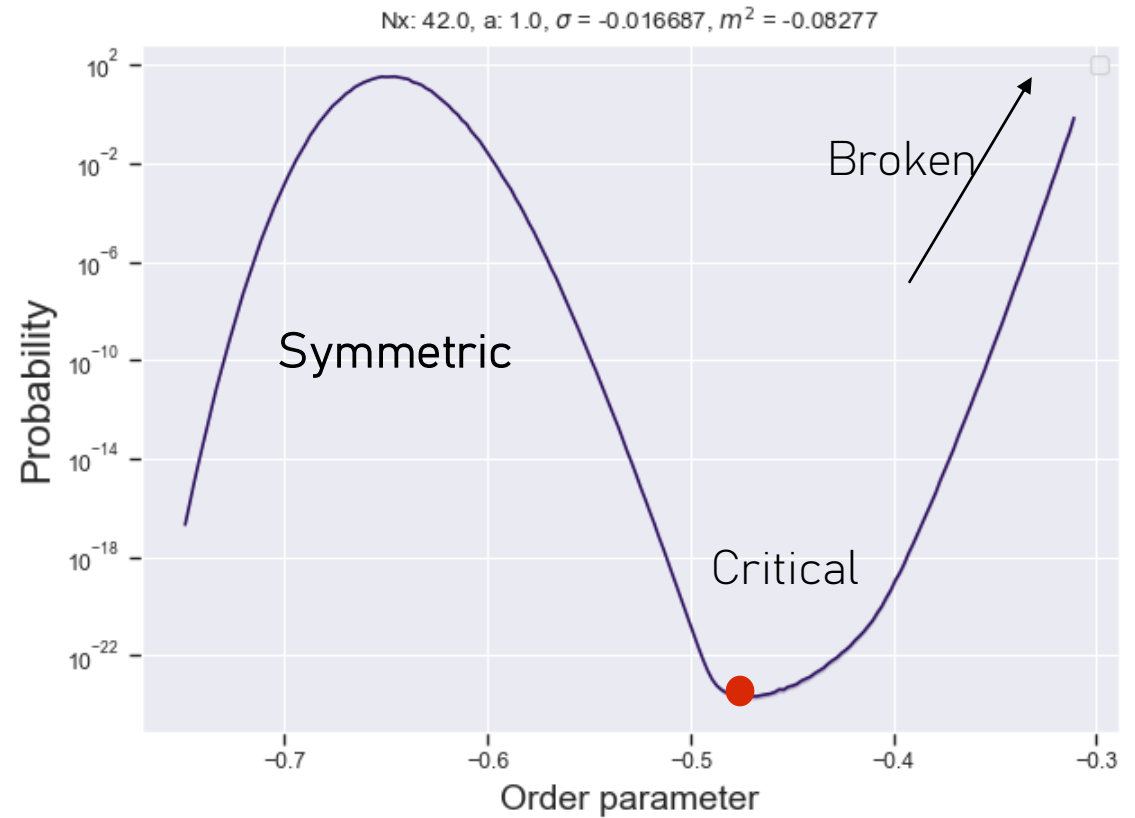
4

Calculate the total nucleation rate, dynamical information × stability info



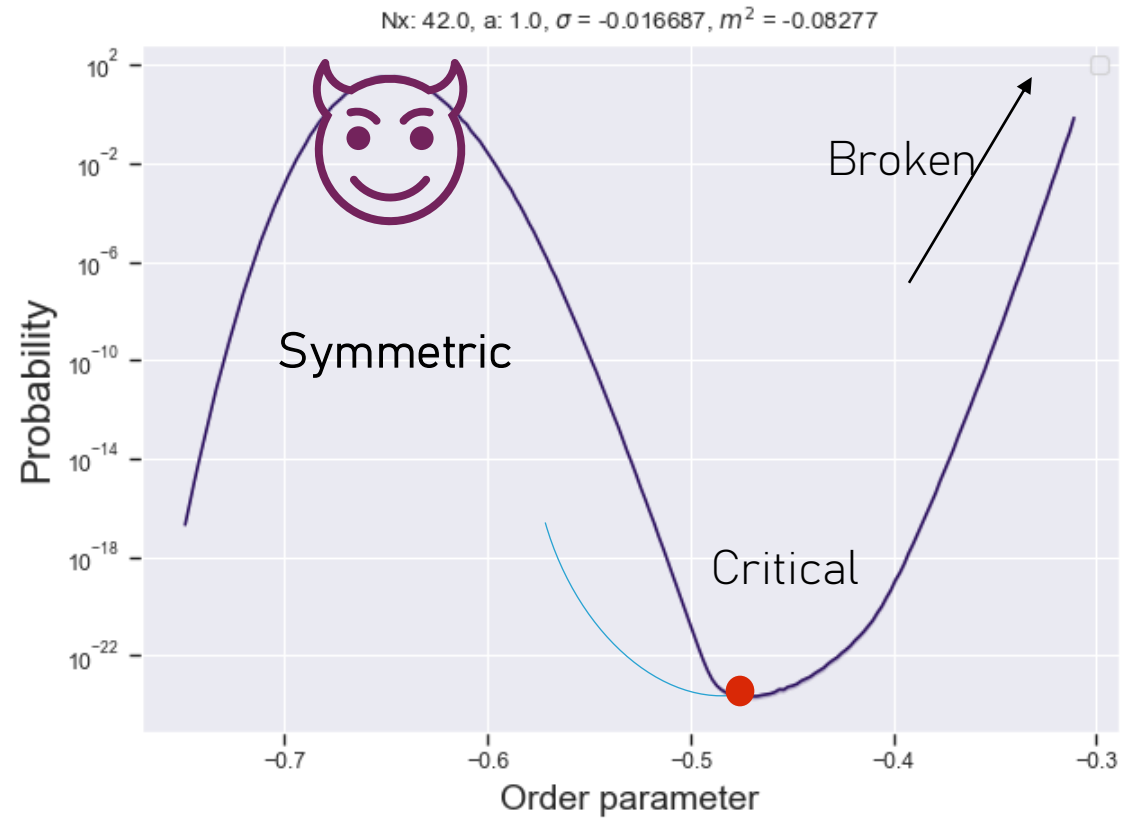
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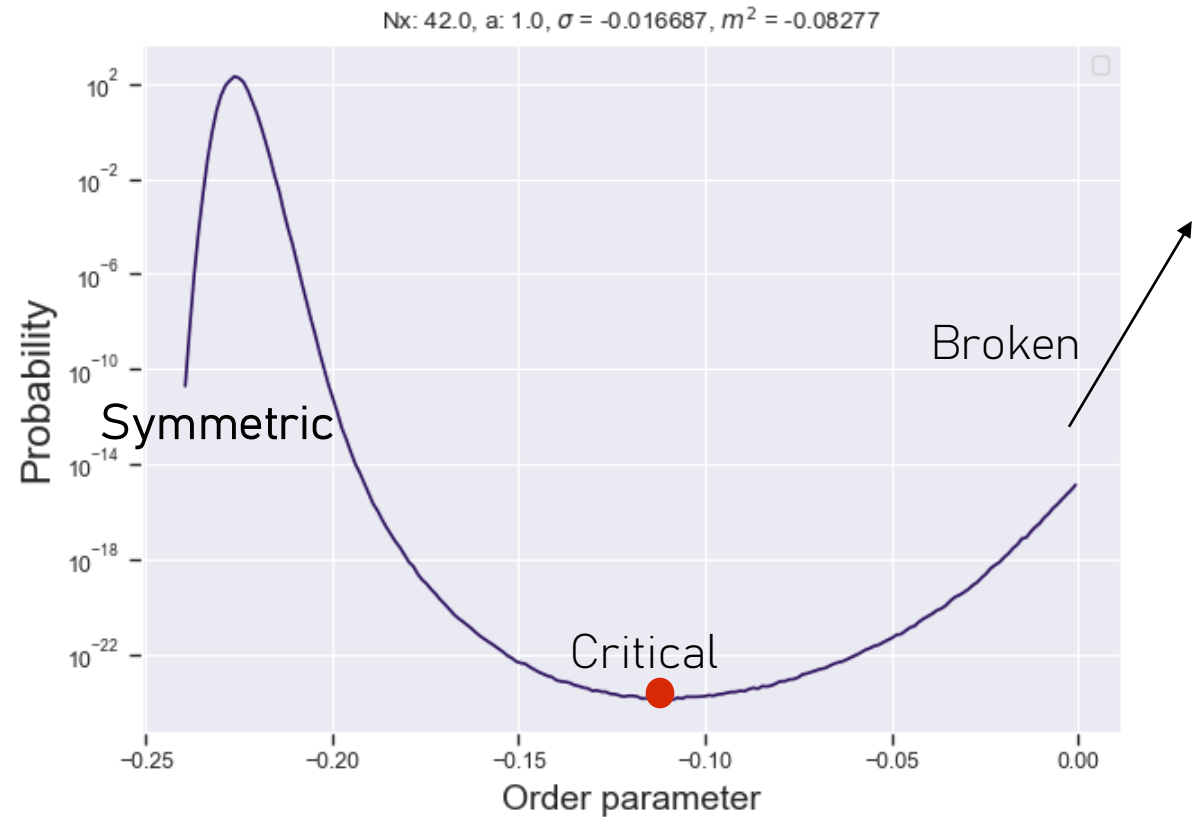
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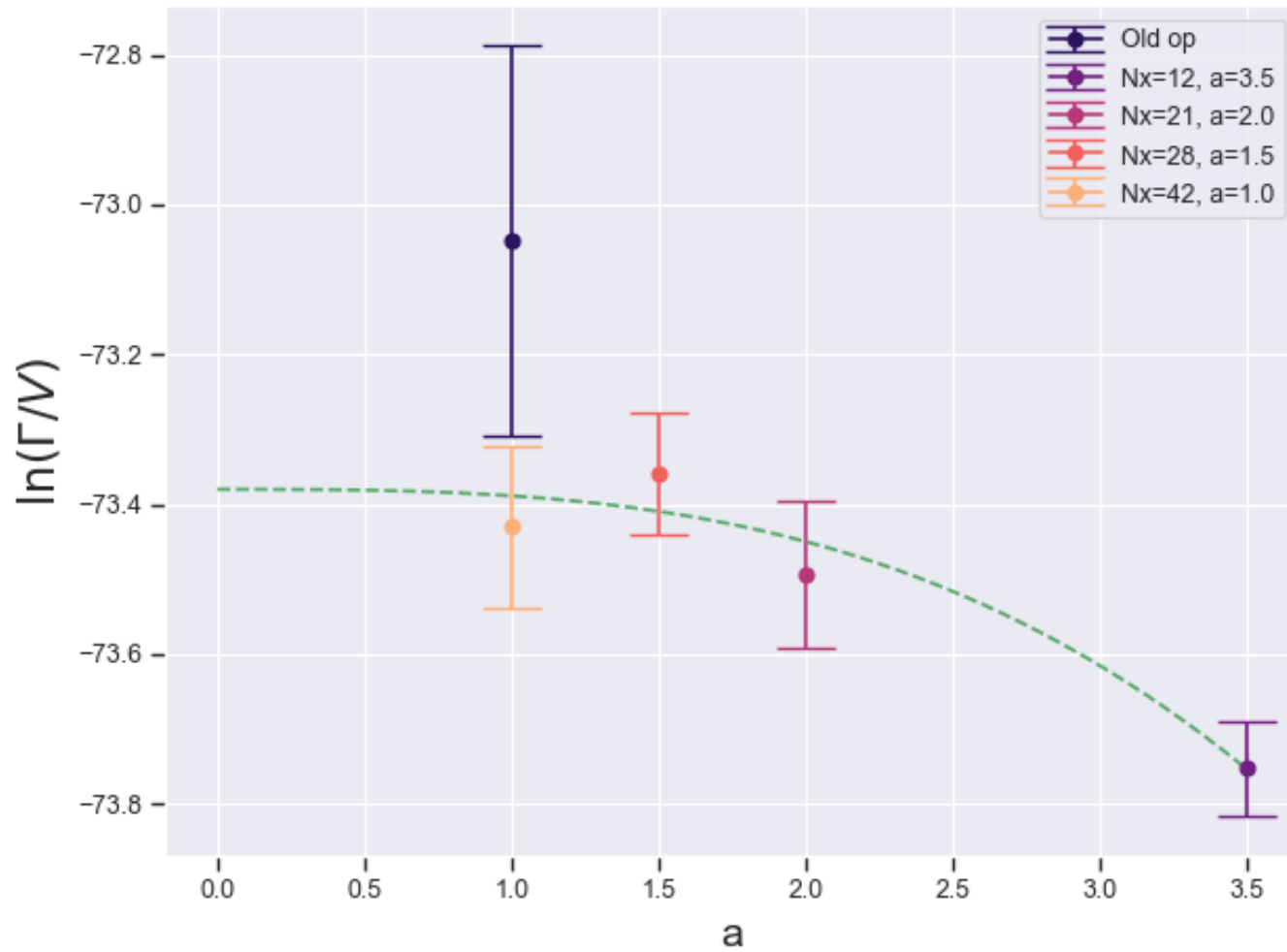
THE PROBLEM WITH LINEAR ORDER PARAMETER

Linear op \rightarrow Quadratic op

$$\frac{1}{V} \sum_i \phi_i \rightarrow \frac{1}{V} \sum_i (\phi_i^2 - \phi \phi_{sym})$$

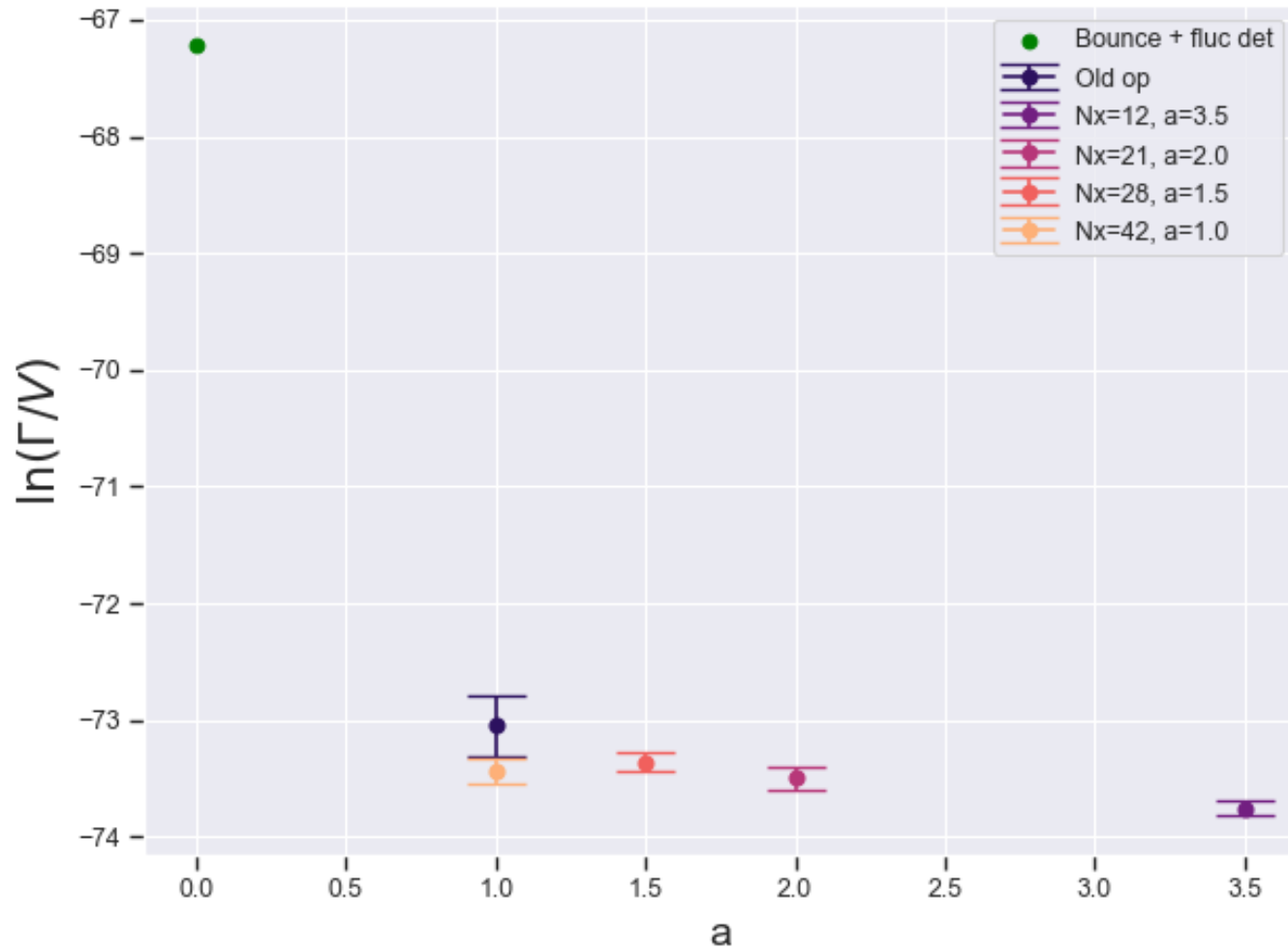


RESULTS



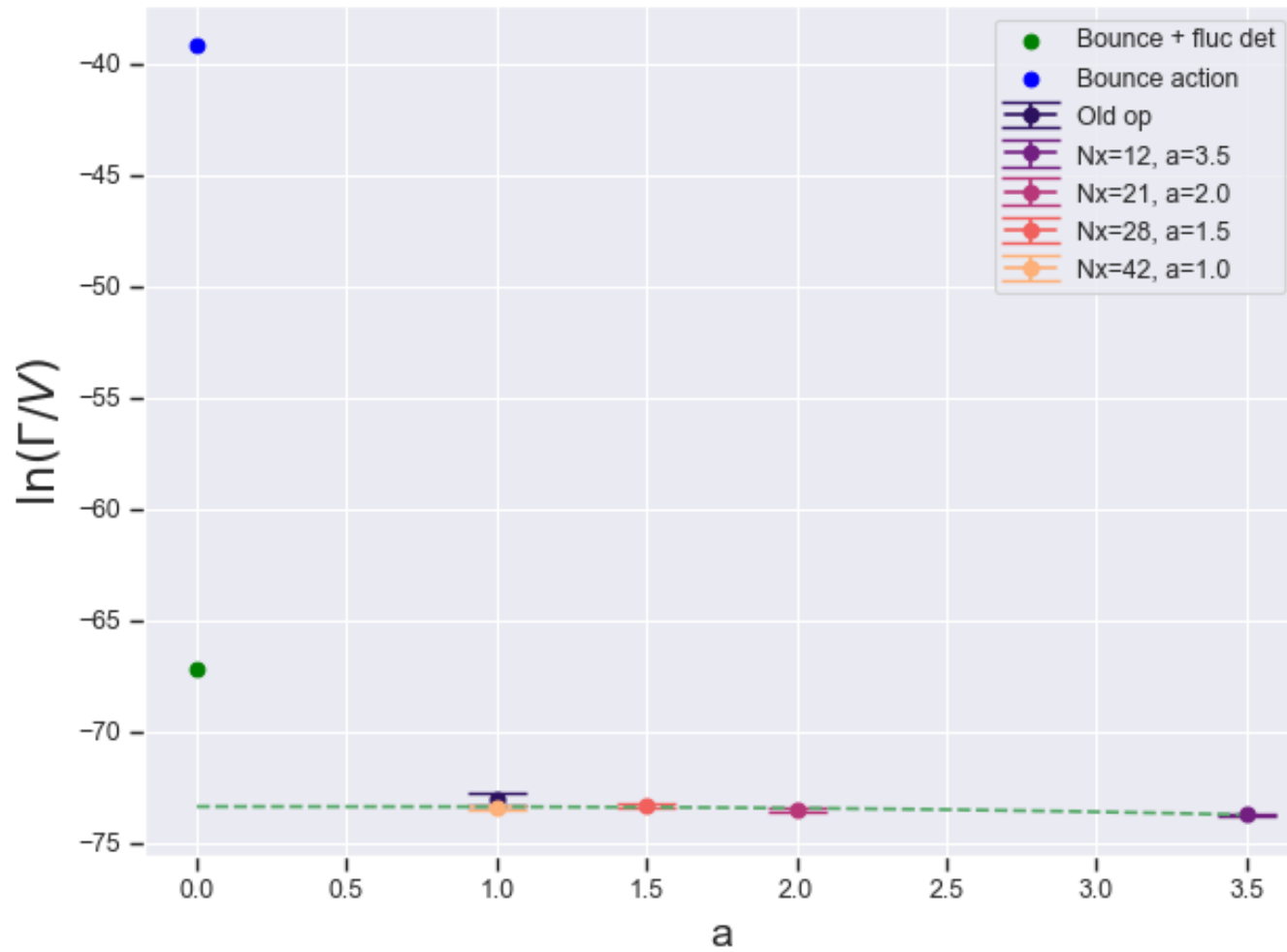
New order parameter runs compared to the old order parameter run

RESULTS



... compared to the bounce action + fluctuation determinant

RESULTS



... and with the bounce action only

CONCLUSIONS

- Allows us to calibrate the uncertainty in PT parameters when obtained from perturbative results
- Our simulations show a substantial suppression of nucleation rate compared to the one loop estimate
- Accurate computations of the nucleation rate are crucial for calculating e.g. the GW power spectrum
- Method and results can be applied to other theories

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One-pumpkin takeaway

There can be large uncertainties in nucleation rates calculated from the bounce action

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Thank you! Questions?

BACKUP

- Volume averaged nucleation rate vs. the perturbative calculation results as a function of temperature T (GeV)
- Lattice spacing fixed, varying physical volume

