

Nordita Winter School 2026

Cosmological Magnetic Fields:
Generation, Observation, and Modeling

Cosmological Magnetic Fields Theory Exercise Session 4: Evolution of magnetic and velocity fields



**UNIVERSITÉ
DE GENÈVE**

Antonino Salvino Midiri

antonino.midiri@unige.ch



**Swiss National
Science Foundation**

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_*$$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = A_V \left(\frac{k}{k_*} \right)^{n_V} \quad k < k_*$$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = A_V \left(\frac{k}{k_*} \right)^{n_V} \quad k < k_* \quad \text{If the velocity field is purely vortical} \rightarrow n_V = 2$$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = A_V \left(\frac{k}{k_*} \right)^{n_V} \quad k < k_*$$

If the velocity field is purely vortical $\rightarrow n_V = 2$
If the velocity field is purely compressional $\rightarrow n_V = 2$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = A_V \left(\frac{k}{k_*} \right)^{n_V} \quad k < k_*$$

If the velocity field is purely vortical $\rightarrow n_V = 2$
If the velocity field is purely compressional $\rightarrow n_V = 0$
For mixed cases $\rightarrow n_V = 0$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = A_V \left(\frac{k}{k_*} \right)^{n_V} \quad k < k_*$$

If the velocity field is purely vortical $\rightarrow n_V = 2$
If the velocity field is purely compressional $\rightarrow n_V = 2$
For mixed cases $\rightarrow n_V = 0$

The relation between peak frequency k_* and length scale L is $k_* = \frac{2\pi}{L}$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = A_V \left(\frac{k}{k_*} \right)^{n_V} \quad k < k_*$$

If the velocity field is purely vortical $\rightarrow n_V = 2$
If the velocity field is purely compressional $\rightarrow n_V = 0$
For mixed cases $\rightarrow n_V = 0$

The relation between peak frequency k_* and length scale L is $k_* = \frac{2\pi}{L}$

We can use the definitions of P_B and P_V to find A_B, A_V

$$\frac{\langle B^2 \rangle}{2} = \frac{1}{2\pi^2} \int_0^{k_*} dk k^2 P_B(k)$$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = A_V \left(\frac{k}{k_*} \right)^{n_V} \quad k < k_*$$

If the velocity field is purely vortical $\rightarrow n_V = 2$
If the velocity field is purely compressional $\rightarrow n_V = 0$
For mixed cases $\rightarrow n_V = 0$

The relation between peak frequency k_* and length scale L is $k_* = \frac{2\pi}{L}$

We can use the definitions of P_B and P_V to find A_B, A_V

$$\frac{\langle B^2 \rangle}{2} = \frac{1}{2\pi^2} \int_0^{k_*} dk k^2 P_B(k) = \frac{A_B}{2\pi^2} \left(\frac{L}{2\pi} \right)^{n_B} \frac{1}{n_B + 3} \left(\frac{2\pi}{L} \right)^{n_B + 3}$$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = A_B \left(\frac{k}{k_*} \right)^{n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = A_V \left(\frac{k}{k_*} \right)^{n_V} \quad k < k_*$$

If the velocity field is purely vortical $\rightarrow n_V = 2$
If the velocity field is purely compressional $\rightarrow n_V = 0$
For mixed cases $\rightarrow n_V = 0$

The relation between peak frequency k_* and length scale L is $k_* = \frac{2\pi}{L}$

We can use the definitions of P_B and P_V to find A_B, A_V

$$\frac{\langle B^2 \rangle}{2} = \frac{1}{2\pi^2} \int_0^{k_*} dk k^2 P_B(k) = \frac{A_B}{2\pi^2} \left(\frac{L}{2\pi} \right)^{n_B} \frac{1}{n_B + 3} \left(\frac{2\pi}{L} \right)^{n_B + 3} \longrightarrow A_B = \frac{n_B + 3}{8\pi} \langle B^2 \rangle L^3$$

Evolution of magnetic and velocity fields

Magnetic field spectra generated during radiation domination have the following large scale behavior

$$P_B(k) = k^{n_B} \frac{n_B + 3}{8\pi(2\pi)^{n_B}} \langle B^2 \rangle L^{3+n_B} \quad k < k_* \quad \text{magnetic fields are purely vortical} \rightarrow n_B = 2$$

The same applies to velocity field spectra, for which

$$P_V(k) = k^{n_V} \frac{n_V + 3}{8\pi(2\pi)^{n_V}} \langle v^2 \rangle L^{3+n_V} \quad k < k_*$$

If the velocity field is purely vortical $\rightarrow n_V = 2$
If the velocity field is purely compressional $\rightarrow n_V = -2$
For mixed cases $\rightarrow n_V = 0$

The relation between peak frequency k_* and length scale L is $k_* = \frac{2\pi}{L}$

We can use the definitions of P_B and P_V to find A_B, A_V

$$\frac{\langle B^2 \rangle}{2} = \frac{1}{2\pi^2} \int_0^{k_*} dk k^2 P_B(k) = \frac{A_B}{2\pi^2} \left(\frac{L}{2\pi} \right)^{n_B} \frac{1}{n_B + 3} \left(\frac{2\pi}{L} \right)^{n_B+3} \longrightarrow A_B = \frac{n_B + 3}{8\pi} \langle B^2 \rangle L^3$$

Evolution of magnetic and velocity fields

By computing the energy on a scale λ

$$B_\lambda^2 = \frac{1}{\pi^2} \int_0^\infty dk \, k^2 P_B(k) e^{-k^2 \lambda^2} = \frac{n_B + 3}{2(2\pi)^{n_B+3}} \Gamma\left(\frac{n_B + 3}{2}\right) \langle B^2 \rangle \left(\frac{L}{\lambda}\right)^{n_B+3}$$

Evolution of magnetic and velocity fields

By computing the energy on a scale λ

$$B_\lambda^2 = \frac{1}{\pi^2} \int_0^\infty dk \, k^2 P_B(k) e^{-k^2 \lambda^2} = \frac{n_B + 3}{2(2\pi)^{n_B+3}} \Gamma\left(\frac{n_B + 3}{2}\right) \langle B^2 \rangle \left(\frac{L}{\lambda}\right)^{n_B+3}$$

We find the following scalings

$$B_\lambda^2 \lambda^{n_B+3} \sim \langle B^2 \rangle L^{n_B+3} \quad \rightarrow \quad \langle B^2 \rangle \sim L^{-(n_B+3)} \quad \rightarrow \quad L \sim \langle B^2 \rangle^{-\frac{1}{n_B+3}}$$

Evolution of magnetic and velocity fields

By computing the energy on a scale λ

$$B_\lambda^2 = \frac{1}{\pi^2} \int_0^\infty dk \, k^2 P_B(k) e^{-k^2 \lambda^2} = \frac{n_B + 3}{2(2\pi)^{n_B+3}} \Gamma\left(\frac{n_B + 3}{2}\right) \langle B^2 \rangle \left(\frac{L}{\lambda}\right)^{n_B+3}$$

We find the following scalings

$$B_\lambda^2 \lambda^{n_B+3} \sim \langle B^2 \rangle L^{n_B+3} \quad \rightarrow \quad \langle B^2 \rangle \sim L^{-(n_B+3)} \quad \rightarrow \quad L \sim \langle B^2 \rangle^{-\frac{1}{n_B+3}}$$

We have seen that in turbulence there is a constant energy transfer rate in the inertial range

Evolution of magnetic and velocity fields

By computing the energy on a scale λ

$$B_\lambda^2 = \frac{1}{\pi^2} \int_0^\infty dk \, k^2 P_B(k) e^{-k^2 \lambda^2} = \frac{n_B + 3}{2(2\pi)^{n_B+3}} \Gamma\left(\frac{n_B + 3}{2}\right) \langle B^2 \rangle \left(\frac{L}{\lambda}\right)^{n_B+3}$$

We find the following scalings

$$B_\lambda^2 \lambda^{n_B+3} \sim \langle B^2 \rangle L^{n_B+3} \quad \rightarrow \quad \langle B^2 \rangle \sim L^{-(n_B+3)} \quad \rightarrow \quad L \sim \langle B^2 \rangle^{-\frac{1}{n_B+3}}$$

We have seen that in turbulence there is a constant energy transfer rate in the inertial range

$$\frac{d\langle v^2 \rangle}{d\eta} \sim \frac{\langle v^2 \rangle}{\tau_{eddy}} \sim \frac{\langle v^2 \rangle}{\frac{L}{\sqrt{\langle v^2 \rangle}}} \sim \frac{\langle v^2 \rangle^{\frac{3}{2}}}{L} \sim \langle v^2 \rangle^{\frac{3}{2} + \frac{1}{n_V+3}} = \langle v^2 \rangle^{\frac{3n_V+11}{2(n_V+3)}}$$

Evolution of magnetic and velocity fields

By computing the energy on a scale λ

$$B_\lambda^2 = \frac{1}{\pi^2} \int_0^\infty dk \, k^2 P_B(k) e^{-k^2 \lambda^2} = \frac{n_B + 3}{2(2\pi)^{n_B+3}} \Gamma\left(\frac{n_B + 3}{2}\right) \langle B^2 \rangle \left(\frac{L}{\lambda}\right)^{n_B+3}$$

We find the following scalings

$$B_\lambda^2 \lambda^{n_B+3} \sim \langle B^2 \rangle L^{n_B+3} \quad \rightarrow \quad \langle B^2 \rangle \sim L^{-(n_B+3)} \quad \rightarrow \quad L \sim \langle B^2 \rangle^{-\frac{1}{n_B+3}}$$

We have seen that in turbulence there is a constant energy transfer rate in the inertial range

$$\frac{d\langle v^2 \rangle}{d\eta} \sim \frac{\langle v^2 \rangle}{\tau_{eddy}} \sim \frac{\langle v^2 \rangle}{\frac{L}{\sqrt{\langle v^2 \rangle}}} \sim \frac{\langle v^2 \rangle^{\frac{3}{2}}}{L} \sim \langle v^2 \rangle^{\frac{3}{2} + \frac{1}{n_V+3}} = \langle v^2 \rangle^{\frac{3n_V+11}{2(n_V+3)}} \quad \rightarrow \quad \langle v^2 \rangle \sim \eta^{-\frac{2(n_V+3)}{n_V+5}}$$

Evolution of magnetic and velocity fields

By computing the energy on a scale λ

$$B_\lambda^2 = \frac{1}{\pi^2} \int_0^\infty dk \, k^2 P_B(k) e^{-k^2 \lambda^2} = \frac{n_B + 3}{2(2\pi)^{n_B+3}} \Gamma\left(\frac{n_B + 3}{2}\right) \langle B^2 \rangle \left(\frac{L}{\lambda}\right)^{n_B+3}$$

We find the following scalings

$$B_\lambda^2 \lambda^{n_B+3} \sim \langle B^2 \rangle L^{n_B+3} \quad \rightarrow \langle B^2 \rangle \sim L^{-(n_B+3)} \quad \rightarrow L \sim \langle B^2 \rangle^{-\frac{1}{n_B+3}}$$

We have seen that in turbulence there is a constant energy transfer rate in the inertial range

$$\frac{d\langle v^2 \rangle}{d\eta} \sim \frac{\langle v^2 \rangle}{\tau_{eddy}} \sim \frac{\langle v^2 \rangle}{\frac{L}{\sqrt{\langle v^2 \rangle}}} \sim \frac{\langle v^2 \rangle^{\frac{3}{2}}}{L} \sim \langle v^2 \rangle^{\frac{3}{2} + \frac{1}{n_V+3}} = \langle v^2 \rangle^{\frac{3n_V+11}{2(n_V+3)}} \quad \rightarrow \langle v^2 \rangle \sim \eta^{-\frac{2(n_V+3)}{n_V+5}}$$
$$\rightarrow L \sim \eta^{\frac{2}{n_V+5}}$$

Evolution of magnetic and velocity fields

By computing the energy on a scale λ

$$B_\lambda^2 = \frac{1}{\pi^2} \int_0^\infty dk \, k^2 P_B(k) e^{-k^2 \lambda^2} = \frac{n_B + 3}{2(2\pi)^{n_B+3}} \Gamma\left(\frac{n_B + 3}{2}\right) \langle B^2 \rangle \left(\frac{L}{\lambda}\right)^{n_B+3}$$

We find the following scalings

$$B_\lambda^2 \lambda^{n_B+3} \sim \langle B^2 \rangle L^{n_B+3} \quad \rightarrow \langle B^2 \rangle \sim L^{-(n_B+3)} \quad \rightarrow L \sim \langle B^2 \rangle^{-\frac{1}{n_B+3}}$$

We have seen that in turbulence there is a constant energy transfer rate in the inertial range

$$\begin{aligned} \frac{d\langle v^2 \rangle}{d\eta} &\sim \frac{\langle v^2 \rangle}{\tau_{eddy}} \sim \frac{\langle v^2 \rangle}{\frac{L}{\sqrt{\langle v^2 \rangle}}} \sim \frac{\langle v^2 \rangle^{\frac{3}{2}}}{L} \sim \langle v^2 \rangle^{\frac{3}{2} + \frac{1}{n_V+3}} = \langle v^2 \rangle^{\frac{3n_V+11}{2(n_V+3)}} \quad \rightarrow \langle v^2 \rangle \sim \eta^{-\frac{2(n_V+3)}{n_V+5}} \\ &\quad \rightarrow L \sim \eta^{\frac{2}{n_V+5}} \end{aligned}$$

Same applies to the magnetic field (from equipartition in the inertial range) $\rightarrow \langle B^2 \rangle \sim \eta^{-\frac{2(n_V+3)}{n_V+5}}$

Evolution of magnetic and velocity fields

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

Evolution of magnetic and velocity fields

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

Evolution of magnetic and velocity fields

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

Exercise no. 1

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

A) Consider fully vortical turbulence. How do $\langle v^2 \rangle$, $\langle B^2 \rangle$, L scale with time?

Exercise no. 1

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

A) Consider fully vortical turbulence. How do $\langle v^2 \rangle, \langle B^2 \rangle, L$ scale with time?

$$n_V = 2 \rightarrow p = \frac{5}{7}, q = \frac{2}{7} \quad \rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

Exercise no. 1

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

B) What happens to the velocity and magnetic field spectra at large scales?

Exercise no. 1

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

B) What happens to the velocity and magnetic field spectra at large scales?

Since $n_B = n_V$ they are both constant

Exercise no. 1

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

C) Consider now the mixed case. How do $\langle v^2 \rangle, \langle B^2 \rangle, L$ scale with time?

Exercise no. 1

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

C) Consider now the mixed case. How do $\langle v^2 \rangle, \langle B^2 \rangle, L$ scale with time?

$$n_V = 0 \rightarrow p = \frac{3}{5} \text{ and } q = \frac{2}{5} \quad \rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-6/5} \quad \rightarrow L \sim \eta^{2/5}$$

Exercise no. 1

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

D) What happens to the velocity and magnetic field spectra at large scales?

Exercise no. 1

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p} \quad p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5} \quad p + q = 1$$

$$\rightarrow L \sim \eta^q$$

At large scales we have

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{-2p} \eta^{(3+n_V)q} \sim \eta^0$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{-2p} \eta^{(3+n_B)q} = \eta^{\frac{2(n_B - n_V)}{n_V + 5}}$$

D) What happens to the velocity and magnetic field spectra at large scales?

P_V stays constant while P_B grows as $\eta^{4/5}$

→ INVERSE CASCADE

Exercise no. 1

Let us now consider the helical case.

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const}$$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p}$$

$$\rightarrow L \sim \eta^q$$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p}$$

$$\rightarrow L \sim \eta^q$$

$$p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5}$$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2p}$$

$$\rightarrow L \sim \eta^q$$

$$p = \frac{n_V + 3}{n_V + 5} \quad q = \frac{2}{n_V + 5}$$

In the scaling laws n_V has to be replaced by $n_V^{\text{equiv}} = -2$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\begin{aligned} \rightarrow \langle v^2 \rangle &\sim \langle B^2 \rangle \sim \eta^{-2p} \\ \rightarrow L &\sim \eta^q \end{aligned} \quad p = \frac{1}{3} \quad q = \frac{2}{3}$$

In the scaling laws n_V has to be replaced by $n_V^{equiv} = -2$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\begin{aligned} \rightarrow \langle v^2 \rangle &\sim \langle B^2 \rangle \sim \eta^{-2p} \\ \rightarrow L &\sim \eta^q \end{aligned} \quad p = \frac{1}{3} \quad q = \frac{2}{3}$$

In the scaling laws n_V has to be replaced by $n_V^{\text{equiv}} = -2$

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{\frac{2(n_V - n_V^{\text{equiv}})}{n_V^{\text{equiv}} + 5}} \sim \eta^{\frac{2(n_V + 2)}{3}}$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{\frac{2(n_B - n_V^{\text{equiv}})}{n_V^{\text{equiv}} + 5}} \sim \eta^{\frac{2(n_B + 2)}{3}}$$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\begin{aligned} \rightarrow \langle v^2 \rangle &\sim \langle B^2 \rangle \sim \eta^{-2p} \\ \rightarrow L &\sim \eta^q \end{aligned} \quad p = \frac{1}{3} \quad q = \frac{2}{3}$$

E) What happens to the velocity and magnetic spectra at large scales in the fully vortical case?

In the scaling laws n_V has to be replaced by $n_V^{equiv} = -2$

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{\frac{2(n_V - n_V^{equiv})}{n_V^{equiv} + 5}} \sim \eta^{\frac{2(n_V + 2)}{3}}$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{\frac{2(n_B - n_V^{equiv})}{n_V^{equiv} + 5}} \sim \eta^{\frac{2(n_B + 2)}{3}}$$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\begin{aligned} \rightarrow \langle v^2 \rangle &\sim \langle B^2 \rangle \sim \eta^{-2p} \\ \rightarrow L &\sim \eta^q \end{aligned} \quad p = \frac{1}{3} \quad q = \frac{2}{3}$$

E) What happens to the velocity and magnetic spectra at large scales in the fully vortical case?

They both grow as $\eta^{8/3}$

In the scaling laws n_V has to be replaced by $n_V^{\text{equiv}} = -2$

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{\frac{2(n_V - n_V^{\text{equiv}})}{n_V^{\text{equiv}} + 5}} \sim \eta^{\frac{2(n_V + 2)}{3}}$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{\frac{2(n_B - n_V^{\text{equiv}})}{n_V^{\text{equiv}} + 5}} \sim \eta^{\frac{2(n_B + 2)}{3}}$$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\begin{aligned} \rightarrow \langle v^2 \rangle &\sim \langle B^2 \rangle \sim \eta^{-2p} \\ \rightarrow L &\sim \eta^q \end{aligned} \quad p = \frac{1}{3} \quad q = \frac{2}{3}$$

E) What happens instead
in the mixed case case?

In the scaling laws n_V has to be replaced by $n_V^{equiv} = -2$

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{\frac{2(n_V - n_V^{equiv})}{n_V^{equiv} + 5}} \sim \eta^{\frac{2(n_V + 2)}{3}}$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{\frac{2(n_B - n_V^{equiv})}{n_V^{equiv} + 5}} \sim \eta^{\frac{2(n_B + 2)}{3}}$$

Exercise no. 1

Let us now consider the helical case. If helicity is conserved we have

$$h = \langle A \cdot B \rangle \sim B^2 L \sim \text{const} \quad \rightarrow -2p + q = 0$$

$$\begin{aligned} \rightarrow \langle v^2 \rangle &\sim \langle B^2 \rangle \sim \eta^{-2p} \\ \rightarrow L &\sim \eta^q \end{aligned} \quad p = \frac{1}{3} \quad q = \frac{2}{3}$$

In the scaling laws n_V has to be replaced by $n_V^{\text{equiv}} = -2$

$$P_V \left(k < \frac{2\pi}{L} \right) \sim \langle v^2 \rangle L^{3+n_V} \sim \eta^{\frac{2(n_V - n_V^{\text{equiv}})}{n_V^{\text{equiv}} + 5}} \sim \eta^{\frac{2(n_V + 2)}{3}}$$

$$P_B \left(k < \frac{2\pi}{L} \right) \sim \langle B^2 \rangle L^{3+n_B} \sim \eta^{\frac{2(n_B - n_V^{\text{equiv}})}{n_V^{\text{equiv}} + 5}} \sim \eta^{\frac{2(n_B + 2)}{3}}$$

E) What happens instead in the mixed case case?

The magnetic spectrum grows as $\eta^{8/3}$ while the velocity spectrum as $\eta^{4/3}$

Summary

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

Summary

@ large scales

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

P_V, P_B constant

Summary

@ large scales

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

P_V, P_B constant

In the non-helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-6/5} \quad \rightarrow L \sim \eta^{2/5}$$

Summary

@ large scales

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

P_V, P_B constant

In the non-helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-6/5} \quad \rightarrow L \sim \eta^{2/5}$$

P_V constant, P_B grows as $\eta^{4/5}$

Summary

@ large scales

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

P_V, P_B constant

In the non-helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-6/5} \quad \rightarrow L \sim \eta^{2/5}$$

P_V constant, P_B grows as $\eta^{4/5}$

In the helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2/3} \quad \rightarrow L \sim \eta^{2/3}$$

Summary

@ large scales

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

P_V, P_B constant

In the non-helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-6/5} \quad \rightarrow L \sim \eta^{2/5}$$

P_V constant, P_B grows as $\eta^{4/5}$

In the helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2/3} \quad \rightarrow L \sim \eta^{2/3}$$

P_V, P_B grow as $\eta^{8/3}$

Summary

@ large scales

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

P_V, P_B constant

In the non-helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-6/5} \quad \rightarrow L \sim \eta^{2/5}$$

P_V constant, P_B grows as $\eta^{4/5}$

In the helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2/3} \quad \rightarrow L \sim \eta^{2/3}$$

P_V, P_B grow as $\eta^{8/3}$

In the helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2/3} \quad \rightarrow L \sim \eta^{2/3}$$

Summary

@ large scales

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

P_V, P_B constant

In the non-helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-6/5} \quad \rightarrow L \sim \eta^{2/5}$$

P_V constant, P_B grows as $\eta^{4/5}$

In the helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2/3} \quad \rightarrow L \sim \eta^{2/3}$$

P_V, P_B grow as $\eta^{8/3}$

In the helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2/3} \quad \rightarrow L \sim \eta^{2/3}$$

P_V grows as $\eta^{4/3}$, P_B grows as $\eta^{8/3}$

Exercise no. 2

@ large scales

In the non-helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-10/7} \quad \rightarrow L \sim \eta^{2/7}$$

P_V, P_B constant

In the non-helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-6/5} \quad \rightarrow L \sim \eta^{2/5}$$

P_V constant, P_B grows as $\eta^{4/5}$

In the helical fully vortical case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2/3} \quad \rightarrow L \sim \eta^{2/3}$$

P_V, P_B grow as $\eta^{8/3}$

In the helical mixed case we have

$$\rightarrow \langle v^2 \rangle \sim \langle B^2 \rangle \sim \eta^{-2/3} \quad \rightarrow L \sim \eta^{2/3}$$

P_V grows as $\eta^{4/3}$, P_B grows as $\eta^{8/3}$

For the 4 cases make a plot of the time evolution of P_B and P_V considering

$$P_B(k) = k^{n_B} \frac{n_B + 3}{8\pi(2\pi)^{n_B}} \langle B^2 \rangle L^{3+n_B} / (1 + k L)^{n_B+m_B} \quad \text{with } m_B = 11/3$$