A roundabout for waves

Anton Akhmerov with Isidora Araya Day, Kostas Vilkelis, Antonio Manesco, Mert Bozkurt, and Valla Fatemi SciPost Phys. 18, 098 (+source code)

Jun 2, 2025, Nordita workshop





The problem

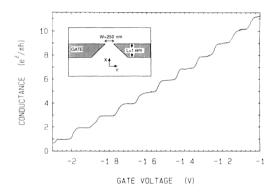




Background: protected transmission

Waves scatter, but sometimes transmission is protected.

 Quantum point contact: transmission protected by smooth potential

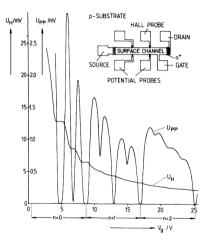


van Wees et al., 1988

Background: protected transmission

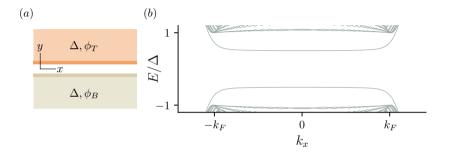
Waves scatter, but sometimes transmission is protected.

- Quantum point contact: transmission protected by smooth potential
- Quantum Hall effect: chiral transmission protected by topology



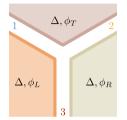
Klitzing et al., 1980

The solution



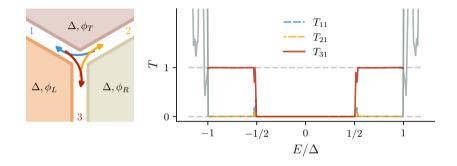
1. Take a short Josepshon junction

The solution



- 1. Take a short Josepshon junction
- 2. Make 3 meet

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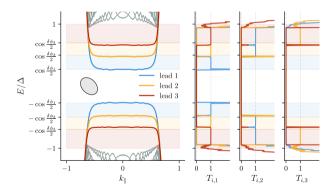
- 1. Take a short Josepshon junction
- 2. Make 3 meet
- 3. Get chiral and quantized transmission

The end. Questions?

Is it really protected?

Works if we:

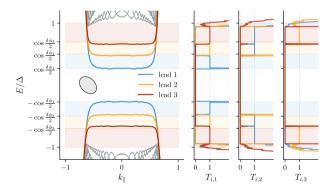
 Make phase differences unequal



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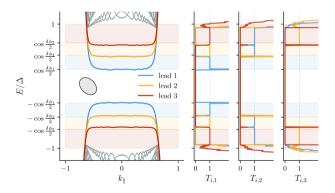
- Make phase differences unequal
- ► Make junction asymmetric



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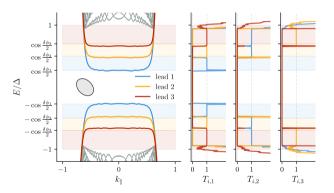
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Works if we:

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- ► Make dispersion anisotropic
- ► ¡Break particle-hole symmetry $H_{ee} \rightarrow 2 \times H_{ee}!$



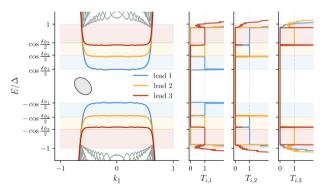
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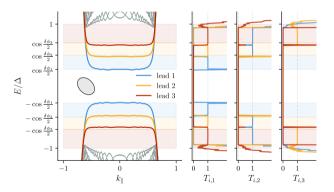
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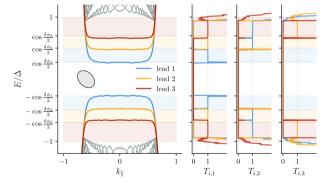
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Breaks if we:

- ► Make $\mu \to \mu(r)$
- Make $\mu \sim \Delta$
- Introduce sharp $< 90^{\circ}$ turns



What is going on?

Foundation

Probing Fermi Sea Topology by Andreev State Transport

Pok Man Tam[®] and Charles L. Kane[®]

Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA

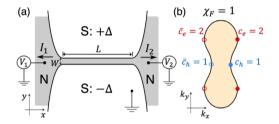
(Received 14 October 2022; accepted 16 January 2023; published 1 March 2023)

We show that the topology of the Fermi sea of a two-dimensional electron gas (2DEG) is reflected in the ballistic Landauer transport *along* a long and narrow Josephson π junction that proximitizes the 2DEG. The low-energy Andreev states bound to the junction are shown to exhibit a dispersion that is sensitive to the Euler characteristic of the Fermi sea (χ_F). We highlight two important relations: one connects the electron or hole nature of Andreev states to the convex or concave nature of Fermi surface critical points, and one relates these critical points to χ_F . We then argue that the transport of Andreev states leads to a quantized conductance that probes χ_F . An experiment is proposed to measure this effect, from which we predict an *I*-V characteristic that not only captures the topology of the Fermi sea in metals, but also resembles the rectification effect in diodes. Finally, we evaluate the feasibility of measuring this quantized response in graphene, InAs and HgTe 2DEGs.

DOI: 10.1103/PhysRevLett.130.096301

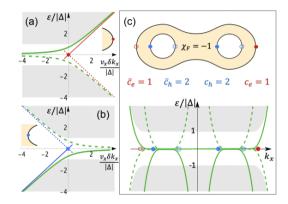
And reev states in a π -junction

 Claim: nonlocal Andreev conductance in a π-junction counts Fermi surfaces



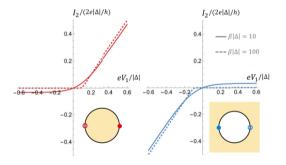
And reev states in a $\pi\text{-junction}$

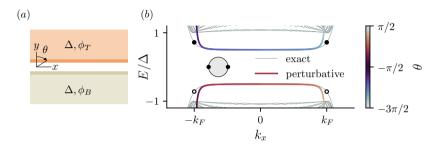
- Claim: nonlocal Andreev conductance in a π-junction counts Fermi surfaces
- Zero modes split near *critical points* and acquire dispersion



And reev states in a π -junction

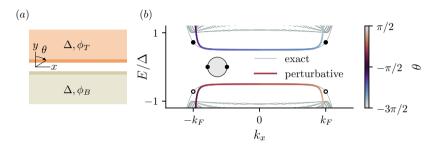
- Claim: nonlocal Andreev conductance in a π-junction counts Fermi surfaces
- Zero modes split near critical points and acquire dispersion
- Adiabaticity prevents backscattering and quantizes charge





Two states at $(k_x, \pm k_y)$ couple due to dispersion nonlinearity $\sim k_y^2$

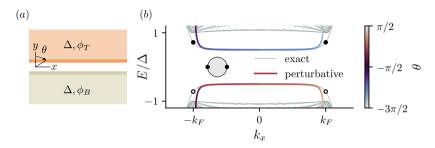
$$H_{\pm} = \Delta \begin{pmatrix} \cos(\delta\phi/2) & \frac{\Delta}{2E_x}\sin^2(\delta\phi/2) \\ \frac{\Delta}{2E_x}\sin^2(\delta\phi/2) & -\cos(\delta\phi/2) \end{pmatrix}$$



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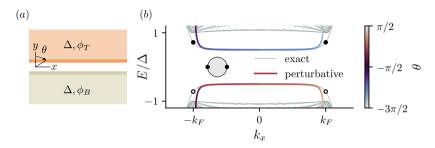
• Positive energies have $\langle k_y \rangle > 0$



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Positive energies have ⟨k_y⟩ > 0
⟨k_y⟩ → 0 ⇒ E → ∞



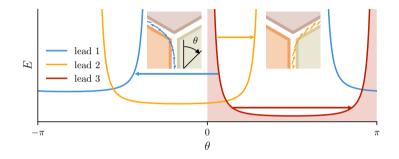
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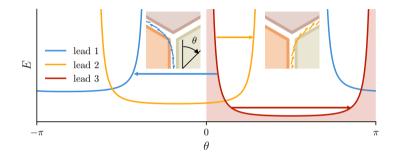
• Positive energies have $\langle k_y \rangle > 0$

 $\blacktriangleright \langle k_y \rangle \to 0 \Rightarrow E \to \infty$

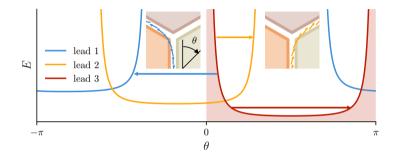
Dispersion forms an energy barrier for momenta!



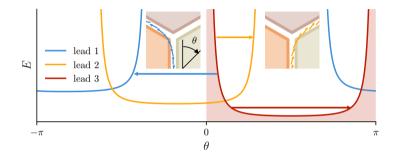
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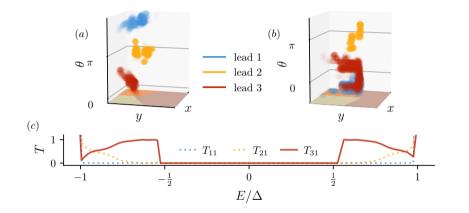


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- \blacktriangleright \Rightarrow only chiral transmission is allowed!

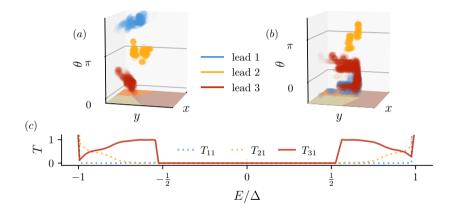


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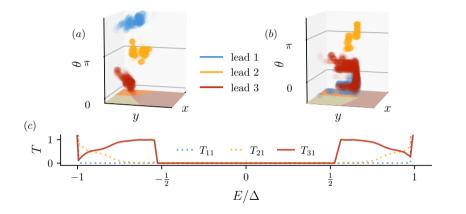
Can we break this pattern?



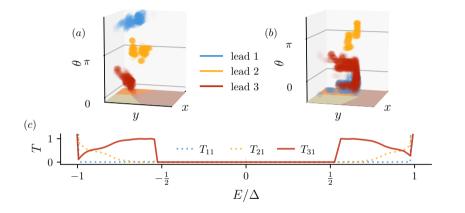
• When turns are $< 90^{\circ}$, transmission is not quantized



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- ► The wave functions start overlapping in the phase space

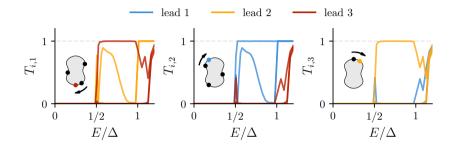


- \blacktriangleright When turns are $<90^\circ,$ transmission is not quantized
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- With smooth turns everything is protected

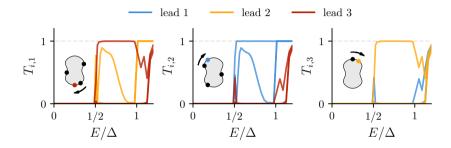


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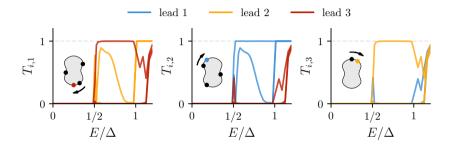
But what about more complex Fermi surfaces?



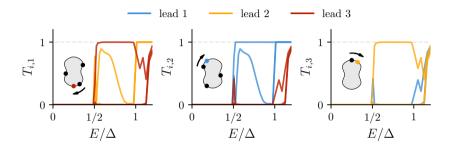
► Peanut-shaped Fermi surface



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- ► Some transmission eigenvalues not quantized...
- However, at least one must stay quantized!



- Peanut-shaped Fermi surface
- ► Some transmission eigenvalues not quantized...
- However, at least one must stay quantized!
- \Rightarrow chiral transport counts Fermi surfaces!

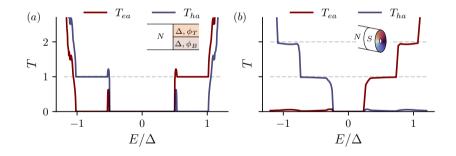
Electrical conductance

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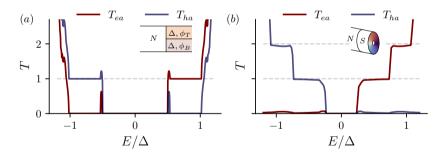
 \blacktriangleright Electrical conductance protected by adiabatically \Rightarrow applies to any ϕ



Electrical conductance

What about electric charge?

- \blacktriangleright Electrical conductance protected by adiabatically \Rightarrow applies to any ϕ
- Same protection applies to a vortex!



- Momentum space separation protects chiral transmission
- Phenomenon generic to gapped multiband systems
- Observation in superconductors possible but hard
- Metamaterials an alternative plaform
- ► A lot of open question about the nature of protection and generality

The end. Questions?

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