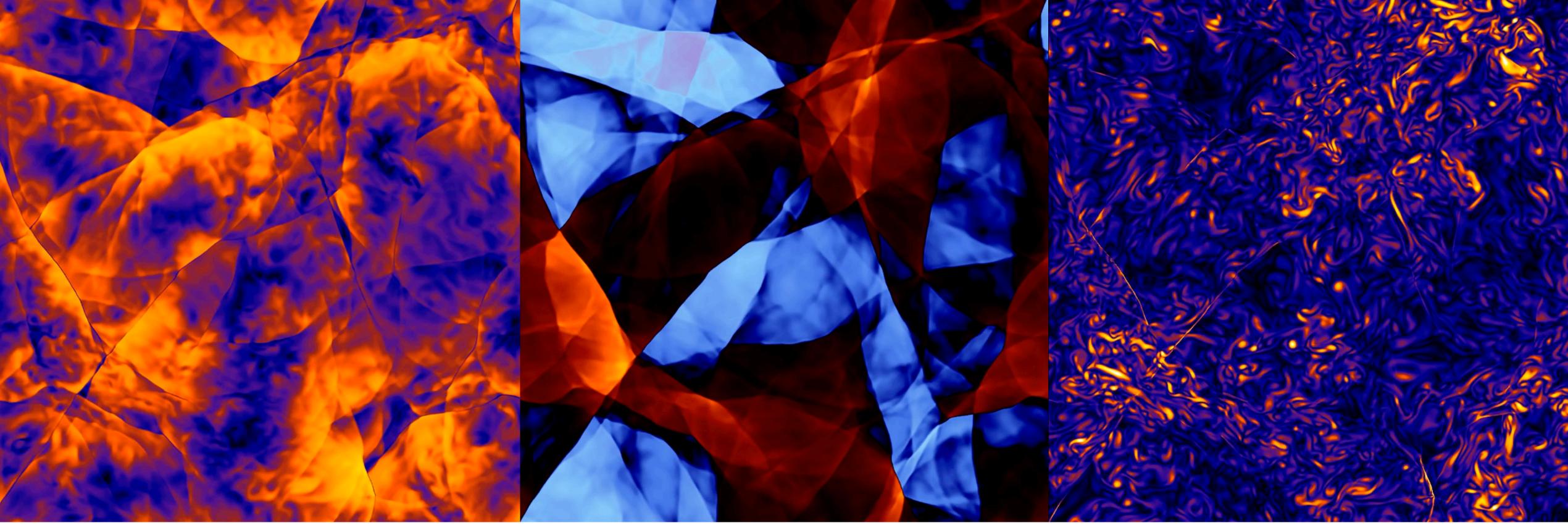
Numerical Simulations of Early Universe Sources of Gravitational Waves - August 7, NORDITA, Sweden





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Work in collaboration with: Alberto Roper Pol, Chiara Caprini, Thomas Konstandin, Ryusuke Jinno, and Henrique Rubira

Based on <u>2508.04263</u>, <u>2409.03651</u>, <u>2209.04369</u> Out today

Gravitational waves

Can we hear them?

Yes!



GW151226

Gravitational waves

Can we hear them?

Yes!

Can we hear GWs from cosmological sources and PTs?

Possibly!



The "sound" of the early universe

Phase transitions

phase transitions

• Phase transitions (PTs) exist in abundance and diversity

• Melting of ice, condensation of snowflakes, **bubbles** in boiling water, ferromagnetism, etc.

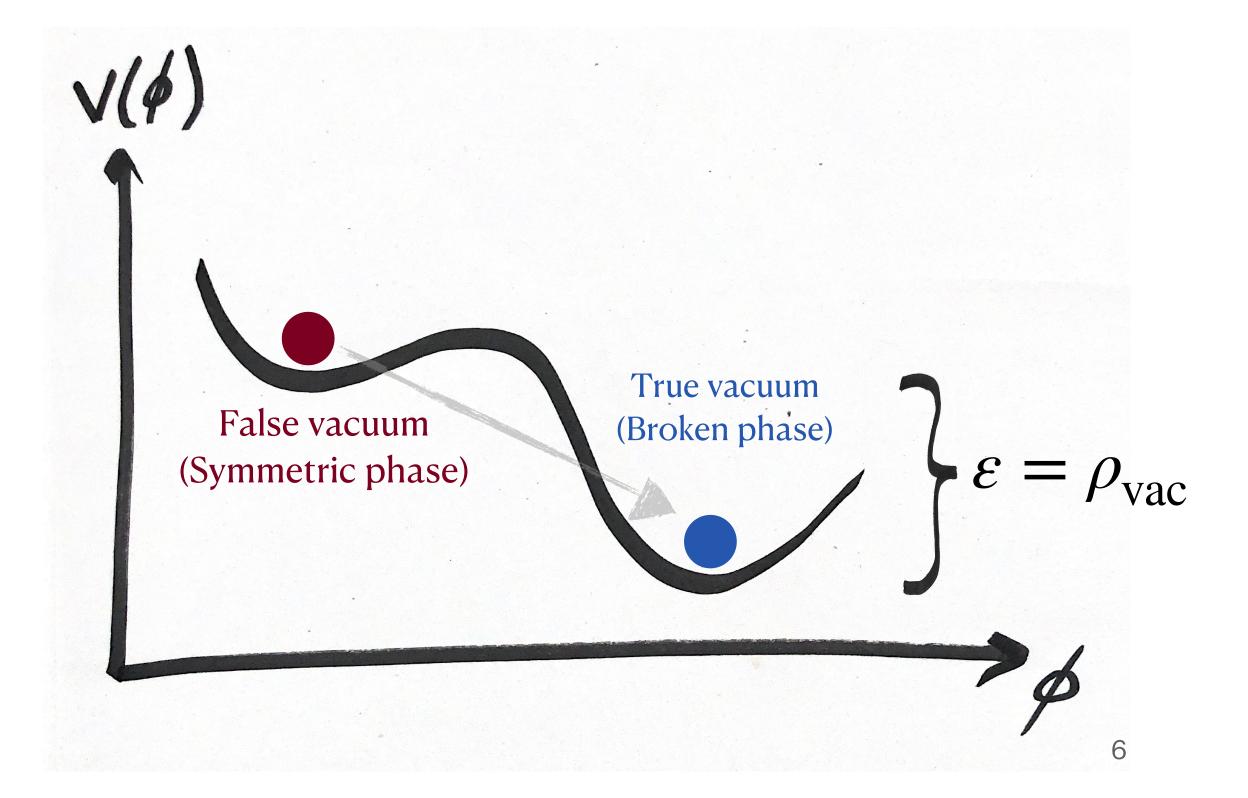
- The list is long, but none pertains to a cosmological context.
- Abundance in diverse physical systems ⇒ cosmological PTs?



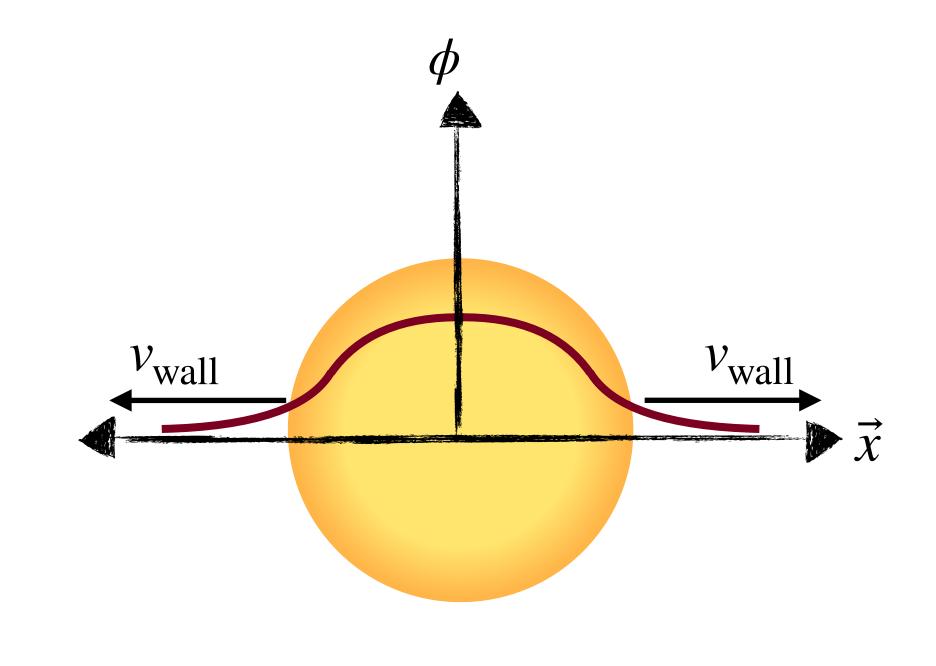
Schematics of a first-order PT

• 1st order PTs proceed through bubble nucleations

Tunneling in field space

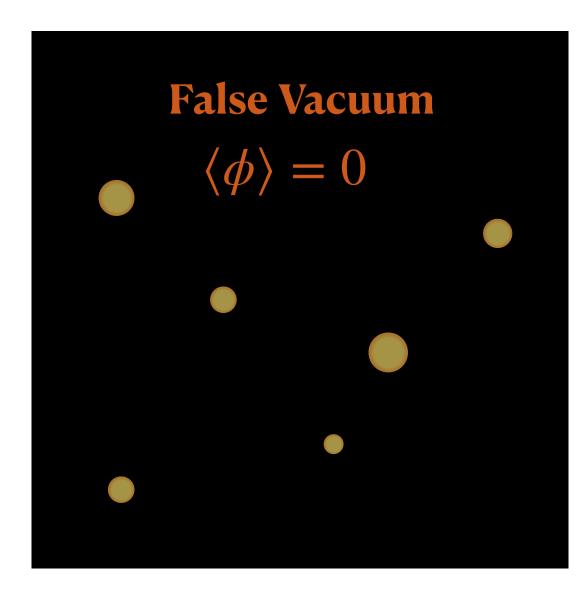


Bubble expansion in physical space



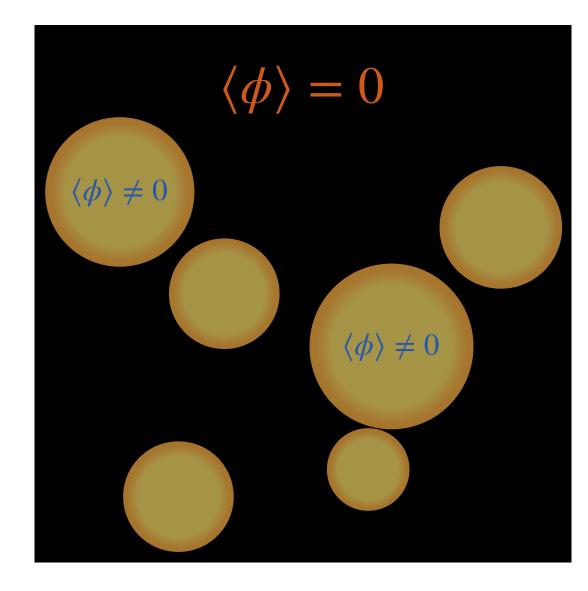
Schematics of a first-order PT

Bubble nucleations



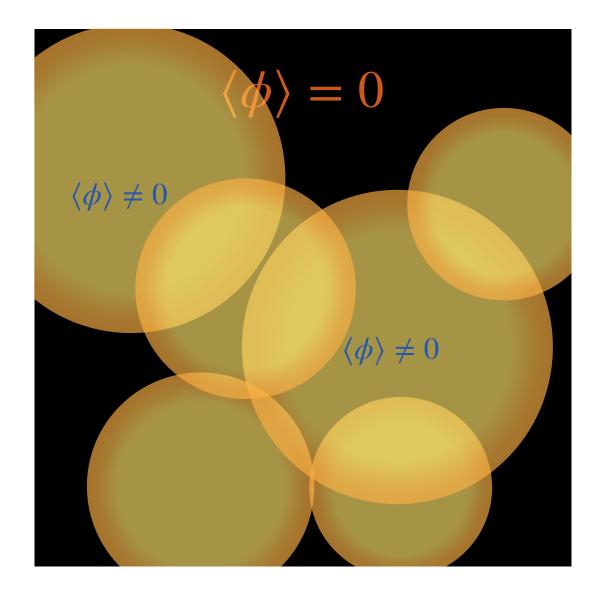
- Universe initially homogeneous
- T drops below T_* and ϕ locally **tunnels** to TV
- Bubbles nucleate and begin expanding

Bubbles expand



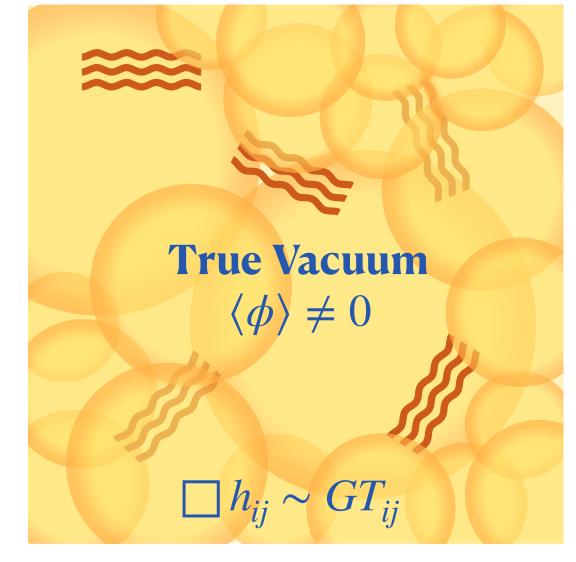
- Expansion proceeds
- Vacuum energy deposits in fluid
- Fluid responds
- Self-similar profiles develop

Bubbles collide



- Bubbles collide
- GW production from anisotropic stress T_{ij}^{plasma}

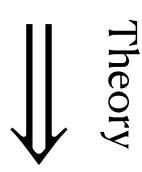
Sound waves



- PT completion
- Long-lasting dynamics (sound waves) source GWs
- Non-linear evolution may produce shocks and turbulence

Particle physics model

- Lagrangian $\mathcal L$
- Thermal field theory
- Effective potential



Simulation

PT parameters

Hydrodynamics

- Character of fluid perturbation
- Linear sound-waves or non-linear evolution
- Detailed evolution of T_{ij}

GW spectrum

• $\Omega_{\text{GW}}(f \mid H_*, \beta, \alpha, v_w)$

PT parameters

- Temperature $T_* : \Longrightarrow H_* \to \text{Horizon size} \to f_{\text{typcial}}$
- Inverse duration β : as in nucleation rate $\Gamma_{\rm nuc} \propto e^{\beta t} \to R_*(\beta)$, $\beta/H \sim \mathcal{O}(100) \Rightarrow$ neglect expansion, $\omega_{\rm peak} \simeq \frac{\beta}{100H} \frac{T}{100{\rm GeV}} {\rm mHz}$
- Strength $\alpha := \rho_{\rm vac}/\rho_{\rm rad}$
- Wall velocity $v_w : \longrightarrow$ energy budget and self-similar profiles

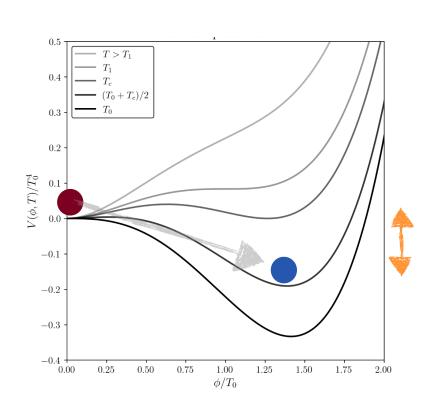
Higgsless simulations

Physical setup

Relativistic hydrodynamics

Thermal

• Perfect fluid:
$$T^{\mu\nu} = (p+\rho)u^{\mu}u^{\nu} - pg^{\mu\nu}$$



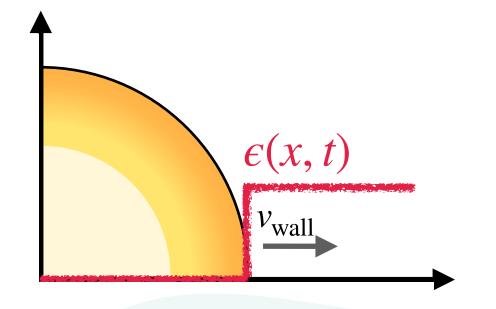
• Bag **E.o.S**.: $\rho = aT^4 + \epsilon$ and. $p = \frac{1}{3}aT^4 - \epsilon$ Vacuum

where
$$\epsilon = \begin{cases} \epsilon_+ \equiv \frac{3\alpha}{4}w, & \text{in the symmetric phase,} \\ 0, & \text{in the broken phase} \end{cases}$$

is the energy difference between the symmetric and broken phases.

• E.o.M. given by total energy-momentum conservation: $\partial_{\mu}T^{\mu\nu}=0$

• 4 dynamical variables: $K^{\mu} \equiv T^{0\mu}$



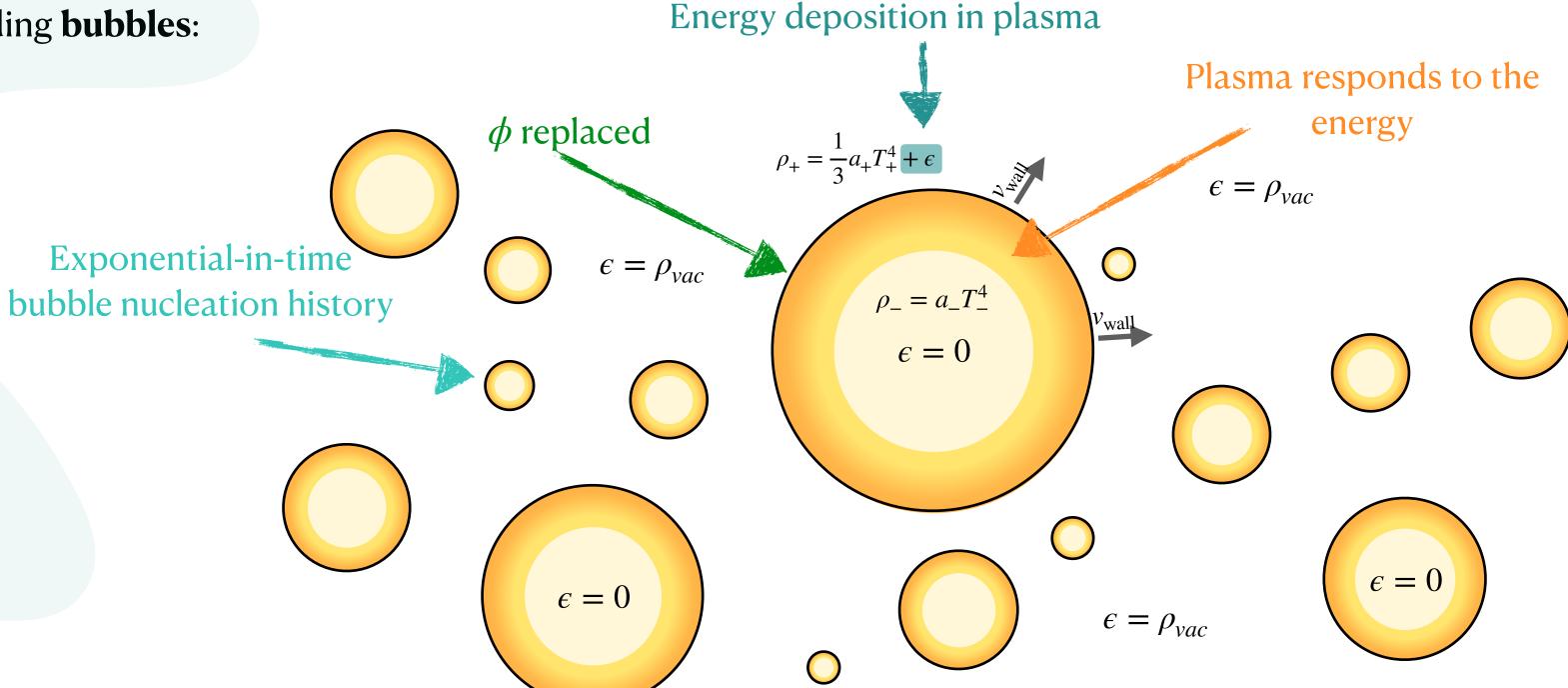
Physical setup

Bubbles and the Equation of State

- Bag E.o.S. suggests natural way to embed the expanding bubbles:
 - Prescribe that $\epsilon = \epsilon(t, \vec{x})$ is a **function** of **space** and **time**
 - Bubbles are radially expanding regions where $\epsilon=0$

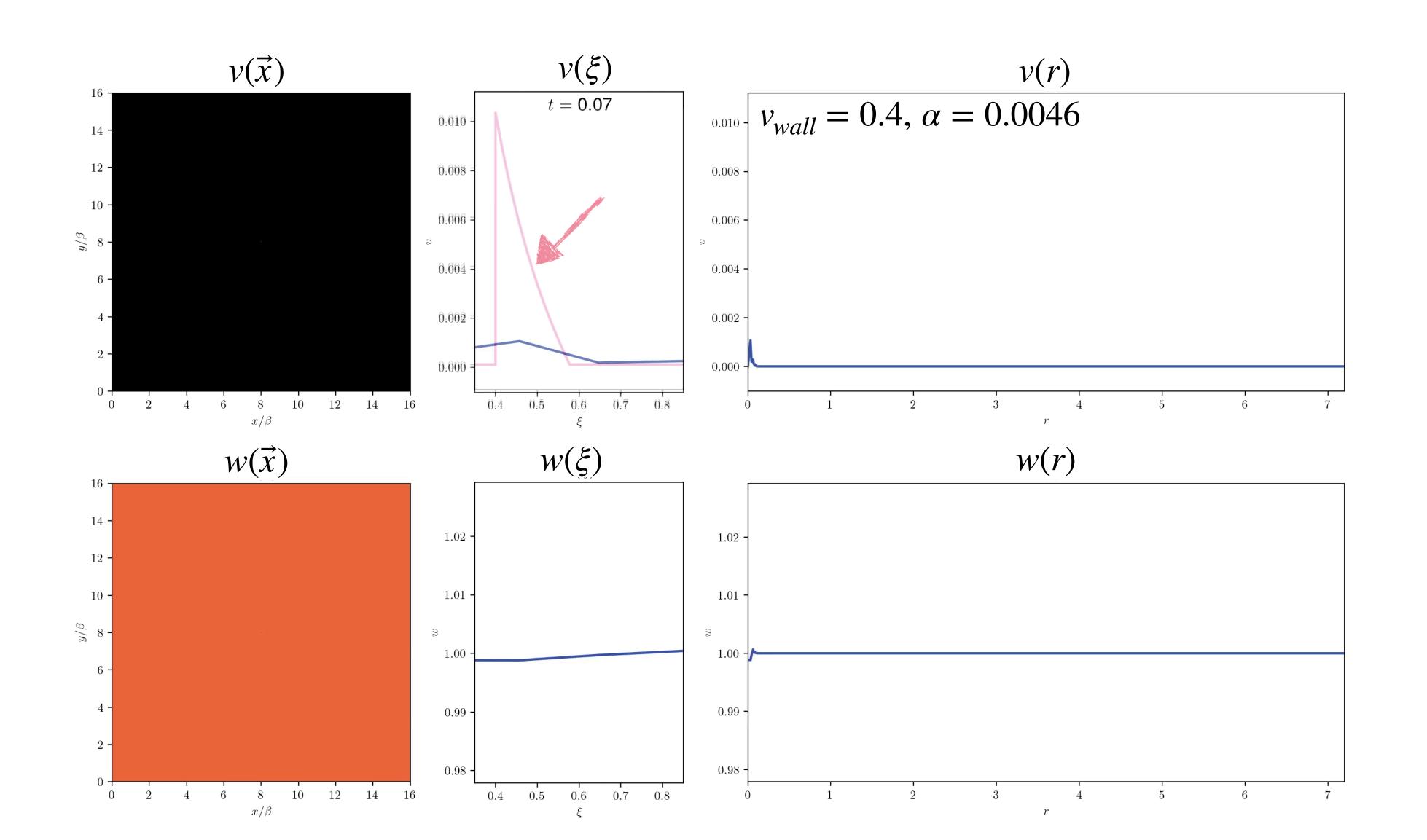
so that
$$e(t, \vec{x} | \alpha) = \begin{cases} 0, & \text{inside bubbles,} \\ \frac{3\alpha}{4}w, & \text{outside bubbles} \end{cases}$$

- To determine $\varepsilon(t, \vec{x})$, we **need** to specify: v_w , α , and an exponential **bubble nucleation history**.
- Nucleation rate $\Gamma(t) \simeq \Gamma_* e^{\beta(t-t_*)}$
- Thus, with these simple 3 ingredients, $\epsilon(t, \vec{x} \mid \alpha)$ is predetermined input to the simulation, effectively replacing $\phi \to \epsilon(t, \vec{x} \mid \alpha)$
- Couples to fluid through $p = \frac{1}{6} \left[-2K^0 4\epsilon(t, \vec{x} \mid \alpha) + \sqrt{(4K^0 4\epsilon(t, \vec{x} \mid \alpha))^2 12K^iK^i} \right]$



as in $T^{ij}\left[K^{\mu}\right] = \frac{K^i K^j}{K^0 + p} + p$

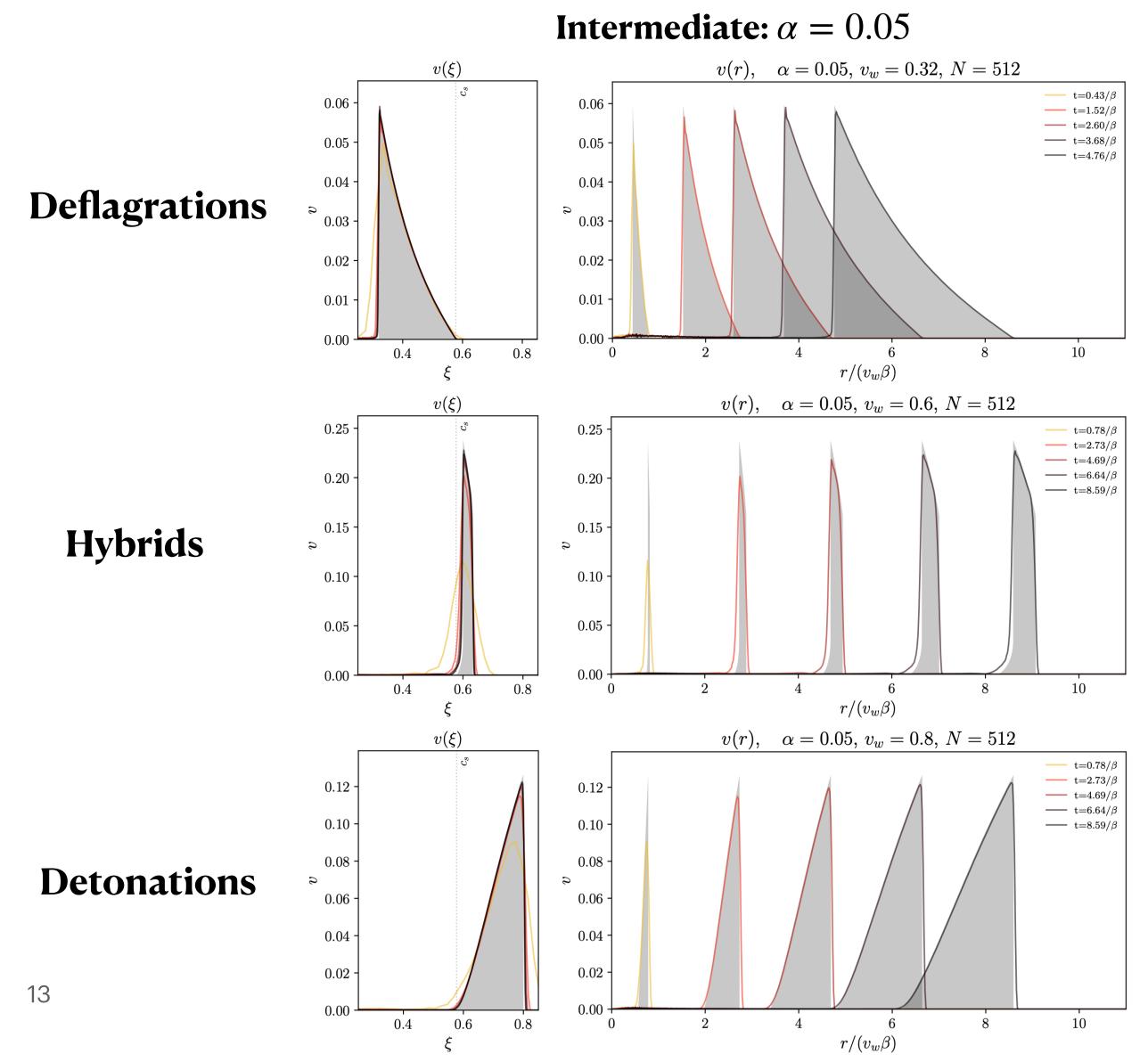
Convergence of self-similar solutions



Convergence of self-similar solutions

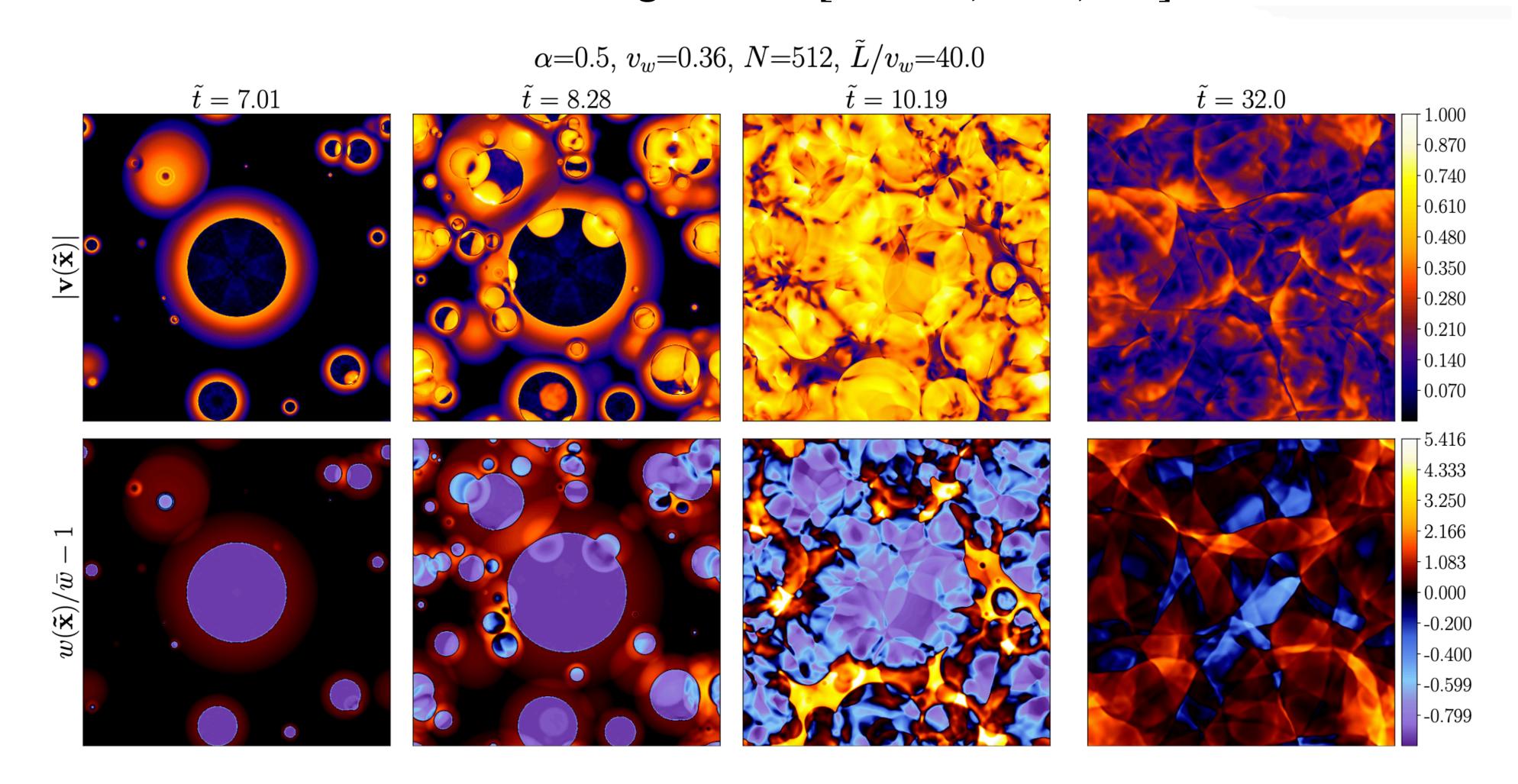
- Remarkable agreement with analytical solutions
- Resolves and handles shocks waves

Convergence sufficiently fast that modest grid size ensures most bubbles in self-similar regime before collision.



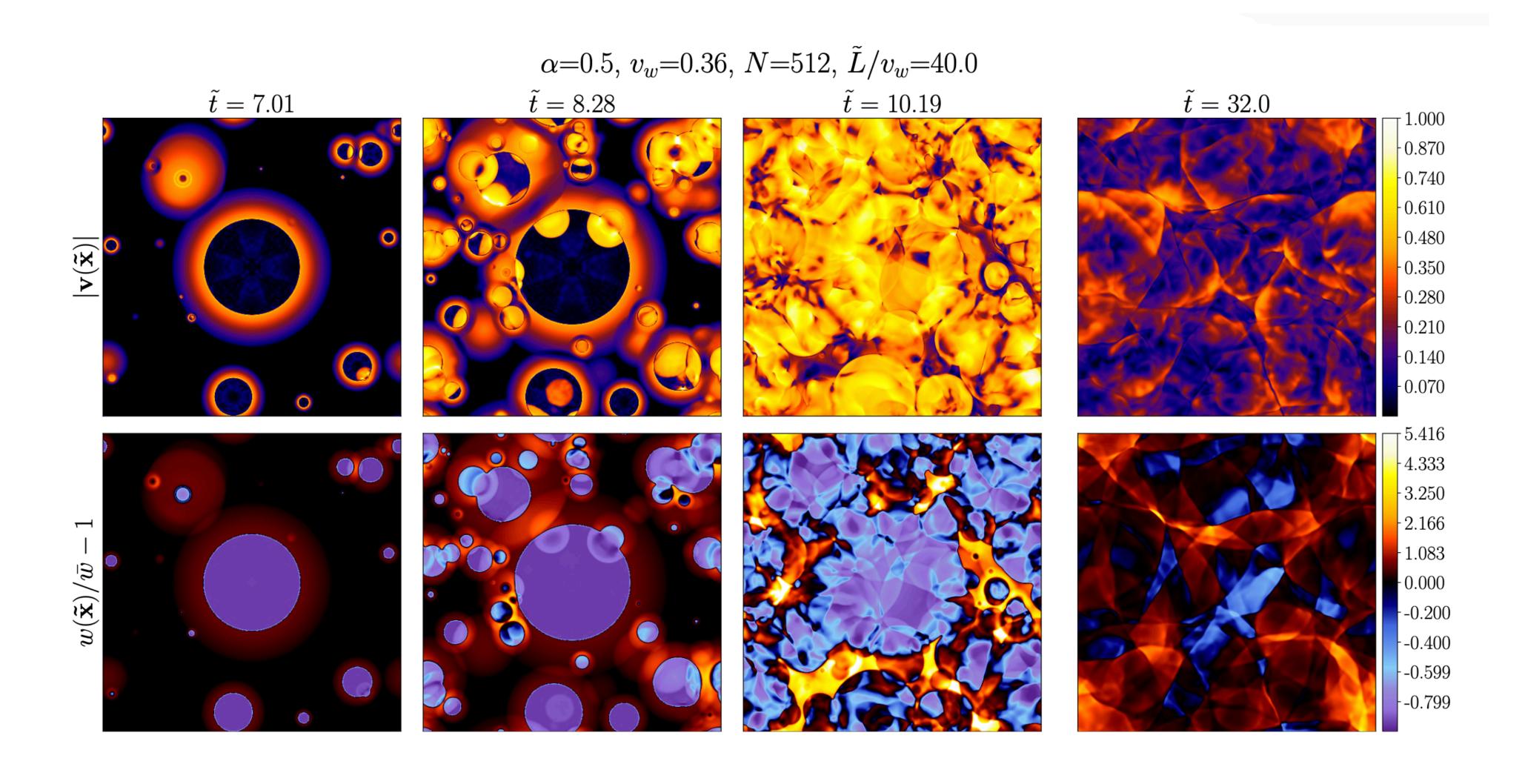
We run simulations for...

...different wall velocities $v_w \in [0.32, 0.8]$ and strengths $\alpha \in [0.0046, 0.05, 0.5]$

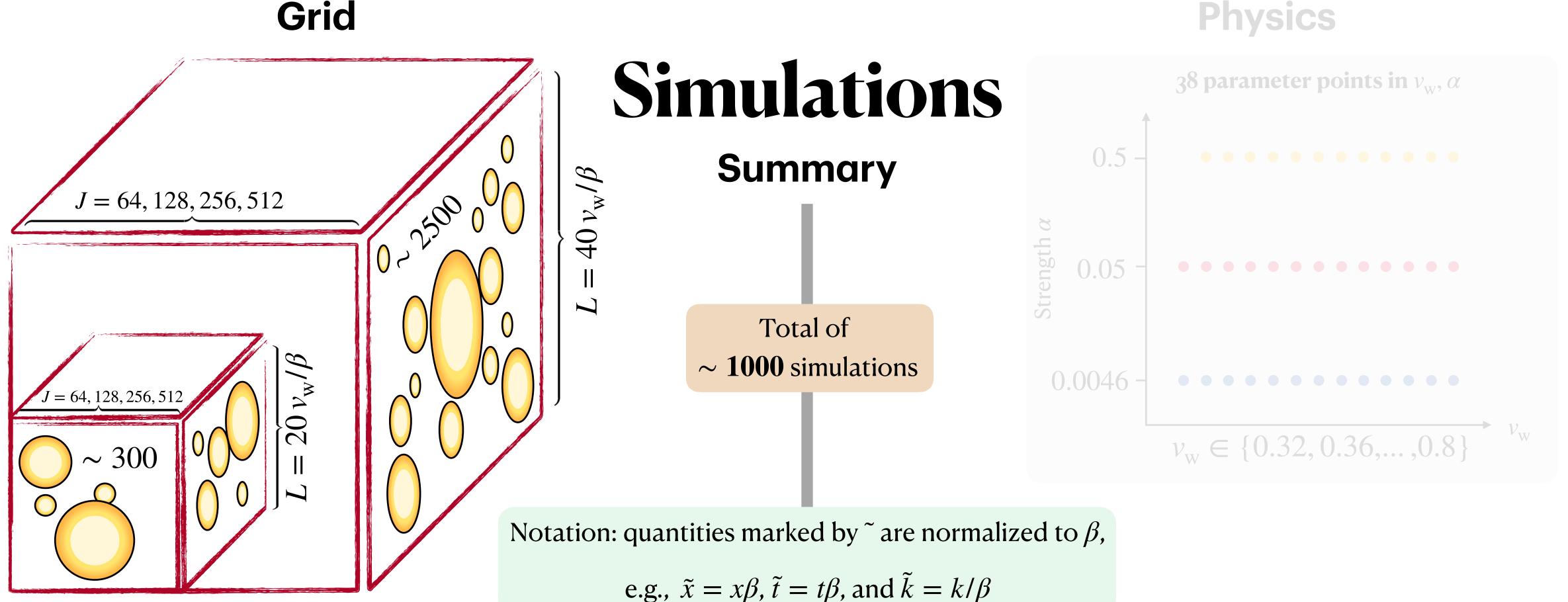


To scan parameter space.

More than 1000 simulations were performed







- **Two box sizes** allows to capture both **IR** (more statistics, ~ 2500 bubbles) and UV (higher resolution per bubble, ~ 300 bubbles)
- A series of resolutions allows to study the convergence of energies and to develop an extrapolation scheme.

- Systematic parameter scan in v_w , α including strong PTs
- Realistic exponential-in-time bubble nucleation histories
- 10 different bubble nucleation histories for $v_w \in \{0.32, 0.6, 0.8\}$ and all α 's to study statistical variance of extracted parameters



J = 64, 128, 256, 512

= 64, 128, 256, 51

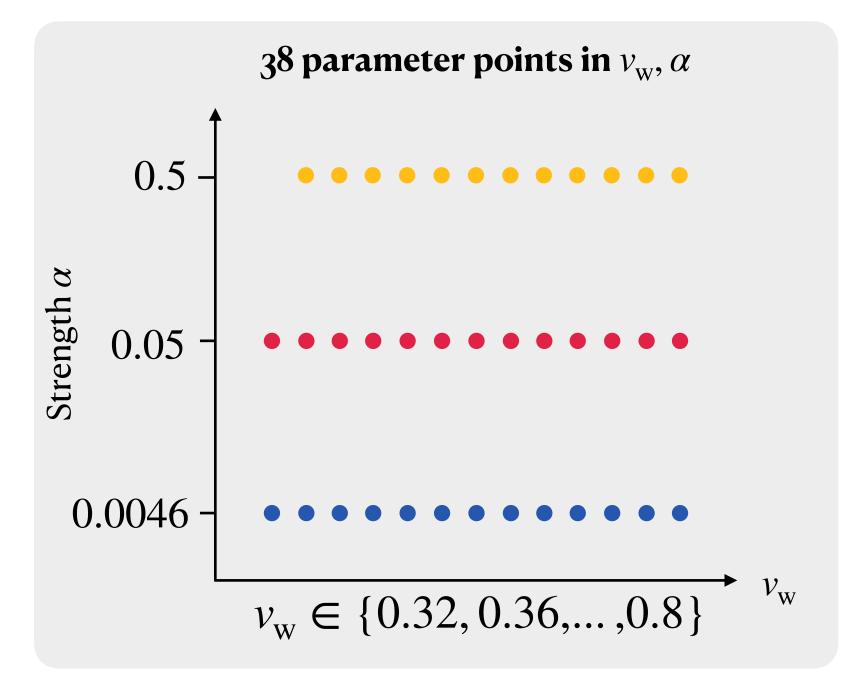
~ 300

Simulations

Summary

Total of ~ 1000 simulations





Notation: quantities marked by $\tilde{\beta}$ are normalized to β ,

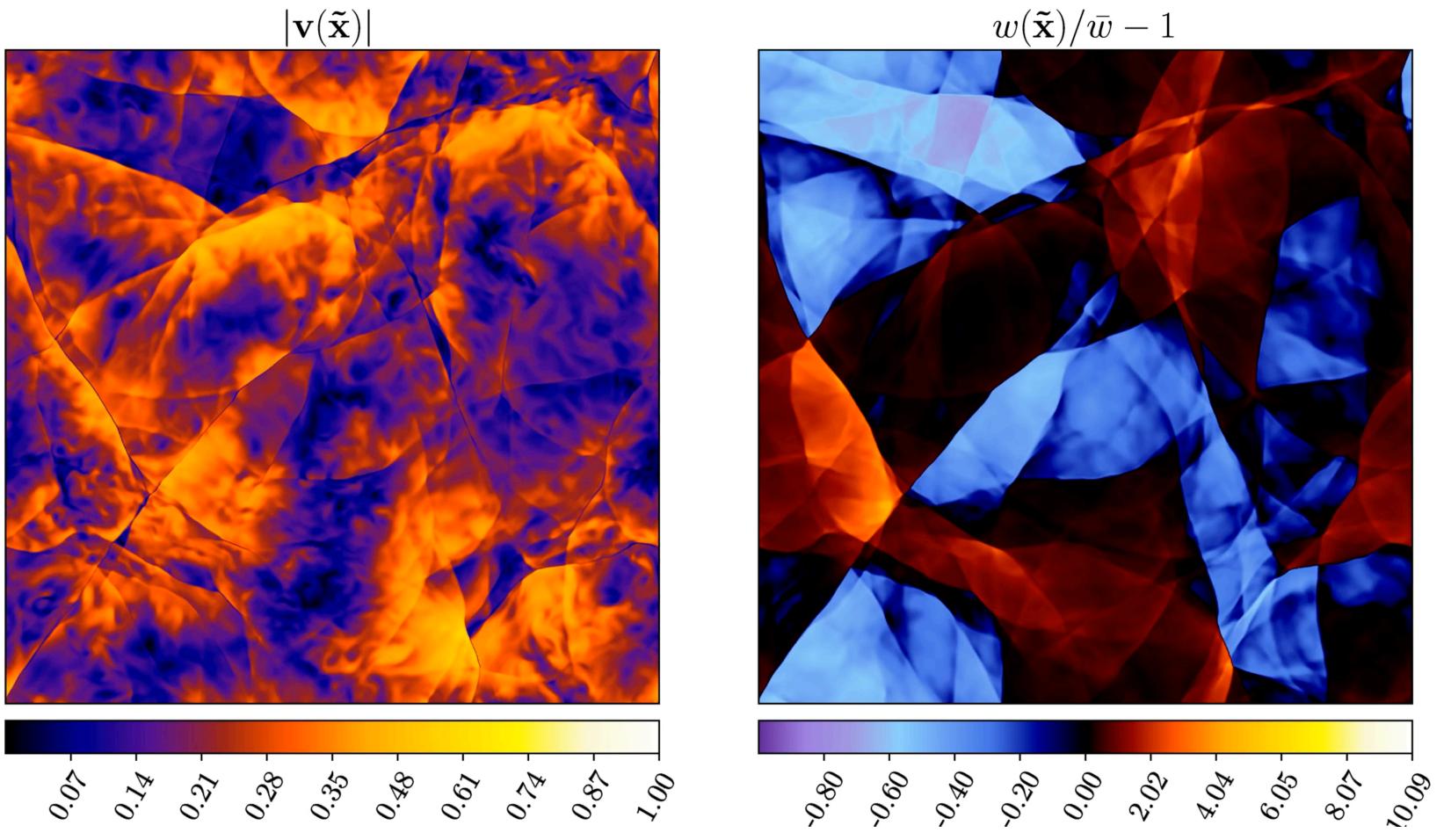
e.g.,
$$\tilde{x} = x\beta$$
, $\tilde{t} = t\beta$, and $\tilde{k} = k/\beta$

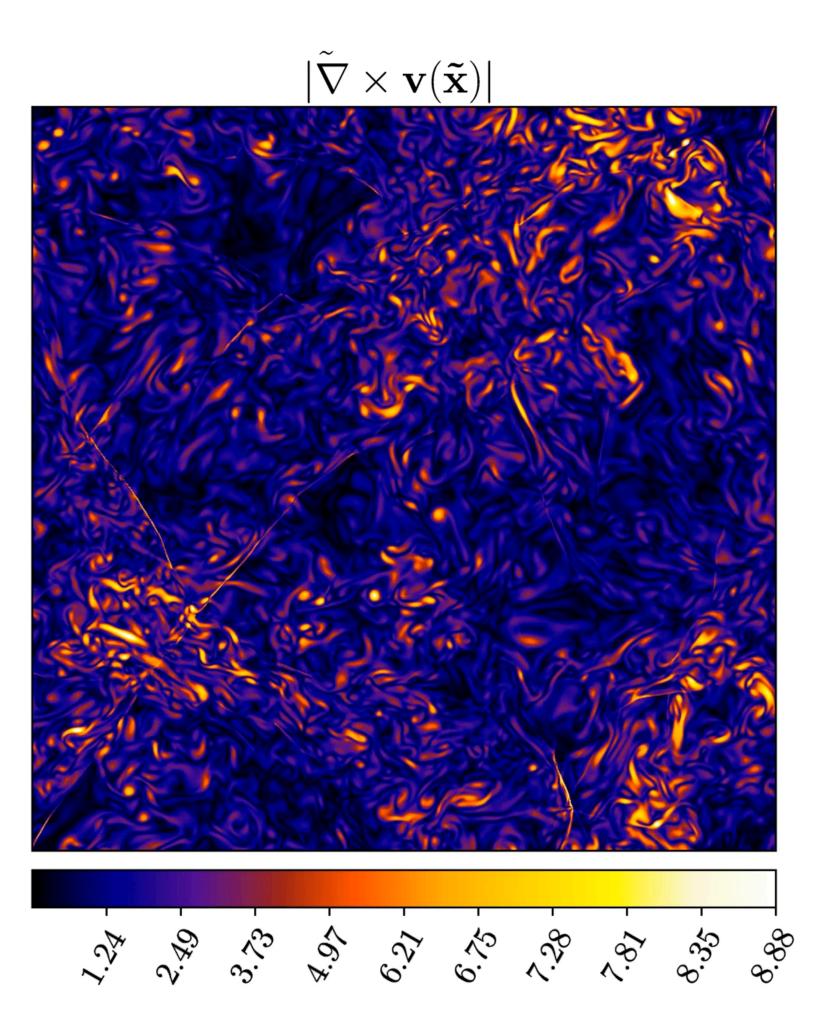
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Example

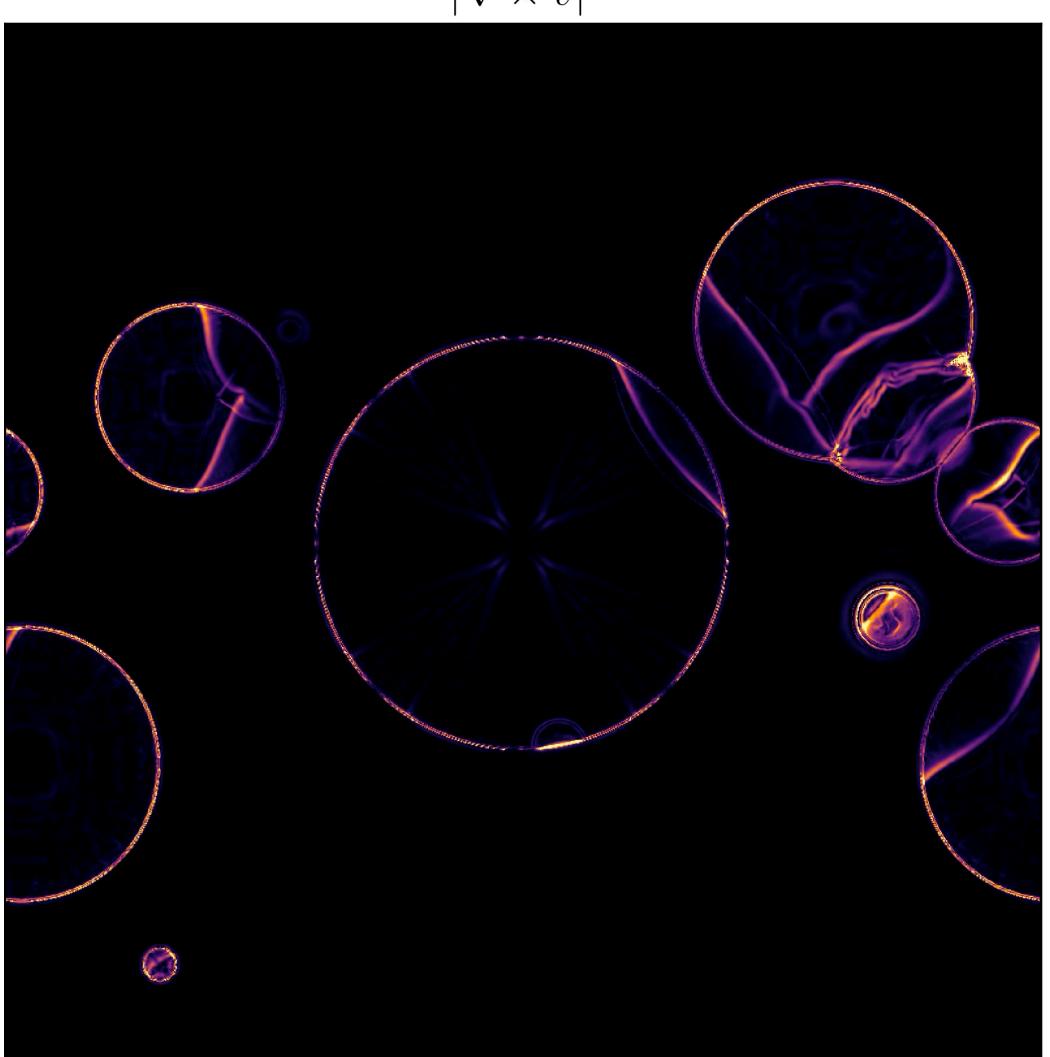
 $\alpha = 0.5, v_w = 0.36, J = 1024, L/v_w = 20.0$ $\tilde{t} = 20.75$



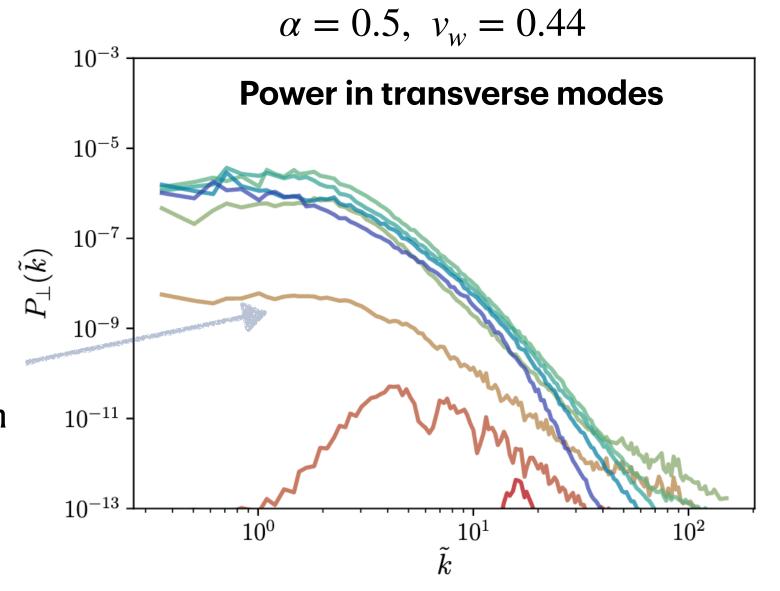


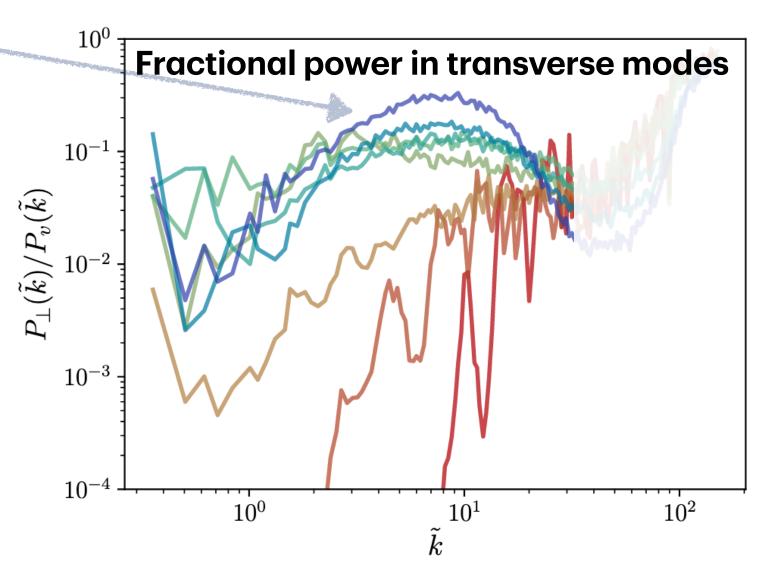
Departure from linearity

|
abla imes v| Vorticity



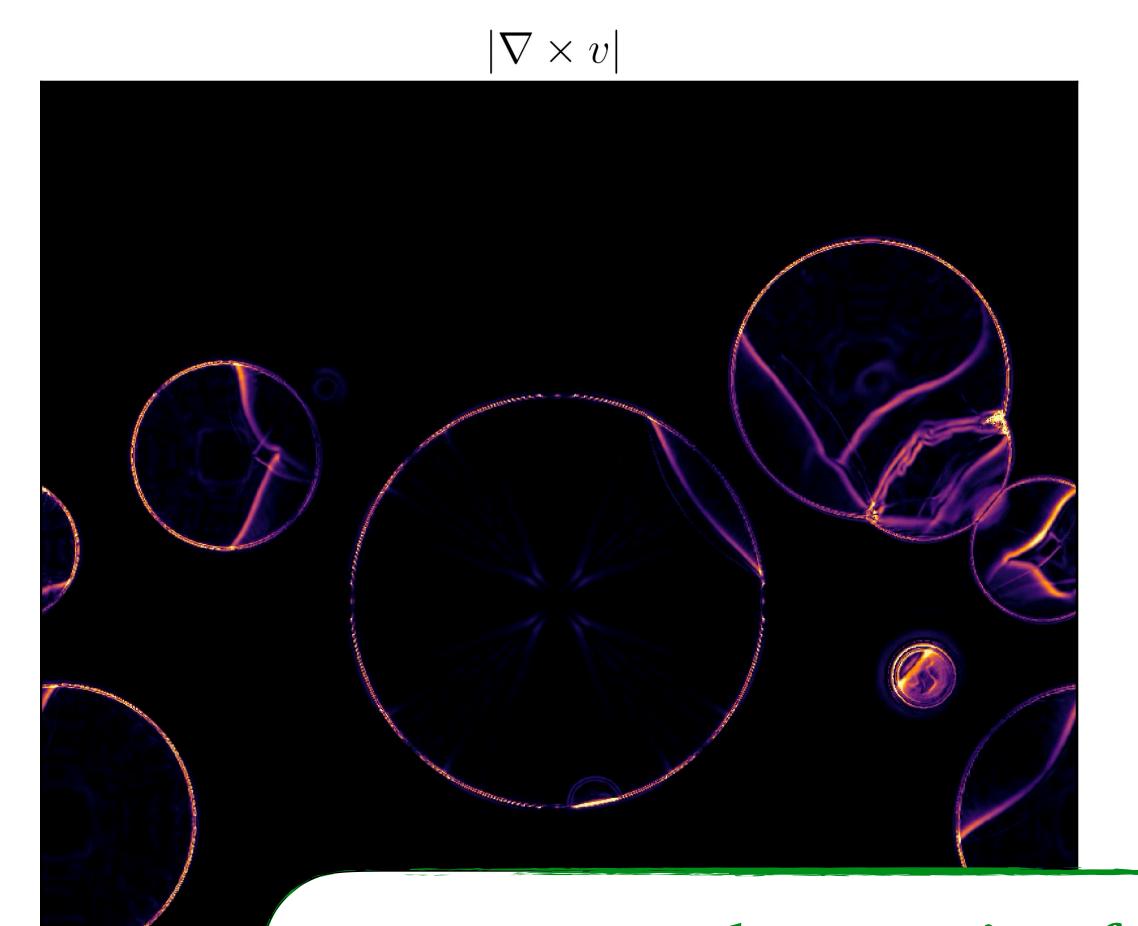
- $P_{\perp}(k)$ initially vanishing
- Transfer of energy from longitudinal to transverse modes
- Transverse velocity spectrum grows with time
- Evolution is non-linear
- 30% of power in transverse modes





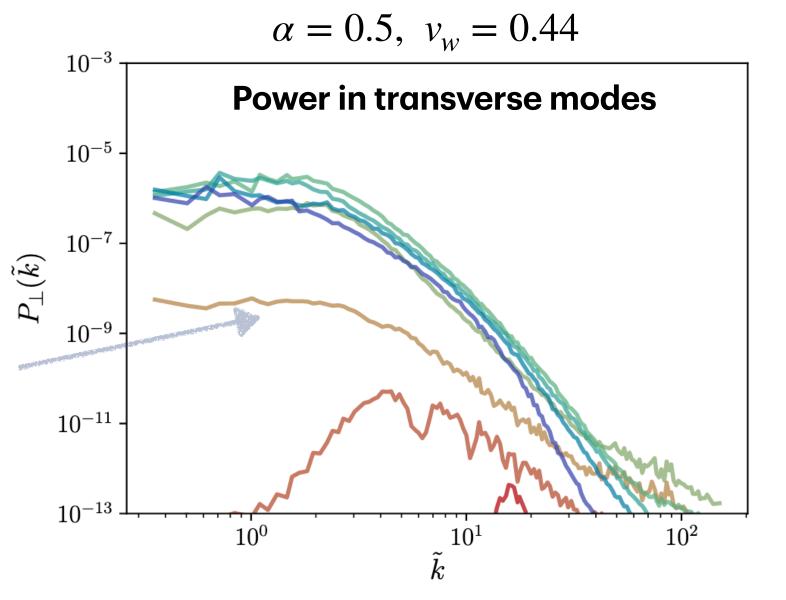
Departure from linearity

Vorticity



0

- $P_{\perp}(k)$ initially vanishing
- Transfer of energy from longitudinal to transverse modes
- Transverse velocity spectrum grows with time
- Evolution is non-linear
- 30% of power in transverse modes



Fractional power in transverse modes

 10^{2}

Testament to the necessity of using numerical simulations to capture realistic fluid evolution and to derive accurate GW predictions.

$$\Omega_{\text{GW}}(\mathbf{k}) = \frac{1}{\rho_{\text{c}}} \frac{d\rho_{\text{GW}}}{d \ln \mathbf{k}} = 3 \mathcal{T}_{\text{GW}} \mathcal{J}(\mathbf{k}) = 3 \mathcal{T}_{\text{GW}} \frac{k}{2} H_*^2 \int_{t_*}^{t_{\text{fin}}} \int_{t_*}^{t_{\text{fin}}} E_{\Pi}(t_1, t_2, \mathbf{k}) \cos \mathbf{k} (t_1 - t_2) dt_1 dt_2 \equiv 3 \mathcal{T}_{\text{GW}} \left(\frac{H_*}{\beta}\right)^2 \mathcal{J}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{q})$$

Complete knowledge of $T_{\mu\nu}$ from simulations

$$\mathcal{F}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{q}) = \frac{\tilde{q}^3}{4\pi^2 \tilde{V} \tilde{\bar{\rho}}_*^2} \int \frac{d\Omega_k}{4\pi} \left[\Lambda_{ij,kl} \tilde{T}_{ij}(\tilde{q}, \tilde{\mathbf{k}}) \tilde{T}_{kl}^*(\tilde{q}, \tilde{\mathbf{k}}) \right]_{\tilde{q} = |\tilde{\mathbf{k}}|}$$

Source GWs from $\tilde{t}_* = 16$ until $\tilde{t}_{fin} = 32$

$$\Omega_{\text{GW}}(\mathbf{k}) = \frac{1}{\rho_{\text{c}}} \frac{\mathrm{d}\rho_{\text{GW}}}{\mathrm{d}\ln\mathbf{k}} = 3\,\mathcal{T}_{\text{GW}} \mathcal{J}(\mathbf{k}) = 3\,\mathcal{T}_{\text{GW}} \frac{k}{2} H_*^2 \int_{t_*}^{t_{\text{fin}}} E_{\Pi}\left(t_1, t_2, k\right) \cos k\left(t_1 - t_2\right) dt_1 dt_2 \equiv 3\,\mathcal{T}_{\text{GW}} \left(\frac{H_*}{\beta}\right)^2 \mathcal{J}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{q})$$

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• $f(t_{-}, k)$ parametrizes the strength of the correlations

UETC $E_{\Pi}(t_1, t_2, k) = 2k^2K^2 \times f(t_-, k)$

from sound waves

- Stationary source \Longrightarrow
- Kinetic energy $\propto K$ is constant
- Correlations captured by $f(t_-, k)$ depending only on $t_- = t_2 t_1$

E.g. 1304.2433 and 2308.12943

Complete knowledge of $T_{\mu\nu}$ from simulations

$$\mathcal{F}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{q}) = \frac{\tilde{q}^3}{4\pi^2 \tilde{V} \tilde{\bar{\rho}}_*^2} \int \frac{d\Omega_k}{4\pi} \left[\Lambda_{ij,kl} \tilde{T}_{ij}(\tilde{q}, \tilde{\mathbf{k}}) \tilde{T}_{kl}^*(\tilde{q}, \tilde{\mathbf{k}}) \right]_{\tilde{q} = |\tilde{\mathbf{k}}|}$$

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• $f(t_{-},k)$ parametrizes the strength of the correlations

UETC
$$E_{\Pi}(t_1, t_2, k) = 2k^2K^2(t_+)f(t_-, k)$$
 • Locally Stationary source \Longrightarrow

from source waves a damped source

- - Kinetic energy $\propto K$ is constant decaying, depending on $t_+ = (t_2 + t_1)/2$
 - Correlations captured by $f(t_-, k)$ depending only on $t_- = t_2 t_1$

Complete knowledge of $T_{\mu\nu}$ from simulations

$$\mathcal{F}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{q}) = \frac{\tilde{q}^3}{4\pi^2 \tilde{V} \tilde{\rho}_*^2} \int \frac{d\Omega_k}{4\pi} \left[\Lambda_{ij,kl} \tilde{T}_{ij}(\tilde{q}, \tilde{\mathbf{k}}) \tilde{T}_{kl}^*(\tilde{q}, \tilde{\mathbf{k}}) \right]_{\tilde{q} = |\tilde{\mathbf{k}}|}$$

Source GWs from $\tilde{t}_* = 16$ until $\tilde{t}_{fin} = 32$

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$$\mathcal{F}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$$

Generalization of GW spectrum parametrization from sound waves

GW efficiency

Integrated kinetic energy

Characteristic scale of fluid perturbations

Spectral shape

GW spectrum from sound waves damped sources

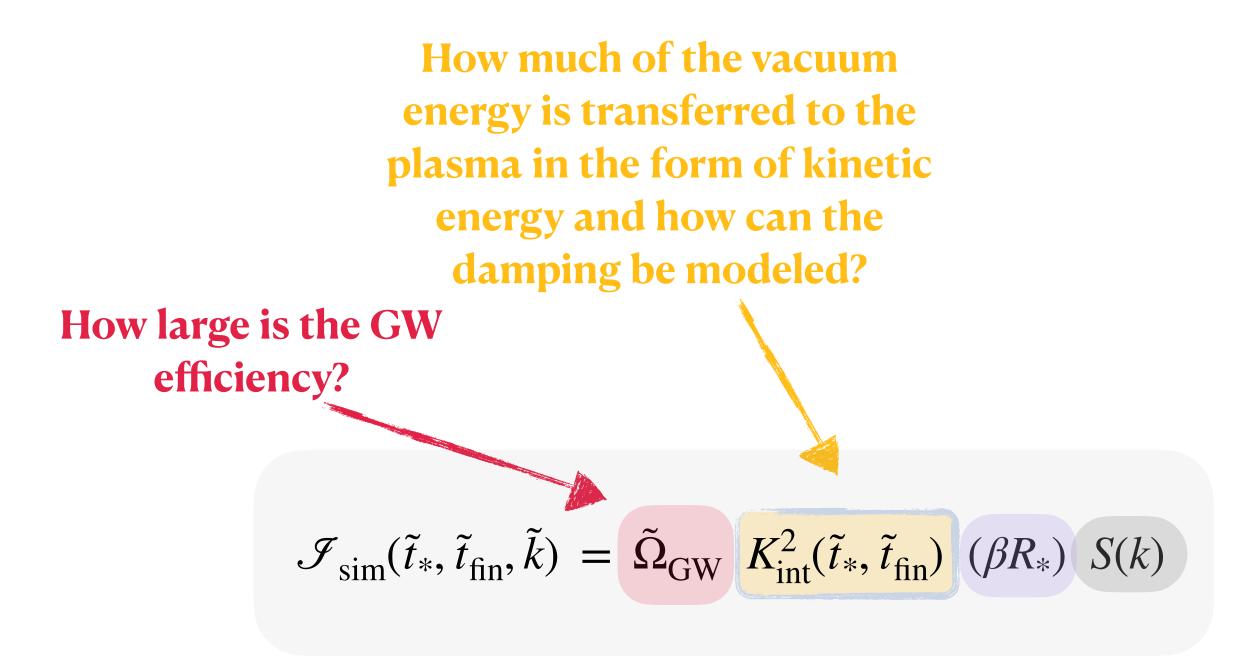
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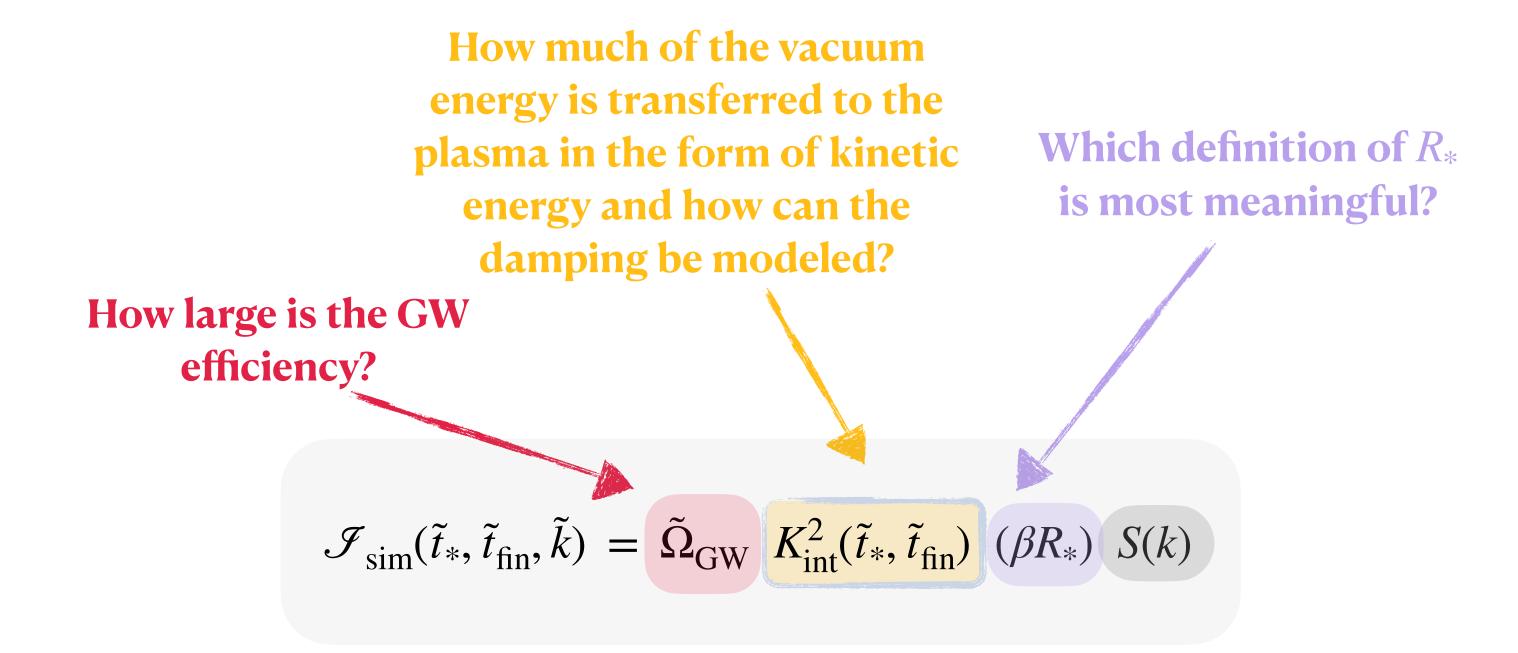
How large is the GW efficiency?

$$\mathcal{F}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} \frac{K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}})}{K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}})} (\beta R_*) S(k)$$

GW spectrum from sound waves damped sources



GW spectrum from sound waves damped sources



GW spectrum from sound waves damped sources

How much of the vacuum energy is transferred to the plasma in the form of kinetic energy and how can the damping be modeled?

Which definition of R_* is most meaningful?

How large is the GW efficiency?

 $\mathcal{F}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$

What is the spectral shape?
In particular, where is the peak, and what are the IR and UV slopes?

Results

$$\mathcal{I}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$$

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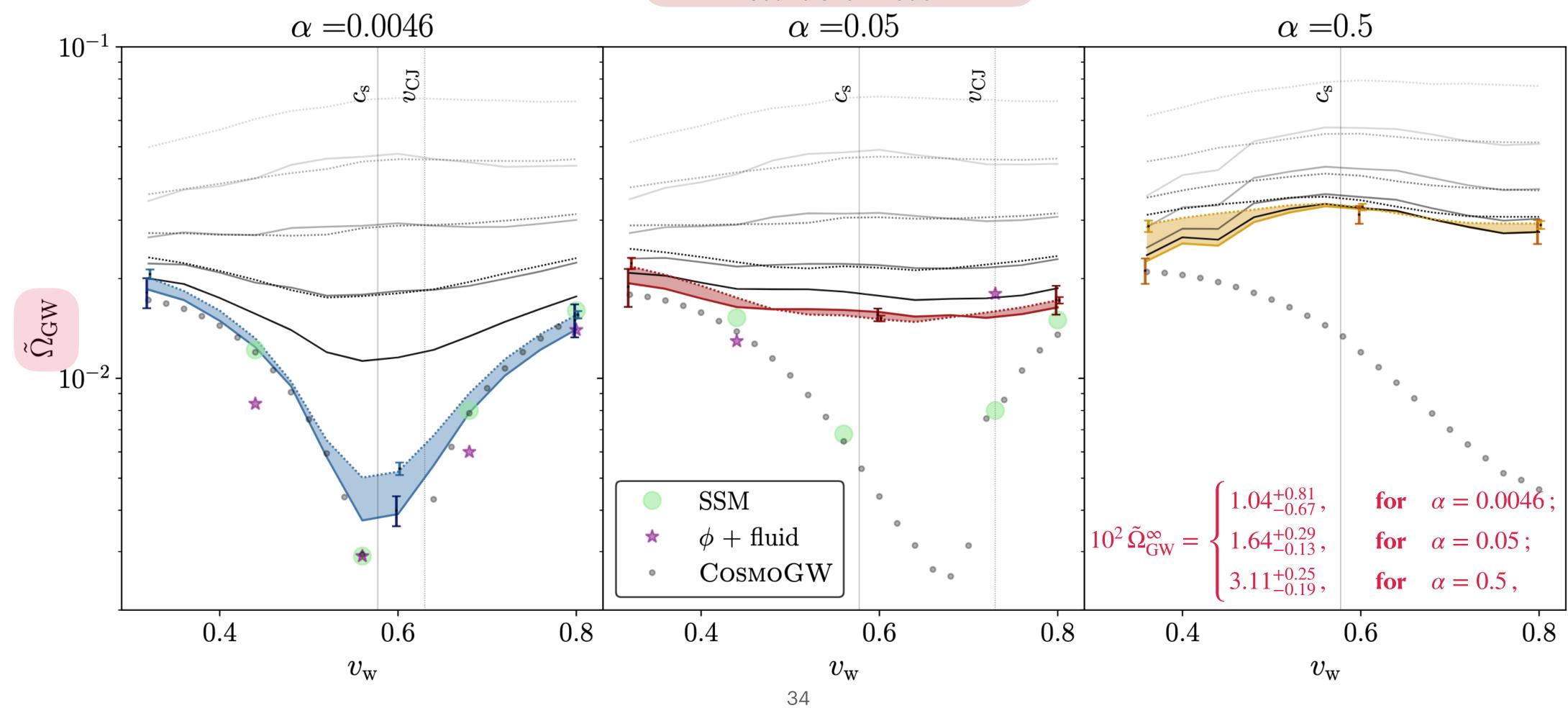
Results GW efficiency

 $\tilde{\Omega}_{\text{GW}} = \frac{\mathcal{I}_{\text{sim}}^{\text{int}}(\tilde{t}_{\text{init}}, \tilde{t}_{\text{end}})}{K_{\text{int}}^2(\tilde{t}_{\text{init}}, \tilde{t}_{\text{end}})(\beta R_*)}$

Weak: Agreement with sound shell model and ϕ + fluid simulations

Interm: Agreement with ϕ + fluid simulations but deviations from sound shell model

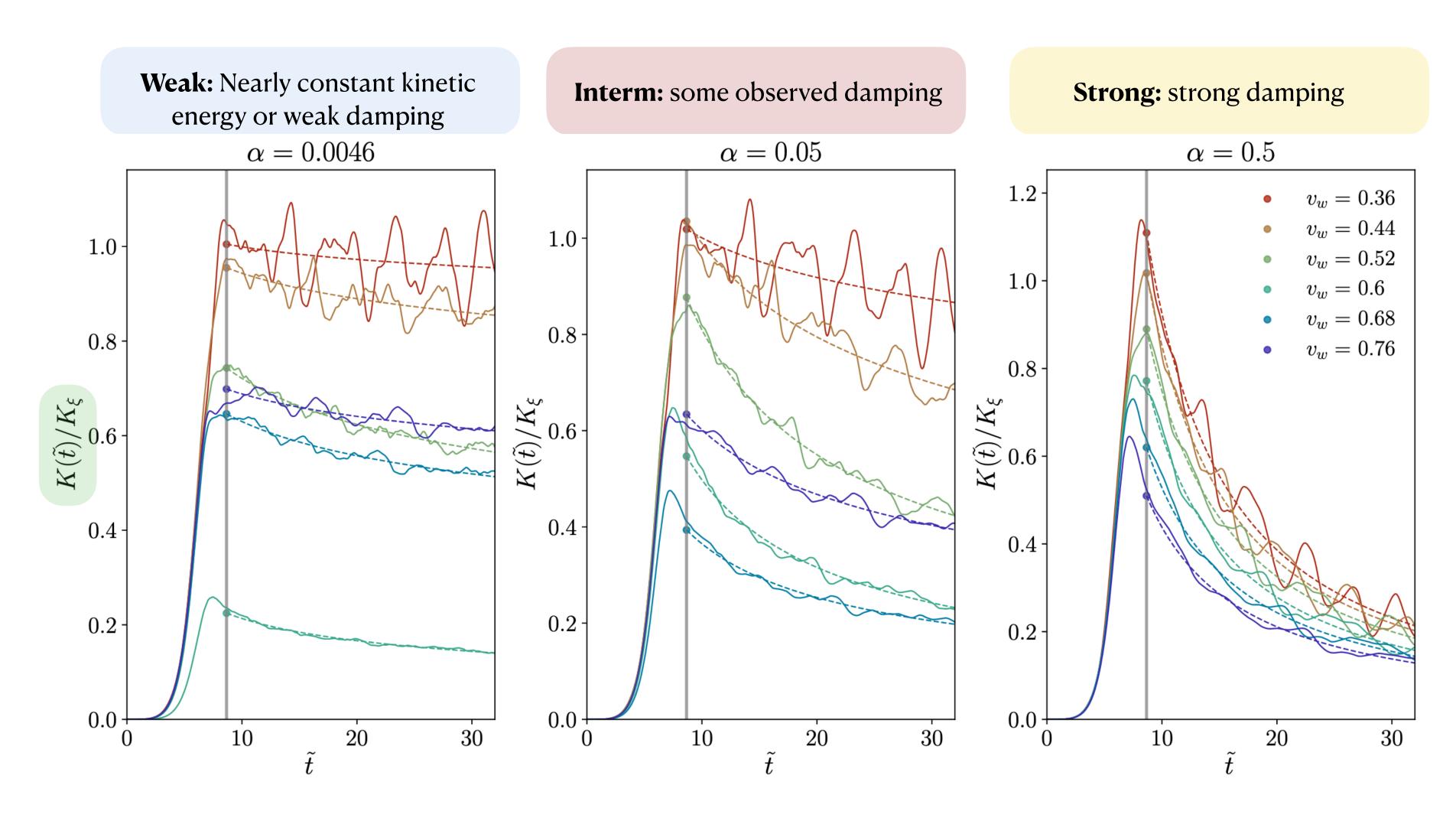
Strong: new prediction



$$\mathcal{F}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$$

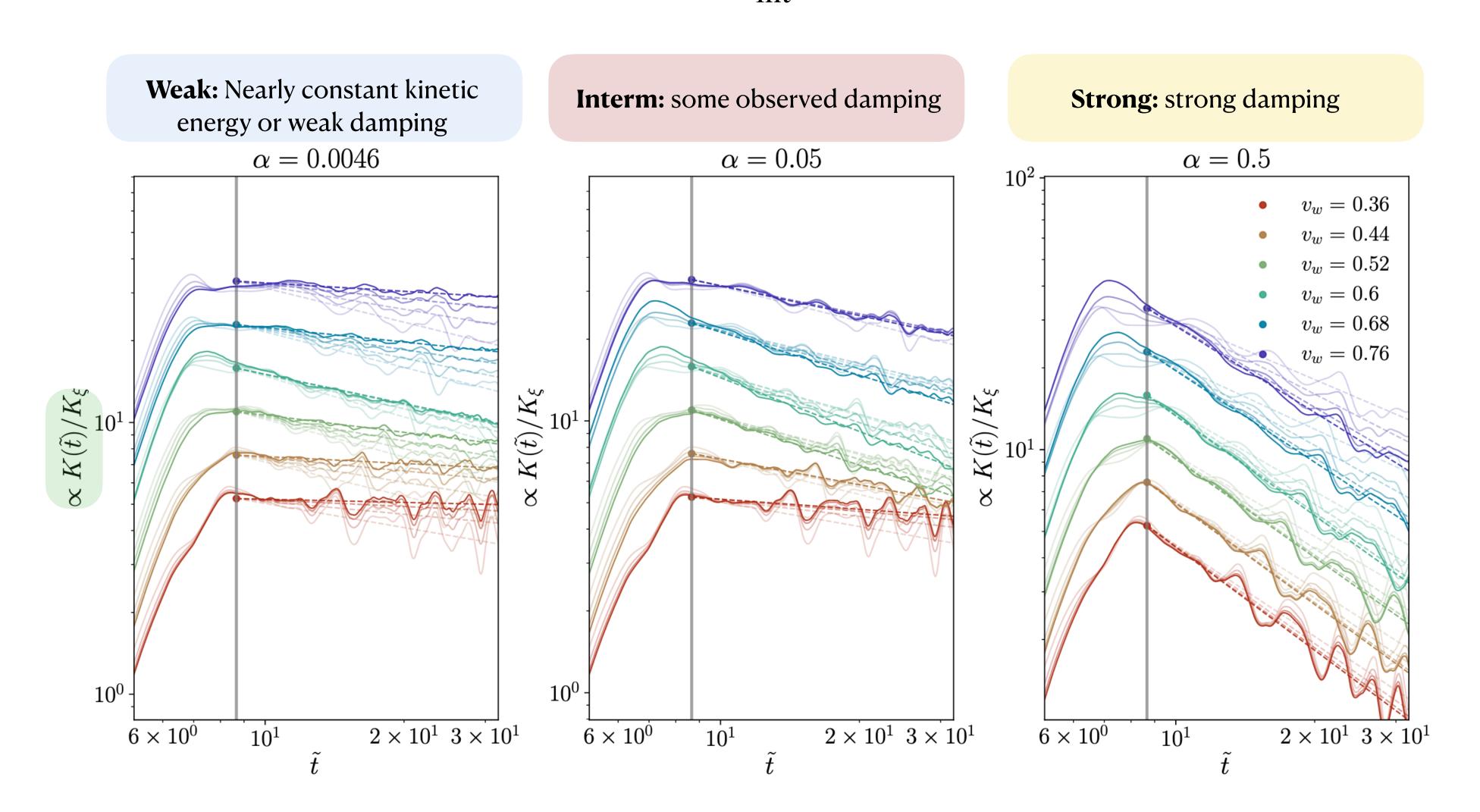
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Results K² int



$$\mathcal{I}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$$

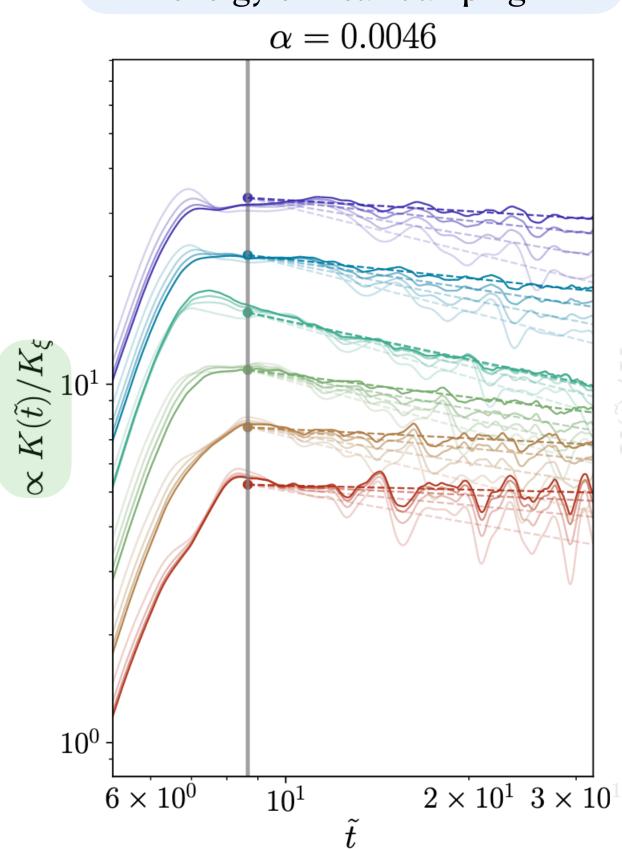
Results K² int



$$\mathcal{I}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$$

Results K² int

Weak: Nearly constant kinetic energy or weak damping



Interm: some observed damping

We fit a **power law**

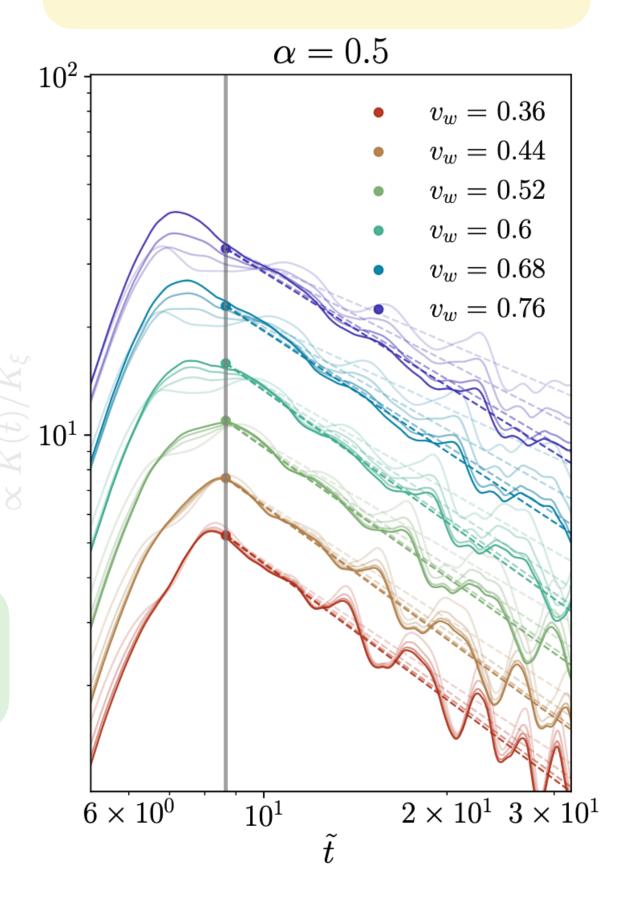
$$K(\tilde{t} > \tilde{t}_0) = K_0 \left(\frac{\tilde{t}}{\tilde{t}_0}\right)^{-b}$$

to the data,

so that

$$K_{\text{int}}^2 = K_0^2 \,\tilde{t}_0 \, \frac{(1 + \tilde{t}_{\text{sw}}/\tilde{t}_0)^{1-2b} - 1}{1 - 2b}$$

Strong: strong damping



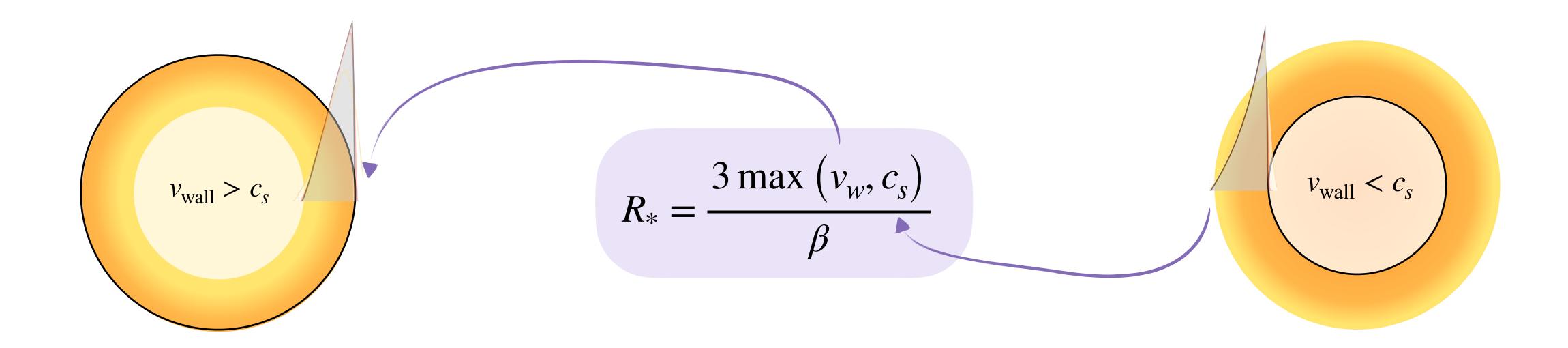
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Detonation

 R_*

Deflagration



- Interpret R_* as the typical size of sound-shells at collision
- This definition reduces the dependence of other parameters on v_w

$$\mathcal{I}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$$

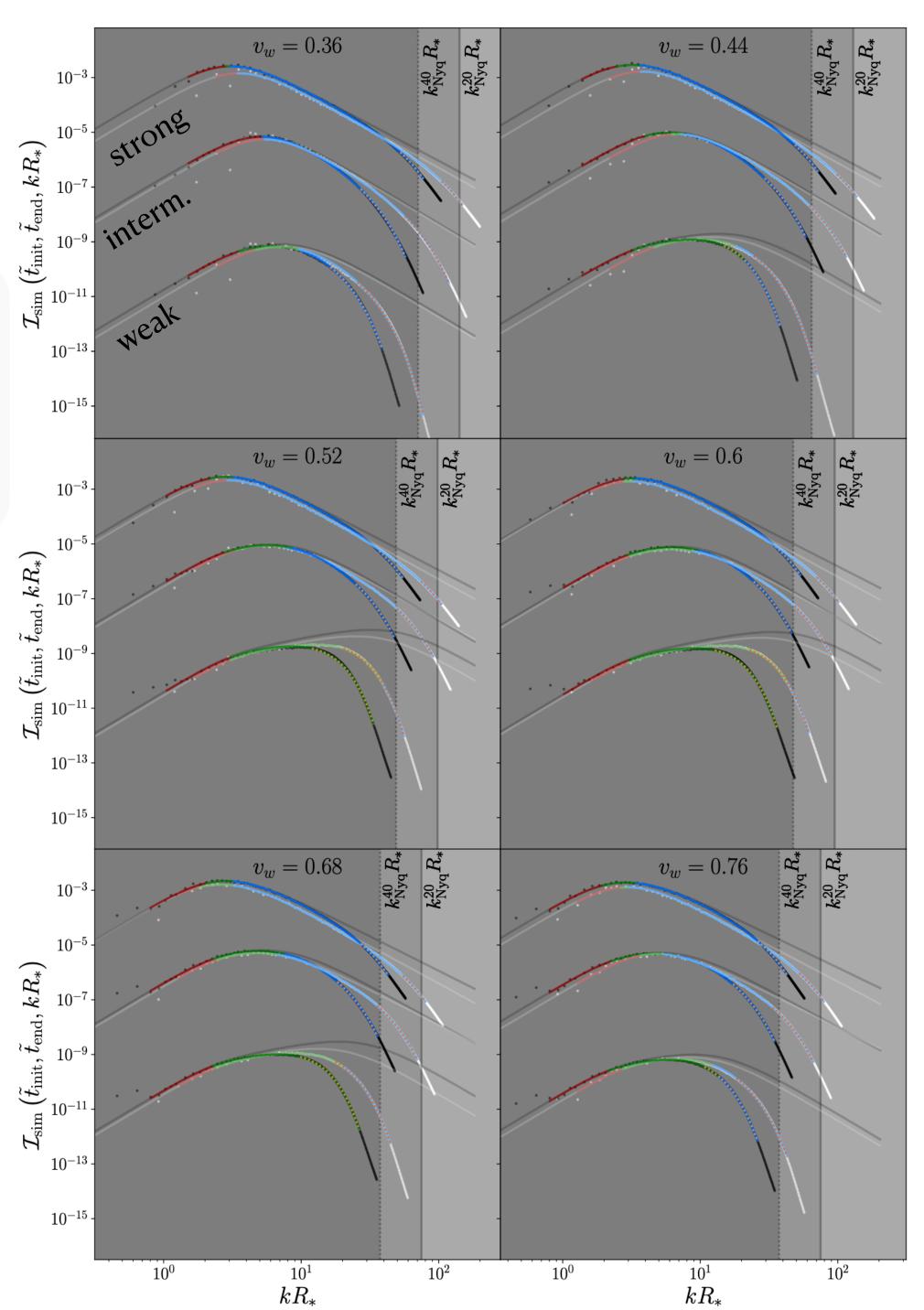
GW spectral shape S(k)

Cubic Linear ~-Cubic
$$S(k) = S_0 \times \left(\frac{k}{k_1}\right)^3 \left[1 + \left(\frac{k}{k_1}\right)^{a_1}\right]^{\frac{-3+1}{a_1}} \left[1 + \left(\frac{k}{k_2}\right)^{a_2}\right]^{\frac{-1+n_3}{a_2}} \times e^{-(k/k_d)^2}$$

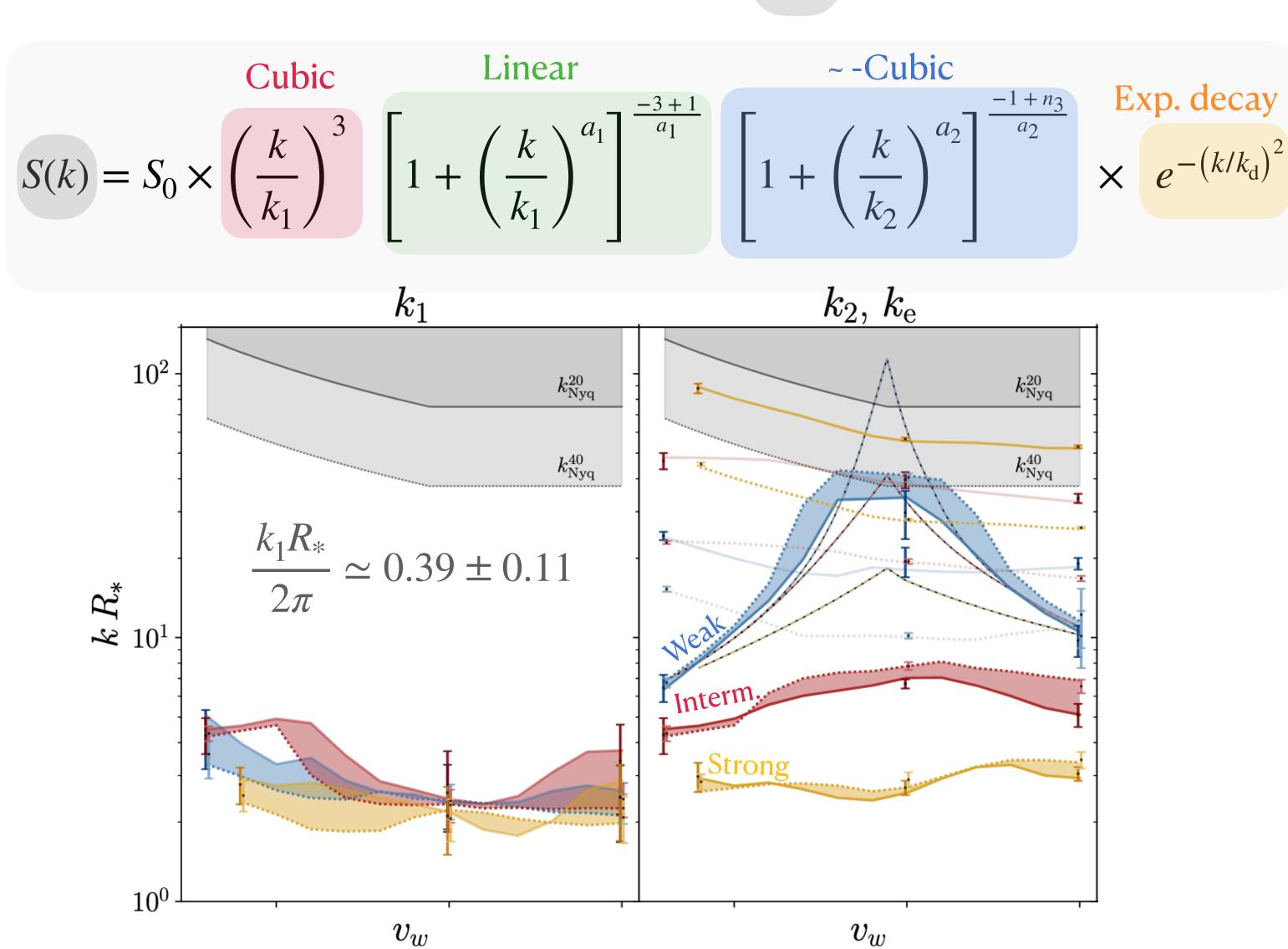
• Double broken power-law with exponential damping in the UV

• Fix IR and Interm. indices and fit S(k) wrt. free parameters

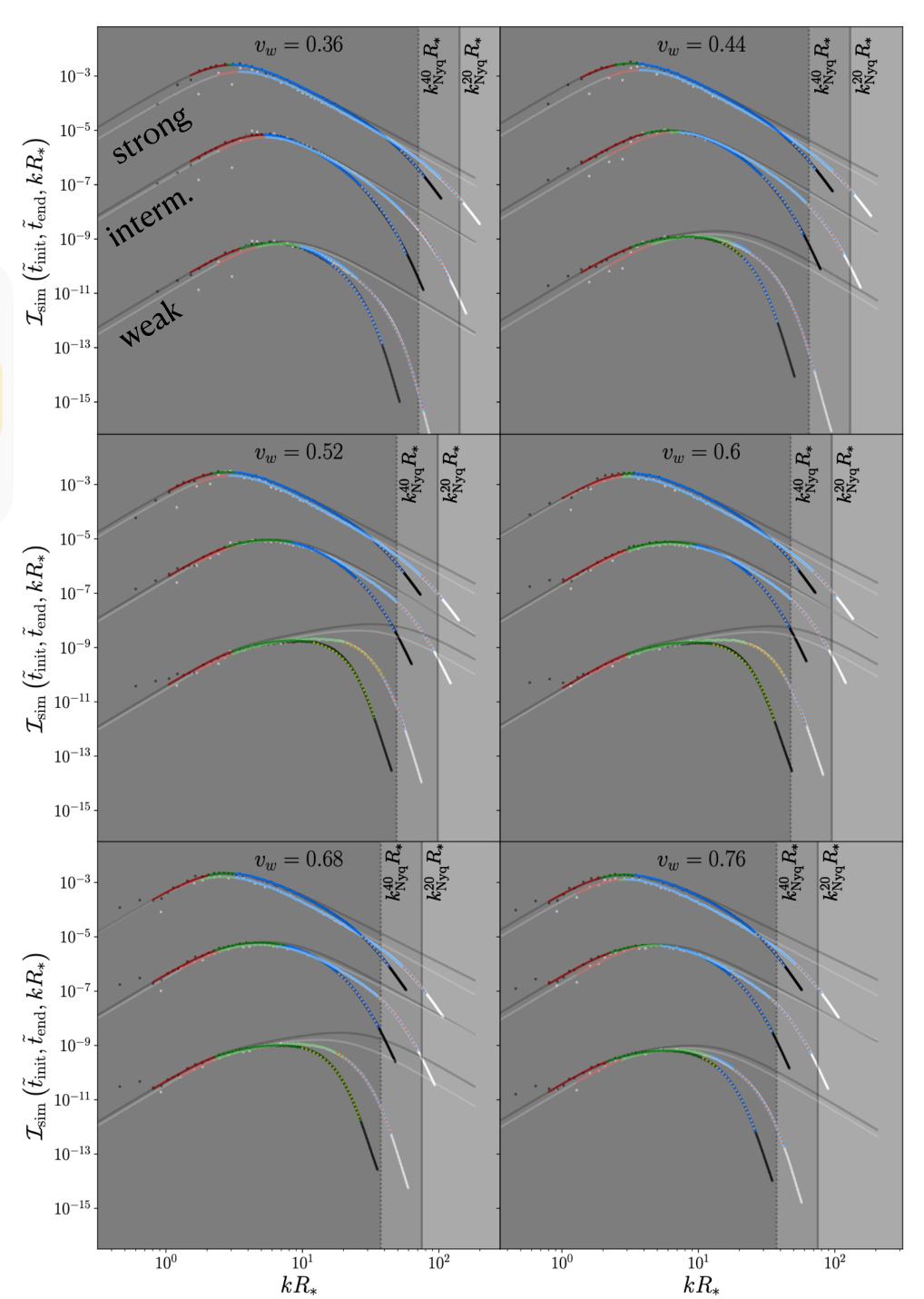
$$\frac{\mathsf{IR}:}{k^3} \longrightarrow k_1 \longrightarrow \frac{\mathsf{Interm}:}{k^1} \longrightarrow k_2 \longrightarrow \frac{\mathsf{UV}:}{k^{n_3}, n_3 \gtrsim -3} \longrightarrow k_d \longrightarrow \frac{\mathsf{Exp.}}{\mathsf{damping}}$$



GW spectral shape S(k)



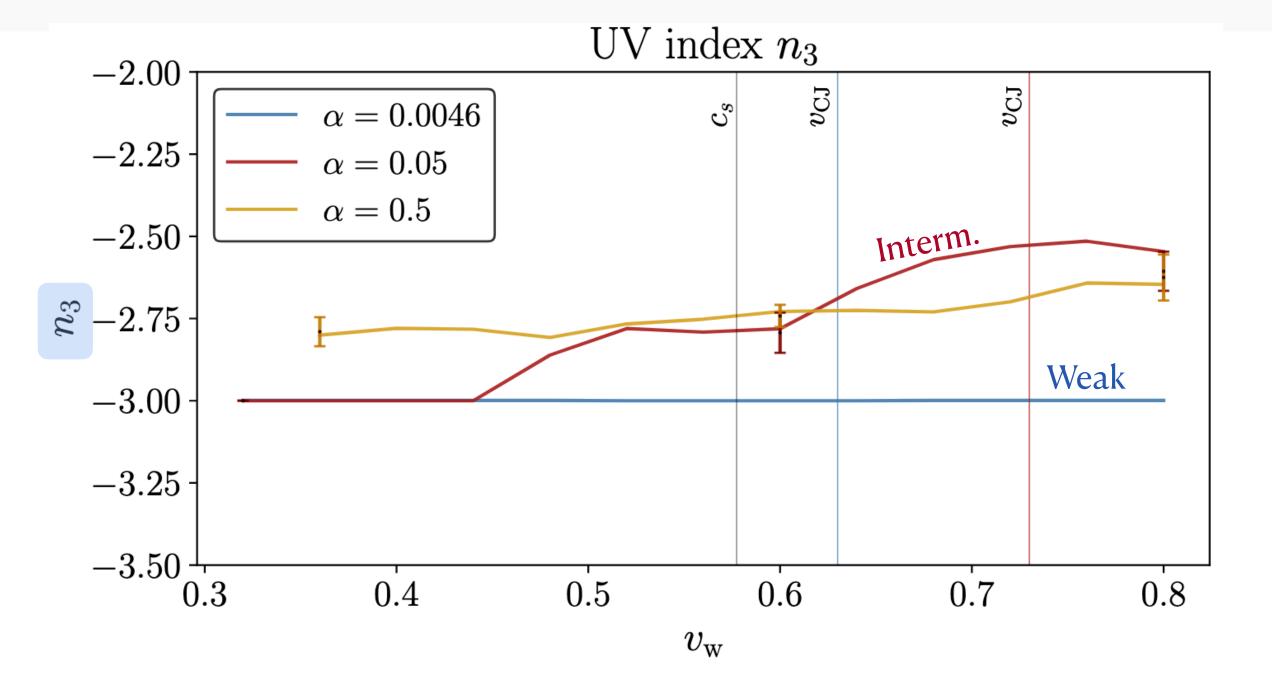
43



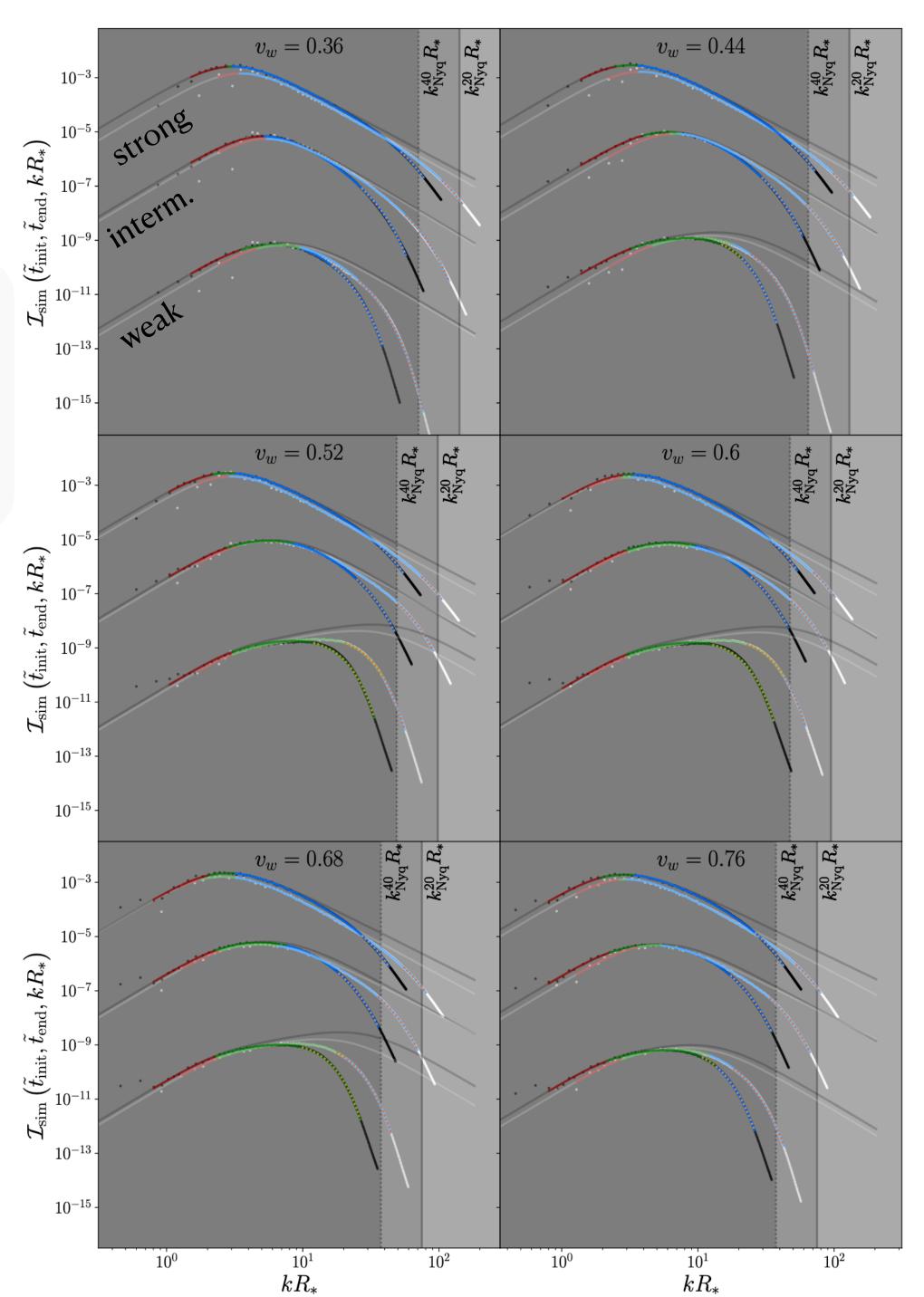
GW spectral shape S(k)

Cubic Linear
$$-\text{Cubic}$$

$$S(k) = S_0 \times \left(\frac{k}{k_1}\right)^3 \left[1 + \left(\frac{k}{k_1}\right)^{a_1}\right]^{\frac{-3+1}{a_1}} \left[1 + \left(\frac{k}{k_2}\right)^{a_2}\right]^{\frac{-1+n_3}{a_2}} \times e^{-(k/k_d)^2}$$



- Departure from -3 indicates a departure from linearity
- Dynamical depth of Strong PTs allow accurate estimation



GW spectrum from sound waves damped sources

How large is the GW efficiency?

$$\mathcal{I}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$$

GW spectrum from sound waves damped sources

$$10^{2} \tilde{\Omega}_{\text{GW}}^{\infty} = \begin{cases} 1.04_{-0.67}^{+0.81}, & \text{for } \alpha = 0.0046; \\ 1.64_{-0.13}^{+0.29}, & \text{for } \alpha = 0.05; \\ 3.11_{-0.19}^{+0.25}, & \text{for } \alpha = 0.5, \end{cases}$$

$$\mathcal{I}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} \frac{K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}})}{K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}})} (\beta R_*) S(k)$$

GW spectrum from sound waves damped sources

How much of the vacuum energy is transferred to the plasma in the form of kinetic energy and how can it be

$$10^{2} \tilde{\Omega}_{\text{GW}}^{\infty} = \begin{cases} 1.04_{-0.67}^{+0.81}, & \text{for } \alpha = 0.0046; \\ 1.64_{-0.13}^{+0.29}, & \text{for } \alpha = 0.05; \\ 3.11_{-0.19}^{+0.25}, & \text{for } \alpha = 0.5, \end{cases}$$

$$\mathcal{F}_{\text{sim}}(\tilde{t}_{*}, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^{2}(\tilde{t}_{*}, \tilde{t}_{\text{fin}}) (\beta R_{*}) S(k)$$

GW spectrum from sound waves damped sources

$$K_{\rm int}^2\left(\tilde{t}_{\rm sw}\right) = \mathcal{K}_0\,\tilde{t}_0\,\frac{(1+\tilde{t}_{\rm sw}/\tilde{t}_0)^{1-2b}-1}{1-2b}$$

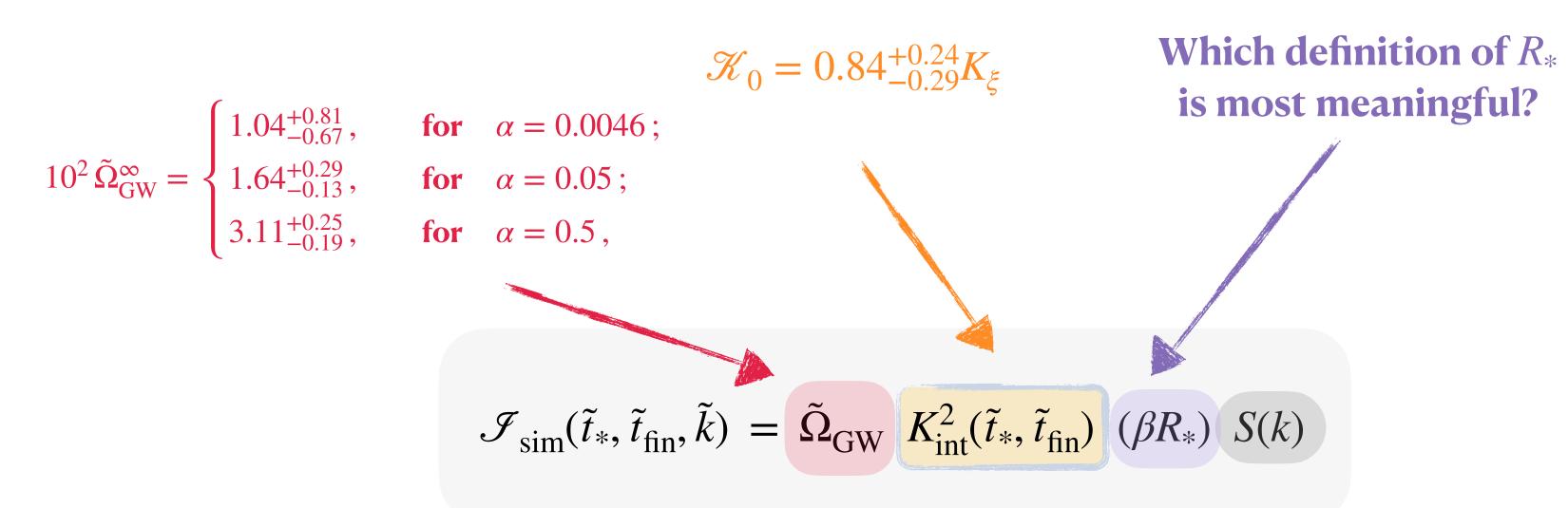
$$\mathcal{K}_0 = 0.84_{-0.29}^{+0.24}K_\xi$$

$$10^2\,\tilde{\Omega}_{\rm GW}^{\infty} = \begin{cases} 1.04_{-0.67}^{+0.81}, & \text{for } \alpha=0.0046\,;\\ 1.64_{-0.13}^{+0.29}, & \text{for } \alpha=0.05\,;\\ 3.11_{-0.19}^{+0.25}, & \text{for } \alpha=0.5\,, \end{cases}$$

$$\mathcal{F}_{\rm sim}(\tilde{t}_*, \tilde{t}_{\rm fin}, \tilde{k}) = \tilde{\Omega}_{\rm GW}\,K_{\rm int}^2(\tilde{t}_*, \tilde{t}_{\rm fin})\,(\beta R_*)\,S(k)$$

GW spectrum from sources damped sources

$$K_{\text{int}}^2(\tilde{t}_{\text{sw}}) = \mathcal{K}_0 \, \tilde{t}_0 \, \frac{(1 + \tilde{t}_{\text{sw}}/\tilde{t}_0)^{1-2b} - 1}{1 - 2b}$$



GW spectrum from sound waves damped sources

$$K_{\rm int}^2\left(\tilde{t}_{\rm sw}\right) = \mathcal{K}_0\,\tilde{t}_0\,\frac{(1+\tilde{t}_{\rm sw}/\tilde{t}_0)^{1-2b}-1}{1-2b}$$

$$\mathcal{K}_0 = 0.84_{-0.29}^{+0.24}K_\xi \qquad R_* = \frac{3\max\left(v_w,c_s\right)}{\beta}$$

$$10^2\,\tilde{\Omega}_{\rm GW}^{\infty} = \begin{cases} 1.04_{-0.67}^{+0.81}, & \text{for } \alpha=0.0046;\\ 1.64_{-0.13}^{+0.29}, & \text{for } \alpha=0.05;\\ 3.11_{-0.19}^{+0.25}, & \text{for } \alpha=0.5, \end{cases}$$

$$\mathcal{F}_{\rm sim}(\tilde{t}_*,\tilde{t}_{\rm fin},\tilde{k}) = \tilde{\Omega}_{\rm GW}\,K_{\rm int}^2(\tilde{t}_*,\tilde{t}_{\rm fin})\,\left(\beta R_*\right)\,S(k)$$

GW spectrum from sound waves damped sources

$$K_{\text{int}}^2(\tilde{t}_{\text{sw}}) = \mathcal{K}_0 \, \tilde{t}_0 \, \frac{(1 + \tilde{t}_{\text{sw}}/\tilde{t}_0)^{1-2b} - 1}{1 - 2b}$$

$$\mathcal{K}_0 = 0.84^{+0.24}_{-0.29} K_{\xi} \qquad \qquad R_* = \frac{3 \max \left(v_w, c_s\right)}{\beta}$$

$$10^2 \tilde{\Omega}_{\rm GW}^{\infty} = \begin{cases} 1.04^{+0.81}_{-0.67}, & \text{for } \alpha = 0.0046; \\ 1.64^{+0.29}_{-0.13}, & \text{for } \alpha = 0.05; \\ 3.11^{+0.25}_{-0.19}, & \text{for } \alpha = 0.5, \end{cases}$$

$$\mathcal{F}_{\text{sim}}(\tilde{t}_*, \tilde{t}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int}}^2(\tilde{t}_*, \tilde{t}_{\text{fin}}) (\beta R_*) S(k)$$

What is the spectral shape? In particular, where is the peak, and what are the IR and UV slopes?

GW spectrum from sources damped sources

$$K_{\mathrm{int}}^{2}\left(\tilde{t}_{\mathrm{sw}}\right) = \mathcal{K}_{0}\,\tilde{t}_{0}\,\frac{(1+\tilde{t}_{\mathrm{sw}}/\tilde{t}_{0})^{1-2b}-1}{1-2b}$$

$$\mathcal{K}_{0} = 0.84_{-0.29}^{+0.24}K_{\xi}$$

$$R_{*} = \frac{3\,\mathrm{max}\left(v_{w},c_{s}\right)}{\beta}$$

$$\frac{k_{1}R_{*}}{2\pi} \simeq 0.39\pm0.11$$

$$S(k) = S_{0}\times\left(\frac{k}{k_{1}}\right)^{3}\left[1+\left(\frac{k}{k_{1}}\right)^{\alpha_{1}}\right]^{\frac{-1+\alpha_{1}}{\alpha_{1}}}\left[1+\left(\frac{k}{k_{2}}\right)^{\alpha_{1}}\right]^{\frac{-1+\alpha_{1}}{\alpha_{2}}}\times e^{-(k/k_{2})^{2}}$$

$$\mathcal{F}_{\mathrm{sim}}(\tilde{t}_{*},\tilde{t}_{\mathrm{fin}},\tilde{k}) = \tilde{\Omega}_{\mathrm{GW}}\,K_{\mathrm{int}}^{2}(\tilde{t}_{*},\tilde{t}_{\mathrm{fin}})\left(\beta R_{*}\right)\,S(k)$$

$$\frac{k_{2}R_{*}}{2\pi} \simeq \begin{cases} 0.49\pm0.024/\Delta_{w}, & \alpha=0.0046\\ 0.93\pm0.13, & \alpha=0.05\\ 0.45\pm0.042, & \alpha=0.5 \end{cases}$$

GW spectrum from sever damped sources



For, as a matter of fact, in radiation-domination...

GW spectrum from sound waves damped sources

including cosmic expansion

... $\partial_{\mu}T^{\mu\nu} = 0$ is conformally invariant



Reinterpret measured K as the coming quantity

$$K_{\rm int}^2 \longrightarrow K_{\rm int,exp}^2 \equiv \left(\beta/H_*\right)^2 \int_{\tilde{\tau}_*}^{\tilde{\tau}_{\rm fin}} \frac{K^2(\tilde{\tau})d\tilde{\tau}}{\tilde{\tau}^2} \qquad \text{Additional damping from cosmic expansion}$$

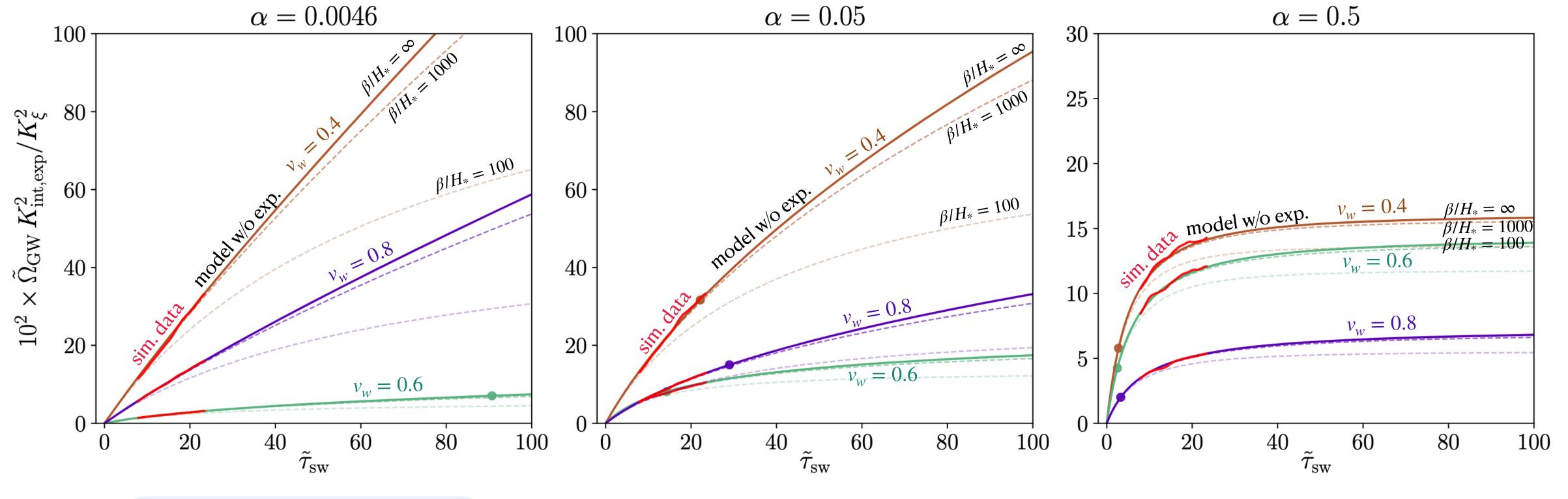
Now modeling expansion

$$\mathcal{I}_{\text{sim}}(\tilde{\tau}_*, \tilde{\tau}_{\text{fin}}, \tilde{k}) = \tilde{\Omega}_{\text{GW}} K_{\text{int,exp}}^2(\tilde{\tau}_*, \tilde{\tau}_{\text{fin}}) (\beta R_*) S(k)$$

GW spectrum from sound waves damped sources

including cosmic expansion

Growth of the GW amplitude w. and w/o cosmic expansion



Weak: Dominated by Hubble damping

Intermediate: Hybrid damping

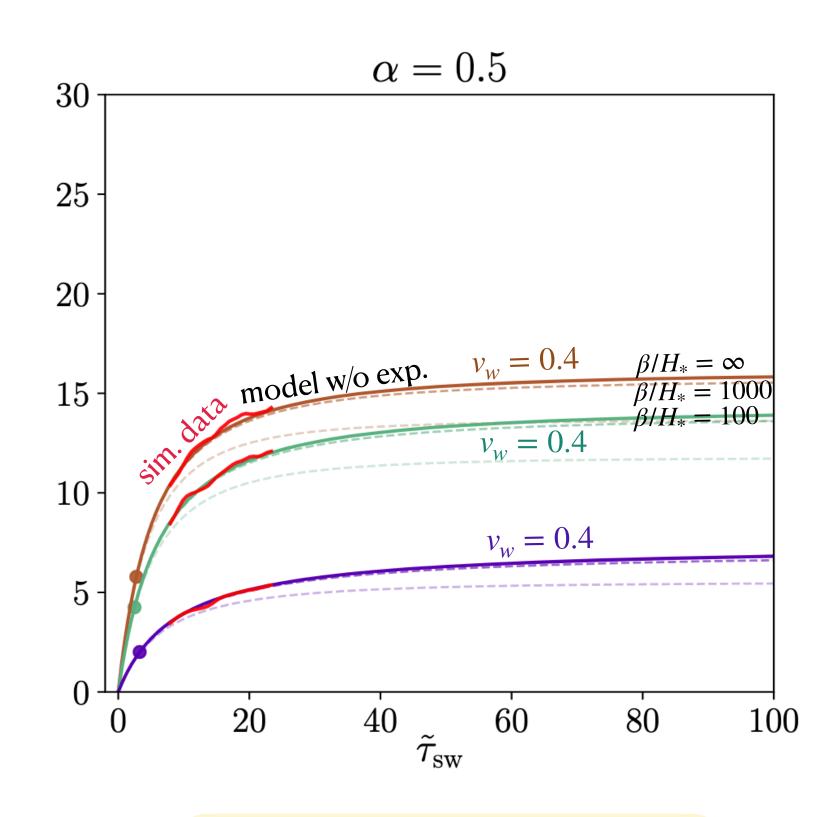
Strong: dominated by fluid damping

GW spectrum from sound waves damped sources

including cosmic expansion

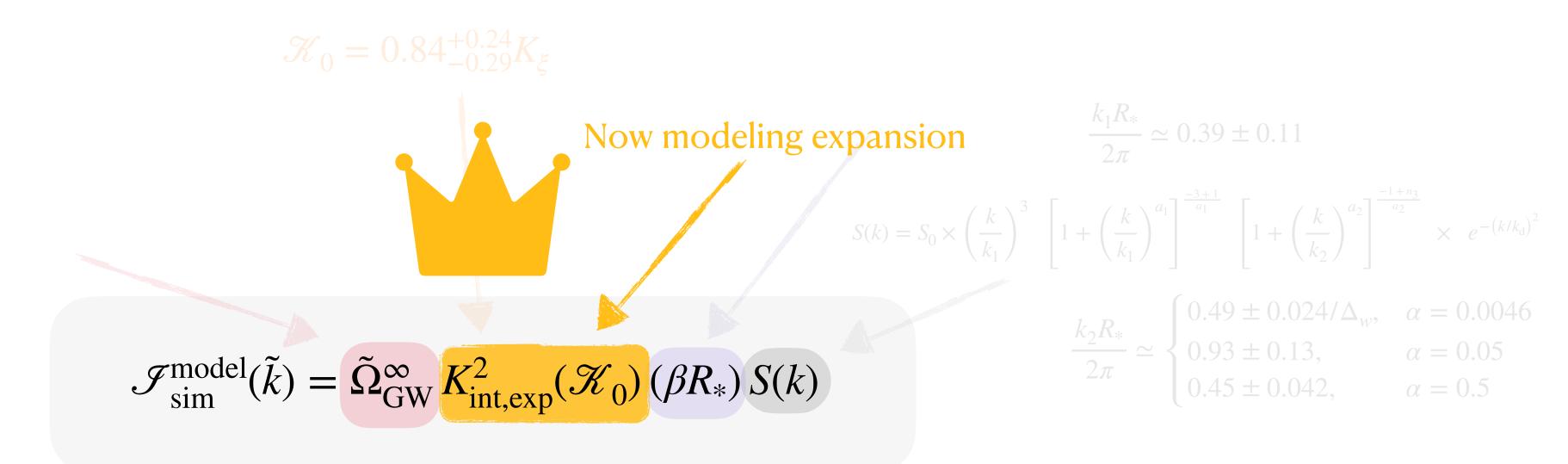
Strong PTs:

- Simulating ~ 5 shock formation times
- Within $\sim 10\%$ of amplitude saturation
- Impact of Hubble expansion small



Strong: dominated by fluid damping

GW spectrum from sound waves damped sources



Main results:

- Template for the GW spectrum motivated by theory and validated on simulation data
- Generalized from SWs to damped sources and including the effect of cosmic expansion,
- Potentially captures GWs from full dynamics, e.g. compressional motion, and turbulence, simultaneously
- Requires specification of the wall velocity v_w , the PT strength α , the ratio β/H_* , and PT temperature T_*

CosmoGW

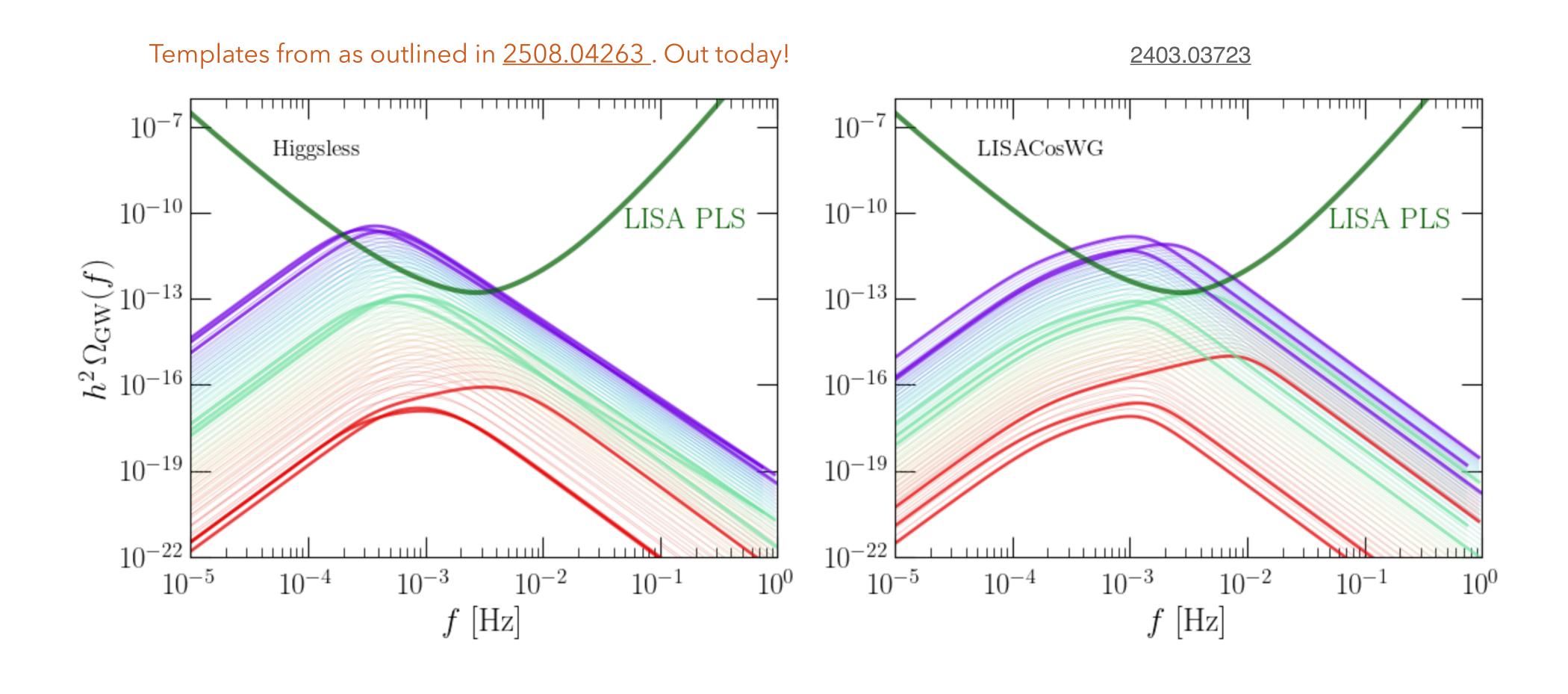
CosmoGW

Making the results readily accessible

- Python package for cosmological sources of GWs developed by Alberto Roper Pol
- Installation: pip install cosmoGW
- Find CosmoGW on Github
- Implemented the results of <u>2409.03651</u> to facilitate community use in e.g. parameter inference studies and detectability forecasts
- New paper summarizing the model <u>2508.04263</u>. Out today!
- Allows interpolating between simulations results to obtain GW spectrum predictions for any PT parameter point
- Takes as input $\{v_w, \alpha, \beta/H_*, N_{\text{shock}}\}$ for source duration $\delta\eta_{\text{fin}} = N_{\text{shock}} R_*/v_f$
- Link to <u>Tutorial</u>

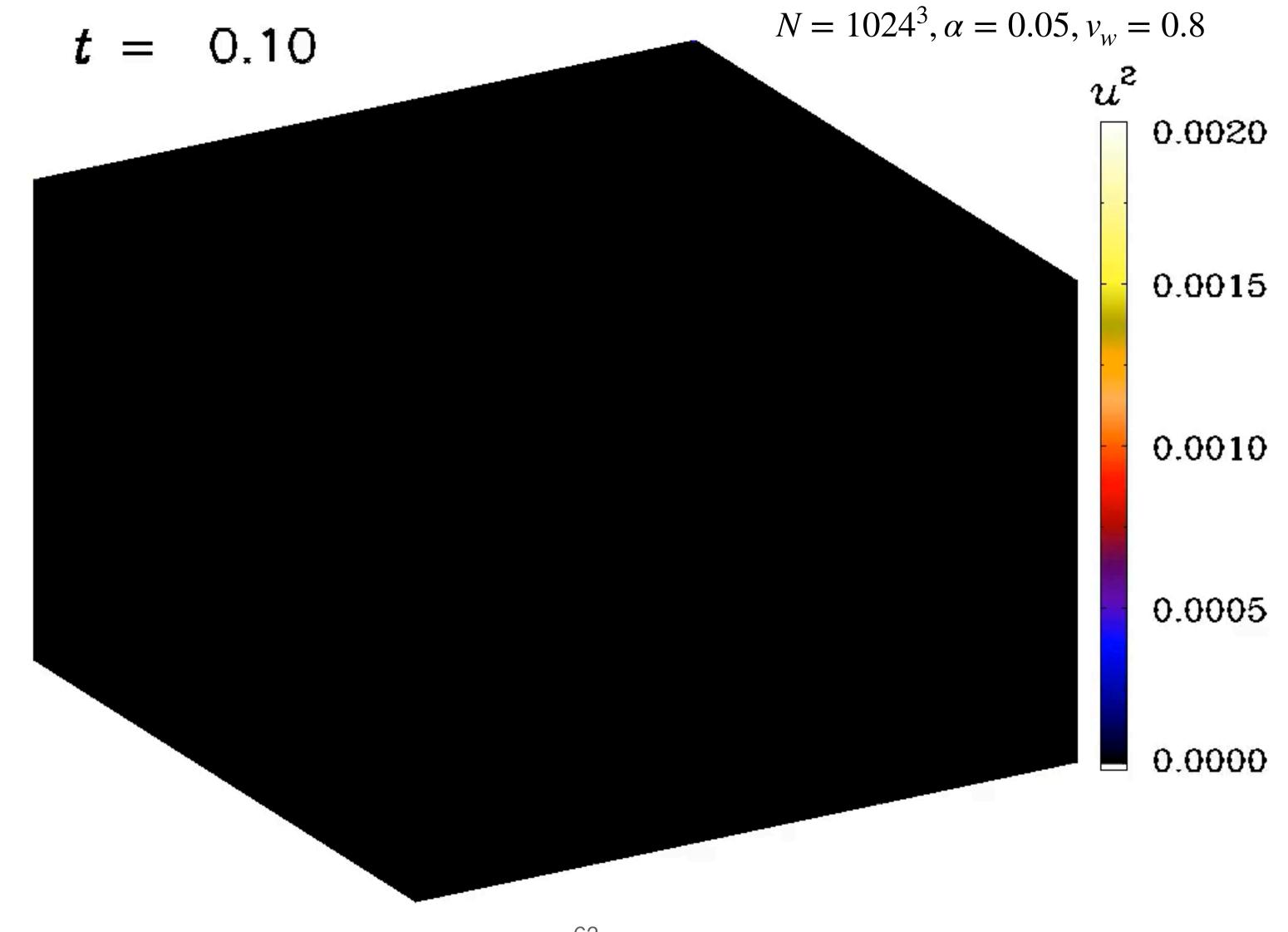
CosmoGW

Making the results readily accessible

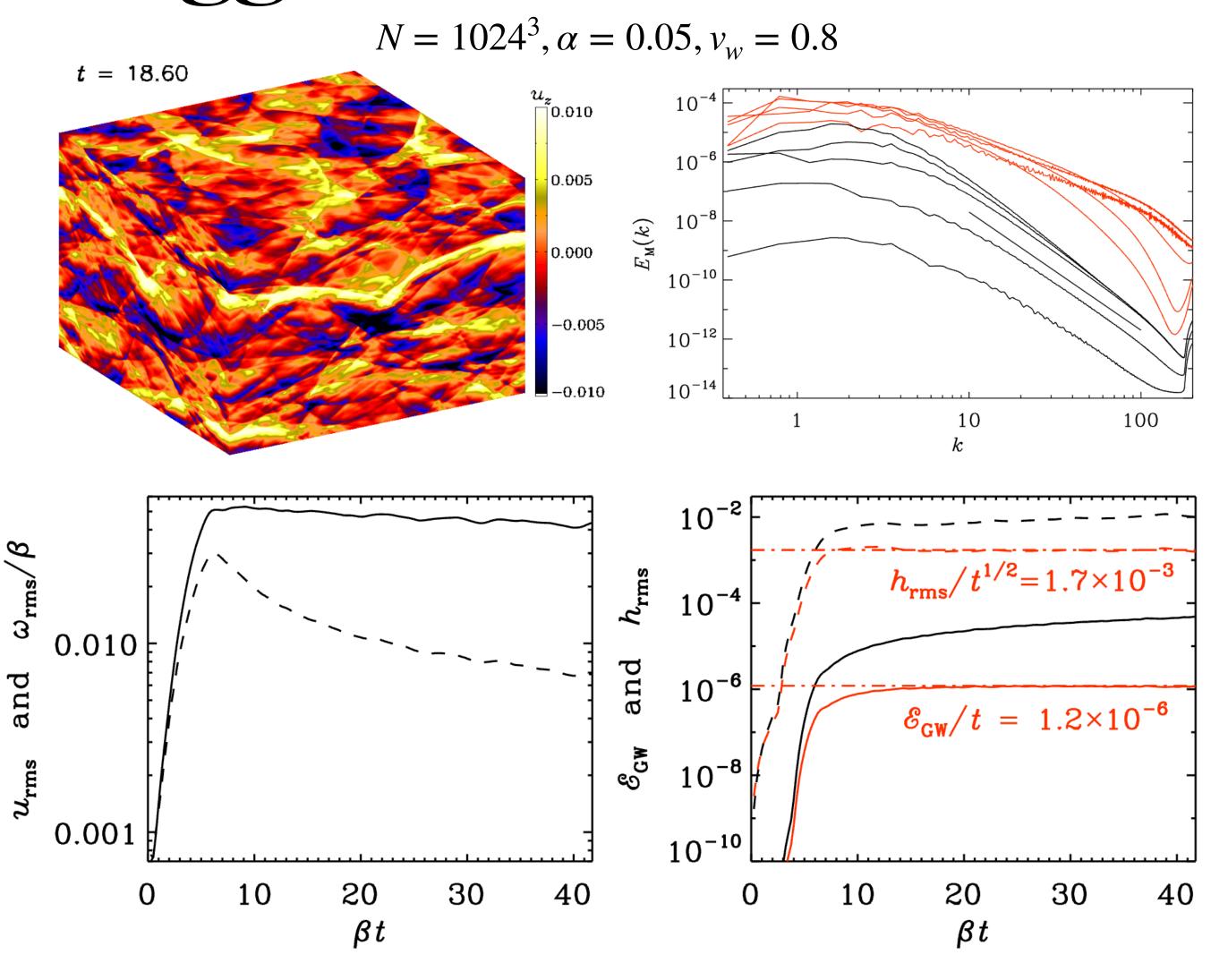


Higgsless in Pencil Code

Higgsless in Pencil Code



Higgsless in Pencil Code



Conclusions and summary

Conclusions and summary

We have...

- Used the novel **Higgsless approach** to simulate the **relativistic hydrodynamics** of 1st-order **cosmological PTs**
- Derived a new **GW-parameterization** generalizing results from sound-waves to **damped sources** accounting also for **cosmic expansion**
- Obtained **GW predictions** based on the **simulation data** for a **large part** of **parameter space** in wall velocity v_w and PT strength α
- Obtained, for the first time, GW predictions for strong PTs
- Established that non-linear evolution dictates the shape and peak location of the GW spectrum, thus rendering full simulations necessary to derive accurate GW predictions
- Captured the saturation of the GW amplitude due to non-linear dynamics
- GW templates available in the Python package CosmoGW
- Provided a tutorial, allowing the community to work interactively with the results and get GW
 predictions for any PT parameter choice

Thank you very much!

Questions?

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