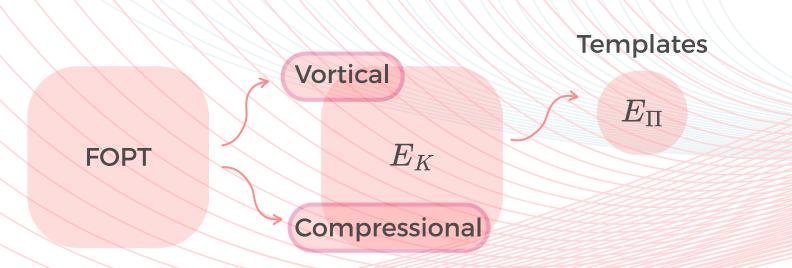


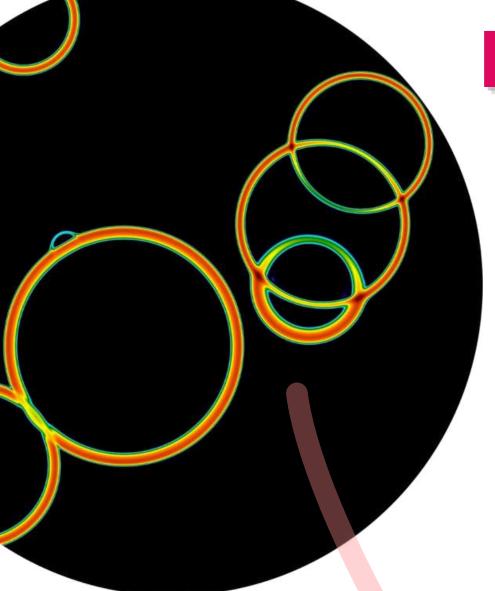
Gravitational wave production: the interplay between vortical and compressional motions

NORDITA (Stockholm) - August 8th, 2025 Numerical Simulations of Early Universe Sources of Gravitational Waves

Madeline SALOMÉ

Part I Introduction From FOPT to GWs' sourcing





First-order Phase Transition

FOPT Bubble expands and turns the Universe from one state to another. When they collides, sound waves are produced (compressional motion in the linear regime).

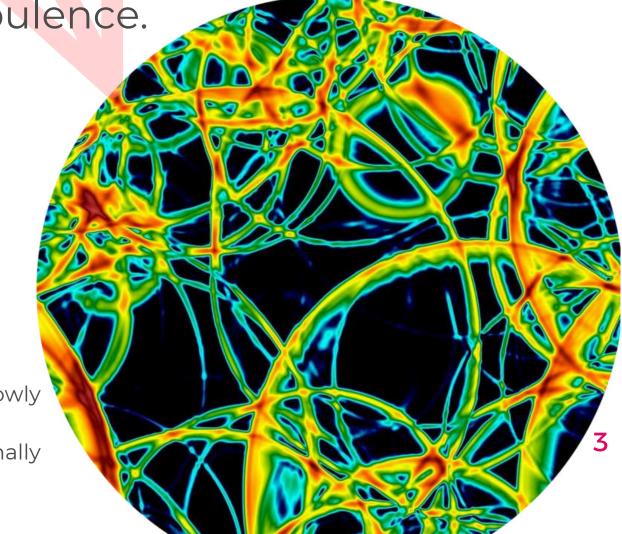
Turbulence The bubble merging possibly induces also vortical motions in the fluid corresponding to magnetohydrodynamic turbulence.

Stress-Energy Tensor

$$T_{ij} = w \gamma^2 u_i u_j - B_i B_j + \partial_i \phi \, \partial_j \phi$$

Publications In preparation.

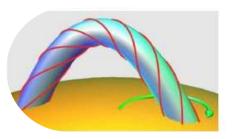
[1] A. Roper Pol, C. Caprini, A. S. Midiri., M. Salomé "Gravitational wave spectrum from slowly decaying sources in the early Universe: constant-in-time and coherent-decay models." [2] A. Roper Pol, M. Salomé. "Gravitational wave production from acoustic and fractionally compressible sources."

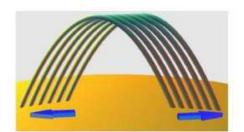


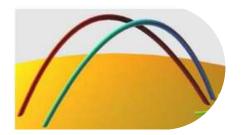
Properties of the field

Velocity Field Longitudinal & normal components + Helicity.

https://www.slideserve.com/mikko/injection-photosph-rique-d-h-licit-magn-tique-implications-sur-l-mergence-du-champ-magn-tique#google_vignette



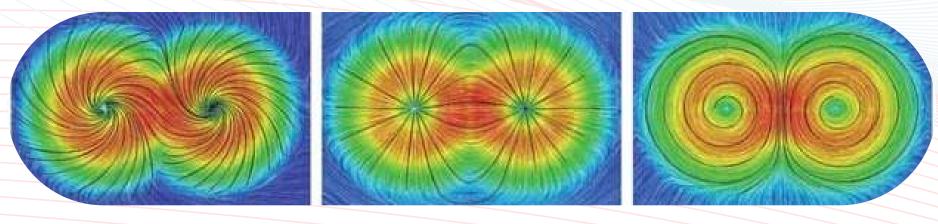




Degree of helicity







Vortical

$$E_\Pi \propto \langle T_{ij}^{TT} T_{ij}^{TT}
angle \sim \int \int \langle u_i u_j
angle \langle u_i u_j
angle \propto \int E_K \int E_K$$

Antonino's talk

GW Equation:

$$(\partial_t^2 - c^2
abla^2) h_{ij}(\mathbf{x},t) = rac{16\pi G}{ac^2} T_{ij}^{TT}(\mathbf{x},t)$$

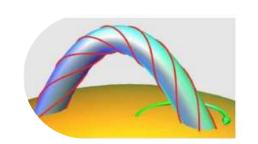
Small perturbation in the spacetime geometry:

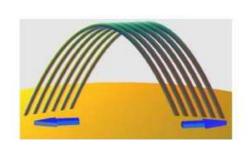
$$ds^2 = a(t)^2 \left[-dt^2 + (\delta_{ij} + h_{ij}) dx^i dx^j
ight]$$

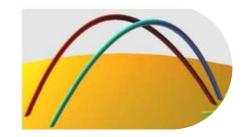
Sources:

$$T_{ij} = w \gamma^2 u_i u_j$$

$$E_{\Pi} \propto \int E_K \int E_K$$







Degree of helicity

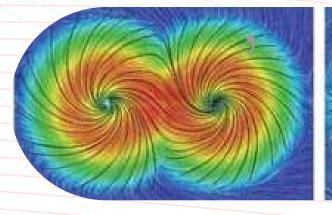
$$E_K^{
m vort}E_K^{
m vort},\, E_K^{
m comp}E_K^{
m comp},\, E_K^{
m vort}E_K^{
m comp}$$

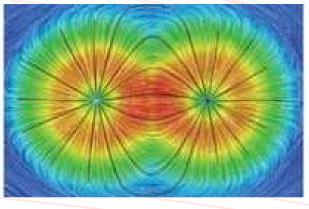


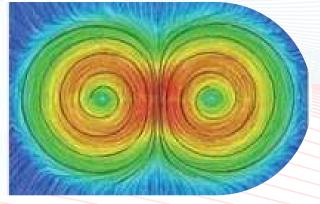
$$E_K^{
m hel}E_K^{
m hel}$$
 $E_K^{
m hel}E_K^{
m comp}$ $E_K^{
m hel}E_K^{
m vort}$



Fraction of compressibility



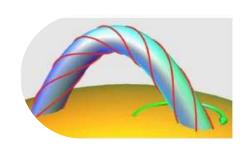


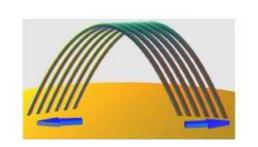


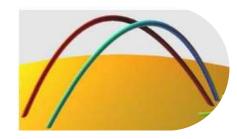
Mix Compressional

Vortical

$$E_{\Pi} \propto \int E_K \int E_K$$







Degree of helicity

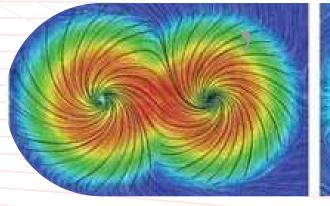
$$E_K^{
m vort}E_K^{
m vort},\, E_K^{
m comp}E_K^{
m comp},\, E_K^{
m vort}E_K^{
m comp}$$

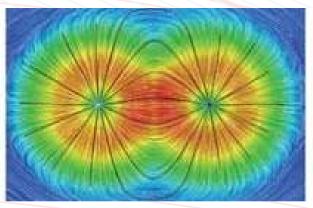


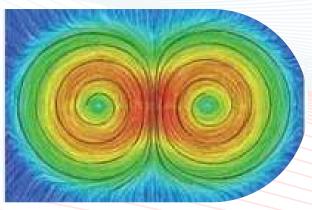


Symmetric part consideration







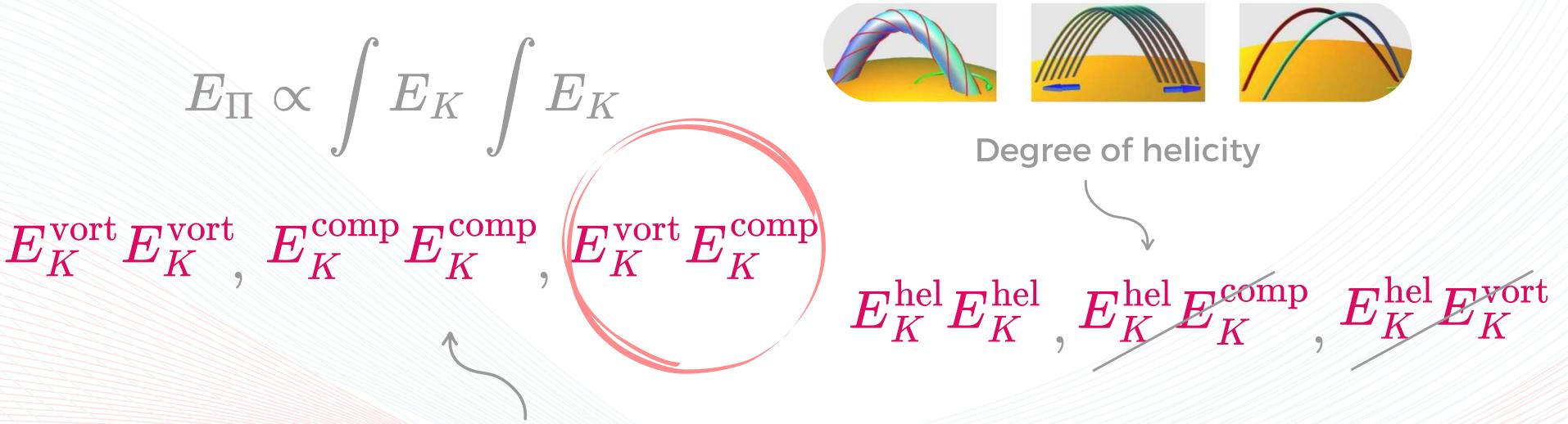


Compressional Mix

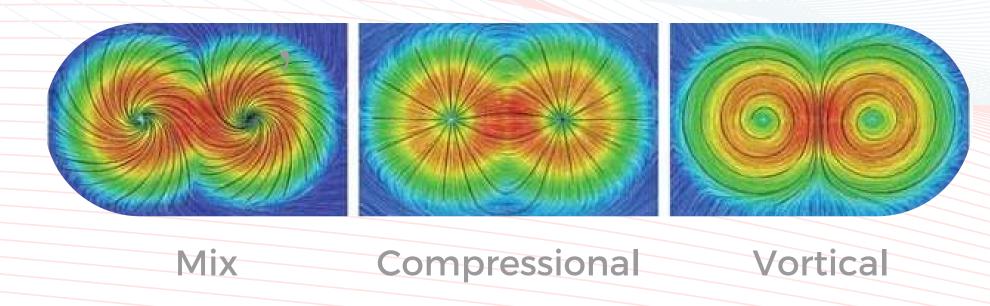
Vortical

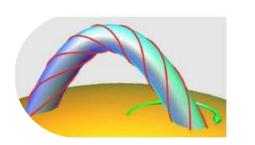
How to combine them to produce GWs?

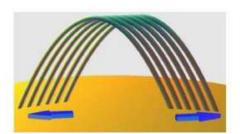
The Unequal-Time Correlator

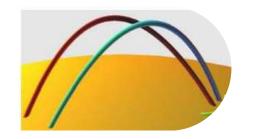


Fraction of compressibility





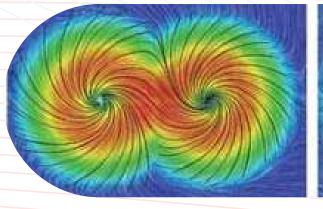




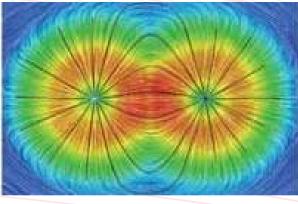
Degree of helicity

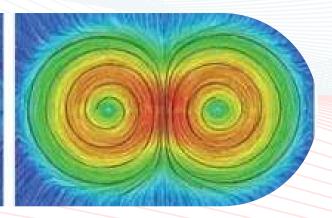
$$(1-q)^2 E_\Pi^{
m vort} + q^2 E_\Pi^{
m comp} + q (1-q) E_\Pi^{
m mix} + arepsilon^2 E_\Pi^{
m hel}$$

Fraction of compressibility



Mix

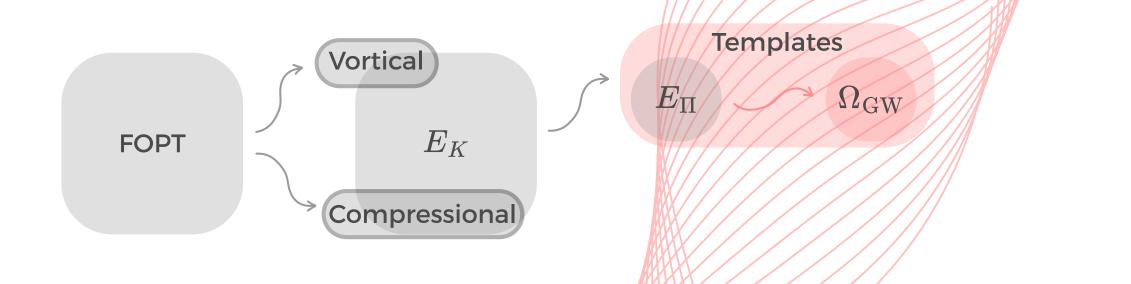




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Part II GW spectrum

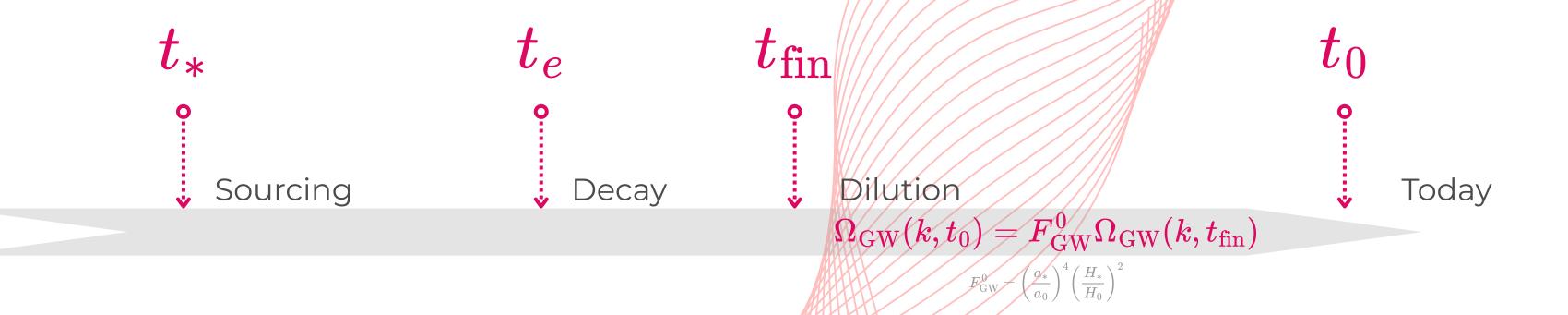
GW energy density spectrum



Our observable

$$\Omega_{
m GW}(k,t) = 3k \int_{t_*}^{\min(t,t_{
m fin})} dt_1 \int_{t_*}^{\min(t,t_{
m fin})} dt_2 rac{E_{
m H}(k,t_1,t_2)}{t_1t_2} \cos\left[k(t-t_1)
ight] \cos\left[k(t-t_2)
ight]$$

Fully vortical motion



Constant-in-time model Antonino's talk

$$E_\Pi(k,t_1,t_2)\simeq E_\Pi(k,t_*)=E_\Pi^*(k)$$

$$oldsymbol{\Omega_{ ext{GW}}(k,t_0) = 3kF_{ ext{GW}}^0E_{ ext{II}}^*(k)oldsymbol{\int_{t_*}^{t_{ ext{fin}}}} rac{dt_1}{t_1} ext{cos}[k(t-t_1)]igg]^2}$$

Fully compressional motion

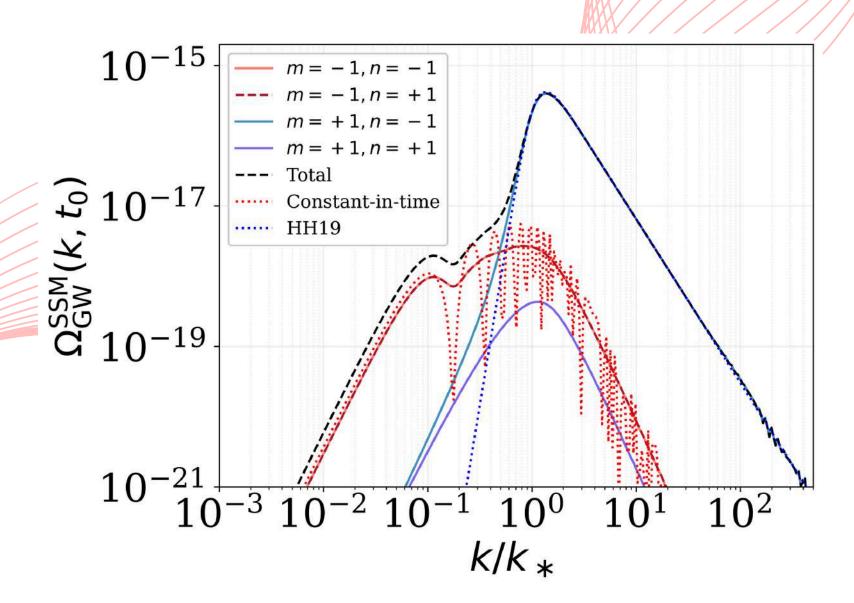
Sound-Shell model

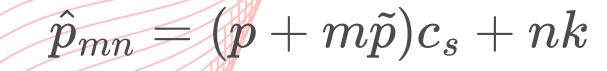
$$E_K(k,t_1,t_2) = E_K^*(k)\cos[kc_s(t_2-t_1)]$$

$$\Delta^2(k,p, ilde{p},t_{
m fin}) = \int_{t_*}^{t_{
m fin}} \int_{t_*}^{t_{
m fin}} rac{dt_1 dt_2}{t_1 t_2} {
m cos}[pc_s(t_2-t_1)] \, {
m cos}[ilde{p}c_s(t_2-t_1)] \, {
m cos}[k(t_2-t_1)]$$

Component analysis

$$\Delta^2(k,p, ilde{p},t_{
m fin}) = rac{1}{4}\sum_{m,n=\pm 1} \Delta^2_{mn}(\hat{p}_{mn})$$

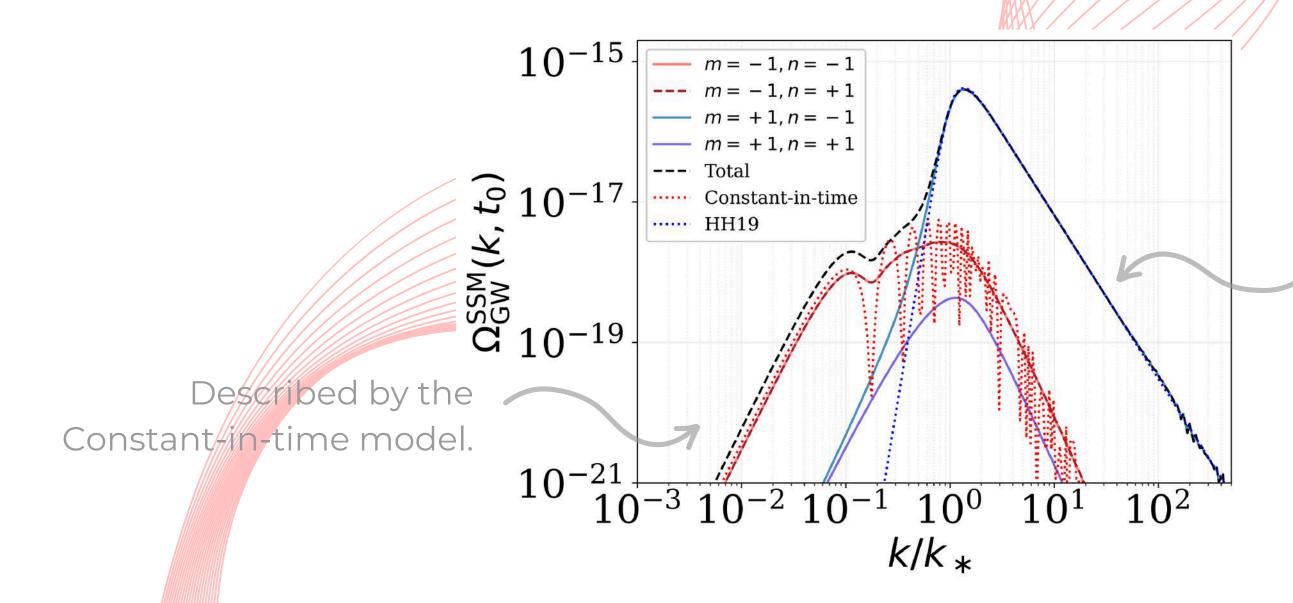




Component analysis

$$\Delta^2(k,p, ilde{p},t_{
m fin}) = rac{1}{4}\sum_{m,n=\pm 1} \Delta^2_{mn}(\hat{p}_{mn})$$

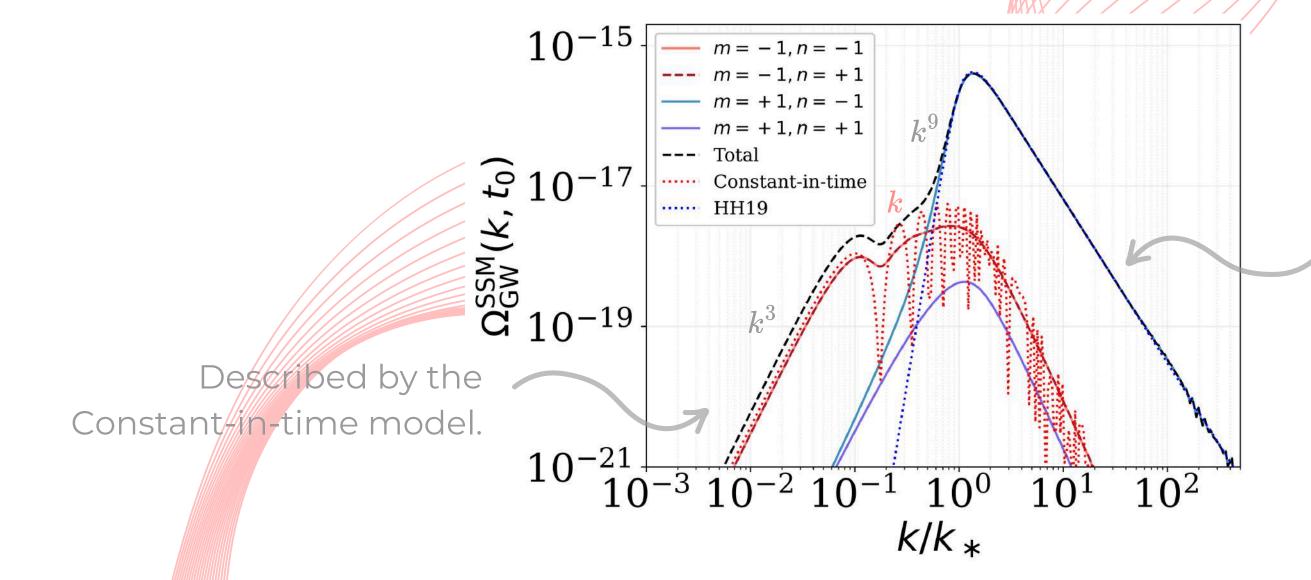




Hindmarsh and Hijazi, 2019

Component analysis

$$\Delta^2(k,p, ilde{p},t_{
m fin}) = rac{1}{4}\sum_{m,n=\pm 1} \Delta^2_{mn}(\hat{p}_{mn})$$



Hindmarsh and Hijazi, 2019

 $\hat{p}_{mn} = (p + m\tilde{p})c_s + nk$

Component analysis

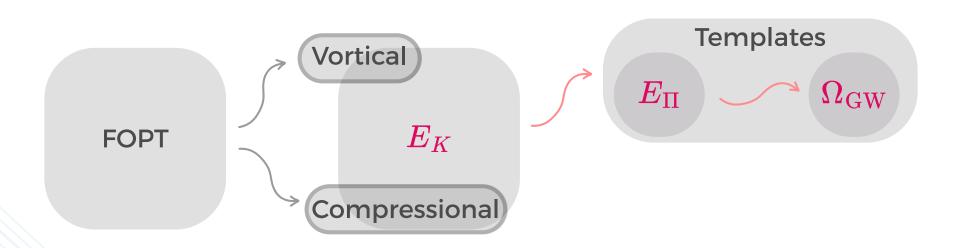
$$\Delta^2(k,p, ilde{p},t_{
m fin}) = rac{1}{4}\sum_{m,n=\pm 1}\Delta^2_{mn}(\hat{p}_{mn})$$

 $\hat{p}_{mn} = (p + m\tilde{p})c_s + nk$

Hindmarsh and Hijazi, 2019

These scales are constructed over time!

CosmoGW Developed by Alberto Roper Pol



Python library Functions for the study of cosmological GW

backgrounds from different sources in the early

Universe (phase transitions).



It includes GW models, postprocessing calculations, numerical computations, plotting routines, and detector sensitivities.

UETC of the anisotropic stresses

Normal component (Vortical)

$$E_{\Pi}^{
m vort}(k) = w^2 rac{k^2}{2} \int_0^{\infty} dp \; E_K^{
m vort}(p) \int_{-1}^1 dz rac{1}{ ilde{p}^4} E_K^{
m vort}(ilde{p}) (1+z^2) (2 ilde{p}^2-p^2(1-z^2))$$

Longitudinal component (Compressional)

$$E_\Pi^{
m comp}(k) = 2 w^2 k^2 \int_0^\infty dp \; p^2 E_K^{
m comp}(p) \int_{-1}^1 dz rac{1}{ ilde{p}^4} E_K^{
m comp}(ilde{p}) (1-z^2)^2$$

Mixed Term

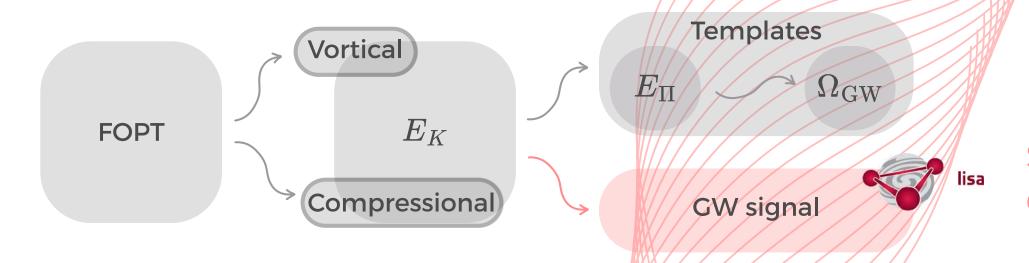
$$E_\Pi^{
m mix}(k) = 2 w^2 k^2 \int_0^\infty dp \; p^2 E_K^{
m vort}(p) \int_{-1}^1 dz rac{1}{ ilde{p}^4} E_K^{
m comp}(ilde{p}) (1-z^4)$$

Helicity

$$E_\Pi^{
m hel}(k) = w^2 rac{k^2}{2} \int_0^\infty dp \, p E_K^{
m hel}(p) \int_{-1}^1 dz rac{1}{ ilde{p}^2} E_K^{
m hel}(ilde{p}) z (k-zp)$$

Available in CosmoGW!

Pencil Code



Simulated by solving GW equations in the Pencil Code.

Pencil Code

A high-order finite-difference code for compressible hydrodynamic flows with magnetic fields.

Spatial Derivative 6th-order, Finite-difference method.

Time Step 3rd-order Runge-Kutta scheme by Williamson (1980).

CosmoGW Read the output files of the Pencil Code's run.

Pencil Code

Type of fields

 $\begin{array}{ll} \text{Magnetic field} & + \sigma \\ \text{Velocity field} & + \sigma \end{array}$

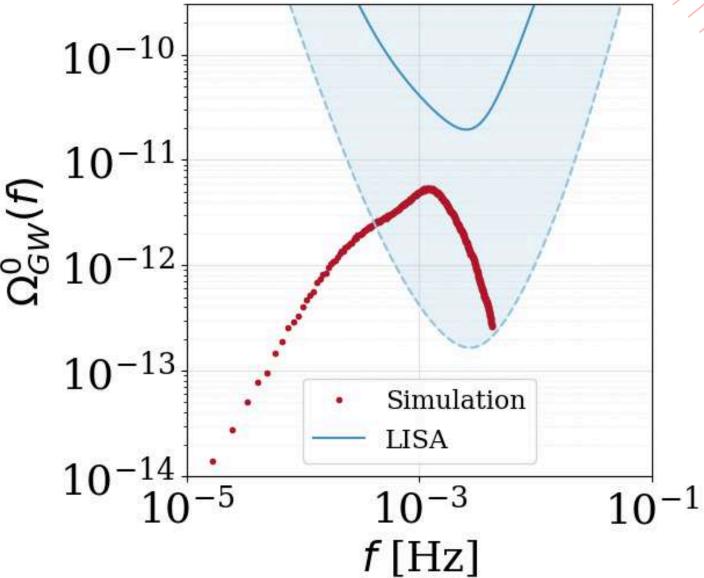
 $arepsilon = rac{2\sigma(1-q)}{1+\sigma^2(1-q)}$

Helicity rate

Compressional rate

al

 $(1-q)^2 E_\Pi^{
m vort} + q^2 E_\Pi^{
m comp} + q (1-q) E_\Pi^{
m mix} + arepsilon^2 E_\Pi^{
m hel}$



Pencil Code

Simulation

 10^{-1}

LISA

 10^{-3}

f[Hz]

Type of fields

Magnetic field
$$+\sigma$$

Velocity field $+\sigma$

 10^{-10}

 10^{-11}

 10^{-13}

S 10-12

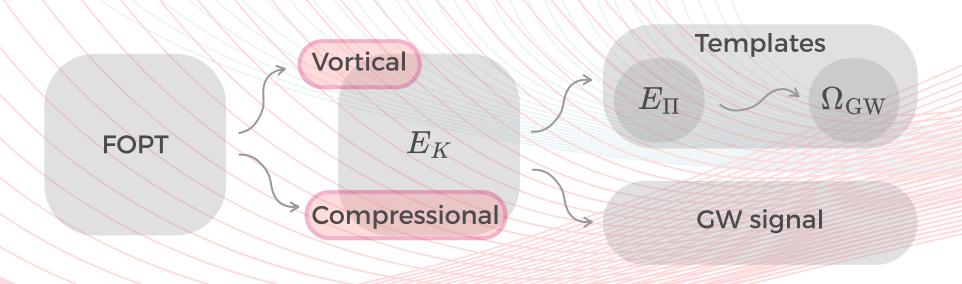
$$(1-q)^2 E_\Pi^{
m vort} + q^2 E_\Pi^{
m comp} + q (1-q) E_\Pi^{
m mix} + arepsilon^2 E_\Pi^{
m hel}$$

Intrument sensibility of LISA.

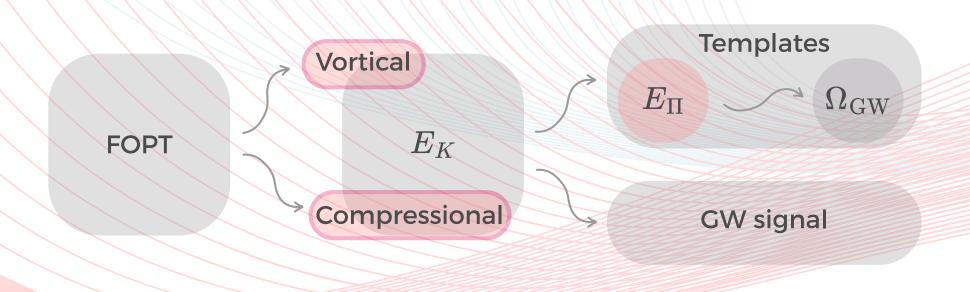
Power Law detection sensibility in LISA data. Signal processing.



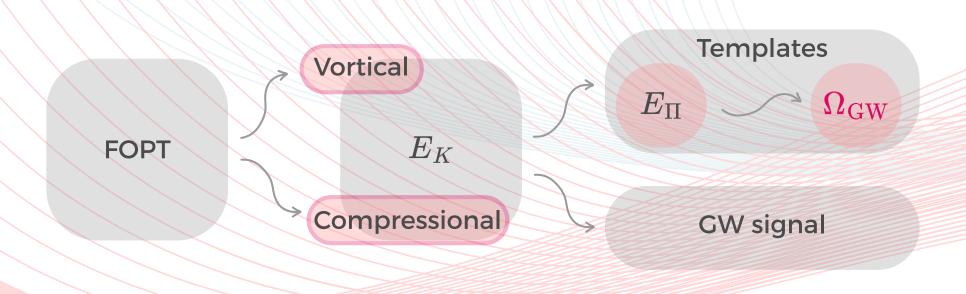
Part III Mixed Term



Part III Mixed Term



Part III Mixed Term



Simulations with the Pencil Code

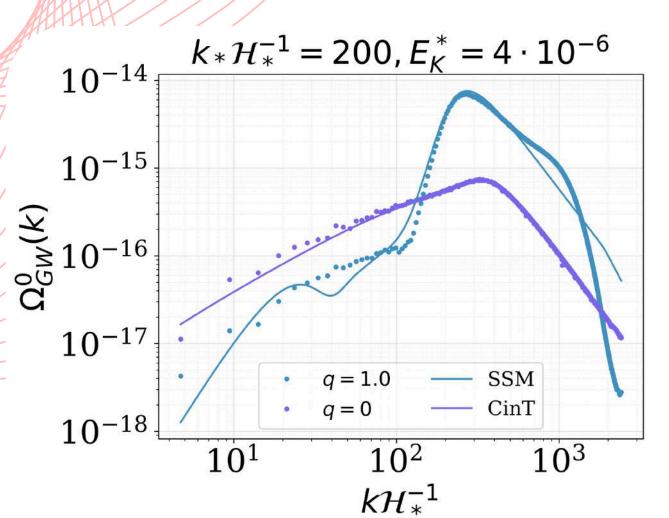
Name	k_*	E_K^*	q	σ	ν	$k_{ m IR}$	Resolution	nt	δt
q0	200	$4\cdot 10^{-6}$	0.0	0.0	10^{-7}	$3/2\pi$	1024^{3}	20000	10^{-4}
q025	200	$4 \cdot 10^{-6}$	0.25	0.0	10^{-7}	$3/2\pi$	1024^{3}	20000	10^{-4}
q05	200	$4 \cdot 10^{-6}$	0.5	0.0	10^{-7}	$3/2\pi$	1024^{3}	20000	10^{-4}
q075	200	$4 \cdot 10^{-6}$	0.75	0.0	10^{-7}	$3/2\pi$	1024^{3}	20000	10^{-4}
q1	200	$4 \cdot 10^{-6}$	1.0	0.0	10^{-7}	$3/2\pi$	1024^{3}	20000	10^{-4}

Table 3 List of Pencil Code simulations for several fractions of compressibility q.

Spectral shape

$k_* \mathcal{H}_*^{-1} = 200, E_K^* = 4 \cdot 10^{-6}$ 10^{-14} 10^{-15} $q \sim 0.7$ 10^{-16} 10^{-17} $q = 1.0 \qquad q = 0.25 \qquad q = 0.25 \qquad q = 0$ $q = 0.75 \qquad q = 0.5$ 10^{1} 10^{1} 10^{2} 10^{3} $k \mathcal{H}_*^{-1}$

Models



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Modelizing the mixed motion spectra

$$\Omega_{\mathrm{GW}}(k,t) = (1-q)^2 \Omega_{\mathrm{GW}}^{\mathrm{vort}}(k,t) + q^2 \Omega_{\mathrm{GW}}^{\mathrm{comp}}(k,t) + q(1-q) \Omega_{\mathrm{GW}}^{\mathrm{mix}}(k,t)$$

Constant-in-time Model

$$E_K^{
m vort}(k,t_1,t_2)=E_K^*(k)$$

$$\Omega_{ ext{GW}}^{ ext{vort}}(k,t_0) = rac{3w^2k^3}{4}F_{ ext{GW}}^0\int_0^\infty dp\, E_K^*(p)\int_{-1}^1 dz\, E_K^*(ilde{p})rac{(1+z^2)}{ ilde{p}^4}(2 ilde{p}^2-p^2(1-z^2))\Delta^2(k,t_{ ext{fin}}) \ \Delta^2(k,t_{ ext{fin}}) = \int_t^{t_{ ext{fin}}}\int_t^{t_{ ext{fin}}}rac{dt_1dt_2}{t_1t_2}\cos[k(t_2-t_1)]$$

Sound-Shell Model

$$E_K^{
m comp}(k,t_1,t_2) = E_K^*(k) \cos[kc_s(t_2-t_1)]$$

$$\Omega_{
m GW}^{
m comp}(k,t_0) = 3w^2k^3F_{
m GW}^0\int_0^\infty dp\, p^2E_K^*(p)\int_{-1}^1 dz\, rac{(1-z^2)^2}{ ilde p^4}E_K^*(ilde p)\Delta^2(k,p, ilde p,t_{
m fin})
onumber \ \Delta^2(k,p, ilde p,t_{
m fin}) = \int_{t_0}^{t_{
m fin}}\int_{t_0}^{t_{
m fin}}rac{dt_1dt_2}{t_1t_2}\cos[pc_s(t_2-t_1)]\cos[ilde pc_s(t_2-t_1)]\cos[k(t_2-t_1)]
onumber \ \Delta^2(k,p, ilde p,t_{
m fin}) = \int_{t_0}^{t_{
m fin}}\int_{t_0}^{t_{
m fin}}rac{dt_1dt_2}{t_1t_2}\cos[pc_s(t_2-t_1)]\cos[ilde pc_s(t_2-t_1)]
onumber \ \Delta^2(k,p, ilde p,t_{
m fin}) = \int_{t_0}^{t_{
m fin}}\int_{t_0}^{t_{
m fin}}rac{dt_1dt_2}{t_1t_2}\cos[pc_s(t_2-t_1)]\cos[ilde pc_s(t_2-t_1)]
onumber \ \Delta^2(k,p, ilde p,t_{
m fin}) = \int_{t_0}^{t_{
m fin}}\int_{t_0}^{t_{
m fin}}rac{dt_1dt_2}{t_1t_2}\cos[pc_s(t_2-t_1)]
onumber \ \Delta^2(k,p, ilde p,t_{
m fin})
onumber \ \Delta^2(k,p, ilde p,t_{
m fin})
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m fin})
onumber \ \Delta^2(k,p,t_{
m fin})
onumber \$$

Modelizing the mixed motion spectra

$$\Omega_{\mathrm{GW}}(k,t) = (1-q)^2 \Omega_{\mathrm{GW}}^{\mathrm{vort}}(k,t) + q^2 \Omega_{\mathrm{GW}}^{\mathrm{comp}}(k,t) + q(1-q) \Omega_{\mathrm{GW}}^{\mathrm{mix}}(k,t)$$

Constant-in-time Model

$$E_K^{
m vort}(k,t_1,t_2)=E_K^*(k)$$

Sound-Shell Model

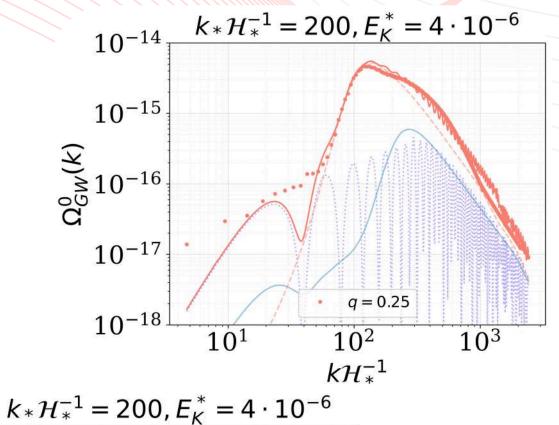
$$E_K^{
m comp}(k,t_1,t_2) = E_K^*(k) \cos[kc_s(t_2-t_1)]$$

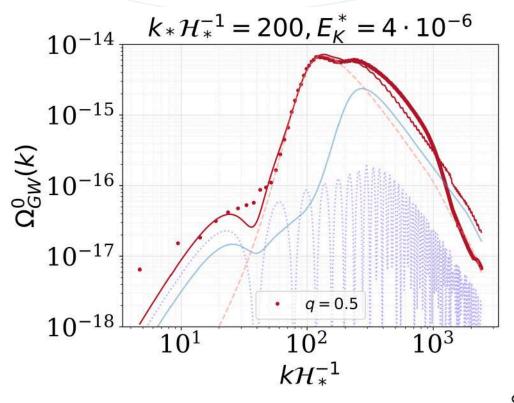
New Model

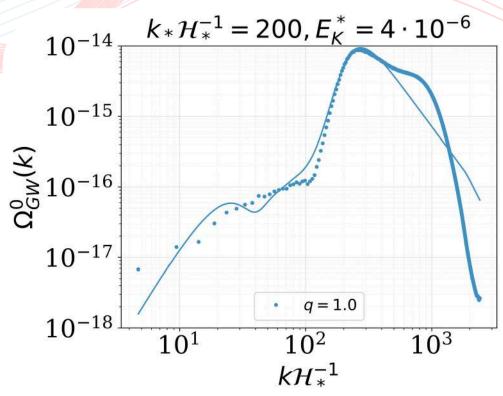
$$\Omega_{
m GW}^{
m mix}(k,t_0) = 3w^2k^3F_{
m GW}^0\int_0^\infty dp\,p^2E_K^*(p)\int_{-1}^1 dz\,rac{(1-z^4)}{ ilde{p}^4}E_K^*(ilde{p})\Delta^2(k, ilde{p},t_{
m fin})
onumber \ \Delta^2(k, ilde{p},t_{
m fin}) = \int_{t_*}^{t_{
m fin}}\int_{t_*}^{t_{
m fin}}rac{dt_1dt_2}{t_1t_2} \!\cos[ilde{p}c_s(t_2-t_1)]\cos[k(t_2-t_1)]
onumber \ \Delta^2(k, ilde{p},t_{
m fin})$$

Modelizing the mixed motion spectra

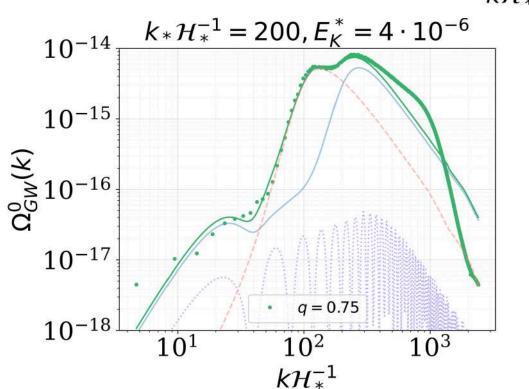
$$\Omega_{\mathrm{GW}}(k,t) = (1-q)^2 \Omega_{\mathrm{GW}}^{\mathrm{vort}}(k,t) + q^2 \Omega_{\mathrm{GW}}^{\mathrm{comp}}(k,t) + q(1-q) \Omega_{\mathrm{GW}}^{\mathrm{mix}}(k,t)$$

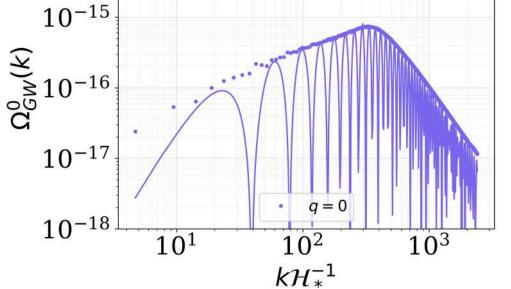






29

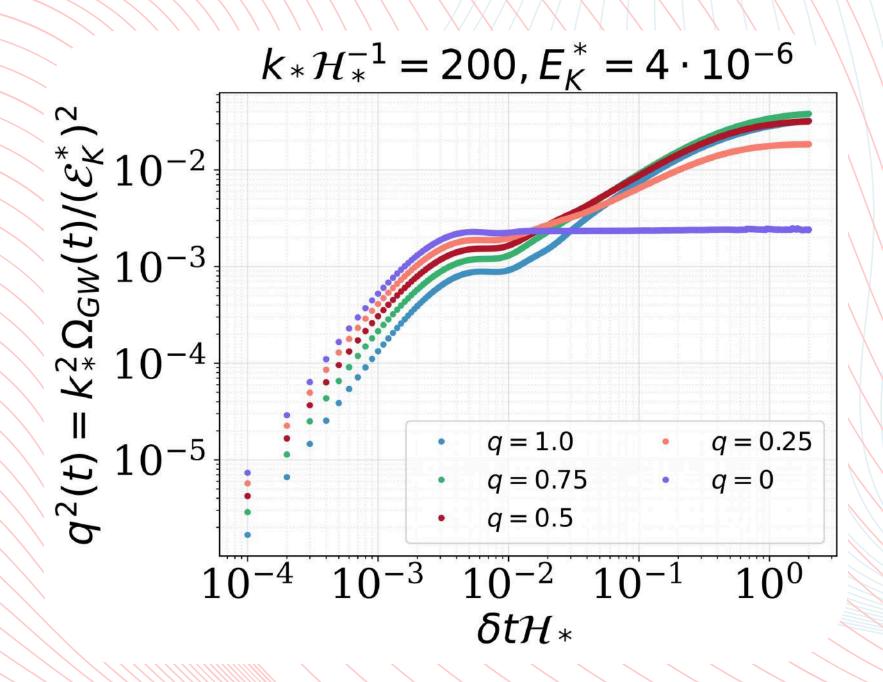


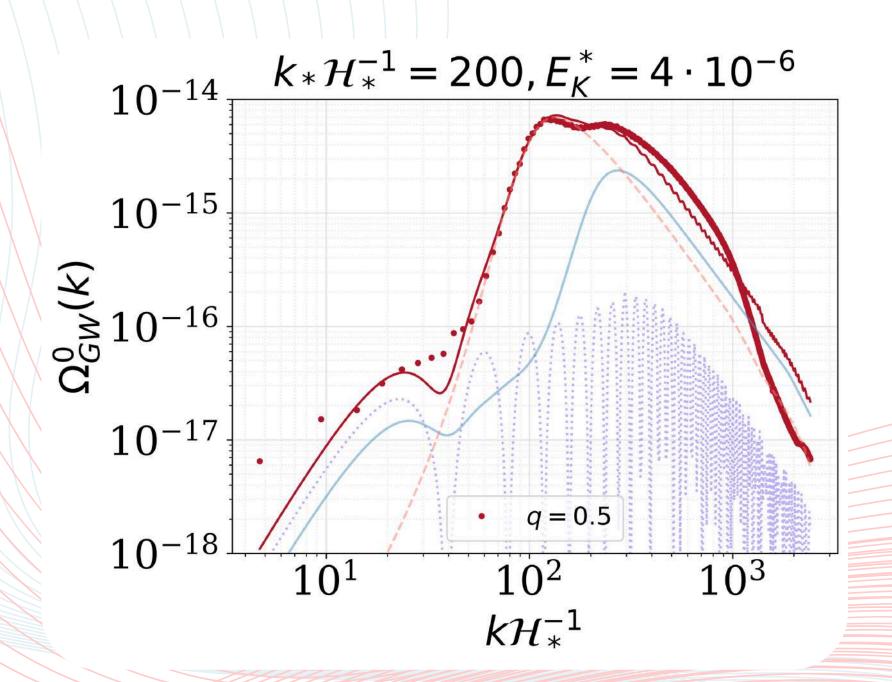


 10^{-14}

[2] A. Roper Pol, M. Salomé. "Gravitational wave production from acoustic and fractionally compressible sources."

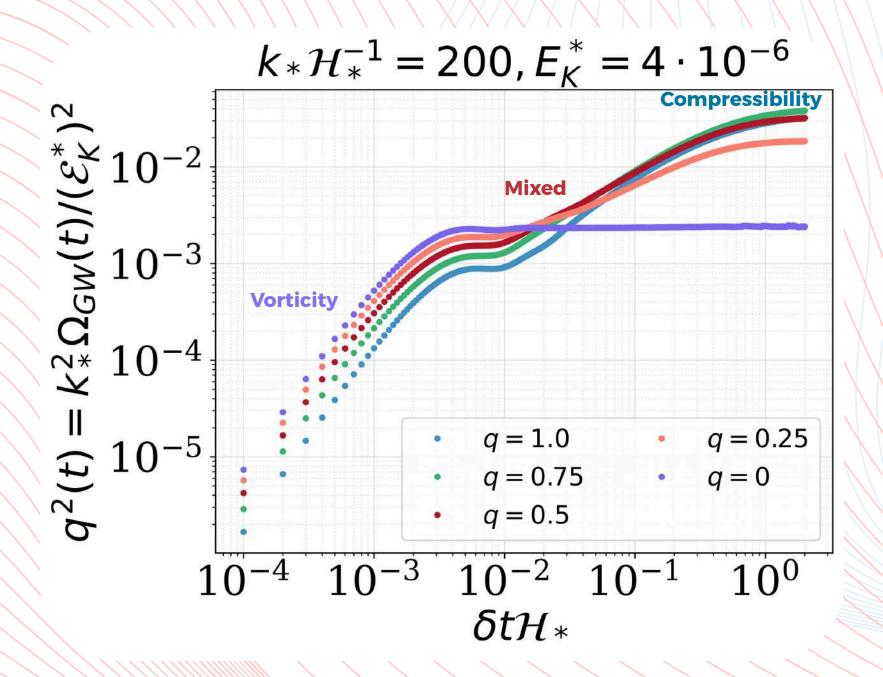
GWs production efficiency Which mode is the most effective for producing GWs?

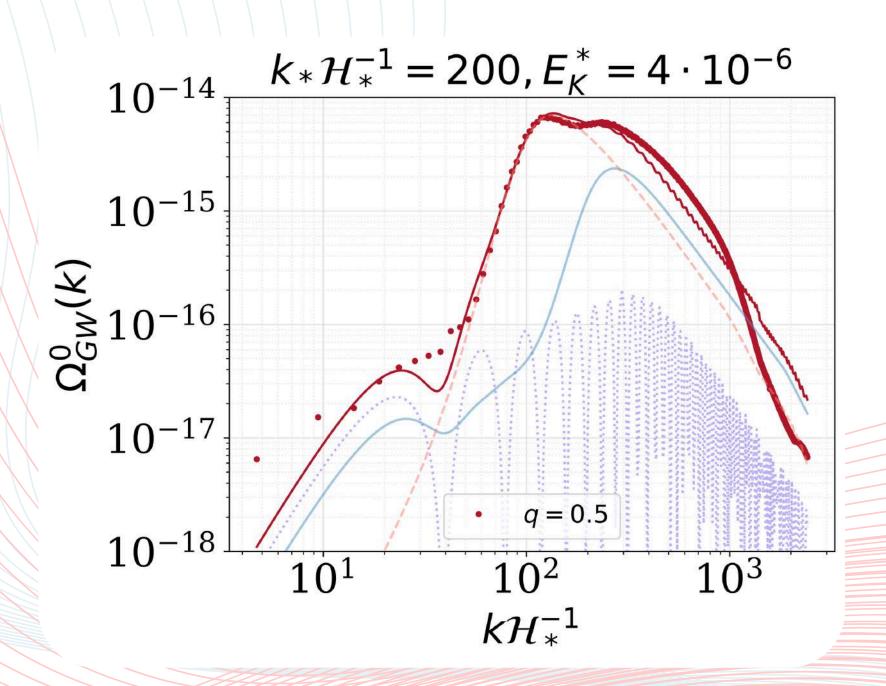




It depends on the sourcing duration and the scale of interest.

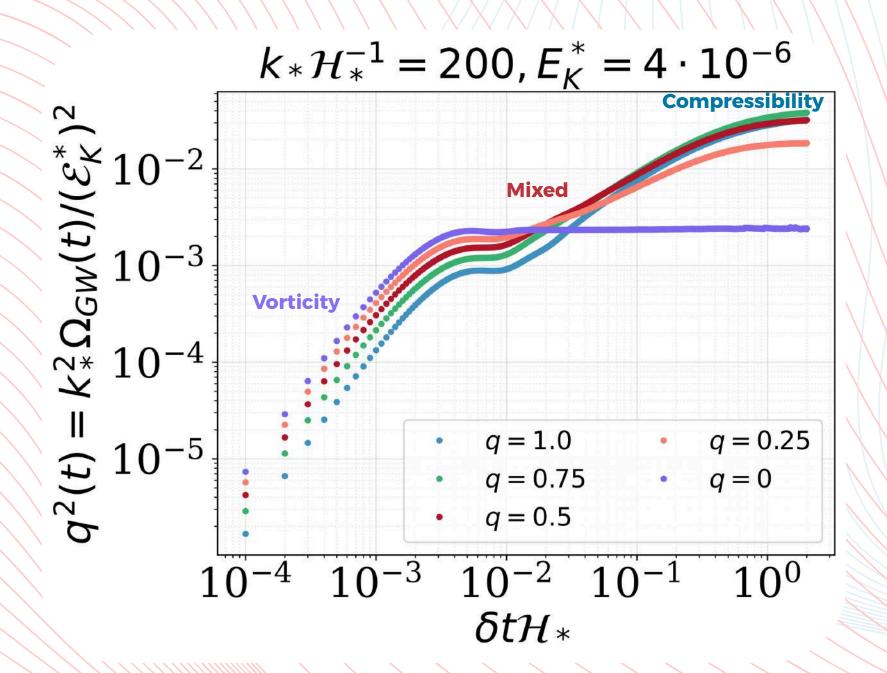
GWs production efficiency Which mode is the most effective for producing GWs?

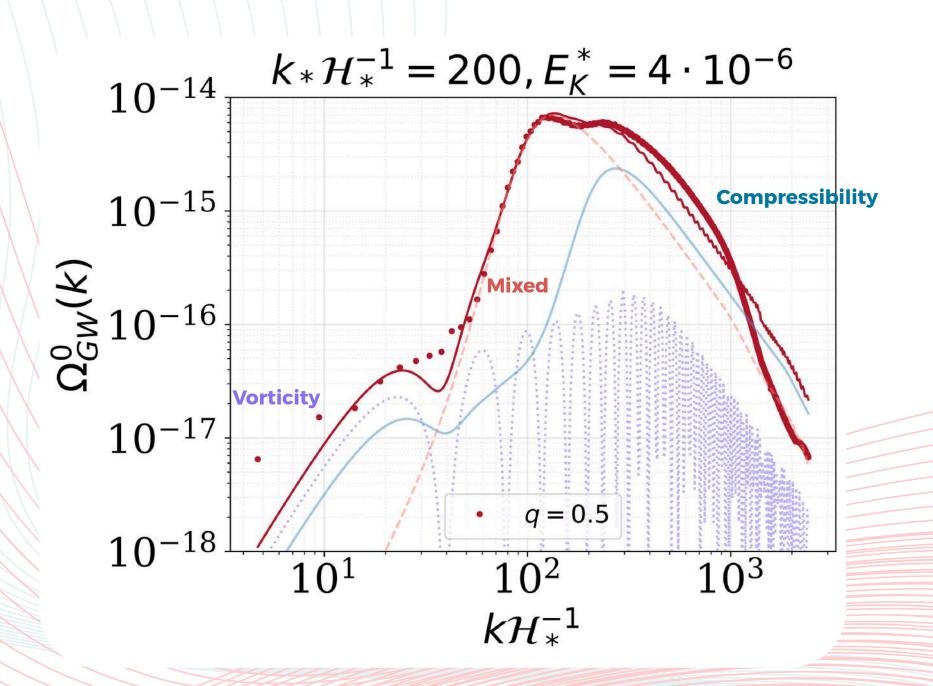




It depends on the sourcing duration and the scale of interest.

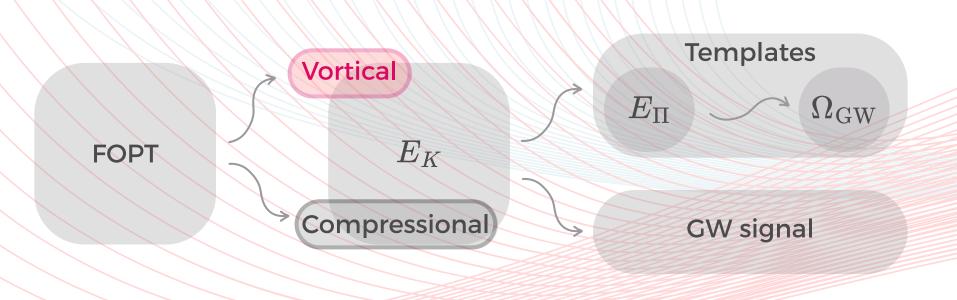
GWs production efficiency Which mode is the most effective for producing GWs?





It depends on the sourcing duration and the scale of interest.

Part IV Vorticity production in FOPT



Vorticity production in FOPT

An overview

$$\omega = \nabla \times \mathbf{u}$$

$$rac{\partial \, \mathbf{u}}{\partial t} = -\mathbf{u} \cdot
abla \mathbf{u} + rac{\mathbf{u}}{3} \psi - rac{3}{4} rac{
abla p}{
ho} + rac{3}{4
ho} \mathbf{J} imes \mathbf{B} + rac{2}{
ho}
abla \cdot (
ho
u \mathbf{S})$$

$$\frac{\partial \omega}{\partial t} + (\mathbf{u} \cdot \nabla)\omega = (\omega \cdot \nabla)\mathbf{u} - \omega(\nabla \cdot \mathbf{u}) \Rightarrow \text{Baroclinic term.} \\ + \frac{3}{4} \frac{1}{\rho^2} \nabla \rho \times \nabla p \Rightarrow \\ + \frac{1}{3} \omega \psi + \frac{\mathbf{u}}{3} \nabla \times \psi \Rightarrow \text{Relativistic term.}$$

$$+\frac{1}{3}\omega\psi+\frac{\mathbf{u}}{3}\nabla\times\psi$$

$$+rac{4}{3}
abla imes \left(rac{\mathbf{J} imes \mathbf{B}}{
ho}
ight)$$
 Magnetic term.

$$+ \nu \nabla^2 \omega + 2\nu \nabla \times (\mathbf{S} \nabla \ln \rho)$$

Viscous term.

Vorticity production in FOPT An overview

$$\frac{\partial \omega}{\partial t} + (\mathbf{u} \cdot \nabla)\omega = (\omega \cdot \nabla)\mathbf{u} - \omega(\nabla \cdot \mathbf{u}) \Rightarrow \text{Baroclinic term.} \qquad p = c_s^2 \rho \\ + \frac{3}{4} \frac{1}{\rho^2} \nabla \rho \times \nabla \rho \\ + \frac{1}{3} \omega \psi + \frac{\mathbf{u}}{3} \nabla \times \psi \\ + \frac{4}{3} \nabla \times \left(\mathbf{J} \times \mathbf{B} \right) \Rightarrow \text{Magnetic term.} \\ + \nu \nabla^2 \omega + 2\nu \nabla \times (\mathbf{S} \nabla \ln \rho)$$

Vorticity production in FOPT An overview

$$\frac{\partial \omega}{\partial t} + (\mathbf{u} \cdot \nabla)\omega = (\omega \cdot \nabla)\mathbf{u} - \omega(\nabla \cdot \mathbf{u}) \Rightarrow \text{Baroclinic term.}$$

$$+ \frac{3}{4} \frac{1}{\rho^2} \nabla \rho \times \nabla \rho$$

$$+ \frac{1}{3} \omega \psi + \frac{\mathbf{u}}{3} \nabla \times \psi$$

$$+ \frac{4}{3} \nabla \times \left(\mathbf{J} \times \mathbf{B} \right) \Rightarrow \text{Magnetic term.}$$

$$+ \nu \nabla^2 \omega + 2\nu \nabla \times (\mathbf{S} \nabla \ln \rho)$$

Vorticity production in FOPT An overview

$$rac{\partial \omega}{\partial t} + (\mathbf{u} \cdot
abla) \omega = rac{\mathbf{u}}{3}
abla imes \psi + 2
u
abla imes (\mathbf{S}
abla \ln
ho)$$

Deepen's talk!

Isak's Talk

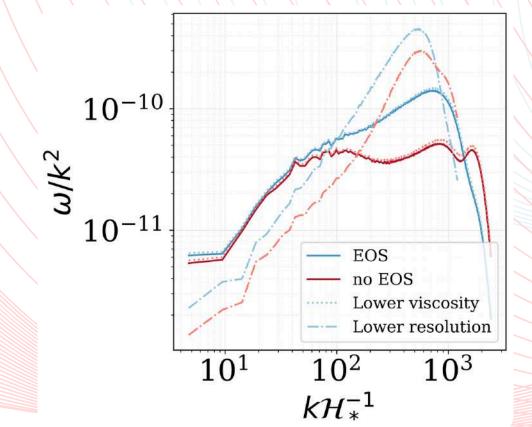
 $q \sim 0.7$

Could be even lower with

magnetic field!

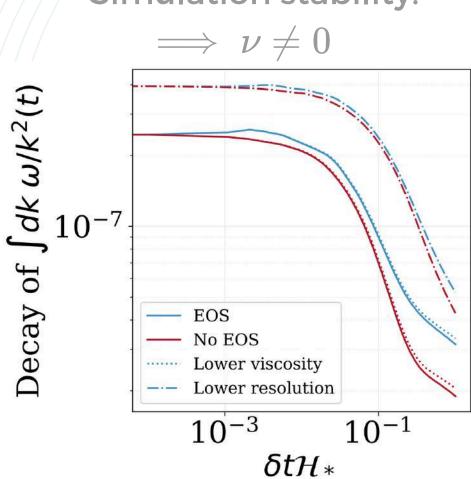
Relativistic term.

Relativistic EOS = EOS Non-rel. EOS = noEOS



Viscous term.

Simulation stability.



Few more words about the Pencil Code

Pencil Code School

1st Pencil Code school on early Universe physics and gravitational waves (Oct. 20-24)

Learning and developing numerical skills applied to early Universe physics using Pencil Code.



Pencil Code School

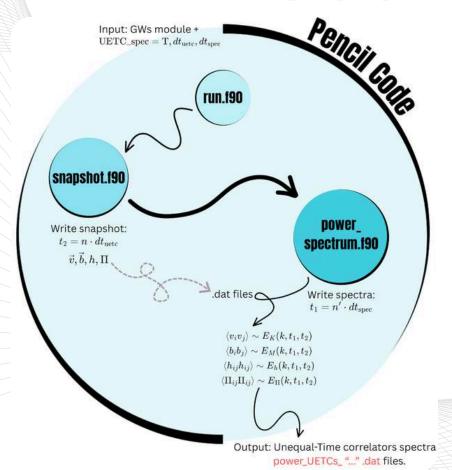
Topics:

- Magnetohydrodynamics of the early Universe,
- Generation and evolution of primordial magnetic fields,
- Chiral magnetohydrodynamics,
- First-order phase transitions,
- Gravitational wave production,
- Axion inflation.

Registration is open and will close on September 10th.

Decorrelation function of the UETCs

UETCs & the Pencil Code



Status To come!

$$E_\Pi(k,t_1,t_2) \sim f(k,t_1) f(k,t_2) \mathcal{D}_{E_\Pi}(k,t_2-t_1) \ f(k,t) = \sqrt{E_\Pi(k,t)}$$

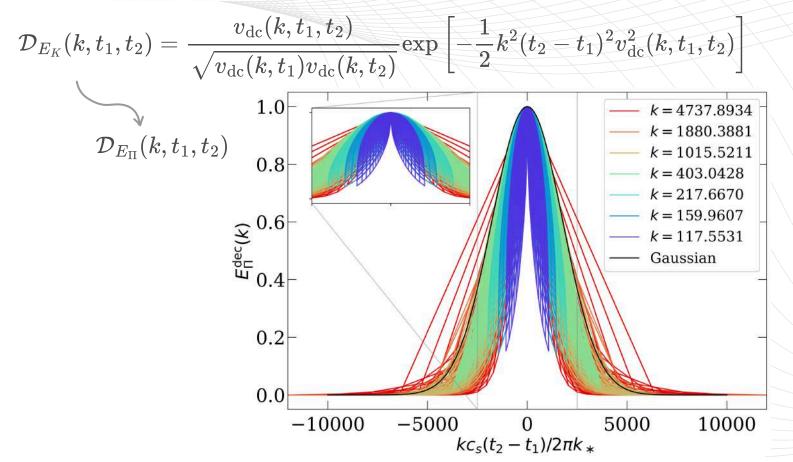
First Results Decorrelation functions.

Kraichnan decorrelation: Gaussian shape is found!

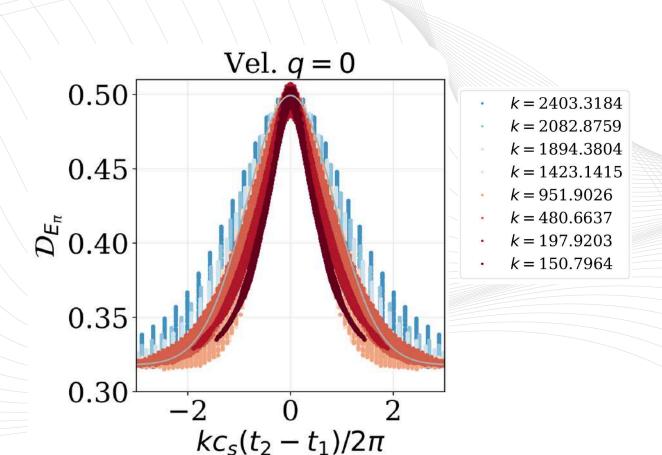
Mark's Talk

J. Correia & M. Hindmarsh & K. Rummukainen & D. J. Weir
"Gravitational waves from strong first order phase transitions."

Analytical description



Simulation



Conclusion

Gravitational wave production: the interplay between vortical and compressional motions.

Main results

Possibilities to source **vorticity** in relativistic primordial plasma leading to a mixed motion: Vortical & Compressional modes.

An additional peak visible in the GW energy density spectrum from **fractionally compressible motion**. A non-negligeable term in some cases.

New generic **model** to describe the GW energy density spectrum by combining the description of each motion: Constant-in-Time model & Sound Shell model.

Helicity part can be added at this generic template.

Thank you for listening!

References

cosmoGW

A. Roper Pol, GitHub project "CosmoGW", https://github.com/AlbertoRoper/CosmoGW

Pencil Code

Pencil Code Collaboration, J. Open Source Software 6, 2807 (2021). The Pencil Code, a modular MPI code for partial differential equations and particles: multipurpose and multiuser-maintained. arXiv: 2009.08231

arXiv:1801.04268 C. Caprini & D.G. Figueroa - Cosmological backgrounds of gravitational waves.

arXiv:2308.12916 R. Sharma & J. Dahl & A. Brandenburg & M. Hindmarsh - Shallow relic gravitational wave spectrum with acoustic peak.

arXiv:1711.03804 A. Brandenburg & T. Kahniashvili & al. - Evolution of hydromagnetic turbulence from electroweak phase transition.

arXiv:2308.12943

A. Roper Pol & S. Procacci & C. Caprini - Characterization of the gravitational wave spectrum from sound waves within the sound shell model.

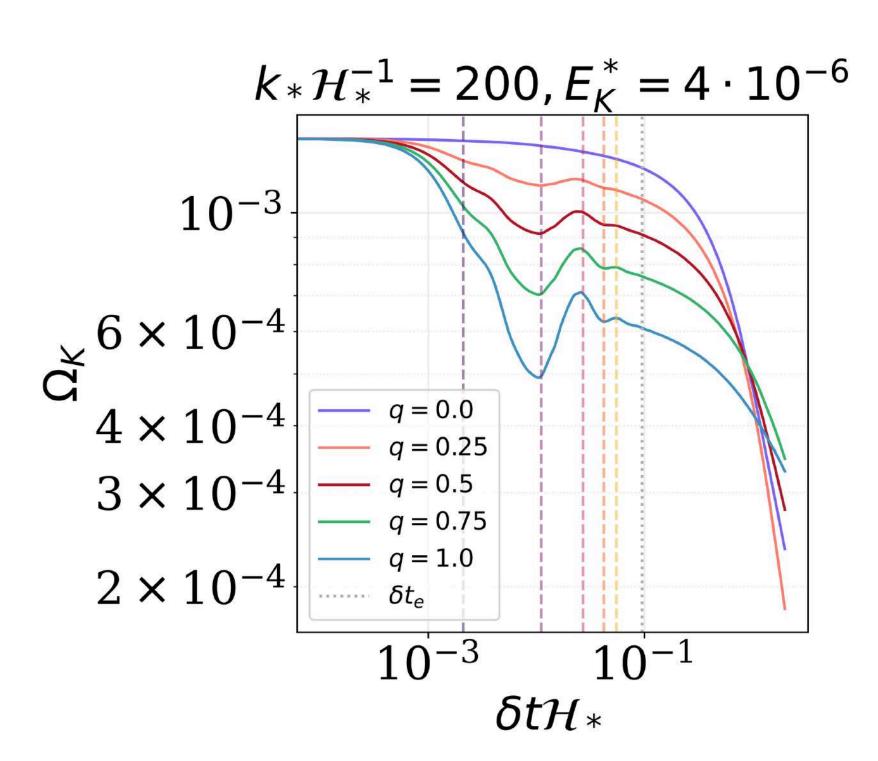
arXiv:2409.01426

L. Giombi & J. Dahl & M. Hindmarsh - Signatures of the speed of sound on the gravitational wave power spectrum from sound waves.

J. Correia & al. - Gravitational waves from strong first order phase transitions.

Slowly decaying sources

Time evolution curve



Vortical description of the decay

$$\Omega_K(t) = \Omega_K^* igg(1 + rac{\delta t}{\delta t_e}igg)^{-p}$$

Exponentially damped cosinus

$$(e^{-A\delta t}\cos[B\cdot 2c_sk_*\delta t]+C)\cdot D+E$$

Slowly decaying sources

Parametrization of the decay

$$\Omega_K(t) = \Omega_K^* imes \left(1 + rac{\delta t}{\delta t_{
m dec}}
ight)^{-p} imes \left(e^{-A\delta t} imes (\cos\left[B \cdot 2c_s k_* \delta t
ight] + C
ight) \cdot D + E
ight)$$

Name	$\delta t_{ m dec}$	p	A	B	C	$\log_{10} D$	$\log_{10} E$
q0	1.8	2.35	-	-	-	-	-
q025	1.8	2.35	170	1.0	0.0	-1.4202	-0.6990
q05	1.8	1.70	130	1.0	0.0	-1.4202	-1.0000
q075	1.8	1.10	110	1.0	0.0	-1.4202	-1.2219
q1	1.8	1.00	90	1.0	0.0	-1.4202	-1.3979

$$p(q) = q^{-0.65}$$
 $A(q) = 90 \cdot q^{-0.5}$
 $E(q) = 0.04 \cdot q^{-1.2}$

 $k_*\mathcal{H}_*^{-1} = 200, E_K^* = 4 \cdot 10^{-6}$ 10^{-3} 3×10^{-4} 2×10^{-4} 10^{-4} 10^{-3} 10^{-2} 10^{-1} 10^{0} $\delta t\mathcal{H}*$

Further studies need to be performed!