TOHOKU UNIVERSITY

COSMIC WISPERS
CA21106

More is Different: Multi-Axion Dynamics Changes Topological Defect Evolution

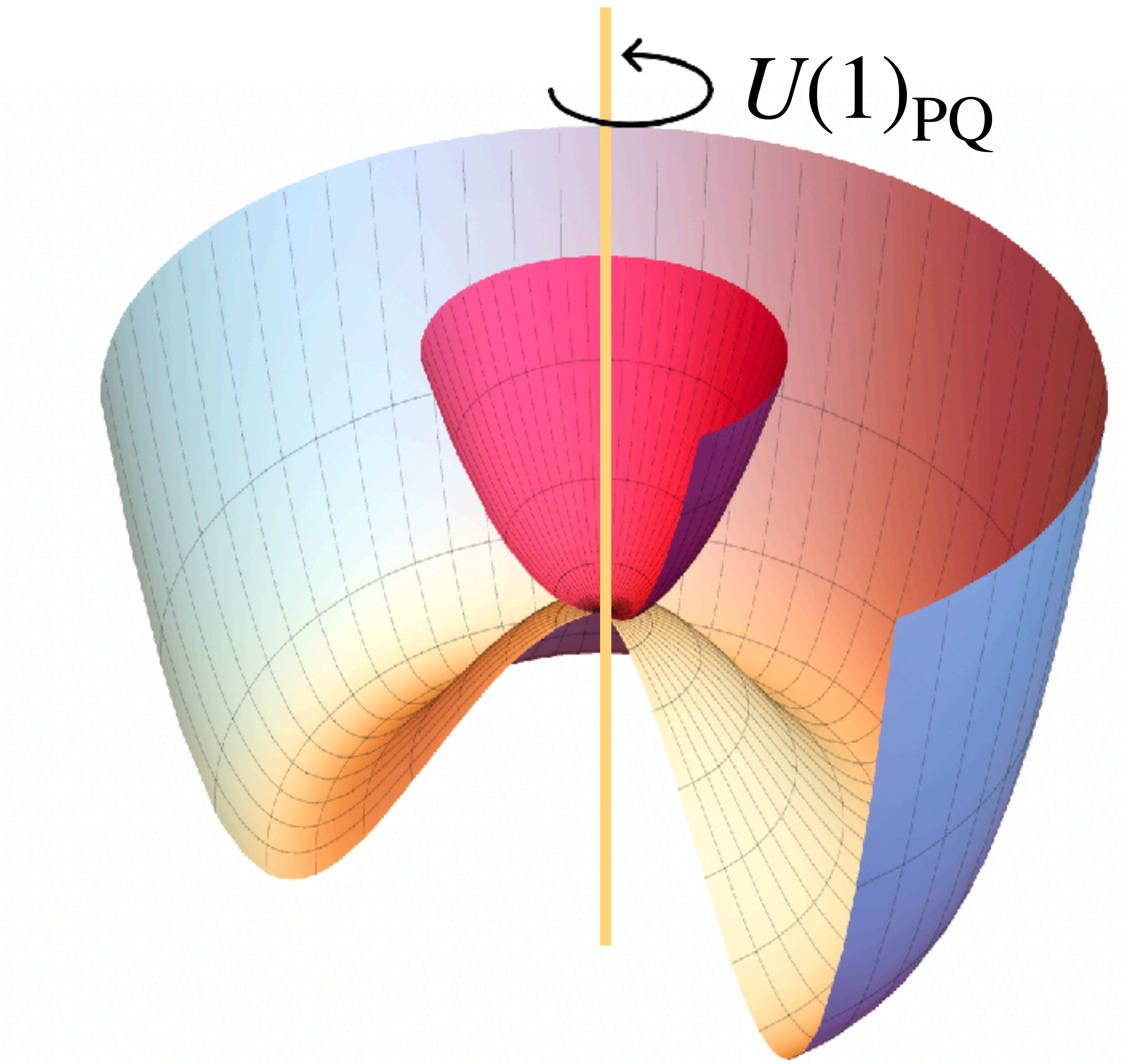
July 2, 2025
@Axions in Stockholm

Fumi Takahashi (Tohoku University)

Based on Lee, Murai, FT and Yin 2409.09749, 2407.09478

QCD axion

The QCD axion is a pseudo Nambu-Goldstone boson associated with SSB of U(1) Peccei-Quinn symmetry.

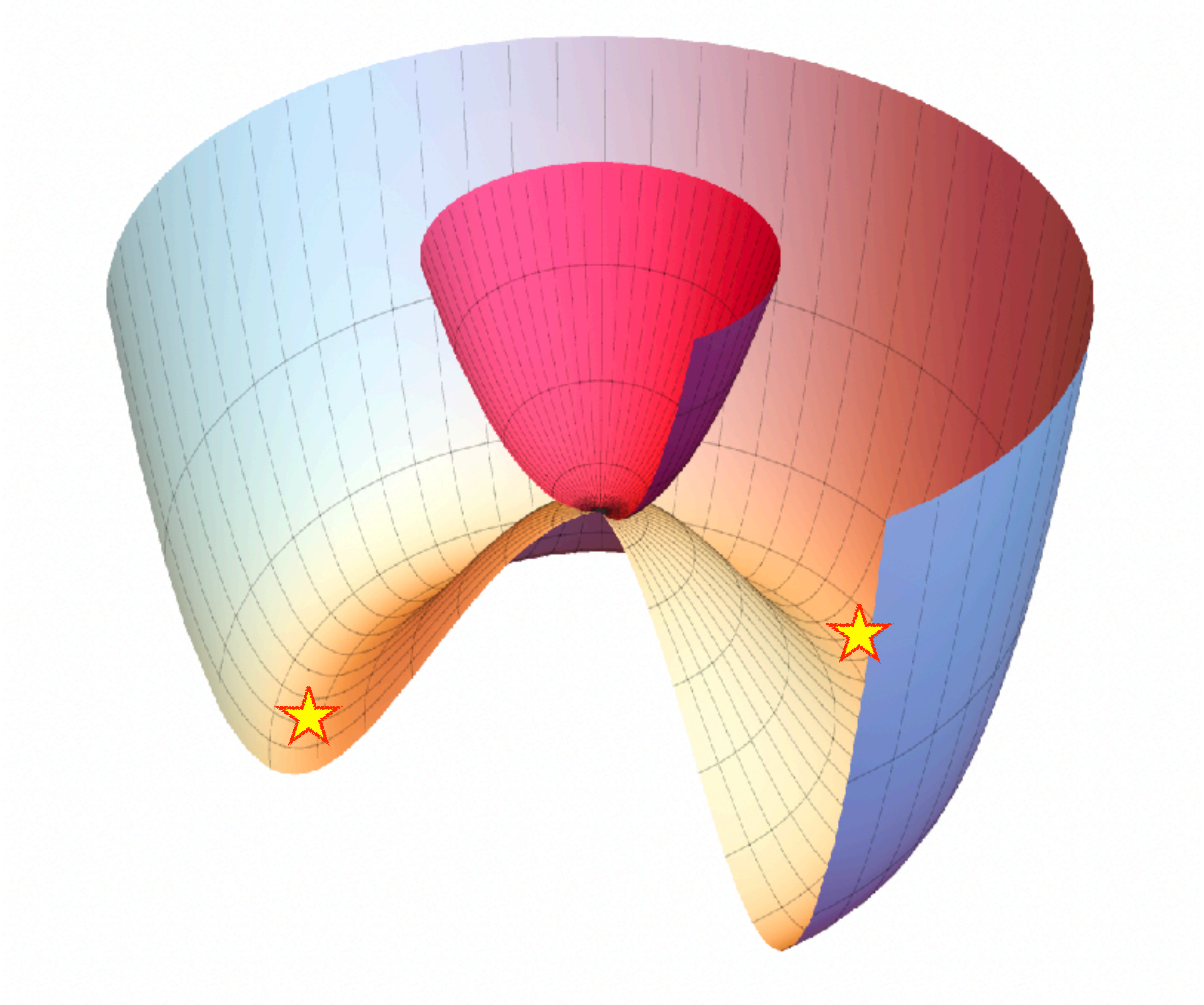


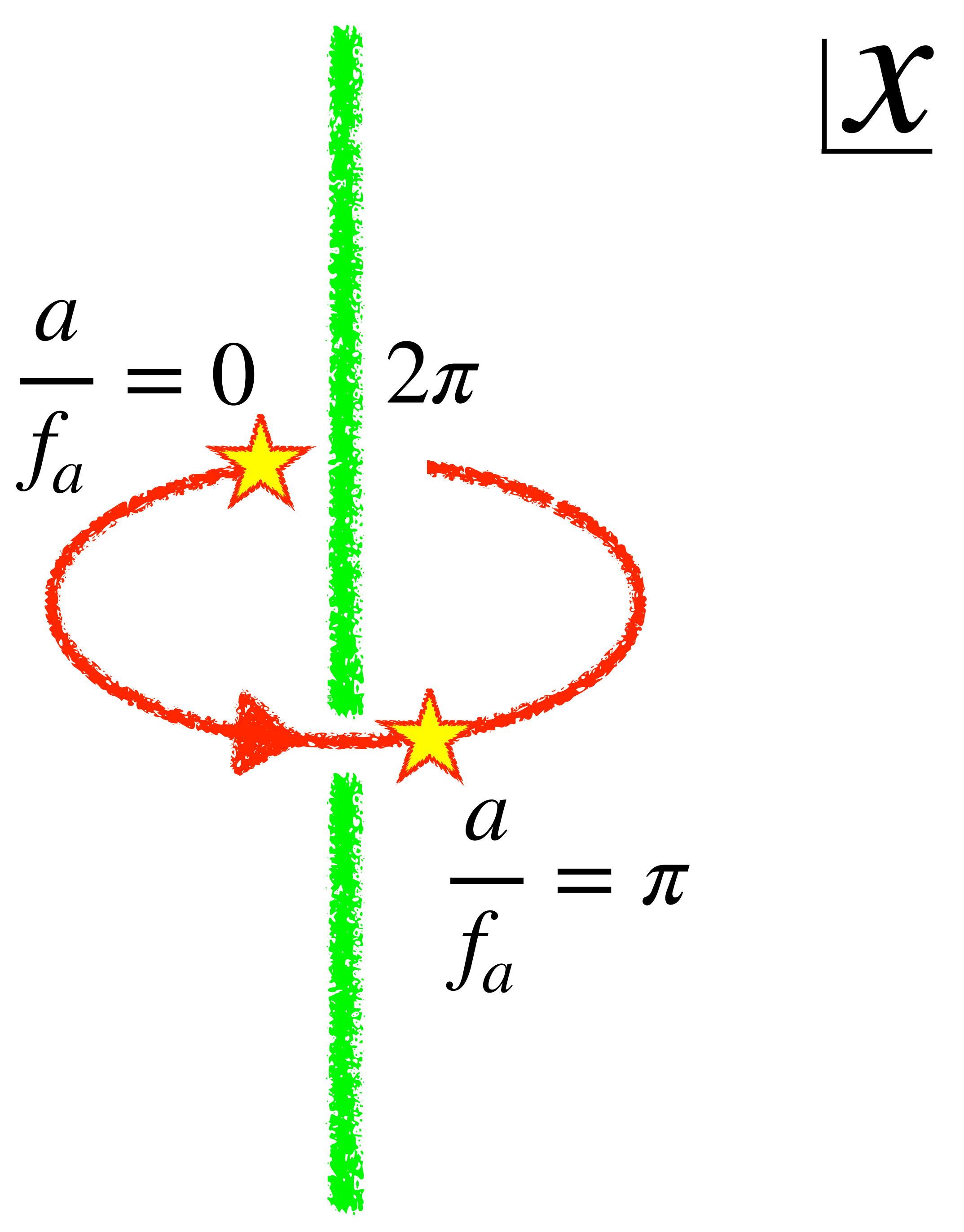
Peccei, Quinn `77, Weinberg `78, Wilczek `78

Axion production from strings/walls

Cosmic strings and domain walls, formed in post-inflationary scenarios, produce axion dark matter.

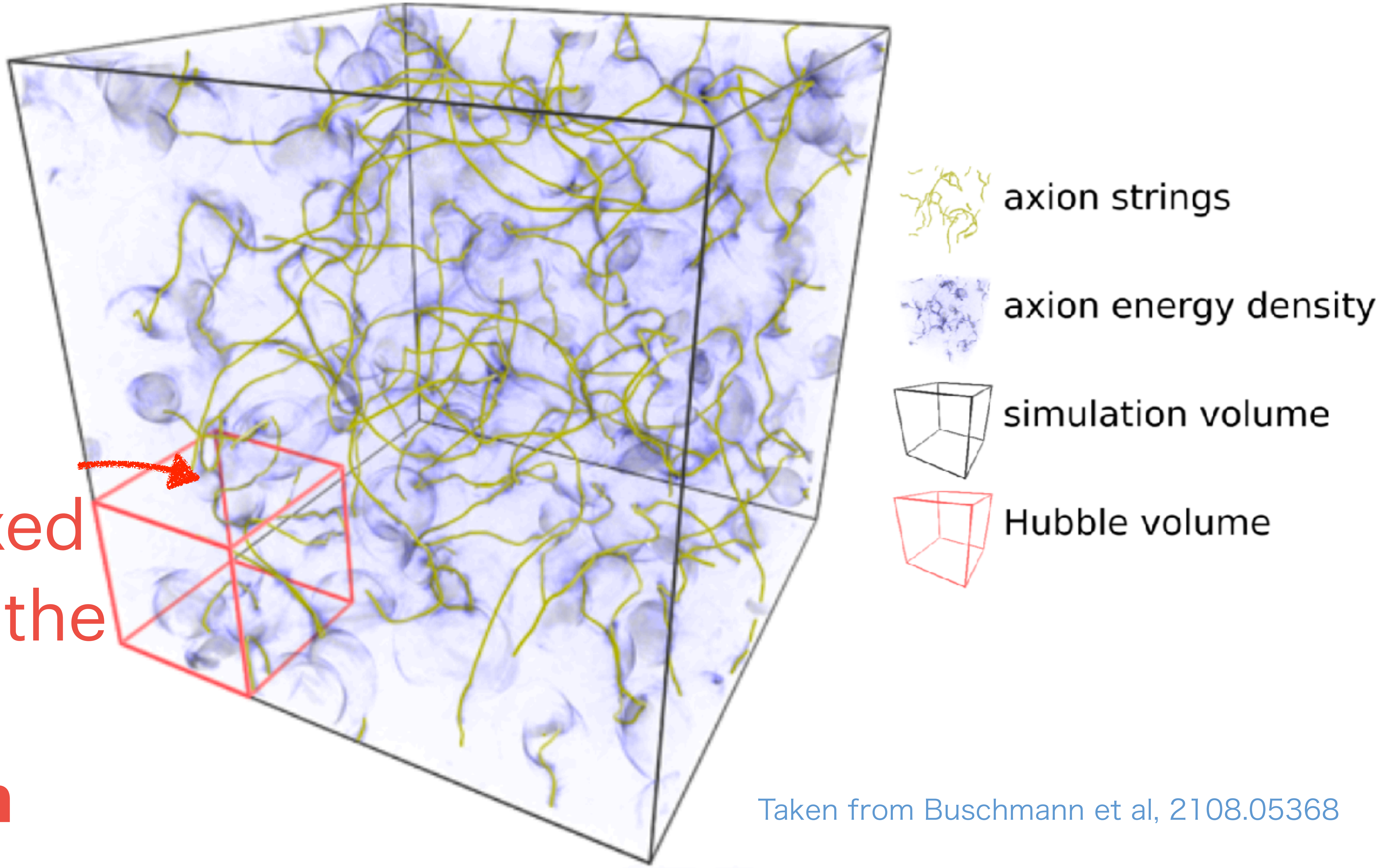
See talks by Long, Kaltschmidt, Gelmini





Cosmic strings

Axion production from strings/walls



(Approximately) fixed numer of strings in the Hubble volume:
Scaling solution

Taken from Buschmann et al, 2108.05368

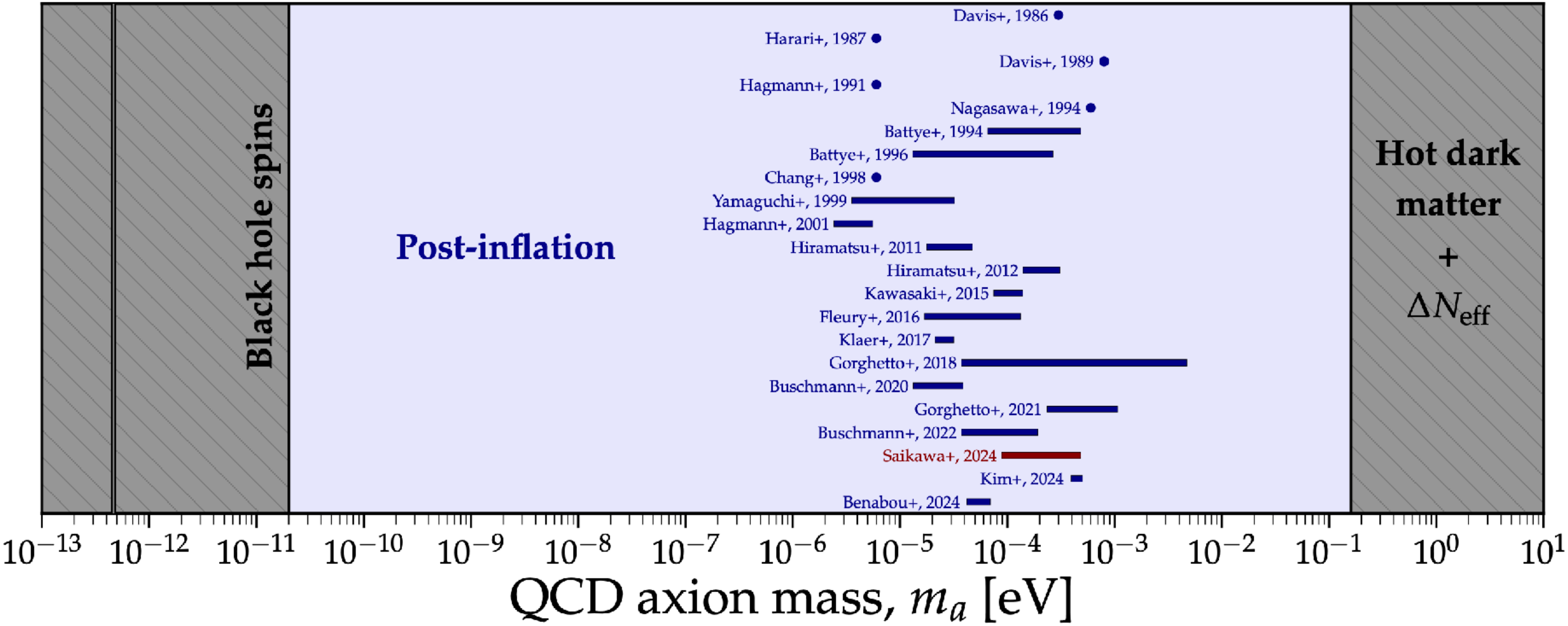
Axion production from strings/walls

$N_{\rm DW} = 1$

Hiramatsu, Kawasaki, Saikawa, Sekiguchi,, 1202.5851

Studying the evolution of strings and domain walls is crucial for predicting the mass of axion dark matter.

Prediction of post-inflationary scenario

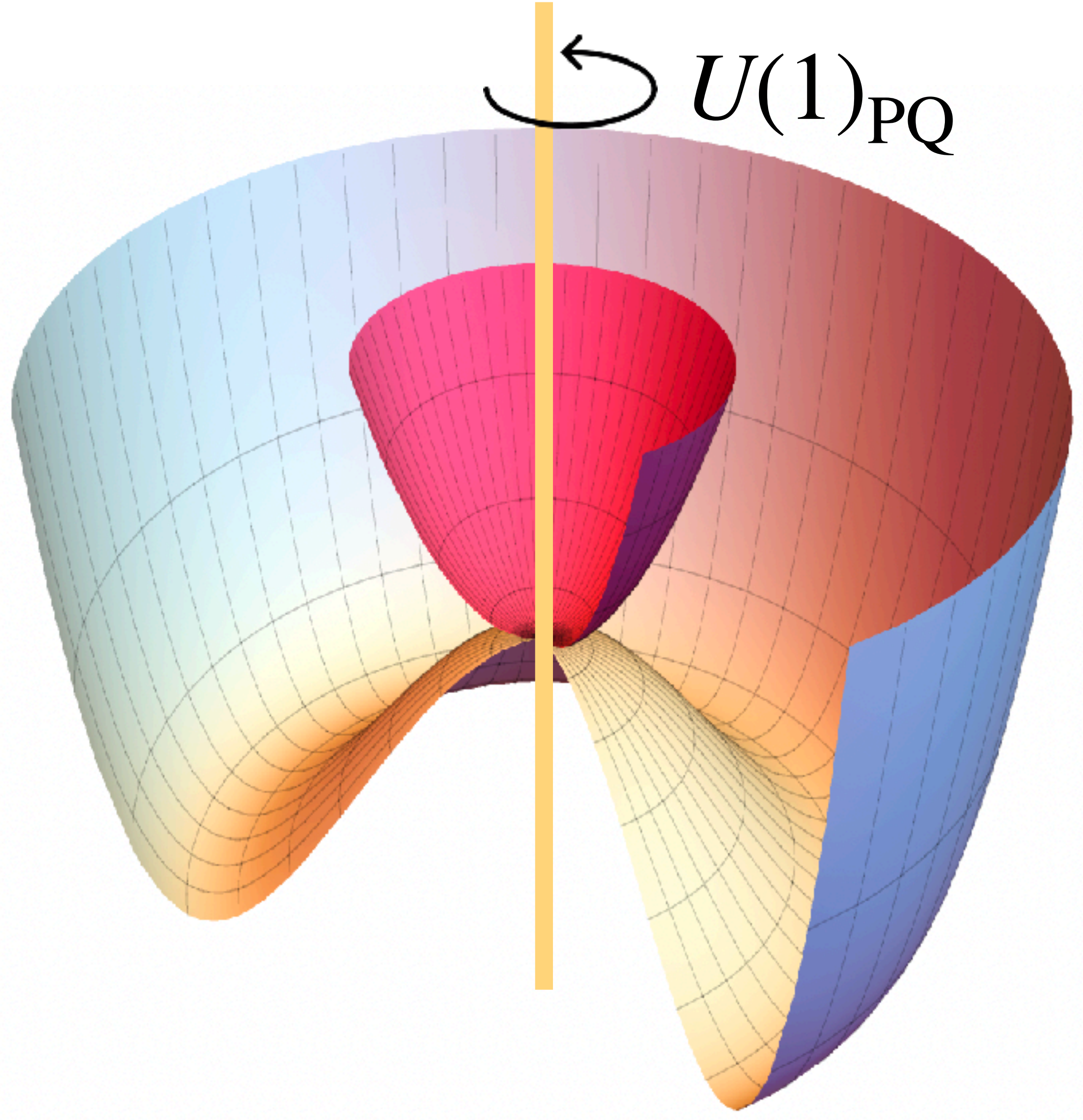


See talk by Kaltschmidt

Kaltschmidt, Redondo, Saikawa and Vaquero, arXiv:2502.02398

Key assumption :

Single PQ scalar model

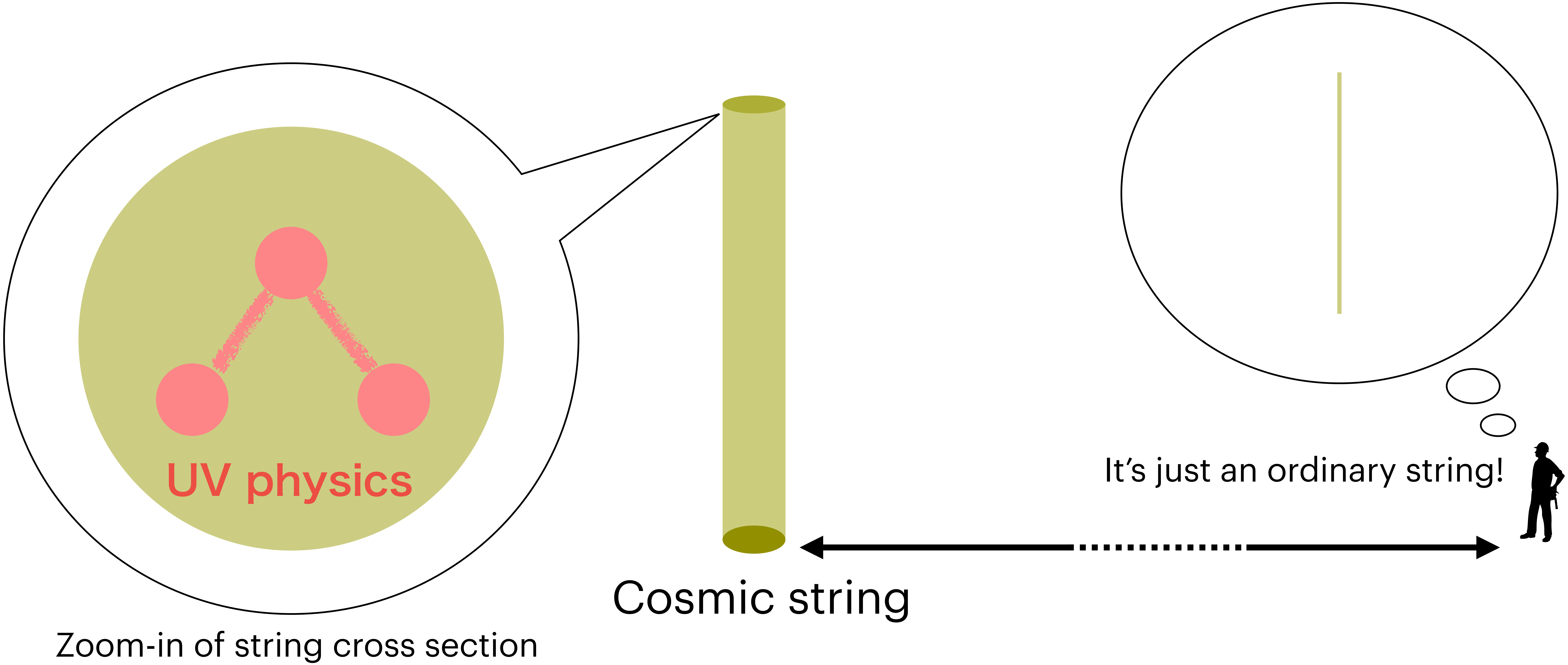


$$\Phi = \frac{f_a + s}{\sqrt{2}} e^{ia/f_a}$$

- Prevalent in high-precision lattice calculations.
- Simplifying assumption or crucial factor?

Origin and breaking of U(1) PQ are unknown!

Is UV physics always confined in the string core?



Two axion system

We introduce two PQ scalars

$$\Phi_1 = \frac{f_1}{\sqrt{2}} e^{i\frac{\phi_1}{f_1}} \quad \text{and} \quad \Phi_2 = \frac{f_2}{\sqrt{2}} e^{i\frac{\phi_2}{f_2}}$$

$\theta_1 \equiv \phi_1/f_1$ and $\theta_2 \equiv \phi_2/f_2$

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$\theta_1 \equiv \phi_1/f_1$ and $\theta_2 \equiv \phi_2/f_2$

and the potential for axions

$$V_1(\phi_1, \phi_2) = \Lambda^4 \left[1 - \cos\left(n_1\frac{\phi_1}{f_1} + n_2\frac{\phi_2}{f_2}\right)\right]$$

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$$V_2(\phi_1, \phi_2) = \Lambda'^4 \left[1 - \cos\left(n'_1 \frac{\phi_1}{f_1} + n'_2 \frac{\phi_2}{f_2} + \alpha\right)\right]$$

$$\Lambda \gg \Lambda'$$

$n_1, n_2, n'_1, n'_2 \in \mathbf{Z}$

with the post-inflationary initial condition.

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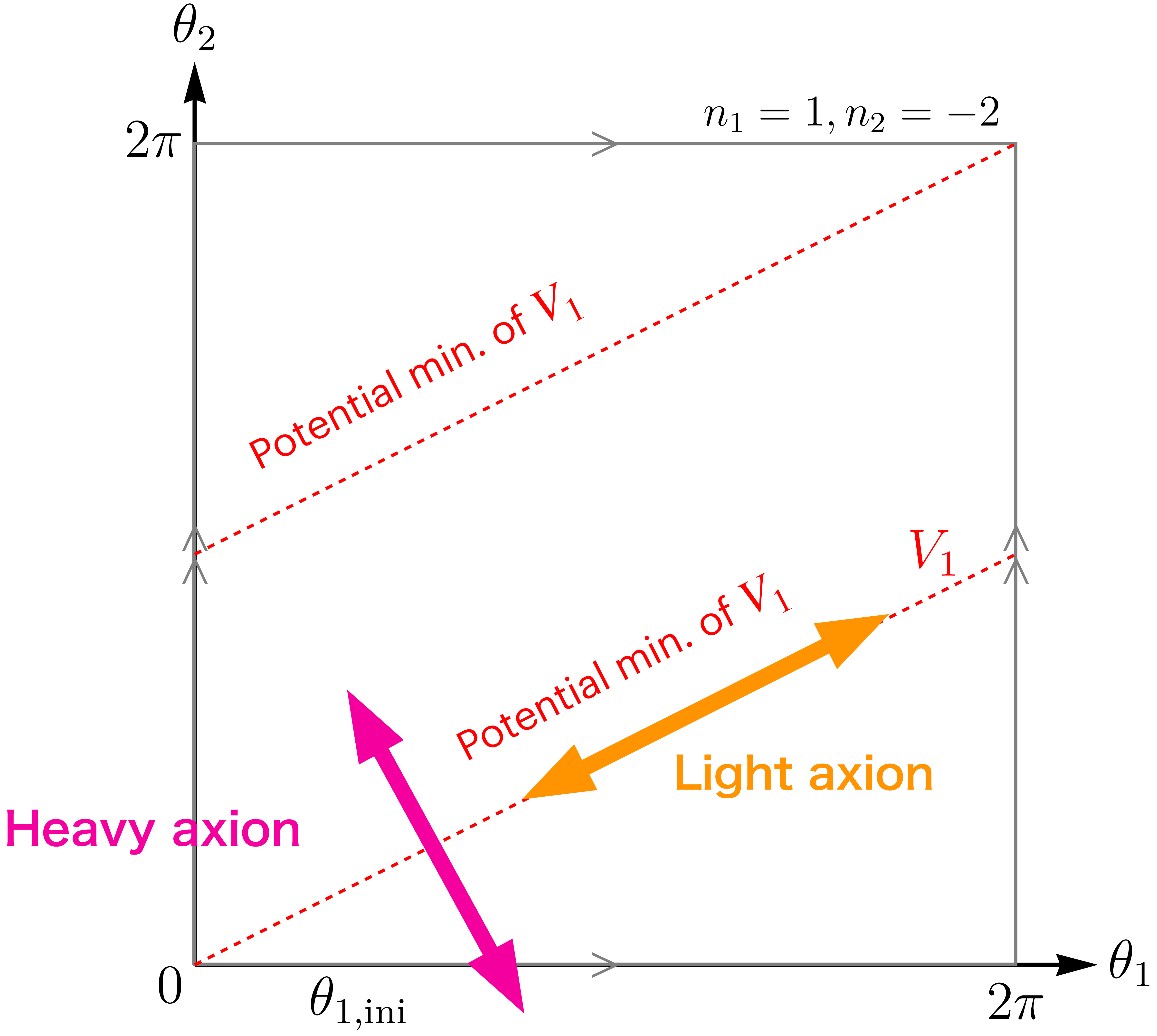
One linear combination of two axions becomes heavy, leaving the orthogonal one (nearly) massless.

$$\phi_{\rm heavy} \propto n_1 \frac{\phi_1}{f_1} + n_2 \frac{\phi_2}{f_2}$$



Two types of strings and multiple DWs

Both ϕ_1, ϕ_2-strings quickly reach the scaling solution, and when V_1 becomes relevant, n_1 (n_2) domain walls appear, attached to the $\phi_1(\phi_2)$-string.

①

②

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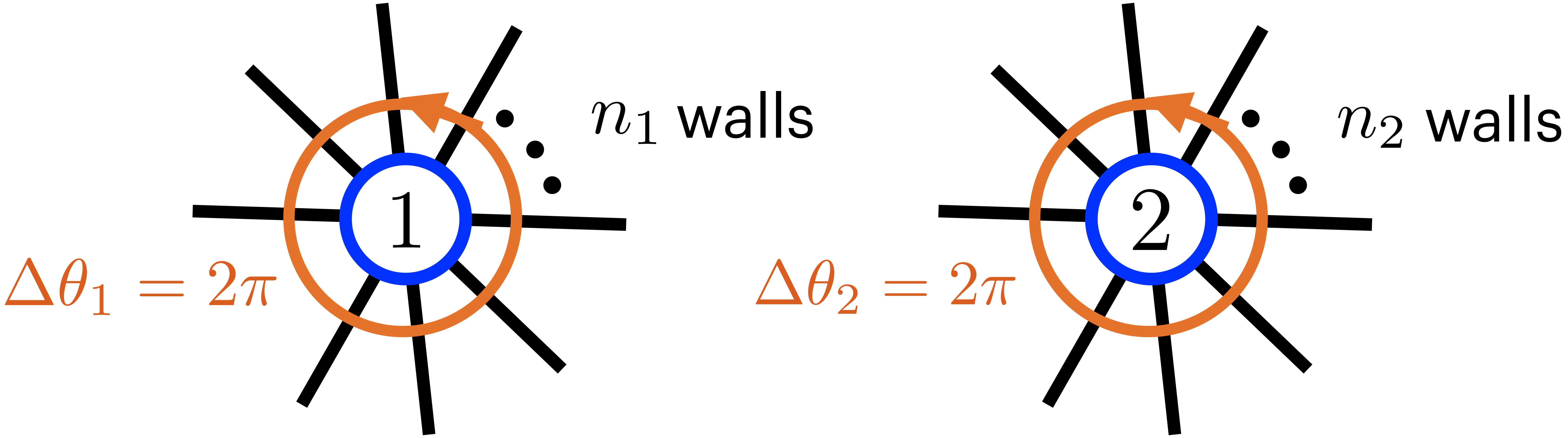

$$V_1(\phi_1, \phi_2) = \Lambda^4 \left[1 - \cos\left(n_1 \frac{\phi_1}{f_1} + n_2 \frac{\phi_2}{f_2}\right)\right] \qquad n_1, n_2 \in \mathbf{Z}$$

Case of $(n_1, n_2) = (2, 1)$

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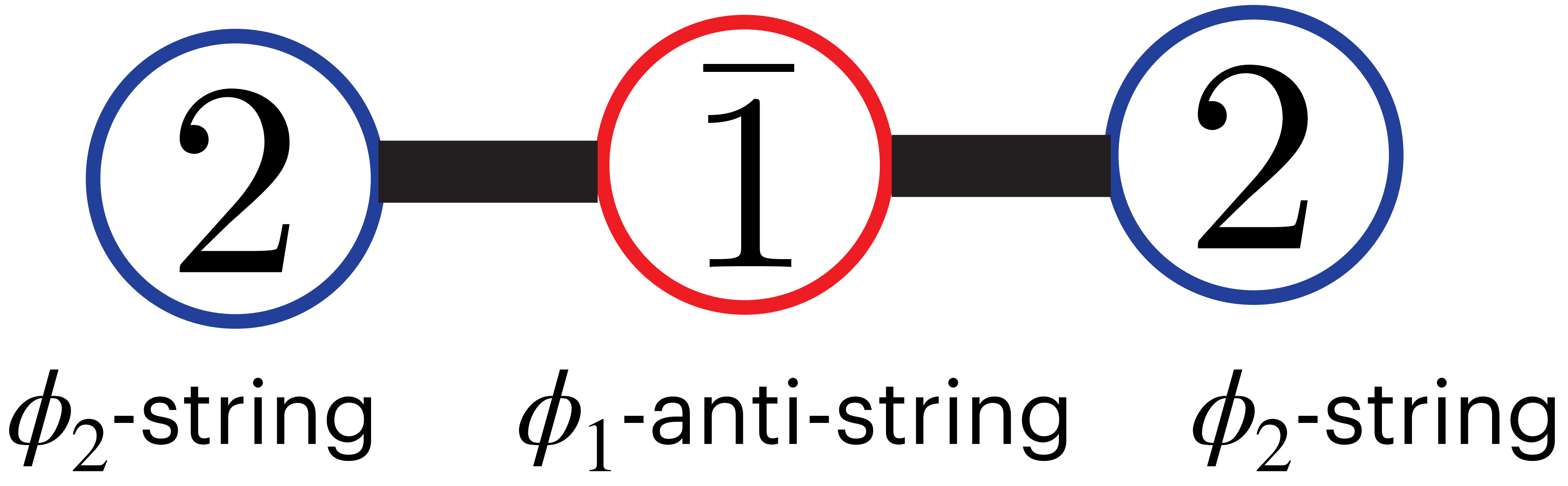

String bundle

(= ordinary cosmic string)

cf. Higaki, Jeong, Kitajima, Sekiguchi and FT, 1606.05552,
See also Eto, Hiramatsu, Saito and Sakakihara, 2309.04248

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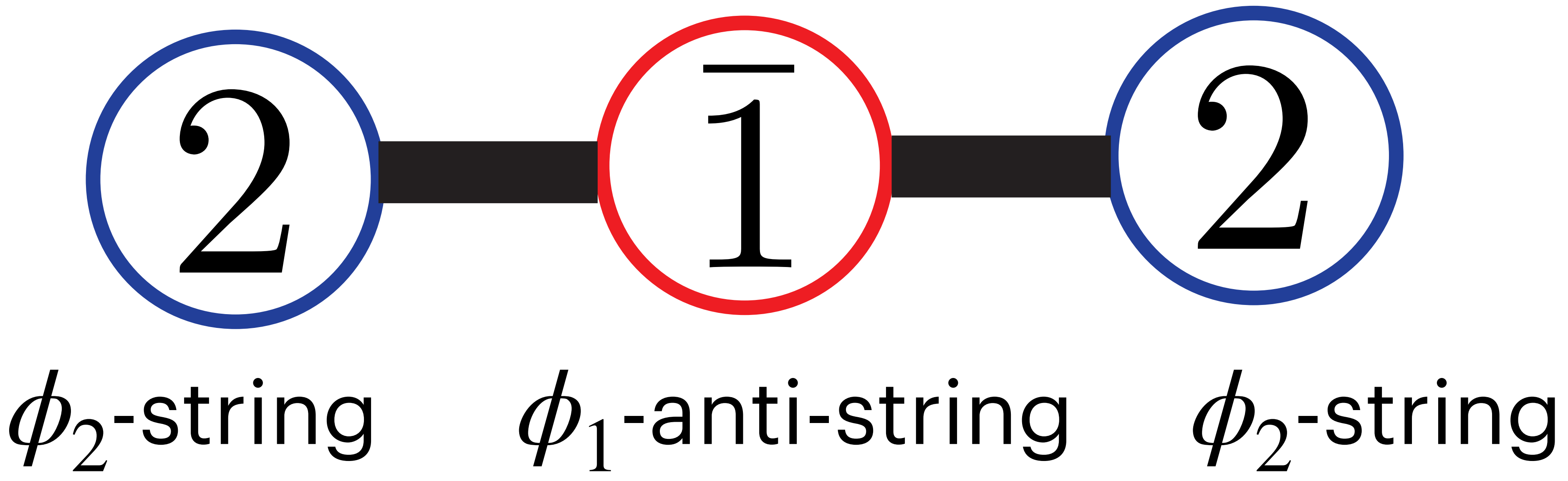

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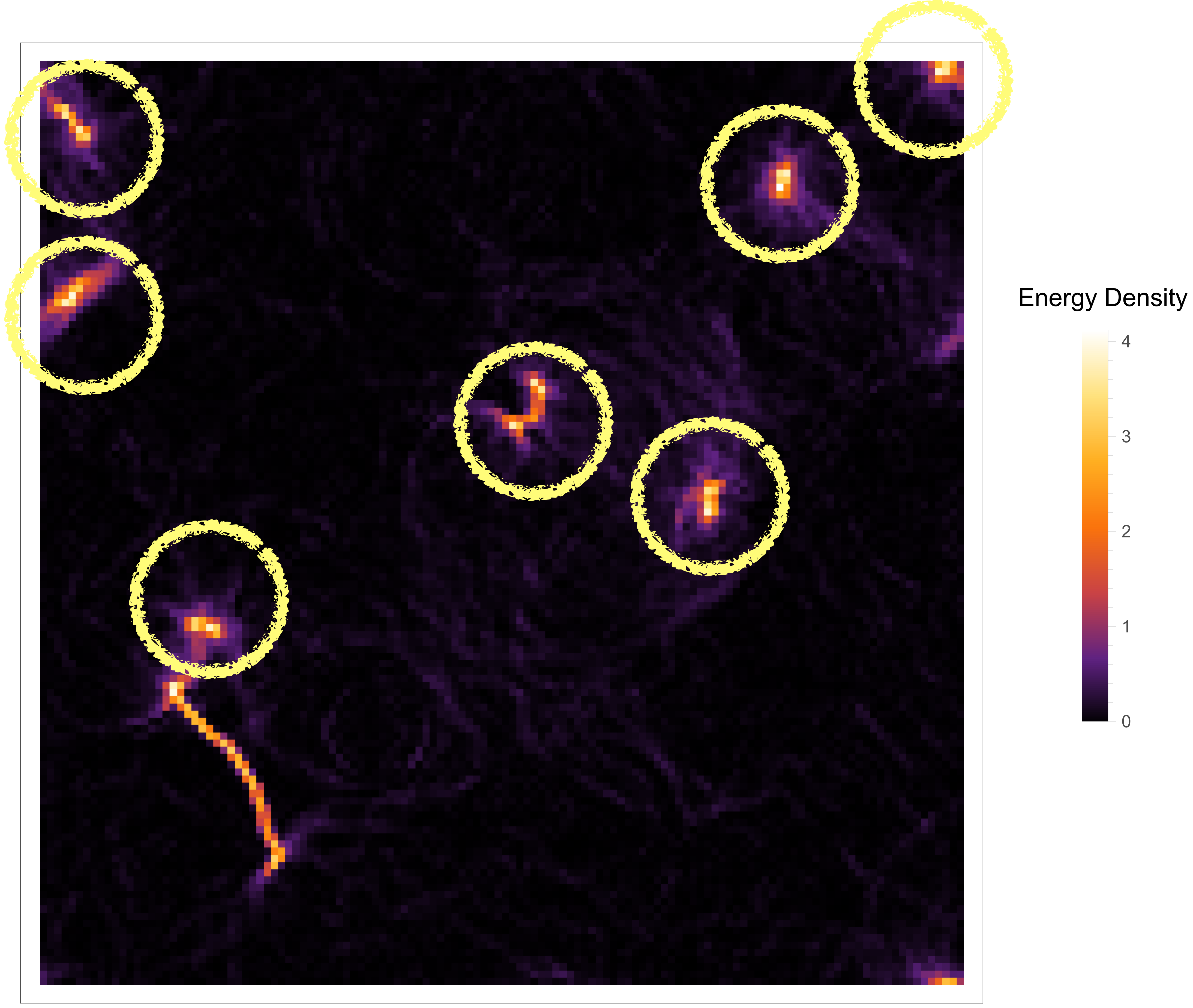

Numerical results (2D)

Lee, Murai, FT and Yin 2409.09749

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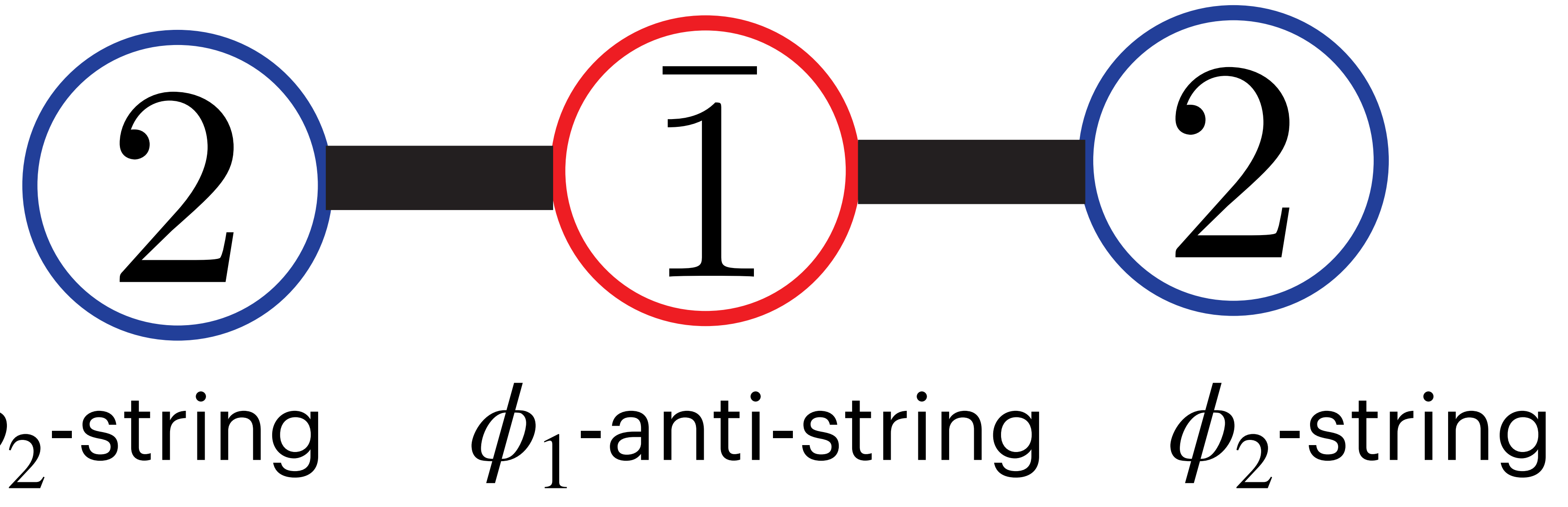

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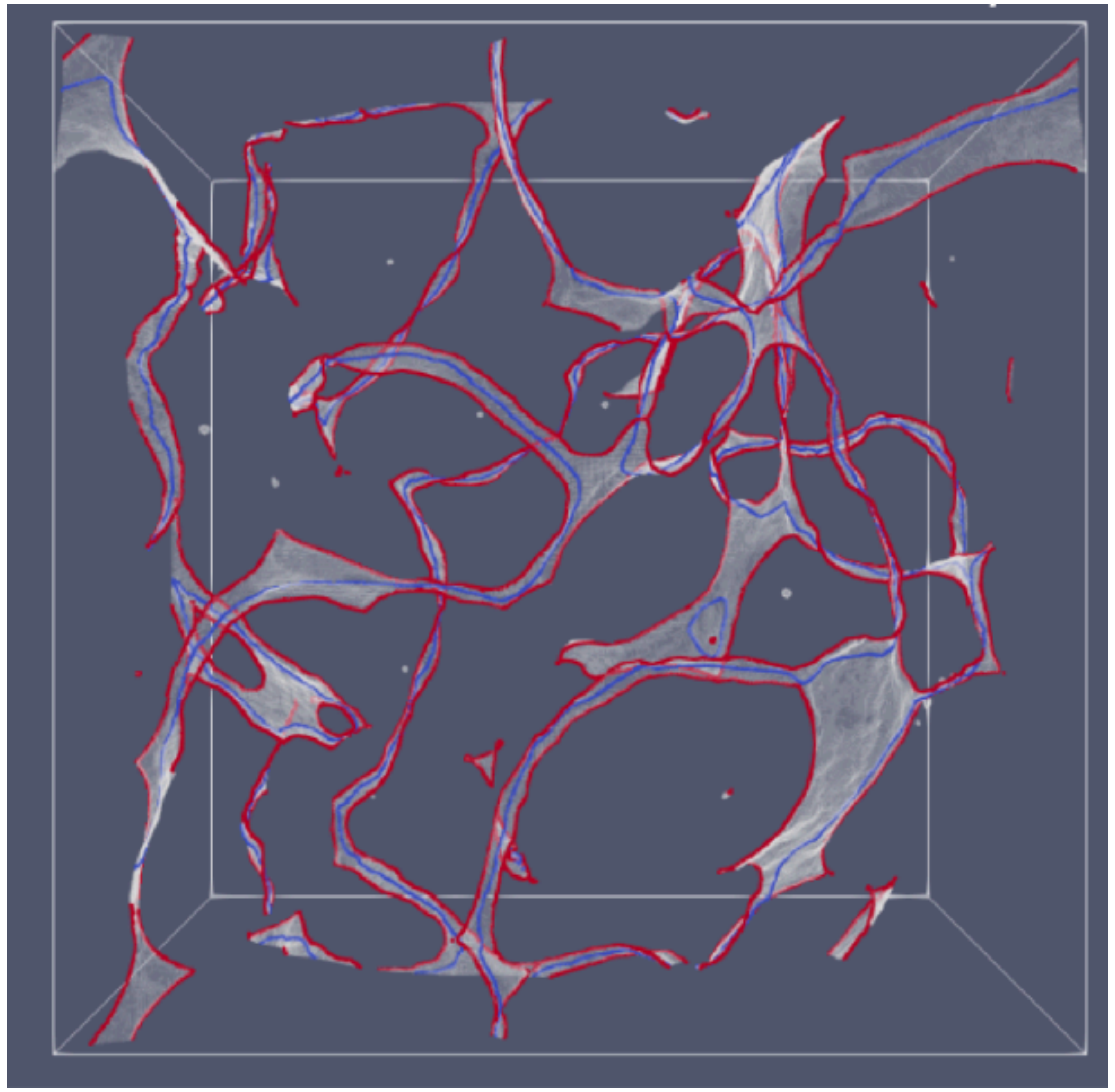

Numerical results (3D)

Lee, Murai, FT and Yin 2409.09749

Case of $(n_1, n_2) = (2, 3)$

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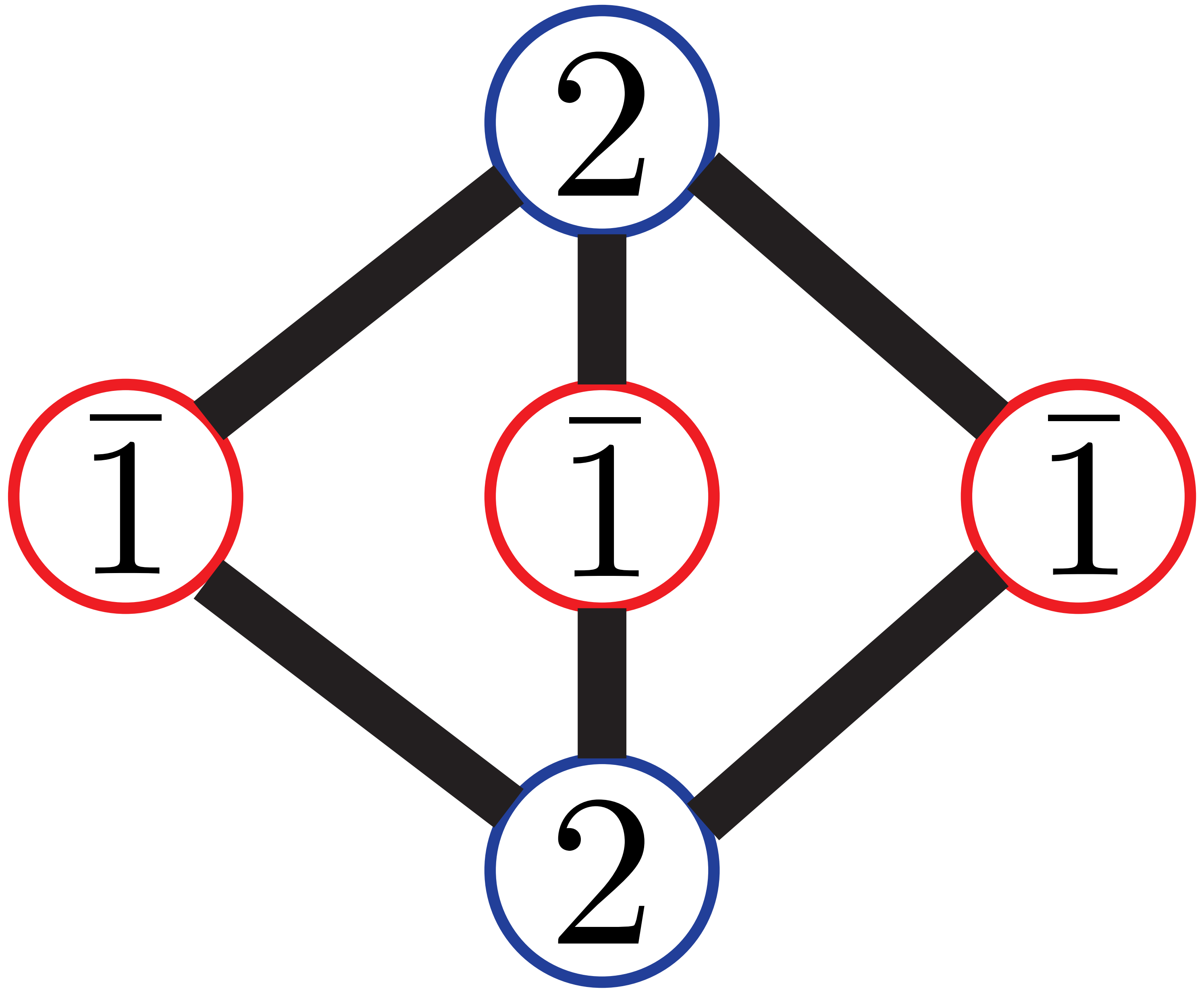

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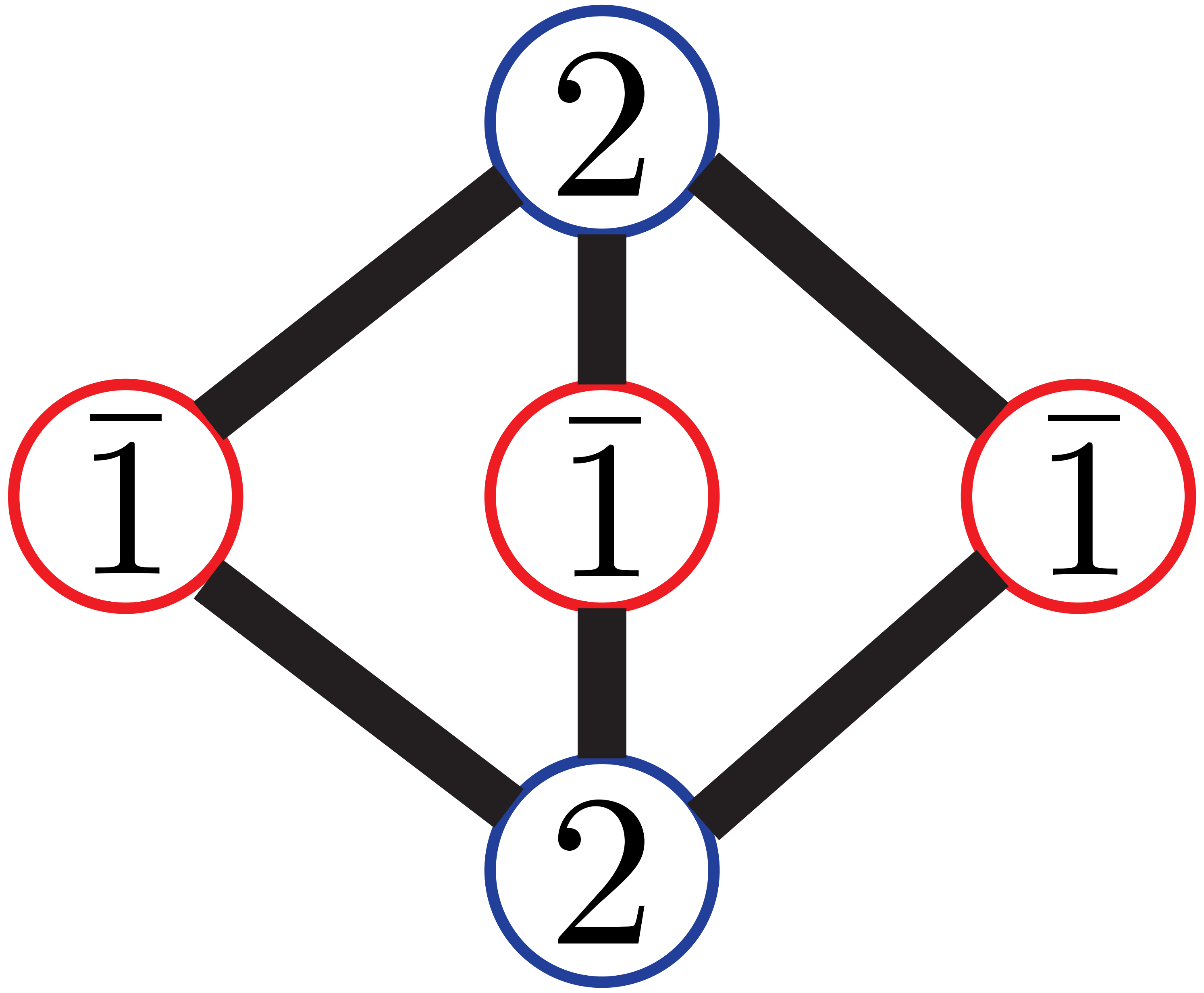

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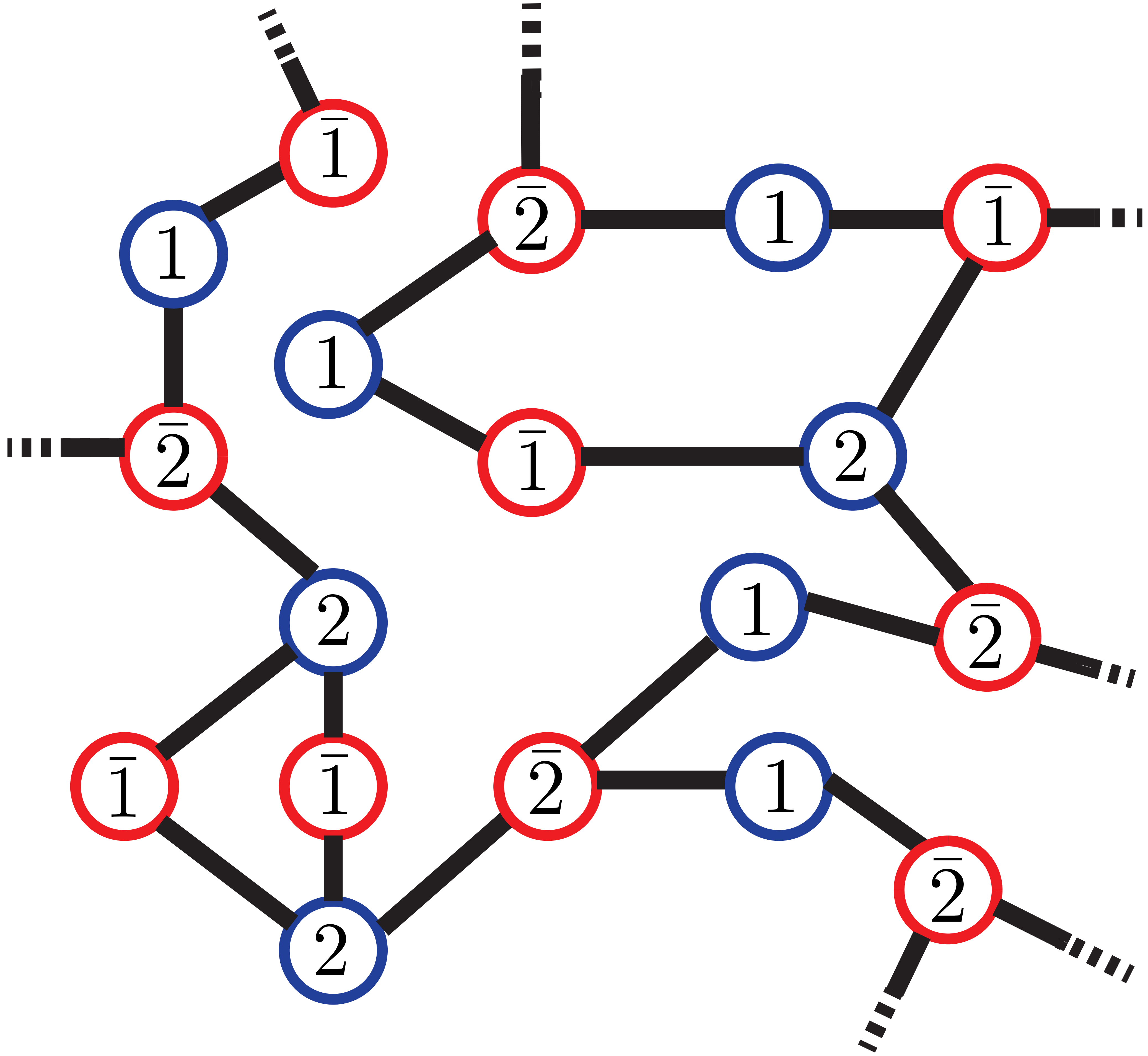
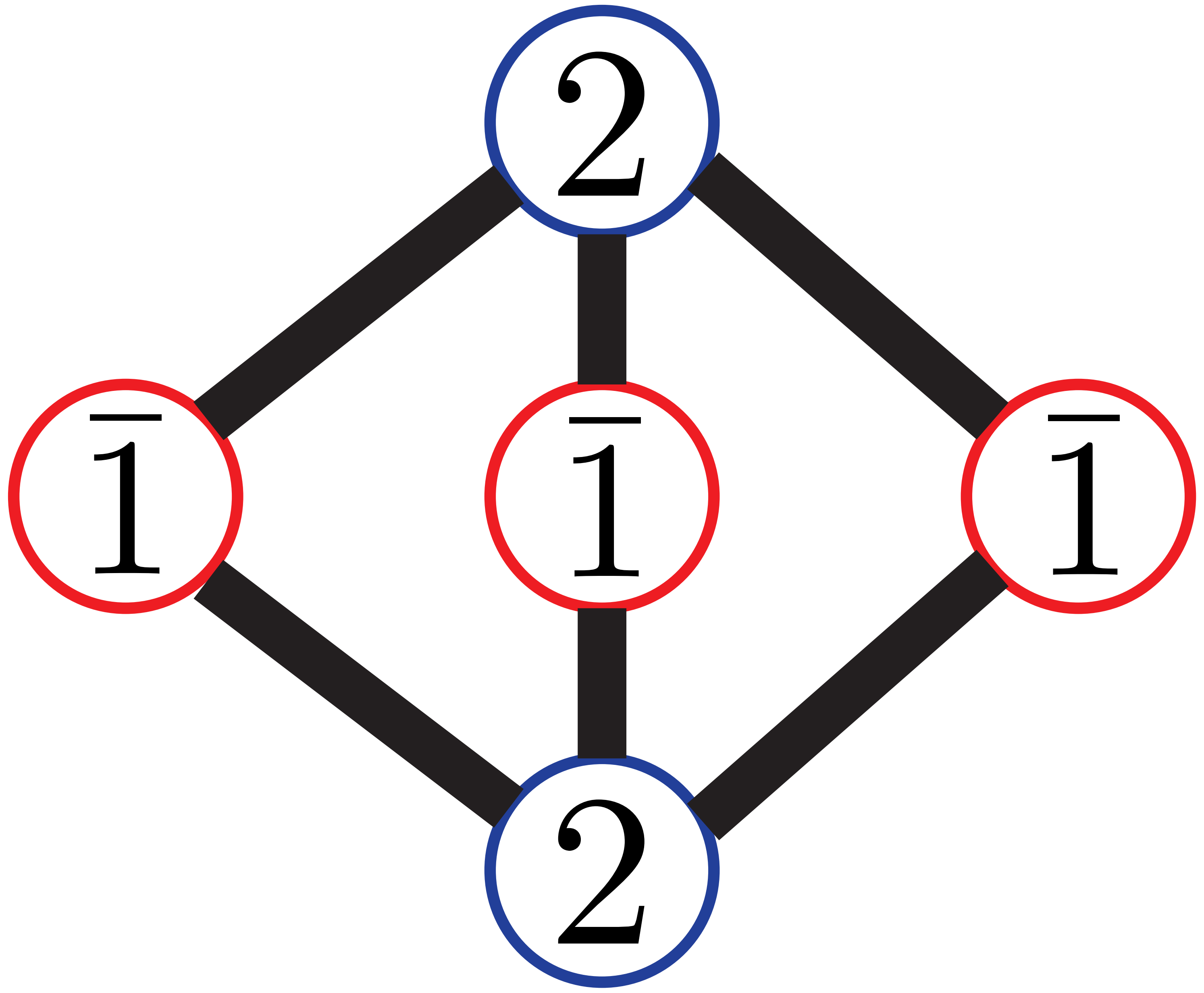


String-wall network

Case of $(n_1, n_2) = (2, 3)$

String-wall network forms in stead of string bundles.



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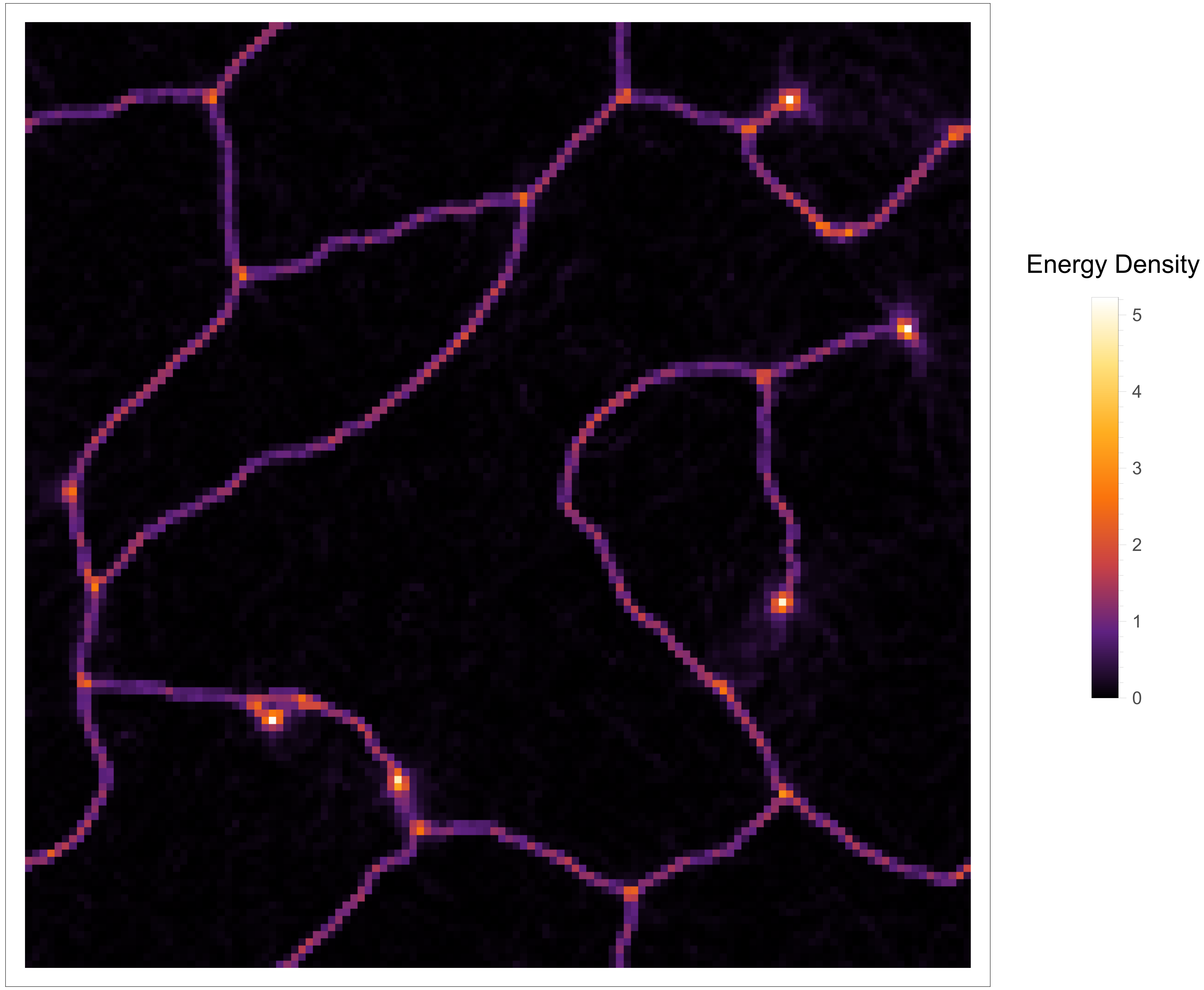

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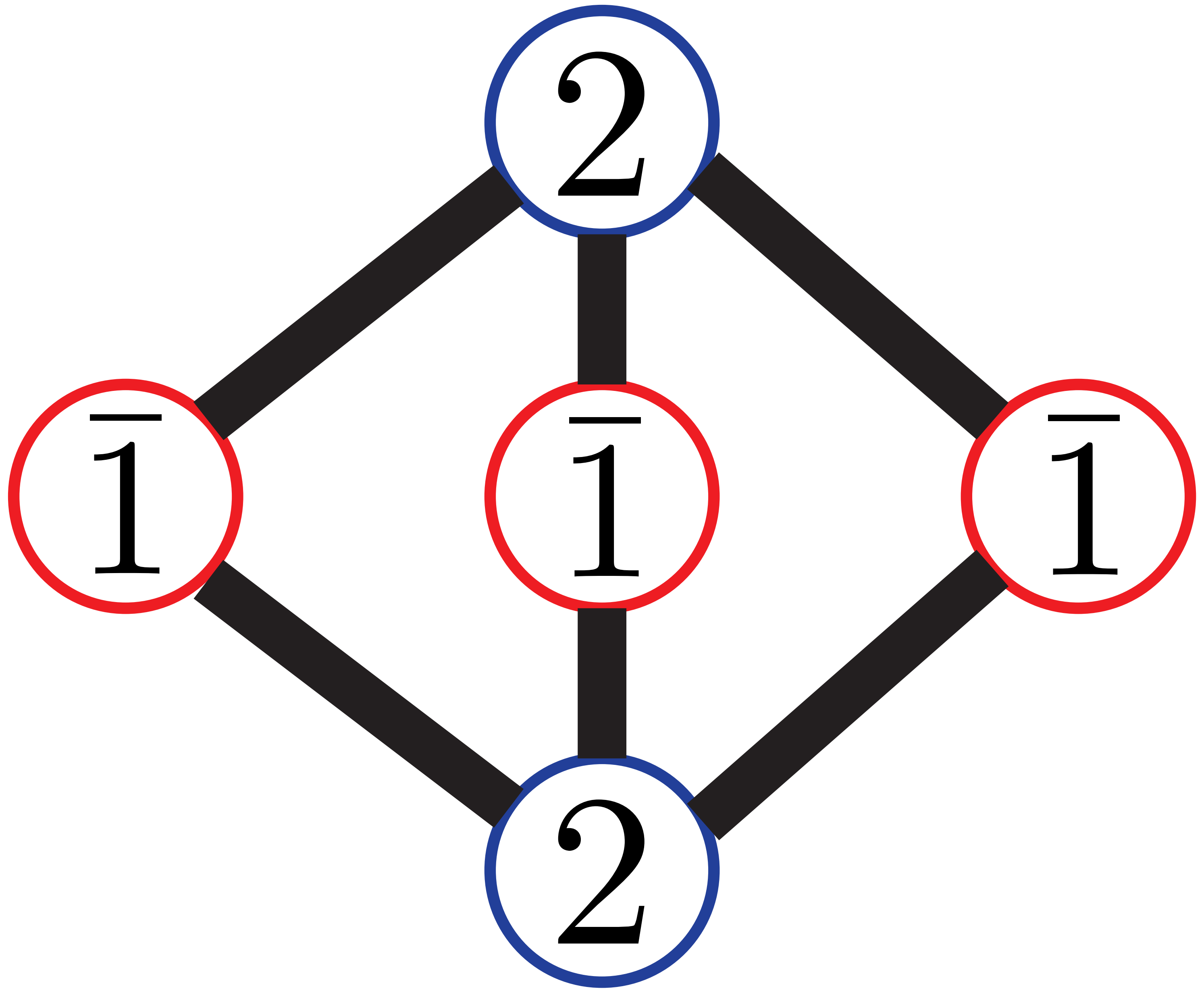

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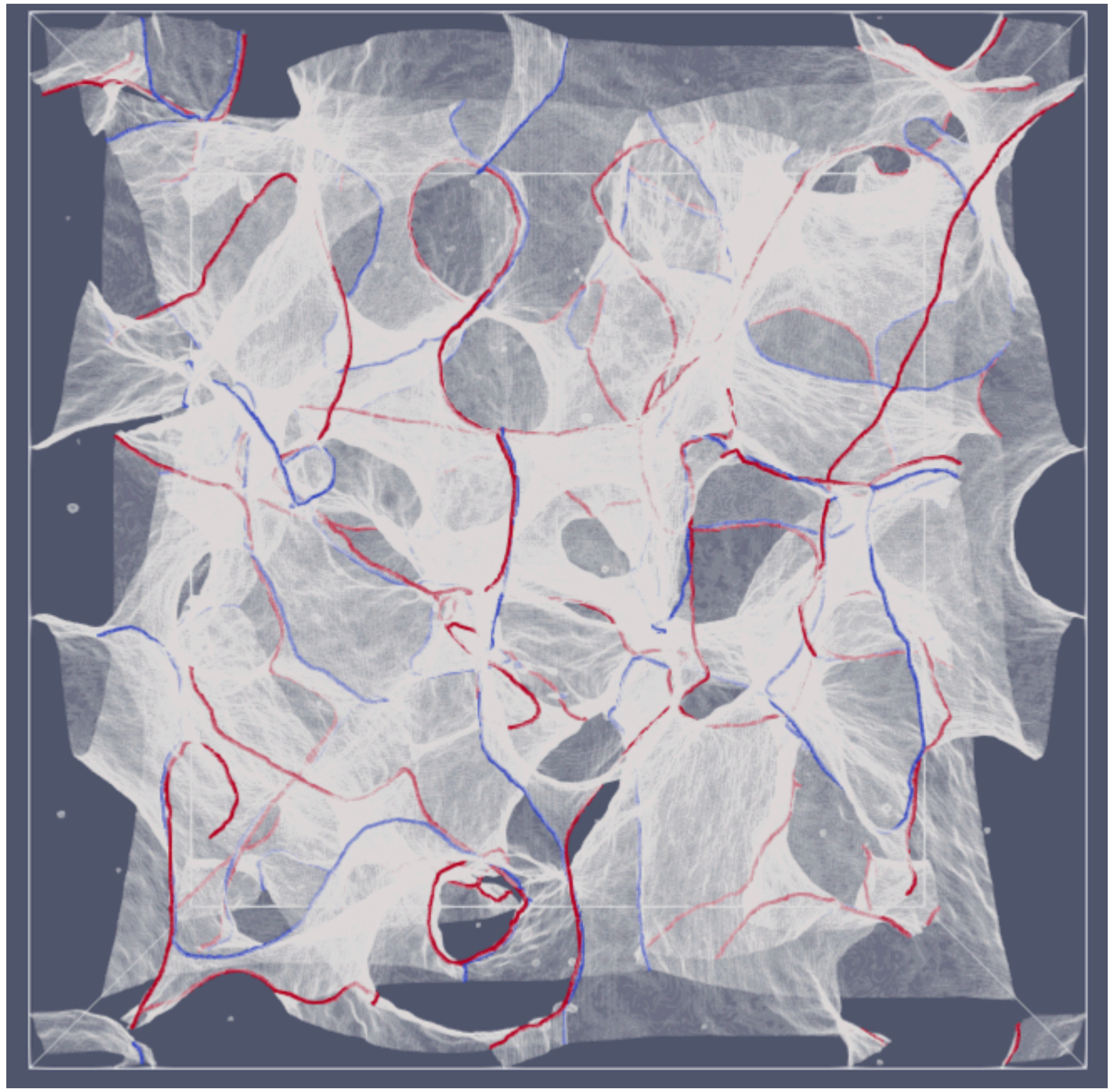
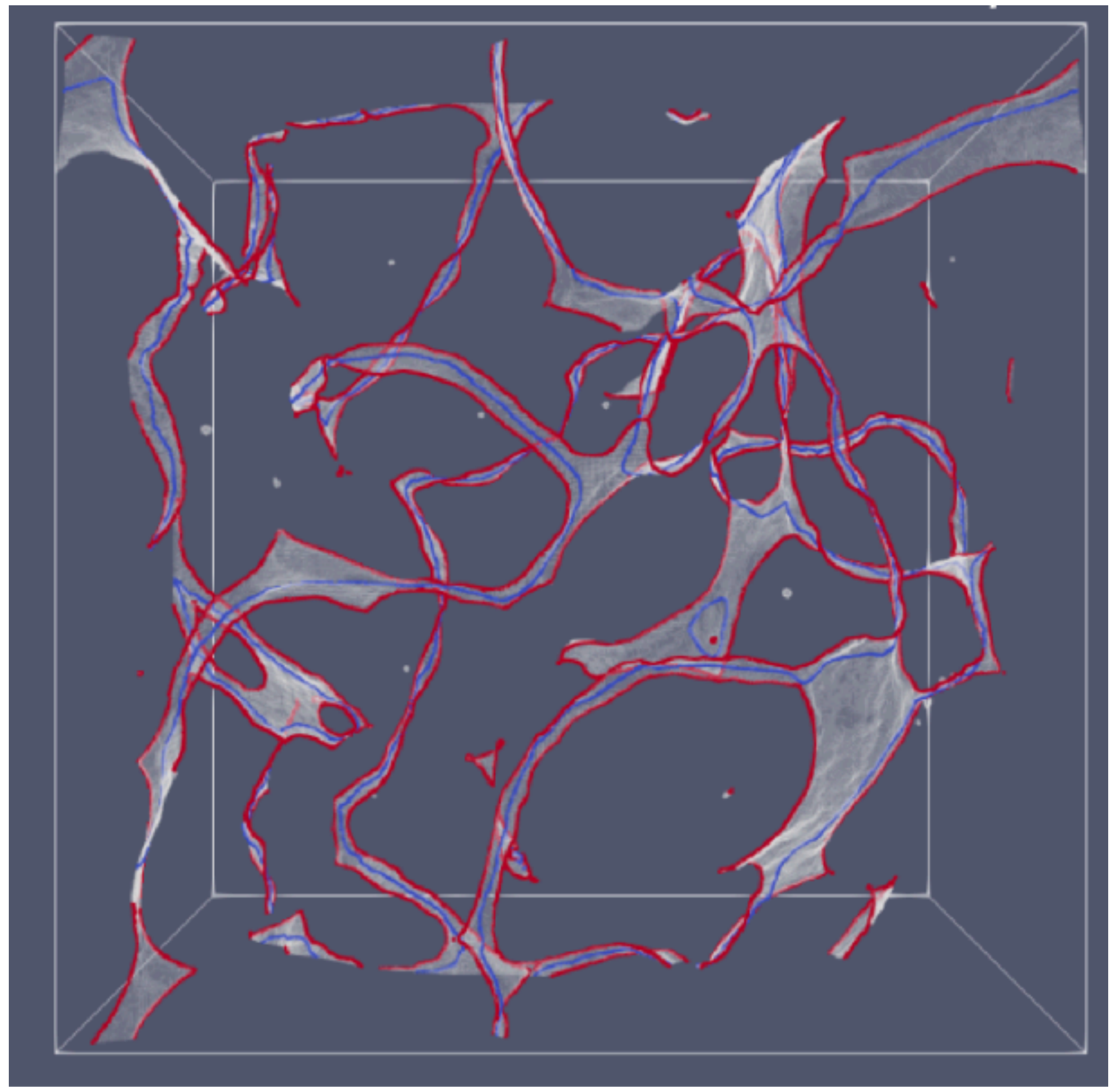


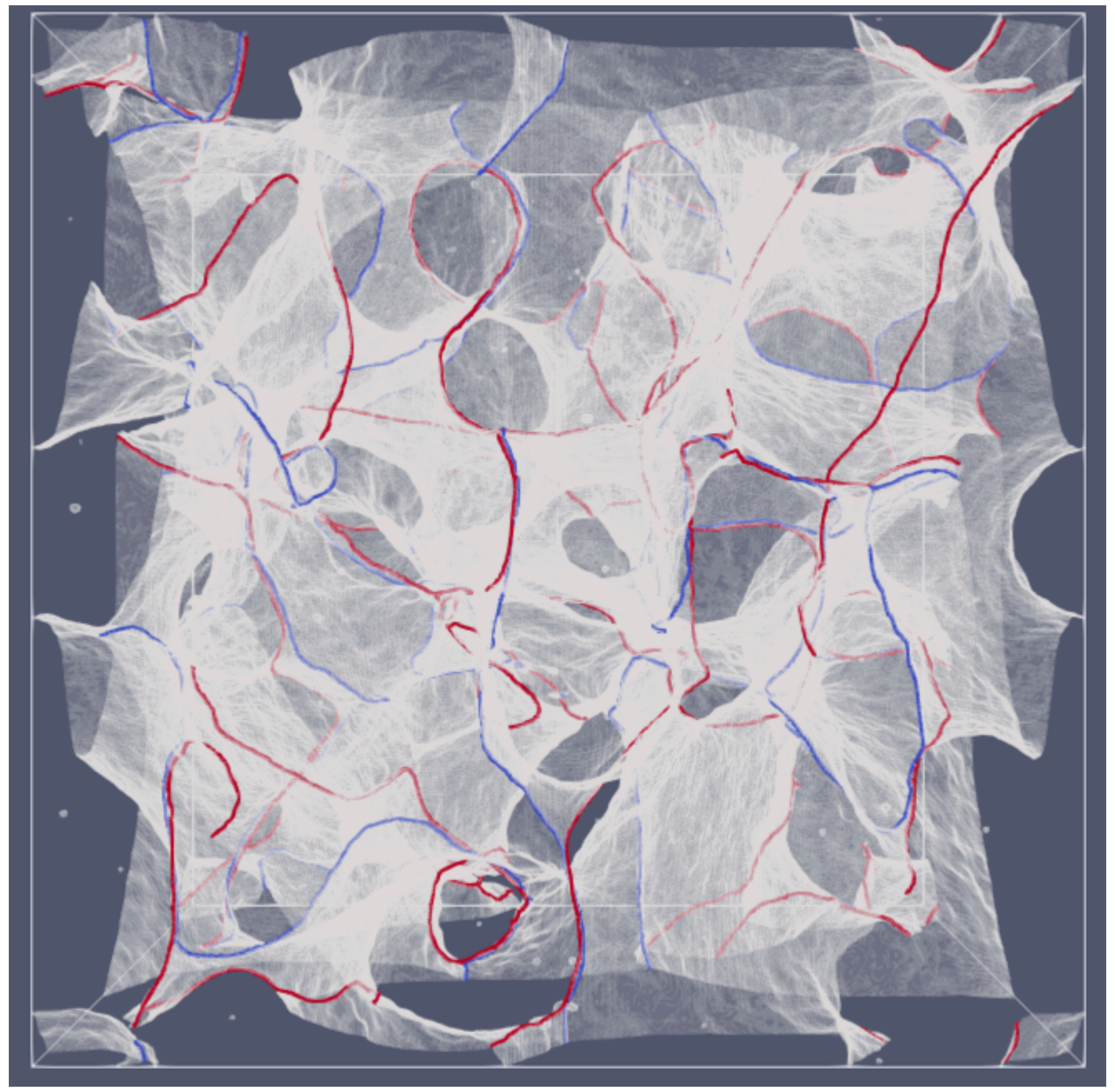
Numerical results (3D)

Coutesy of Junseok Lee

In the post-inflationary scenario, a string-wall network forms instead of ordinary cosmic strings if $n_1, n_2 \geq 2$.



$n_1 = 1$ or $n_2 = 1$



$n_1, n_2 \geq 2.$

These heavy axion DWs also produce a large amount of GWs.

Higaki, Jeong, Kitajima, Sekiguchi and FT, 1606.05552

Mixed initial conditions

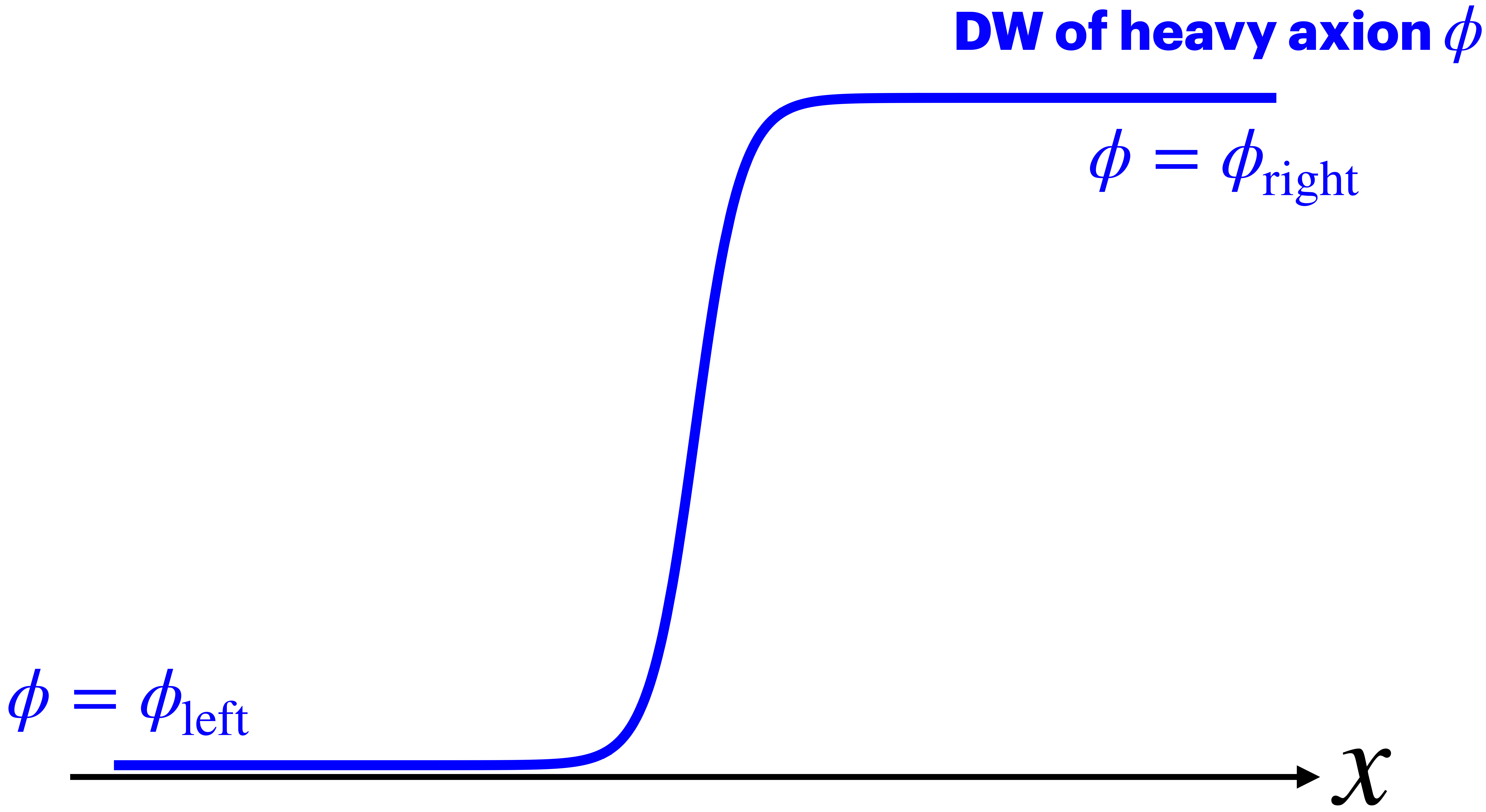
We may impose a mixed ”pre-post” initial condition, i.e.,

pre-inflationary initial condition for ϕ_1

post-inflationary initial condition for ϕ_2

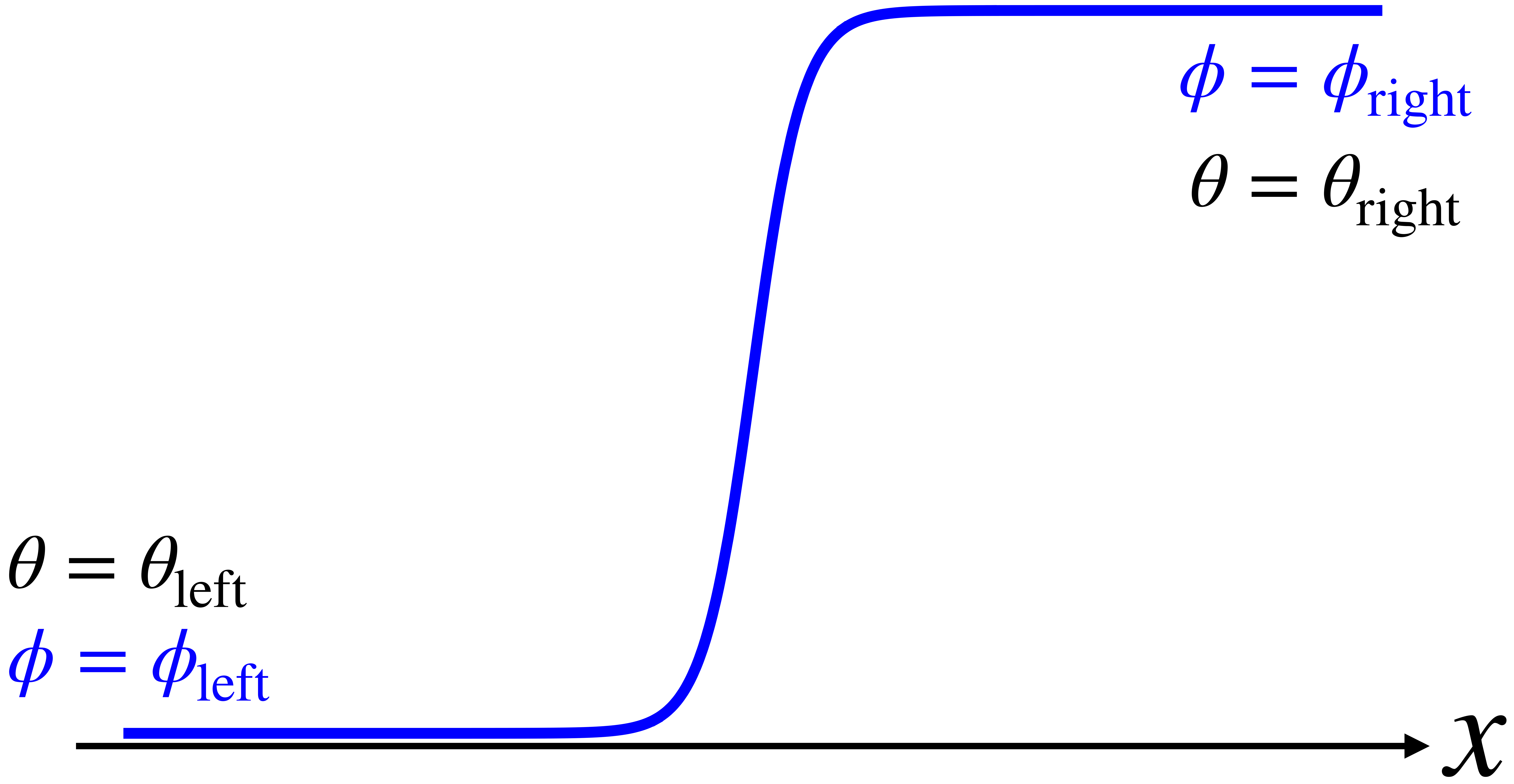
Then, no strings bundles are formed, and string-wall network of ϕ_2 remains if $n_2 \geq 2$ (even if $n_1 = 1$).

Mixed initial conditions makes string-wall network formation more likely.


DW of heavy axion ϕ
$\phi = \phi_{\rm right}$
$\phi = \phi_{\rm left}$
x

Couple ϕ to gluons: $\phi G\tilde{G}$

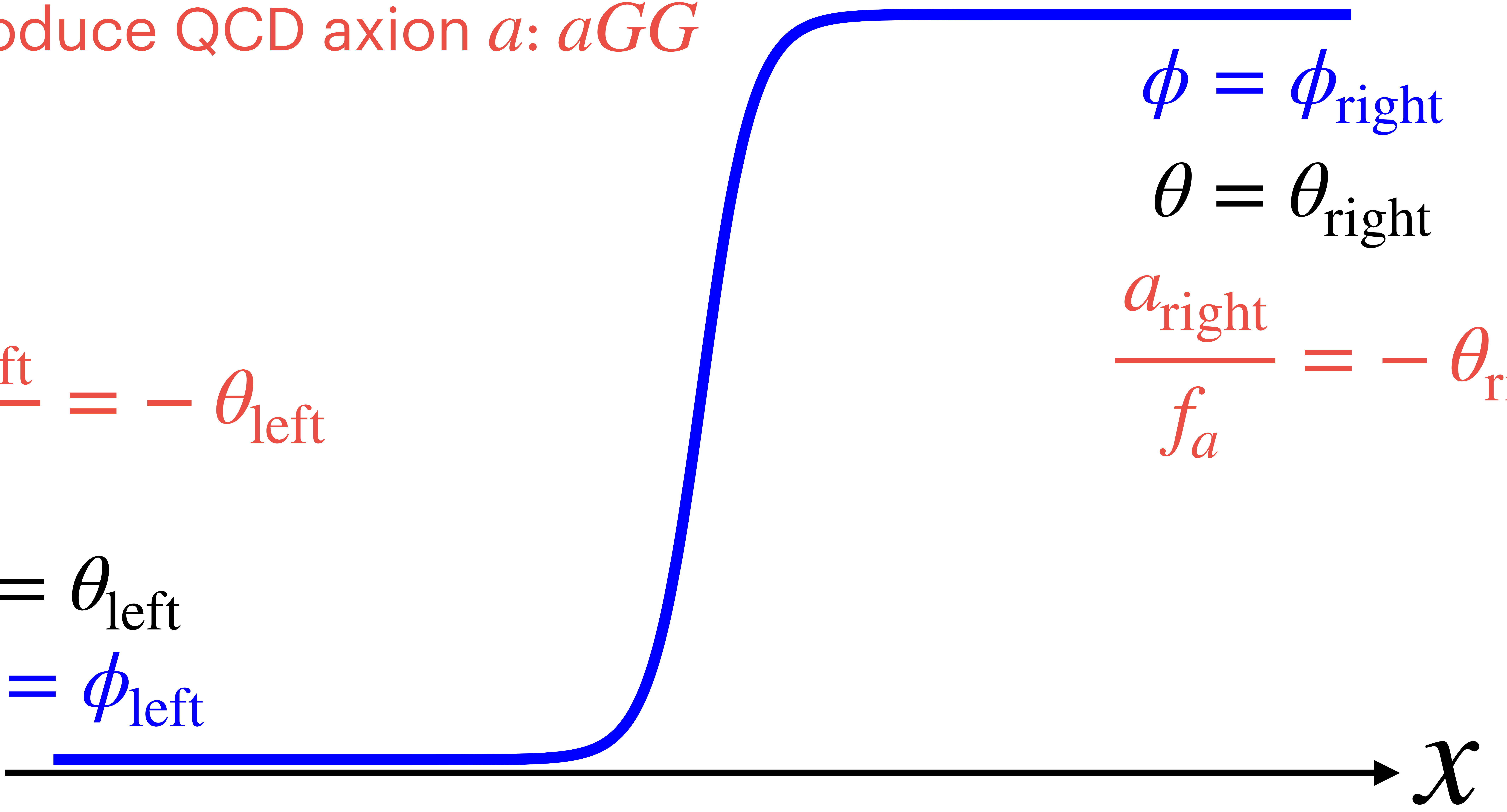
DW of heavy axion ϕ



Couple ϕ to gluons: $\phi G\tilde{G}$

Introduce QCD axion a: $aG\tilde{G}$

DW of heavy axion ϕ



$\phi = \phi_{\text{right}}$

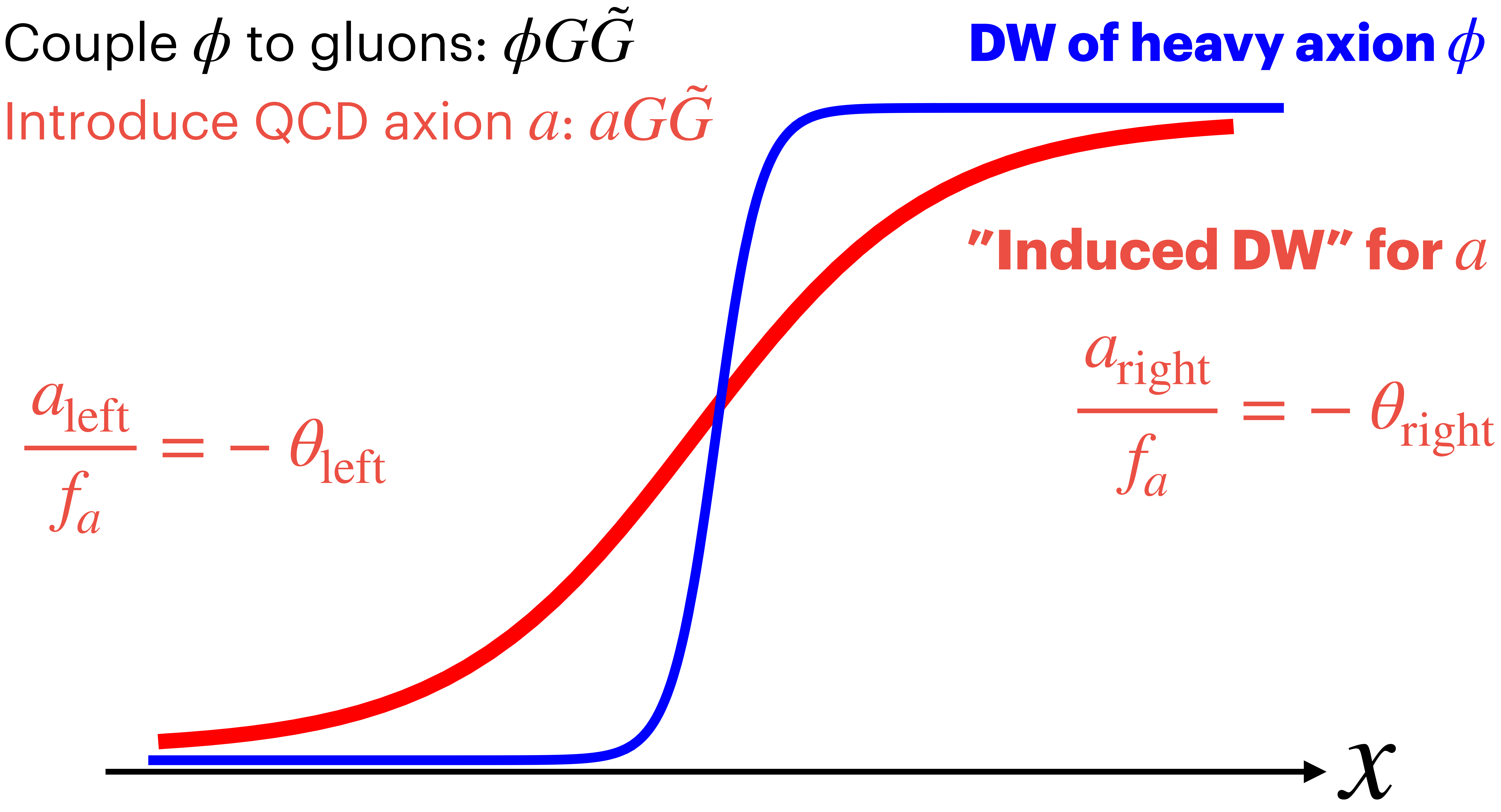
$\theta = \theta_{\text{right}}$

$$\frac{a_{\text{right}}}{f_a} = -\theta_{\text{right}}$$

$$\frac{a_{\text{left}}}{f_a} = -\theta_{\text{left}}$$

$\theta = \theta_{\text{left}}$

$\phi = \phi_{\text{left}}$



Couple ϕ to gluons: $\phi G\tilde{G}$
Introduce QCD axion a: $aG\tilde{G}$
DW of heavy axion ϕ
"Induced DW" for a
$\frac{a_{\rm left}}{f_a} = -\theta_{\rm left}$
$\frac{a_{\rm right}}{f_a} = -\theta_{\rm right}$
x

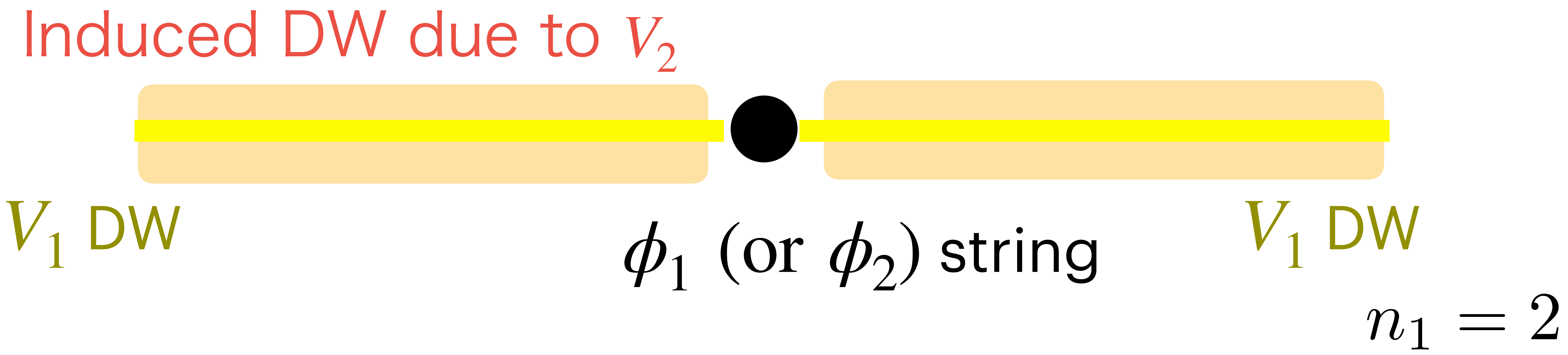
"Induced DW" formation due to V_2

Now we consider the DW formation due to $V_2 \ll V_1$.

$$V_1(\phi_1, \phi_2) = \Lambda^4 \left[1 - \cos\left(n_1 \frac{\phi_1}{f_1} + n_2 \frac{\phi_2}{f_2}\right)\right]$$

$$V_2(\phi_1, \phi_2) = \Lambda'^4 \left[1 - \cos\left(n_1' \frac{\phi_1}{f_1} + n_2' \frac{\phi_2}{f_2}\right)\right] \qquad \Lambda' \ll \Lambda$$

Both V_1 and V_2 can be minimized in any domains, and so there is no potential bias at the minimum.


Induced DW due to V_2
V_1 DW
ϕ_1 (or ϕ_2) string
V_1 DW
$n_1 = 2$

The string-wall network persists even after V_2 is included, since induced DWs appear and the domains remain degenerate in energy.

Thus, even a minimal extension with two PQ scalars leads to the cosmological DW problem if n_1 , n_2 ≥ 2.

The simplest solution to the DW problem is to introduce a potential bias along the heavy axion, leaving the PQ mechanism intact. (I'll present an alternative approach shortly.)

Induced domain walls of QCD axion

We consider induced domain walls of the QCD axion a, arising from its mixing with a heavy axion ϕ.

Pre-inflationary condition for the QCD axion $A = \frac{f_a}{\sqrt{2}} e^{i \frac{a}{f_a}}$,

Post-inflationary condition for the heavy axion $\Phi = \frac{f_\phi}{\sqrt{2}} e^{i \frac{\phi}{f_\phi}}$

The axion potential:

$$V(a,\phi) = \chi(T)\left[1-\cos\left(\frac{a}{f_a}+\frac{\phi}{f_\phi}\right)\right] + \Lambda^4\left[1-\cos\left(2\frac{\phi}{f_\phi}\right)\right]$$

$$\chi(T) \equiv \frac{m_a^2(T) f_a^2}{N_a^2} = \begin{cases} \chi_0 & (T < T_{\rm QCD}) \\ \chi_0 \left(\frac{T}{T_{\rm QCD}}\right)^{-n} & (T \geq T_{\rm QCD}) \end{cases} ,$$

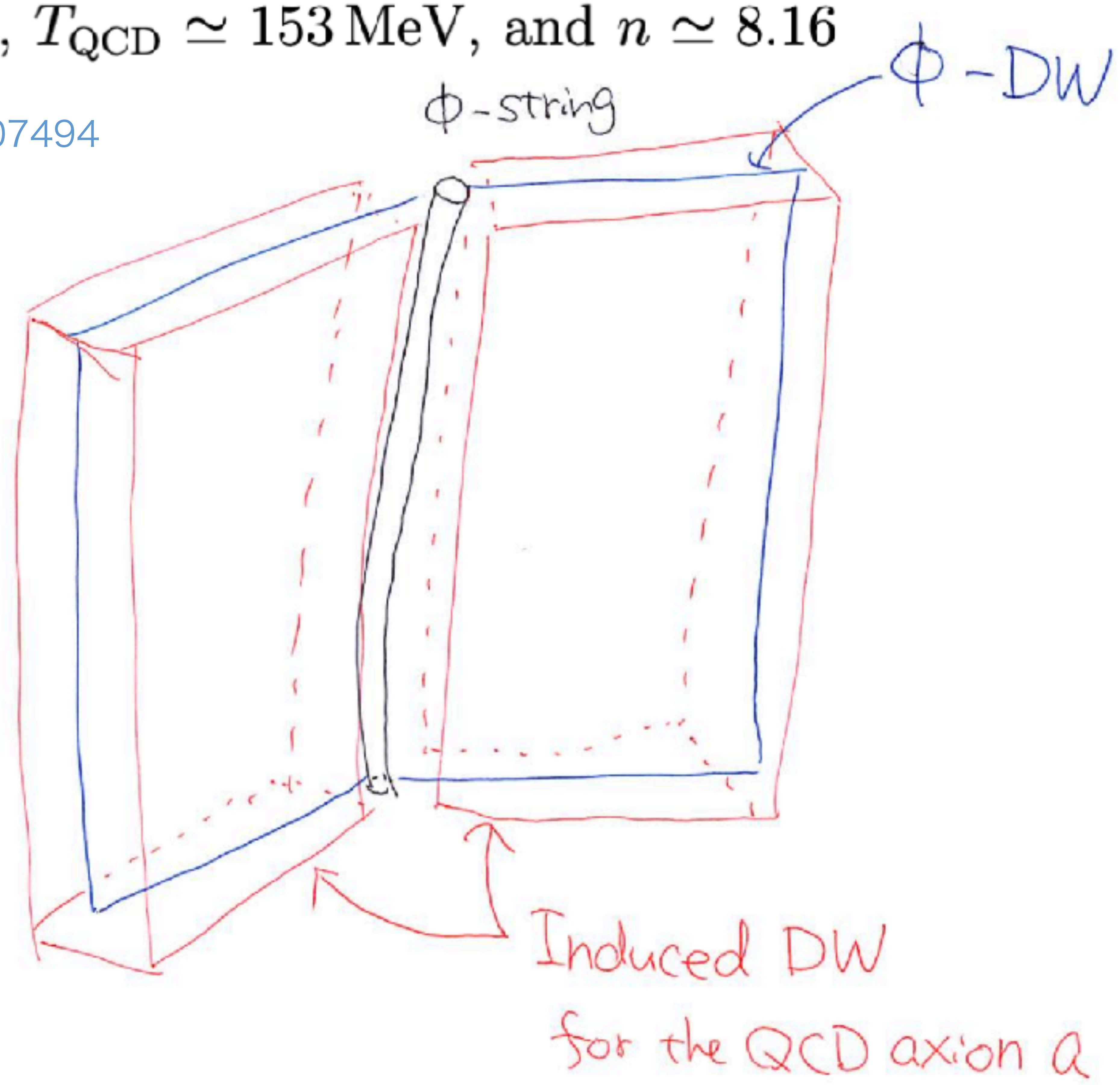
$\chi_0 \simeq (75.6\,{\rm MeV})^4$, $T_{\rm QCD} \simeq 153\,{\rm MeV}$, and $n \simeq 8.16$

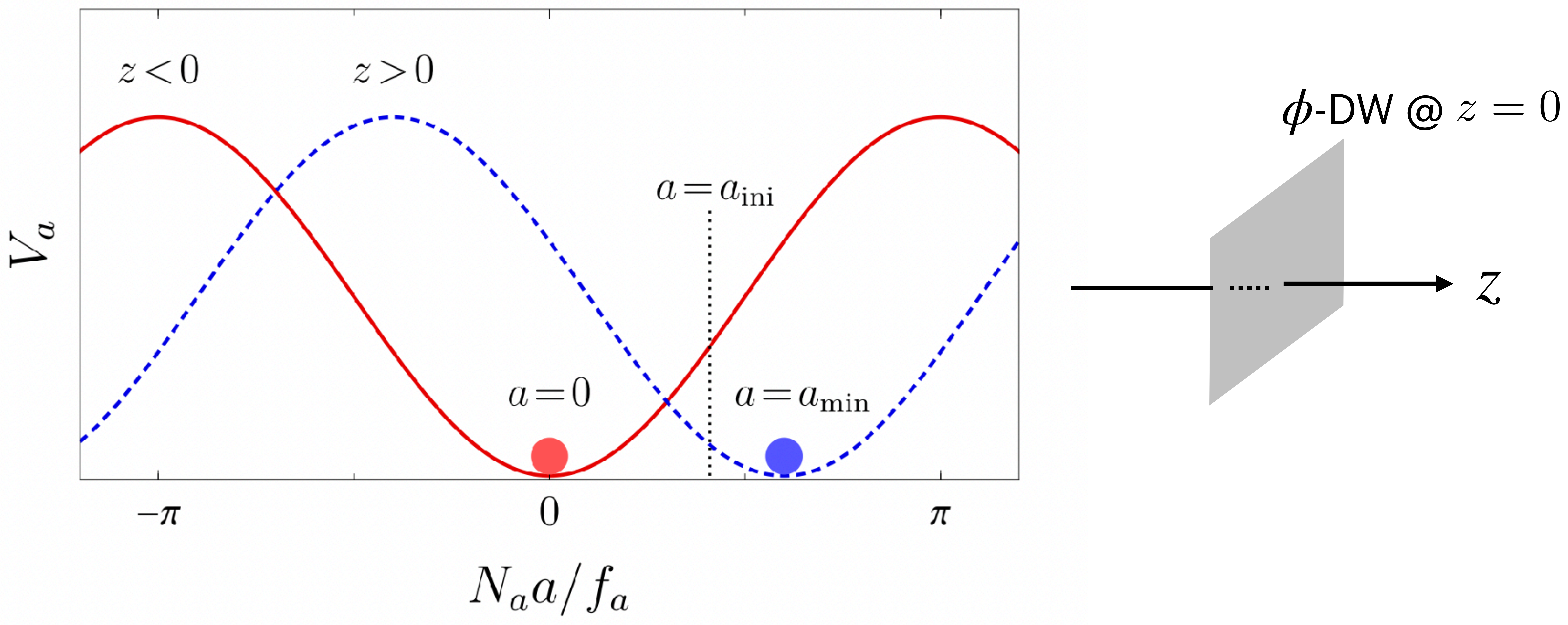
A. Borsanyi et al, 1606.07494

For simplicity we assume a hierachy in the mass and the tension:

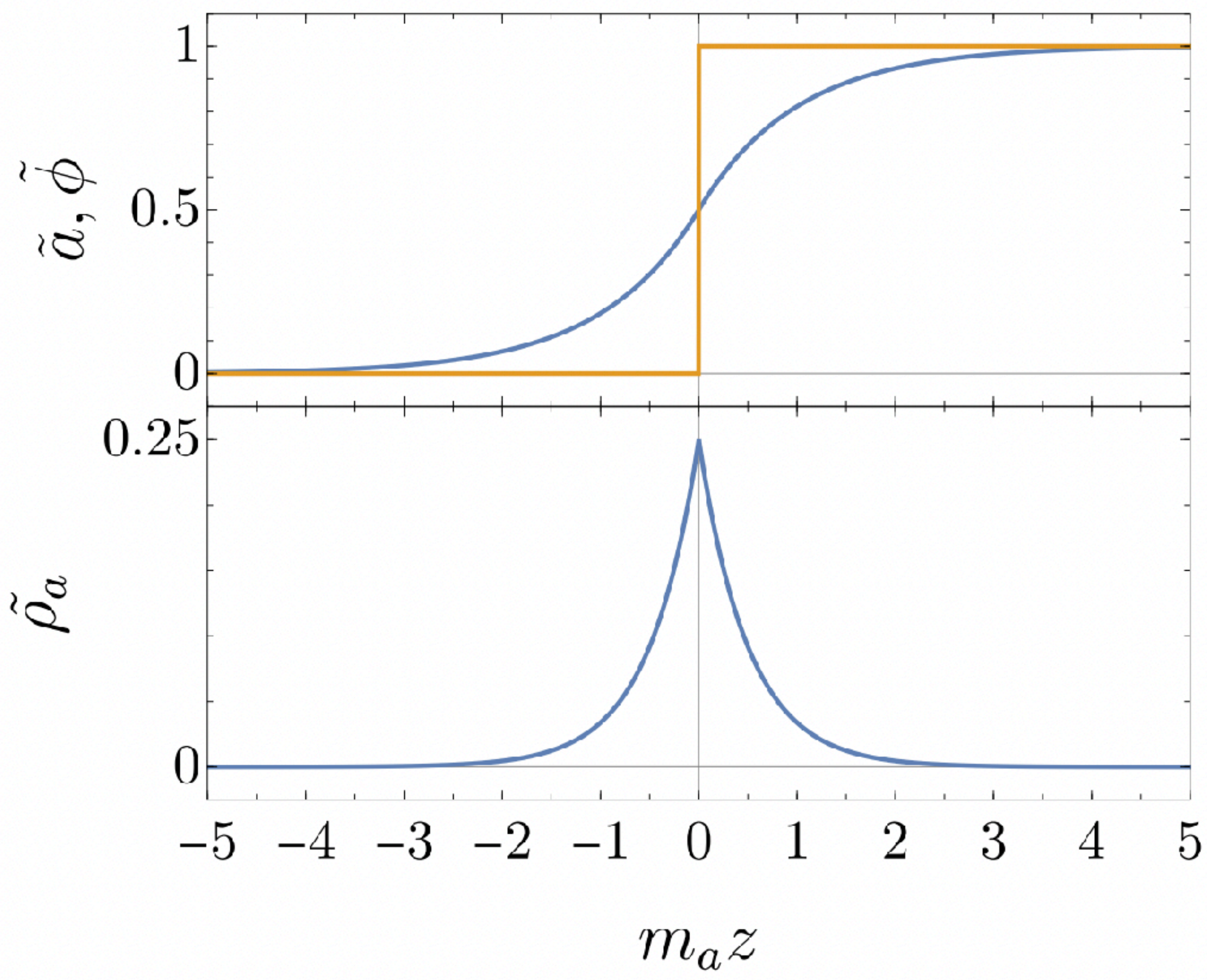
$$m_\phi \sim \frac{\Lambda^2}{f_\phi} \gg m_{a0} = m_a(T=0)$$

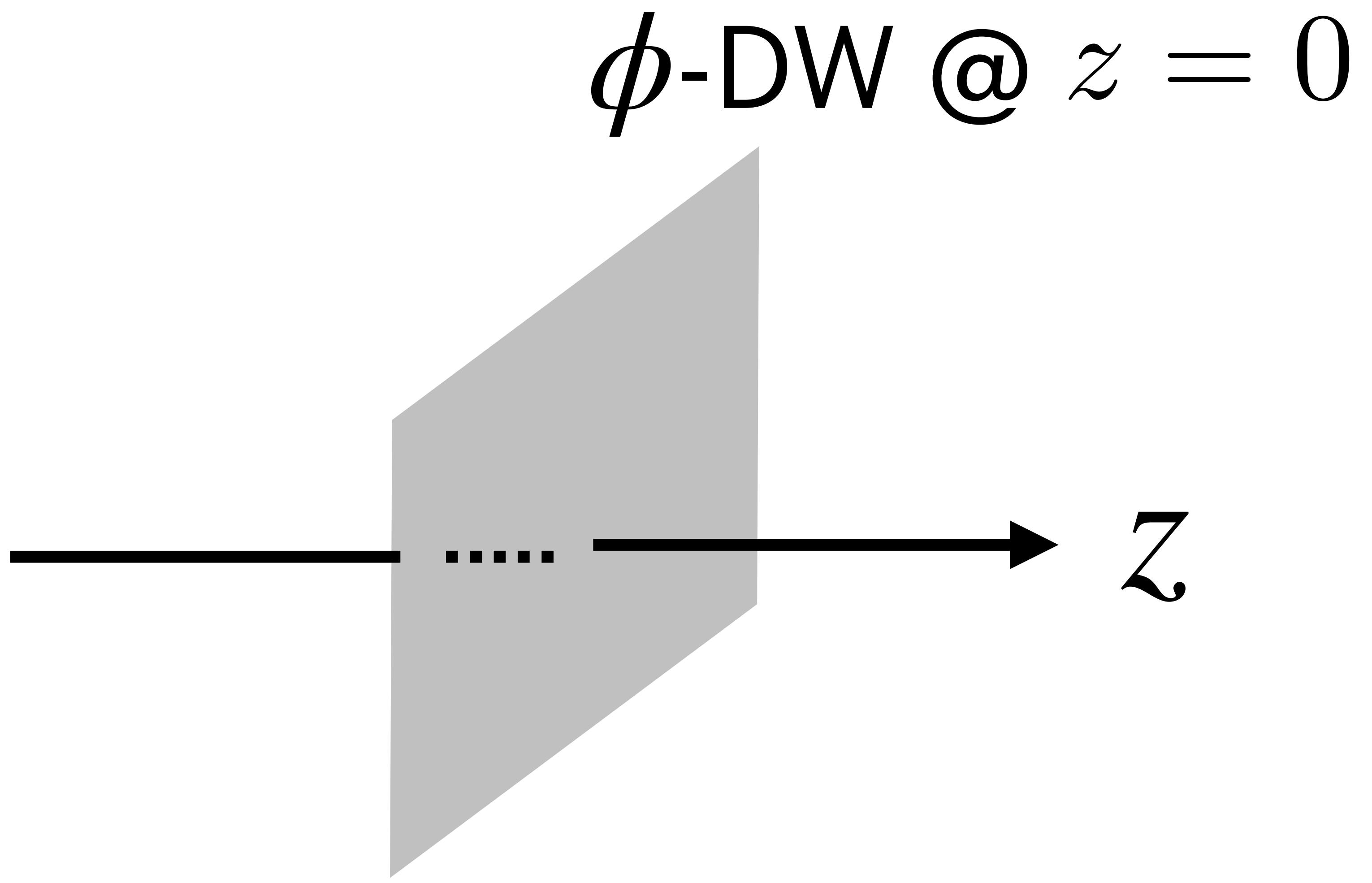
$$m_\phi f_\phi^2 \gg m_{a0} f_a^2$$





$z<0$
$z>0$
$a=a_{\rm ini}$
$a=0$
$a=a_{\rm min}$
V_a
$-\pi$
0
π
$N_a a/f_a$
ϕ-DW @ $z=0$
z



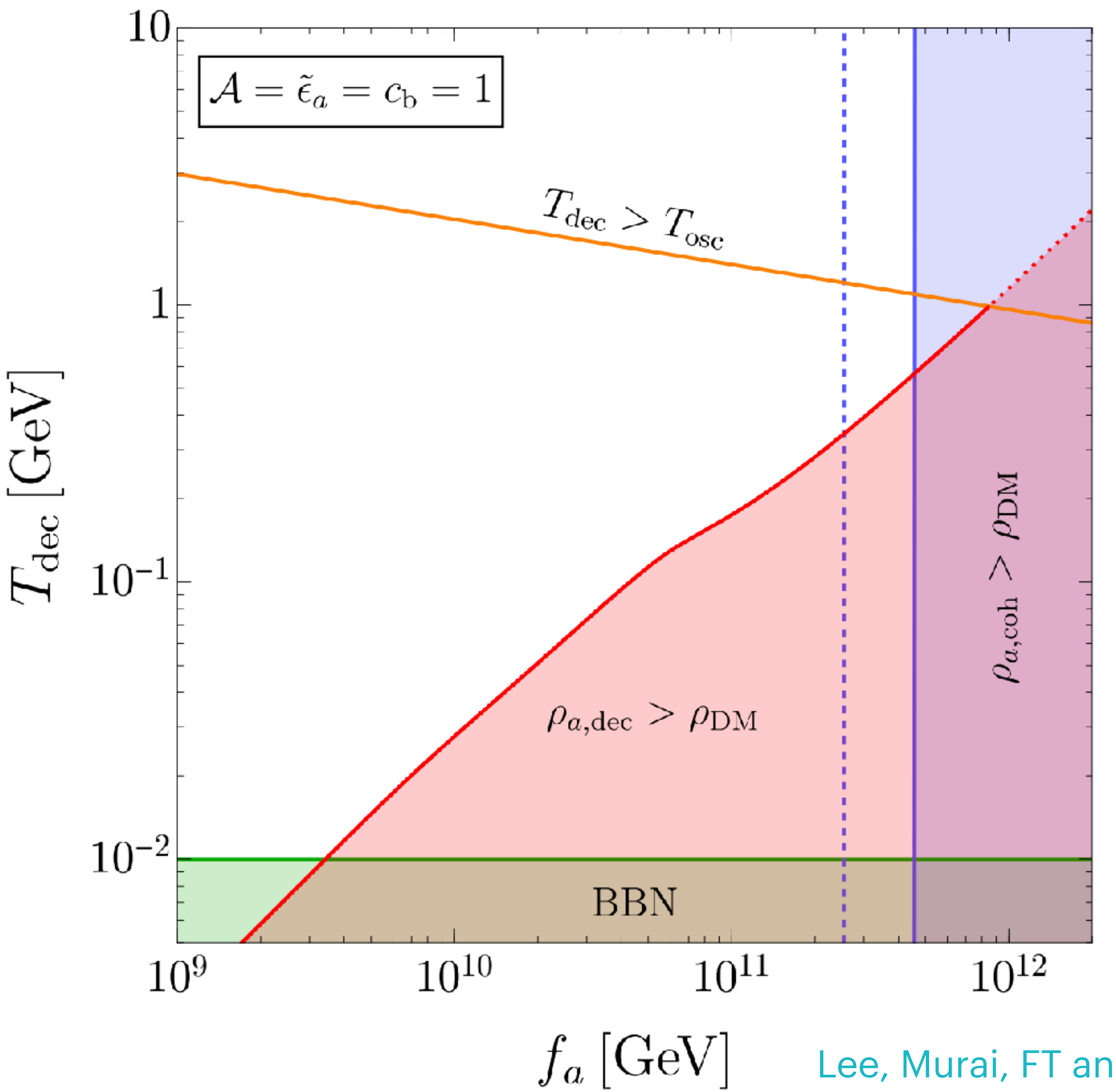


$\sigma_a \sim m_a a_{\min}^2 < 8 m_a f_a^2$

The tension of induced DWs is always smaller than that of ordinary DWs.

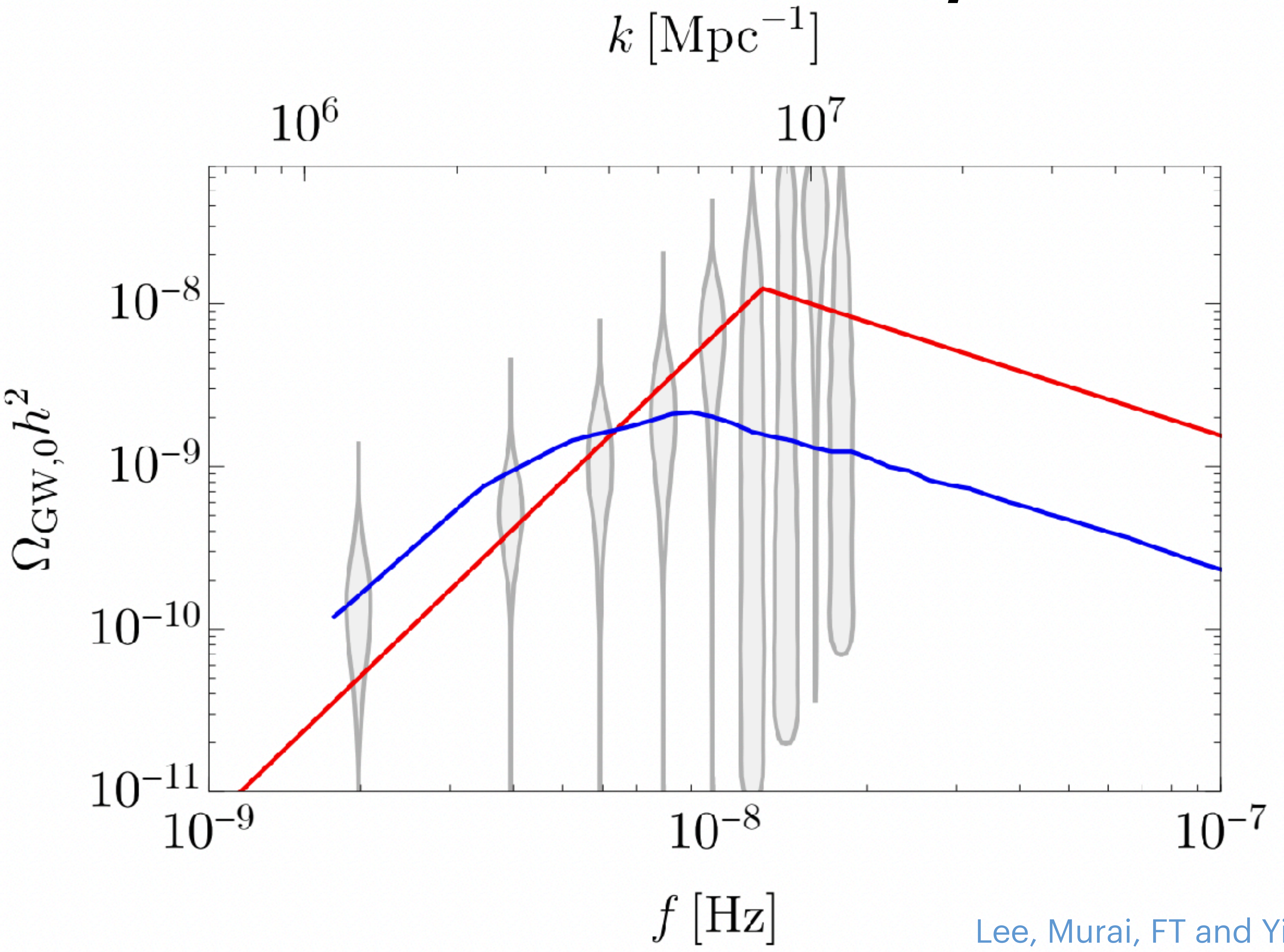
Decay temperature of the ϕ-DWs

(We introduce a small potential bias along ϕ-direction to make DWs unstable)



Lee, Murai, FT and Yin 2407.09478

GW spectrum from the ϕ-DW collapse



Red:
Analytical estimate with
$\sigma_\phi = 6 \times 10^{15}\,\mathrm{GeV}^3$
$T_{\mathrm{dec}} = 100\,\mathrm{MeV}$

Blue:
Numerical estimate with
$\sigma_\phi = 1.3 \times 10^{15}\,\mathrm{GeV}^3$
$T_{\mathrm{dec}} = 150\,\mathrm{MeV}$

Kitajima, Lee, Murai, FT and Yin 2306.17146

Lee, Murai, FT and Yin 2407.09478

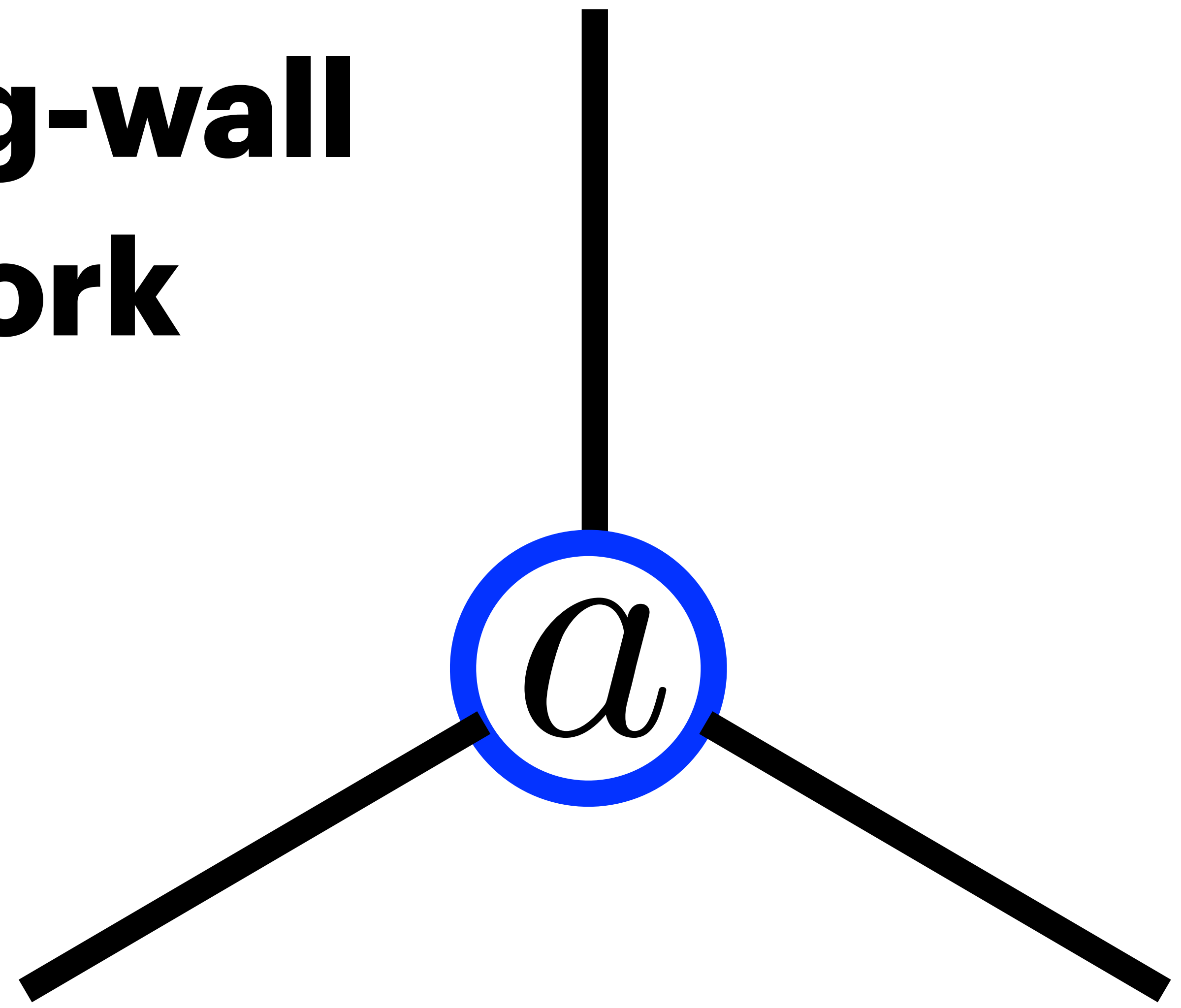
New solution to the QCD axion DW problem

Lee, Murai, FT and Yin 2507.xxxxx, FT and Yin 2012.11576

$$\mathcal{L} = -N_{\rm DW}\frac{g_s^2}{32\pi^2}\frac{a}{v_a}G\tilde{G}$$

$N_{\rm DW} = 3 \text{ or } 6$ for DFSZ axion

String-wall network



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$$\mathcal{L} = -N_{\rm DW}\frac{g_s^2}{32\pi^2}\frac{a}{v_a}G\tilde{G} \quad \Rightarrow \quad \mathcal{L} = -\frac{g_s^2}{32\pi^2}\left(N_{\rm DW}\frac{a}{v_a}+\frac{\varphi}{v_\varphi}\right)G\tilde{G}$$

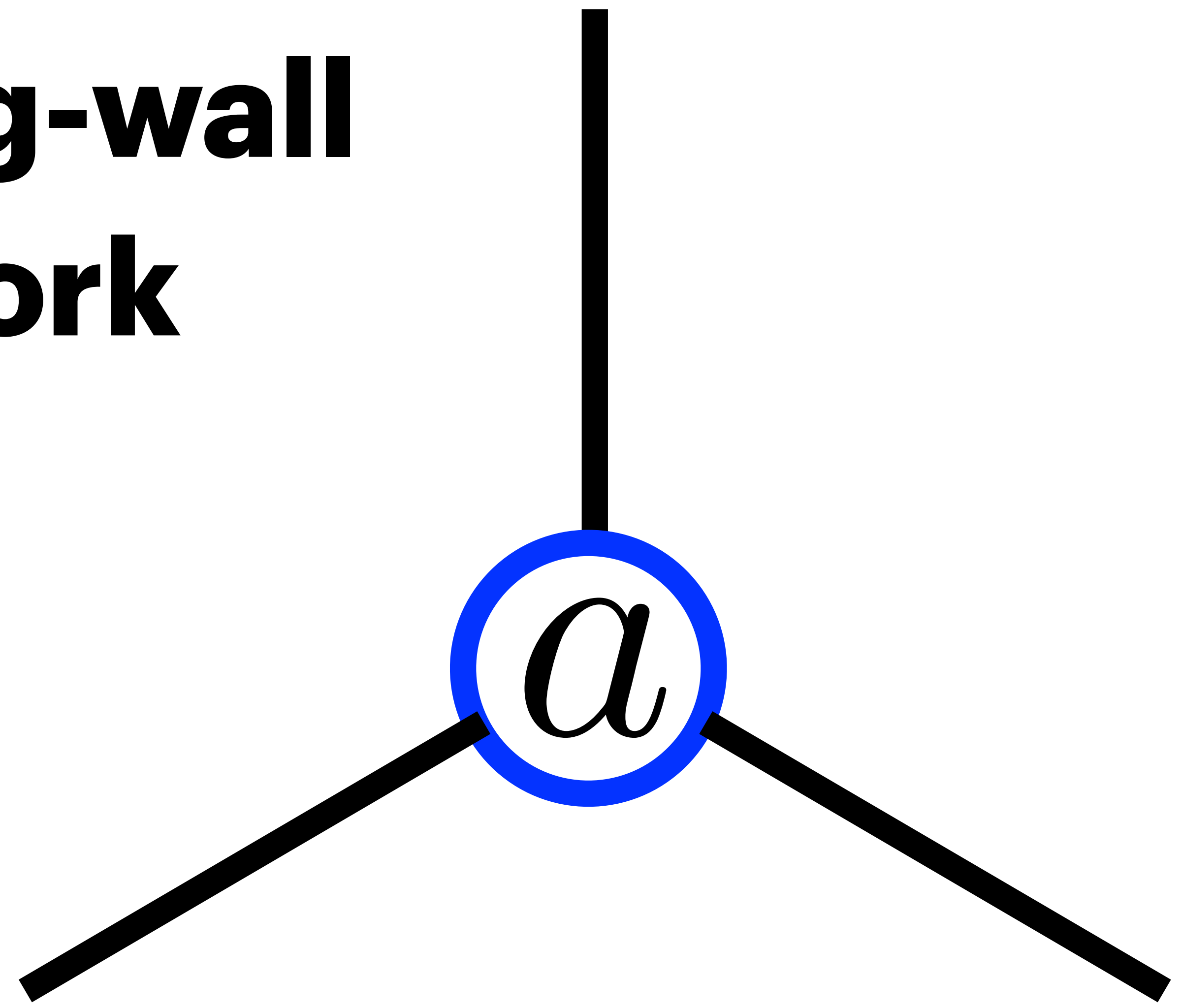
$N_{\rm DW} = 3$ or 6 for DFSZ axion

φ : Massless (or very light) axion

e.g. KSVZ axion

(different from a)

String-wall network



New solution to the QCD axion DW problem

Lee, Murai, FT and Yin 2507.xxxxx, FT and Yin 2012.11576

$$\mathcal{L} = -N_{\rm DW}\frac{g_s^2}{32\pi^2}\frac{a}{v_a}G\tilde{G} \Rightarrow \mathcal{L} = -\frac{g_s^2}{32\pi^2}\left(N_{\rm DW}\frac{a}{v_a}+\frac{\varphi}{v_\varphi}\right)G\tilde{G}$$

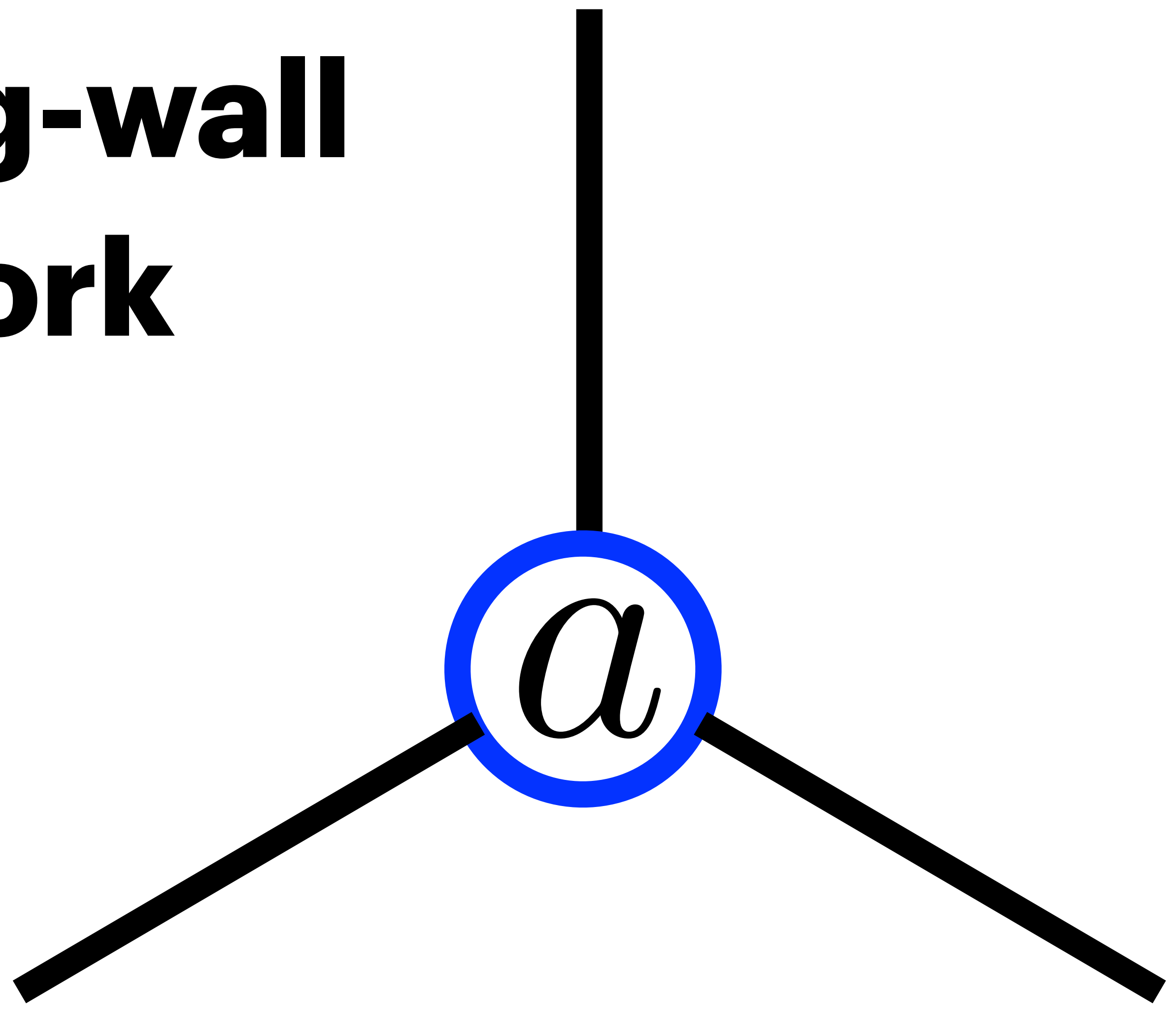
$N_{\rm DW} = 3 \text{ or } 6$ for DFSZ axion

φ : Massless (or very light) axion

e.g. KSVZ axion

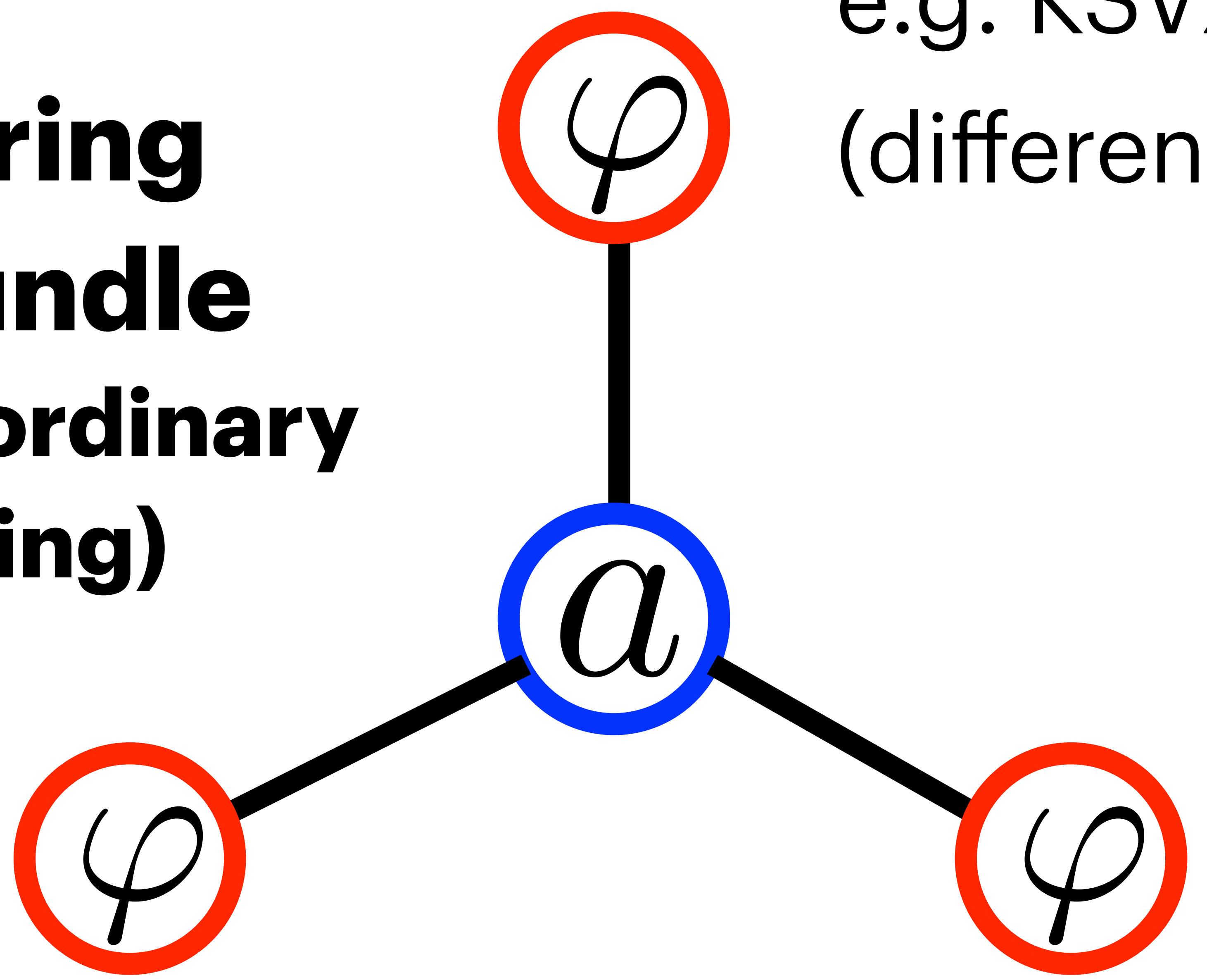
(different from a)

String-wall network



String bundle

(= ordinary string)



New solution to the QCD axion DW problem

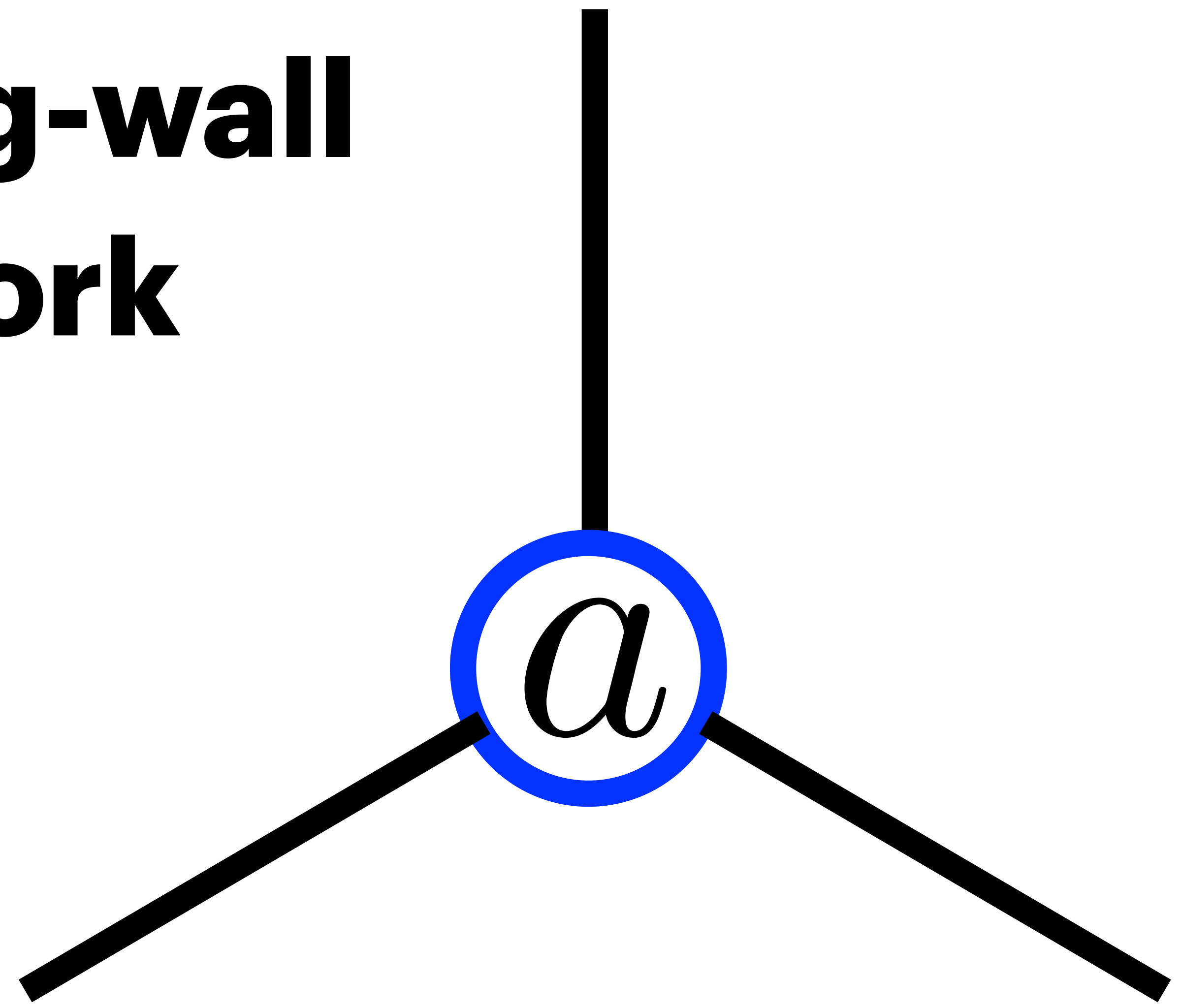
Lee, Murai, FT and Yin 2507.xxxxx, FT and Yin 2012.11576

Adding another light axion can solve the DW problem by converting the string-wall network into string bundes.

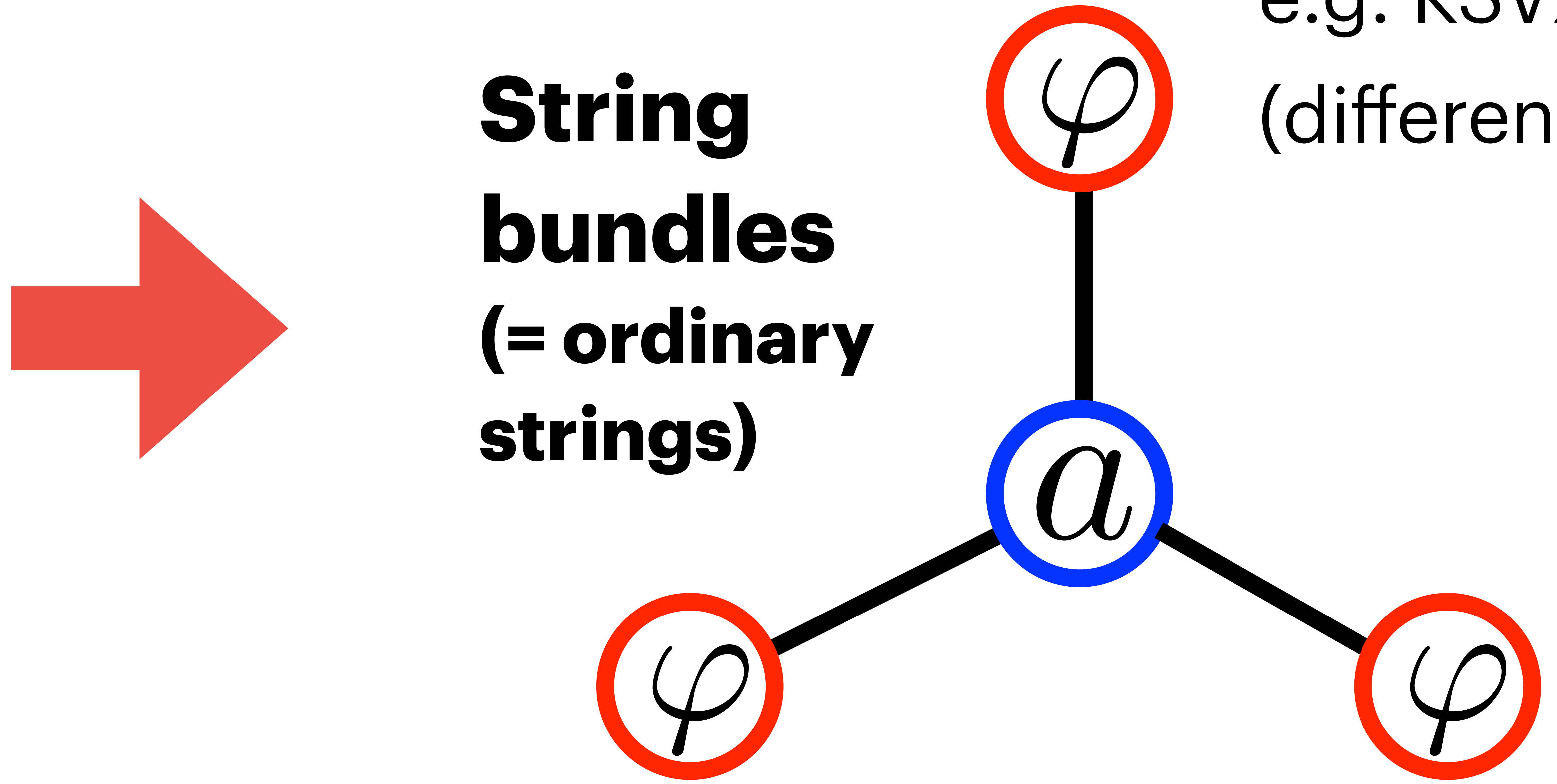
$N_{\rm DW} = 3 \text{ or } 6$ for DFSZ axion

φ : Massless (or very light) axion
e.g. KSVZ axion
(different from a)

String-wall network



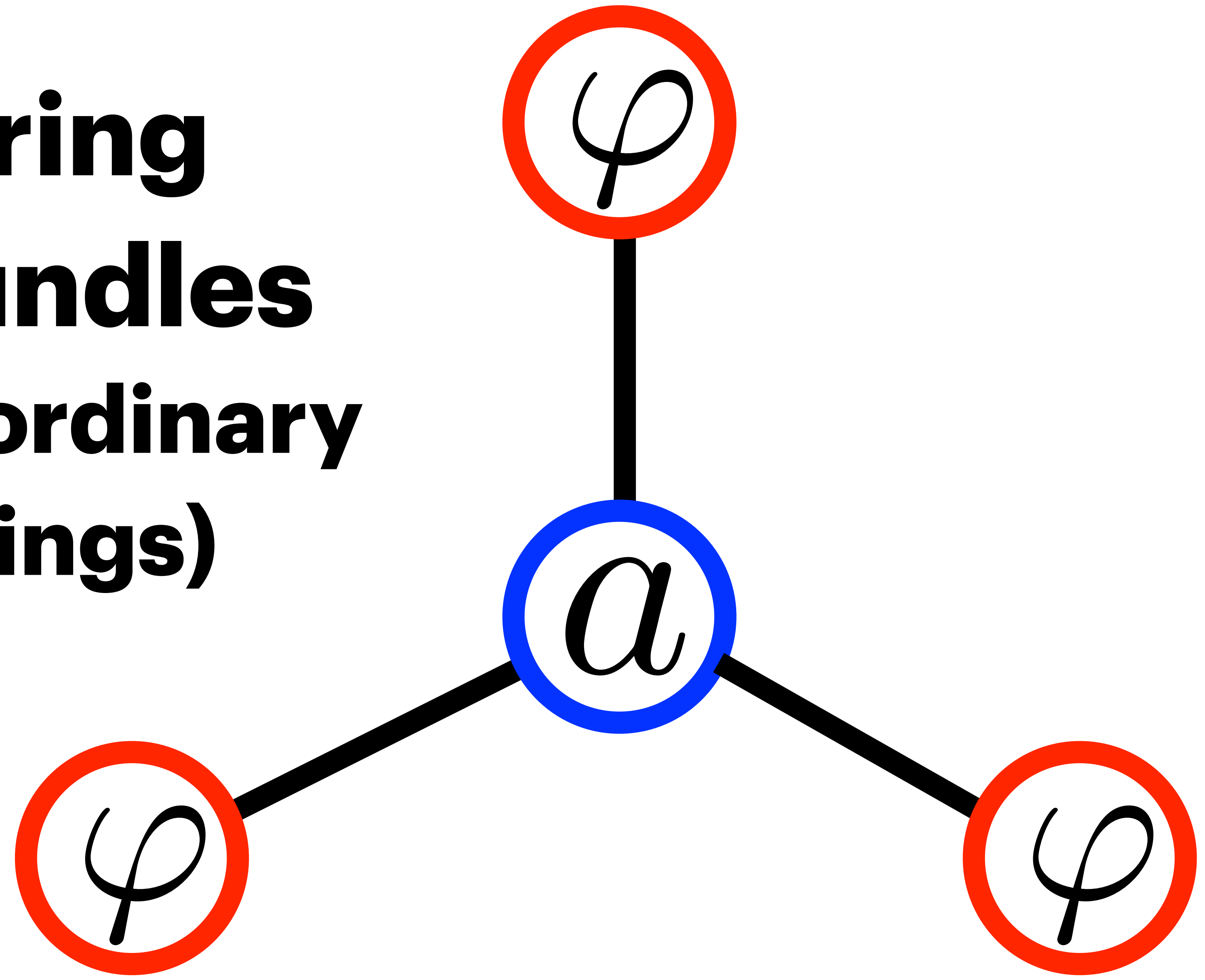
String bundles
(= ordinary strings)



New solution to the QCD axion DW problem

Lee, Murai, FT and Yin 2507.xxxxx, FT and Yin 2012.11576

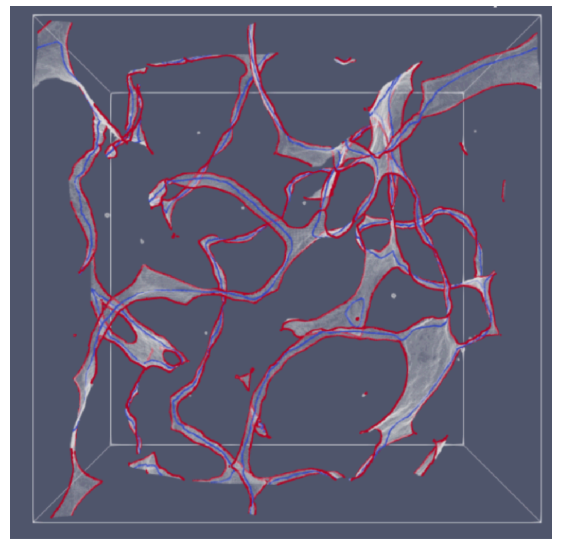
String bundles (= ordinary strings)



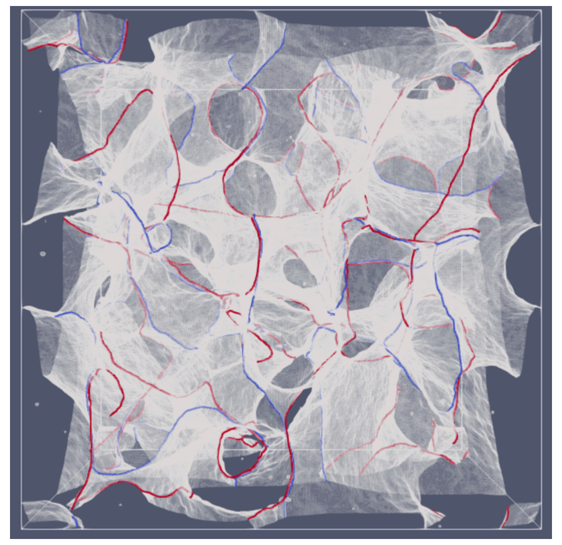
Unlike the usual QCD axion strings, the string bundles are stable and contribute to the cosmic birefringence!

See talks by Andrew Long and Eiichiro Komatsu for CB

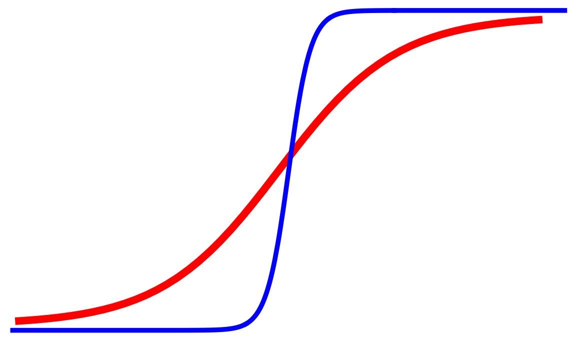
Summary



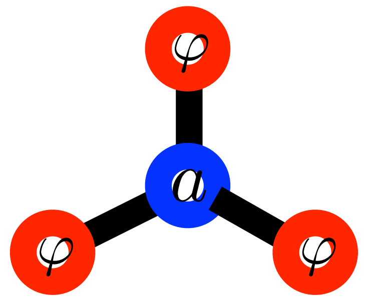
- The origin and breaking of U(1) PQ are unknown.



- A minimal extension from one to two PQ scalars often leads to stable DWs with large tension.
- Their decay can produce GWs and QCD axion DM even for small f_a.



- Adding another light axion can solve the QCD axion DW problem, leading to stable string bundles.



In multi-axion systems, cosmology changes drastically —
More is different!

Back-up slides

Symmetry breaking pattern

$$U(1)_{\phi_1} \times U(1)_{\phi_2} = U(1)_H \times U(1)_L$$

$$\xrightarrow{V_1} Z_{\frac{n_1^2+n_2^2}{d}} \times U(1)_{L=\mathrm{PQ}}$$

$$\xrightarrow{V_2} Z_{\frac{n_1^2+n_2^2}{d}} \times Z_{\frac{|n_1 n_2' - n_2 n_1'|}{d}} = N_{\mathrm{DW}}$$

$$V_1(\phi_1, \phi_2) = \Lambda^4 \left[1 - \cos\left(n_1 \frac{\phi_1}{f_1} + n_2 \frac{\phi_2}{f_2}\right)\right] \qquad d = \gcd(n_1, n_2)$$

$$V_2(\phi_1, \phi_2) = \Lambda'^4 \left[1 - \cos\left(n_1' \frac{\phi_1}{f_1} + n_2' \frac{\phi_2}{f_2} + \alpha\right)\right] \qquad \Lambda \gg \Lambda'$$

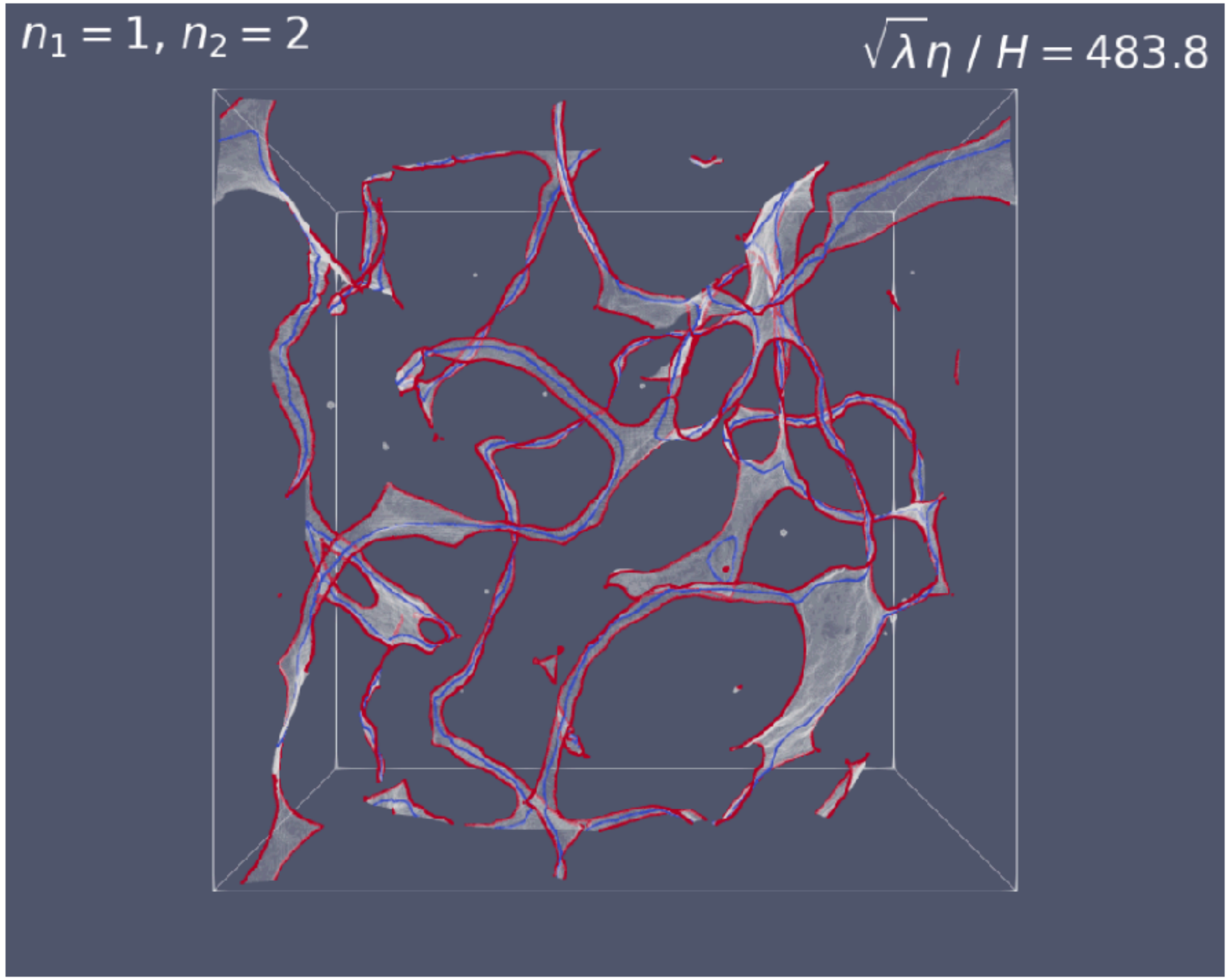
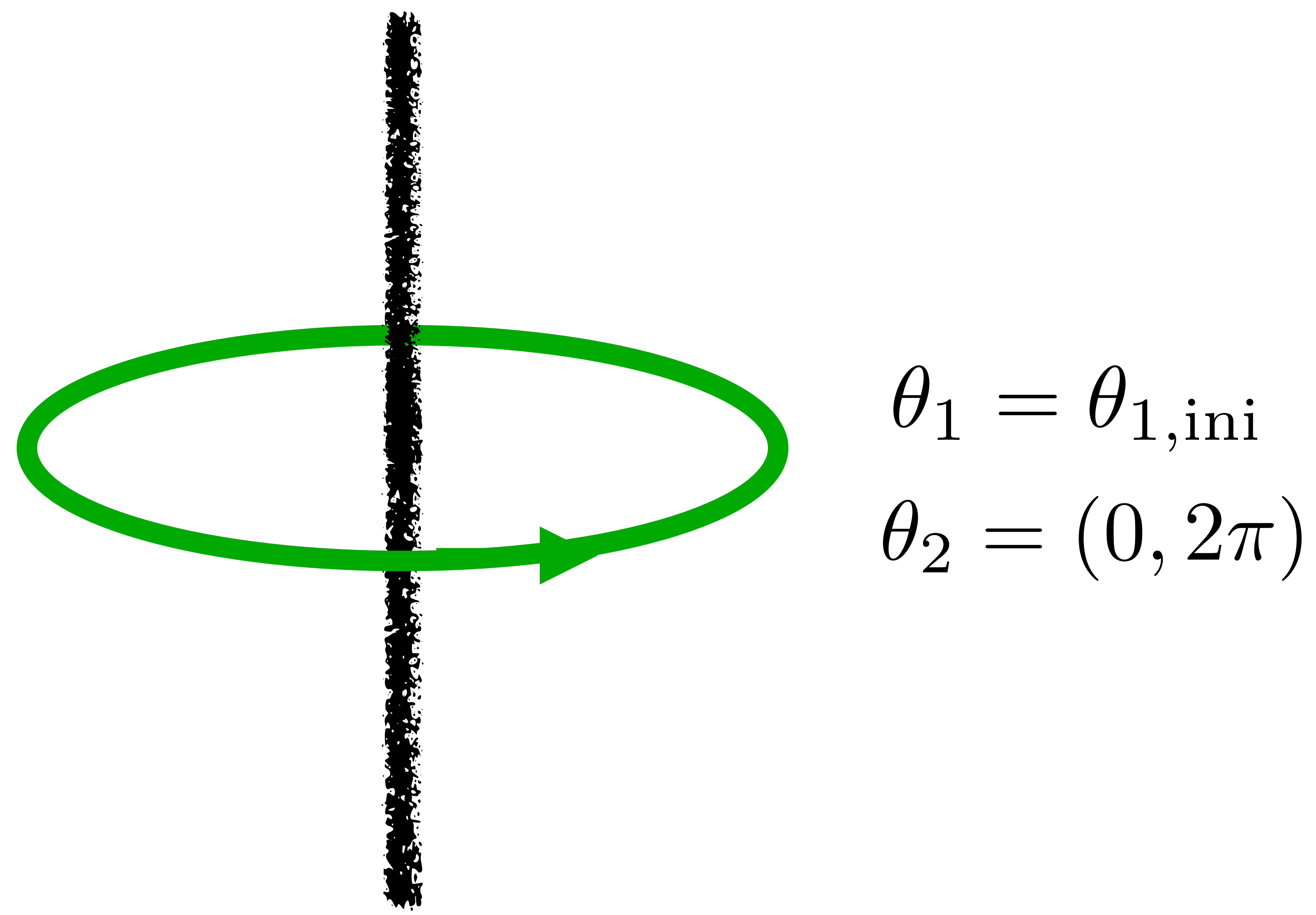
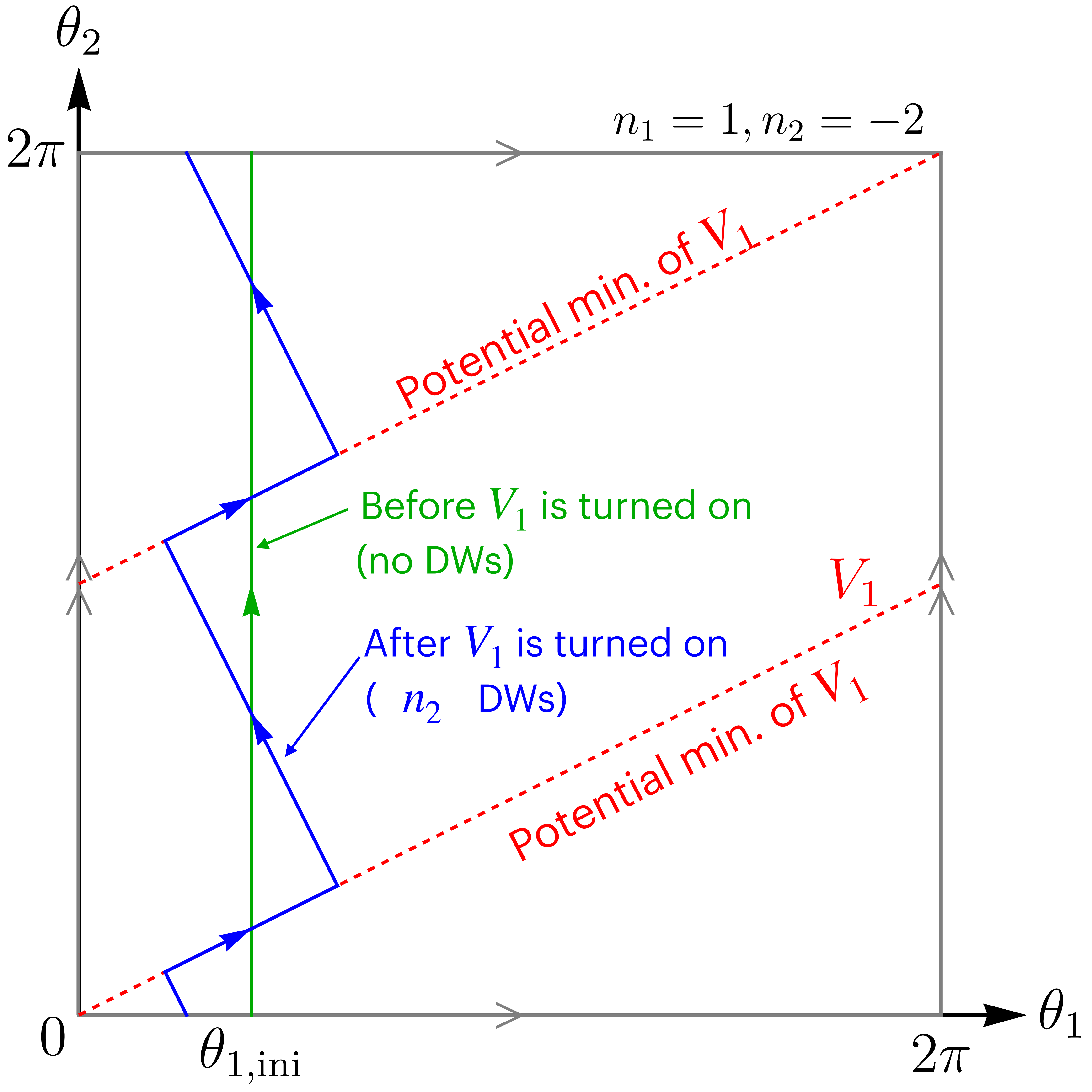


FIG. 8. The snapshots of strings and DWs at the time when $H^{-1} = 483.8/\sqrt{\lambda}\eta$ for $(n_1, n_2) = (1, 2)$ are shown. The blue (red) lines correspond to 1-strings (2-strings) and the semi-transparent white surfaces correspond to heavy axion DWs. One side of the box is the size of two Hubble horizons.

Pre-inflationary for ϕ_1, post-inflationary for ϕ_2 with $(n_1, n_2) = (1, -2)$ and $(n_1', n_2') = (1, -3)$.

Green vertical line: field configuration around a ϕ_2 string in the absence of V_1 and V_2.

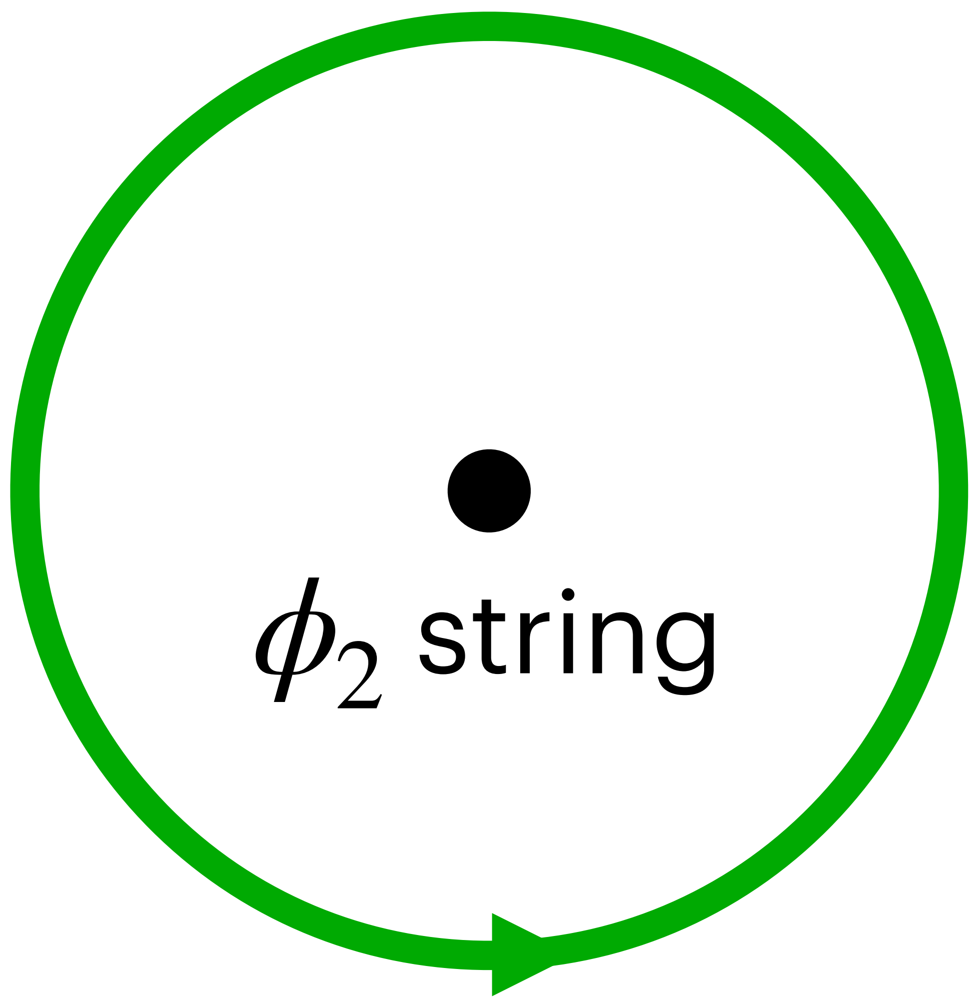




Before V_2 is turned on.

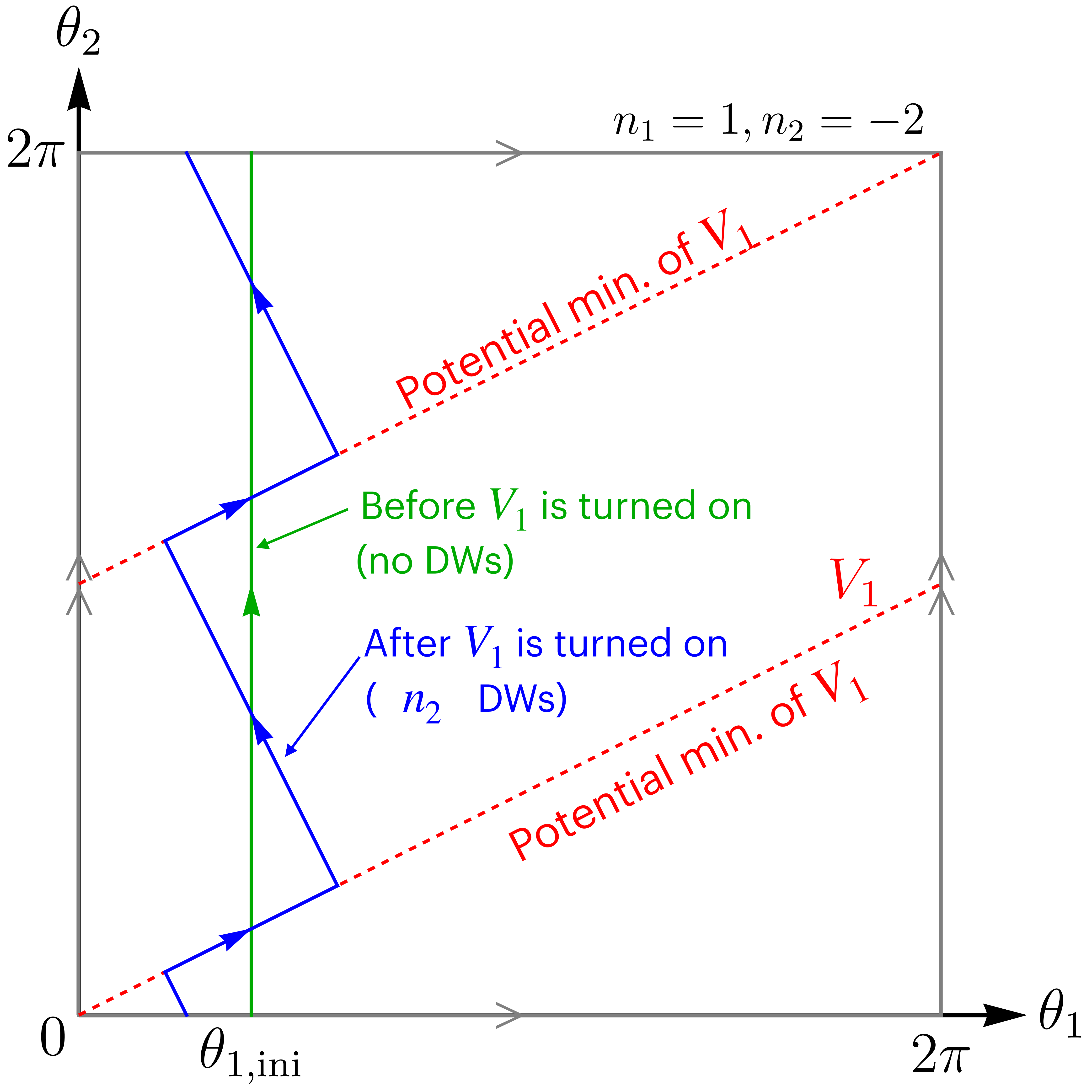
Pre-inflationary for ϕ_1, post-inflationary for ϕ_2 with $(n_1, n_2) = (1, -2)$ and $(n_1', n_2') = (1, -3)$.

Green vertical line: field configuration around a ϕ_2 string in the absence of V_1 and V_2.



$\theta_1 = \theta_{1,\mathrm{ini}}$

$\theta_2 = (0, 2\pi)$

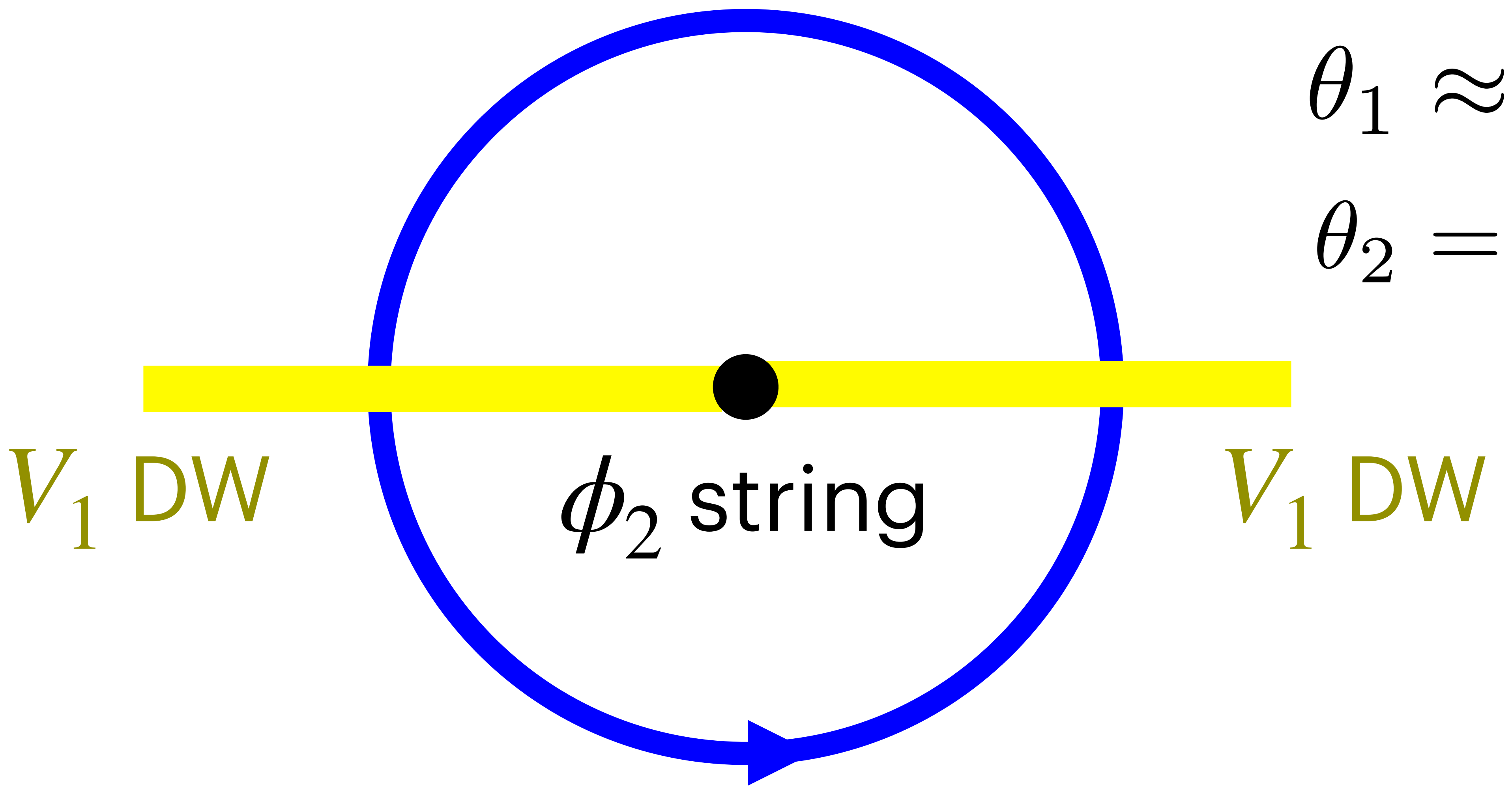


Before V_2 is turned on.

Pre-inflationary for ϕ_1, post-inflationary for ϕ_2 with $(n_1, n_2) = (1, -2)$ and $(n_1', n_2') = (1, -3)$.

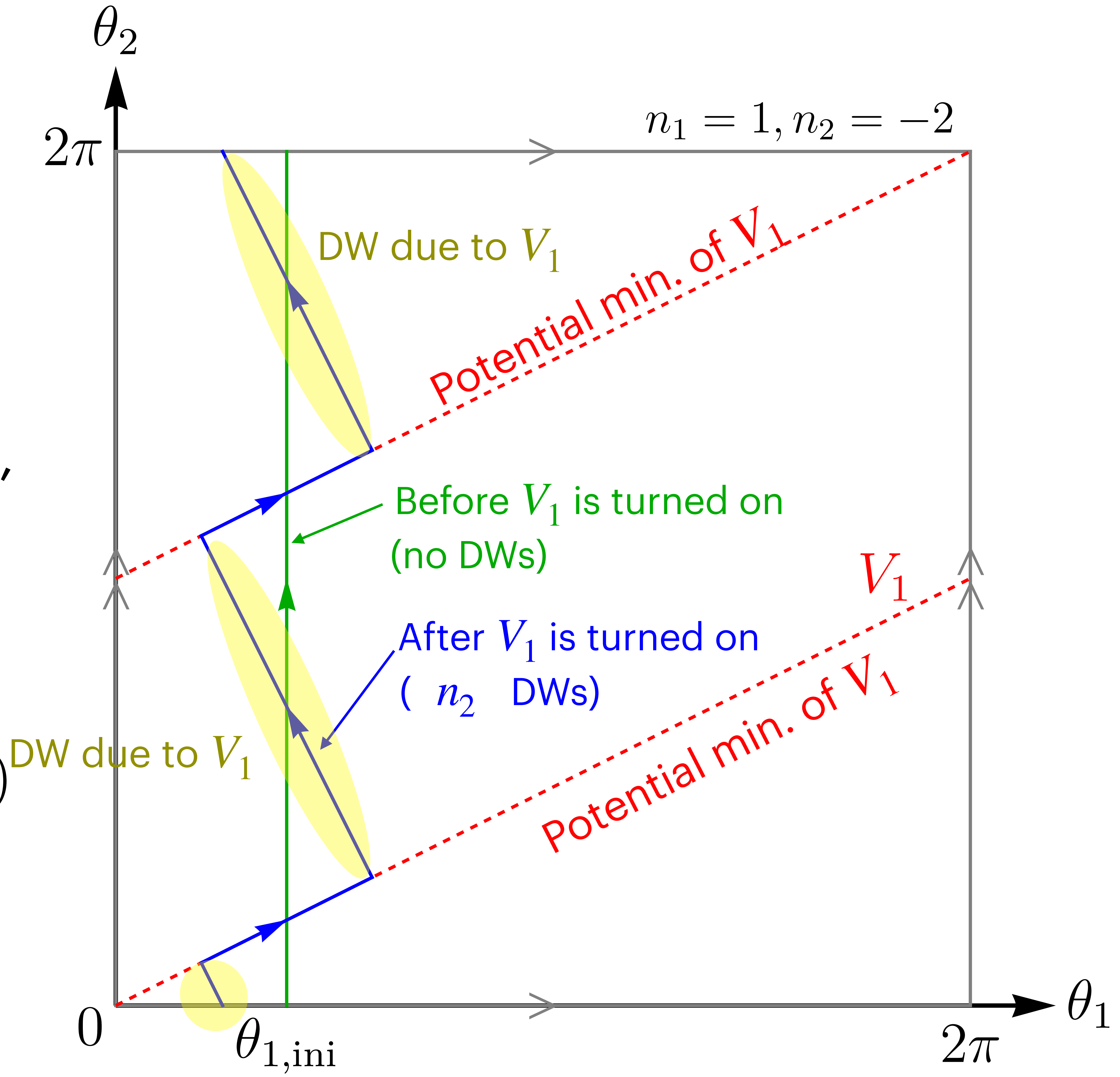
Green vertical line: field configuration around a ϕ_2 string in the absence of V_1 and V_2.

Blue line: ϕ_2 string and two DWs with V_1 present, and V_2 absent.

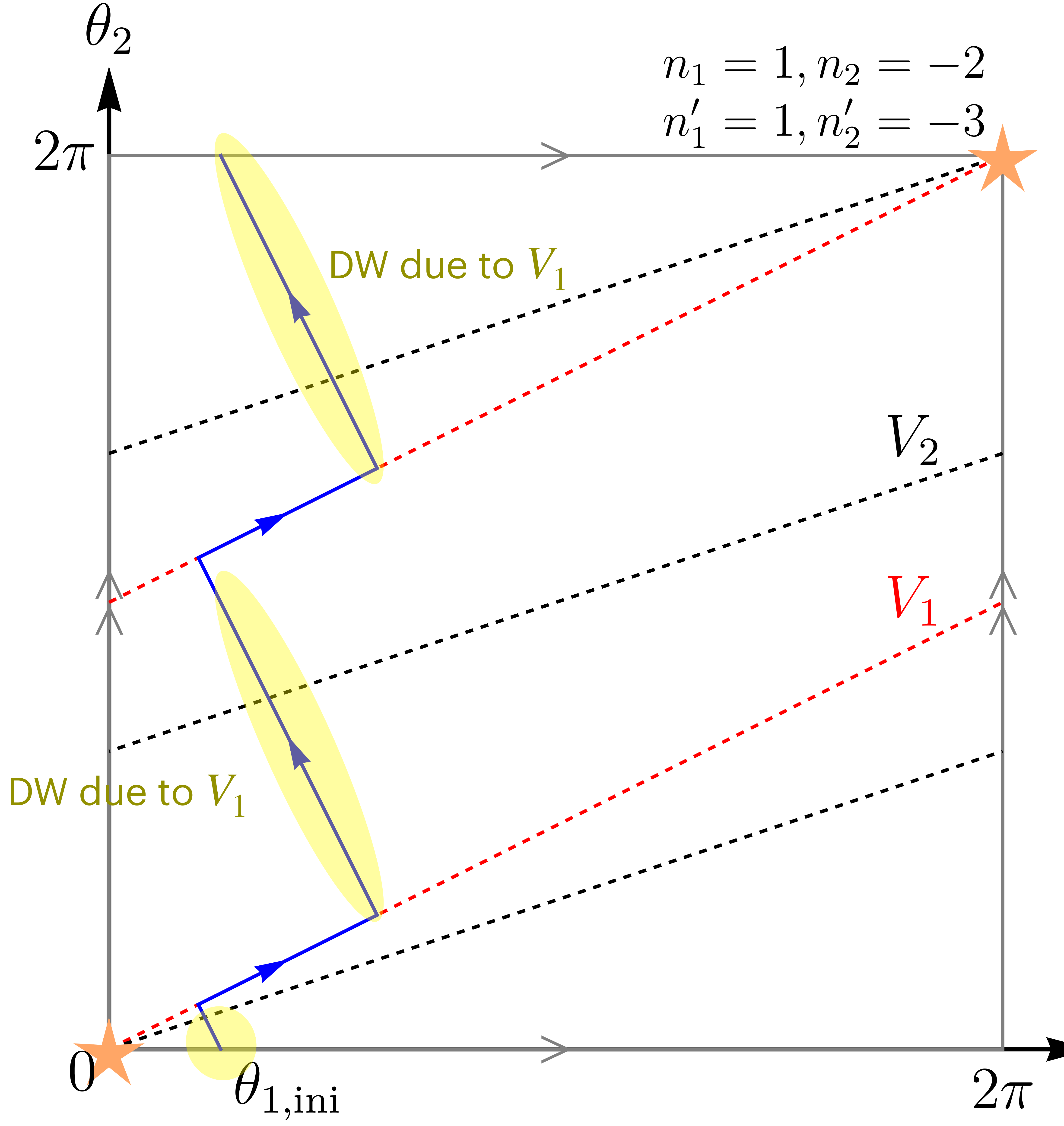


$\theta_1 \approx \theta_{1,\mathrm{ini}}$

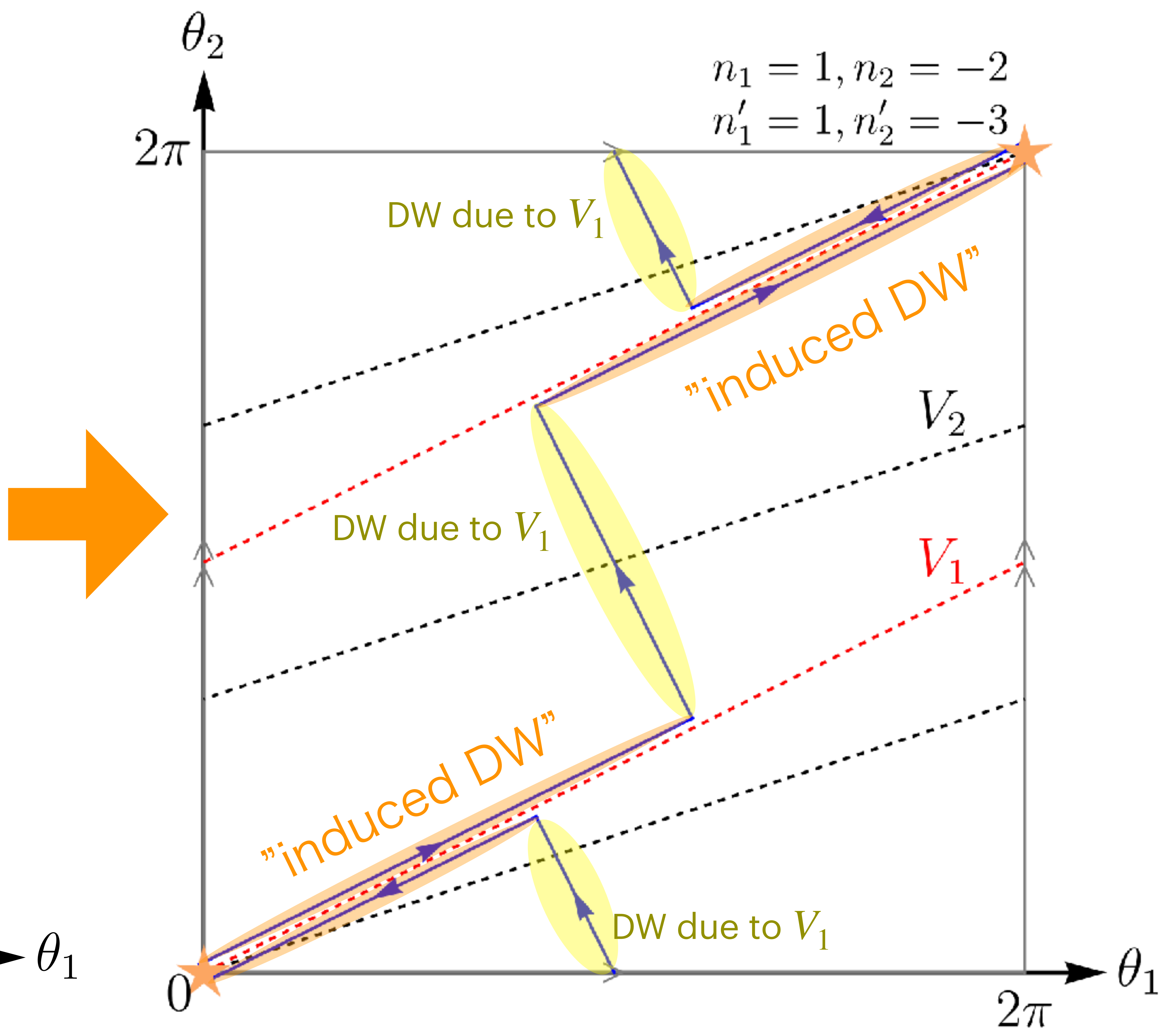
$\theta_2 = (0, 2\pi)$



Before V_2 is turned on.



Before V_2 is turned on.



After V_2 is turned on.

TABLE I. Fate of the string-wall network in the post-post scenario. We consider cases where at most one of n_1, n_2, n_1', n_2' is zero. Here, we assume $|n_1| \geq |n_2|$ for convenience. For $|n_1| < |n_2|$, this table should read with $n_1 \leftrightarrow n_2$. $D \equiv |n_1 n_2' - n_2 n_1'|$. d is the greatest common divisor of $(|n_1|, |n_2|)$, with $d = n_2$ for $n_1 = 0$ and $d = n_1$ for $n_2 = 0$. $N_{\rm DW} \equiv D/d$. Columns V_1 and V_2 show topological defects formed or decayed when each potential becomes effective. We assume that the effects of V_1 are more significant than those of V_2.

n_1	n_2	$N_{\rm DW}$	V_1	V_2	figure	
1	0	-	decay of ϕ_1 strings	ϕ_2 strings with $\lvert n_2' \rvert$ DWs	-	
≥ 2	0	-	ϕ_1 network & ϕ_2 strings	stable network	-	
≥ 1	1	0	θ_L string bundles	θ_L string bundles	Figs. 2 and 8	-
		1		decay of network		Fig. 7
		≥ 2		stable network		-
≥ 2	≥ 2	-	stable network	stable network with induced DWs or collapse by transient bias	Fig. 2	-

TABLE II. Fate of the string-wall network in the pre-post scenario. Other definitions and column descriptions are the same as in Table I.

$\lvert n_2 \rvert$	D	n_2'/n_2	V_1	V_2	figure
0	-	-	ϕ_2 strings without DWs	ϕ_2 strings with $\lvert n_2' \rvert$ DWs	-
1	-	-	decay of ϕ_2 strings	no defects	-
≥ 2	0	integer	ϕ_2 strings with $\lvert n_2 \rvert$ DWs	stable network	Fig. 5
		non-integer		network collapse by bias	Fig. 6
	$\neq 0$	-		stable network with induced DWs* or collapse by transient bias	Figs. 3 and 4

* If $n_1 = 0$ and n_2'/n_2 is an integer, the stable network does not accompany induced DWs.