Strings, gravity, and nonlocality

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Based on:

hep-th/0604072; hep-th/0703116; arXiv:0705.1816, w. Gross and Maharana; arXiv:0705.2197

Cosmology, strings, and phenomenology Nordita, Stockholm A plausible viewpoint: the black hole information paradox is of comparable importance to the paradox of the classical instability of matter

Why a paradox? Apparently must abandon a cherished principle of physics:

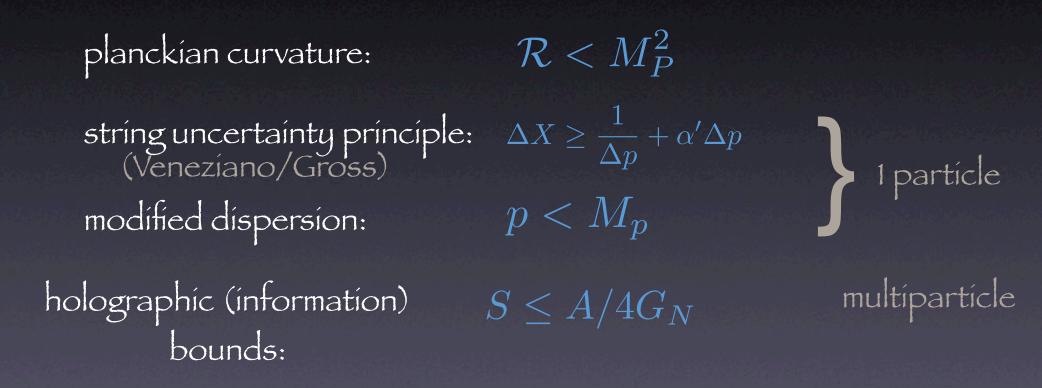
- unitarity and energy conservation (QM violated)
 stability (remnants)
- macroscopic locality (information escapes)

...widespread belief

But, if nonlocality:

1) What is the mechanism 2) How is Hawking's argument evaded 3) Where does GR+local QFT fail? - what is the correspondence limit for new physics?

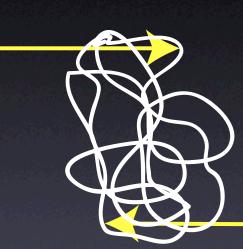
Some existing proposals for the correspondence limit:



Will explore some issues that seem to shed some light on these questions.

Begin with: The role of strings

If there is nonlocality in string theory, would expect to find evidence of it in high-energy scattering. Does string extendedness provide the mechanism for nonlocality?



What does this have to do w/BH formation? (Does it prevent? Or is this BH formation?) (Q's: Strominger, Gross, ...; string spreading - Susskind)

Long strings? $L \sim E/M_s^2$

String uncertainty principle? $\Delta X \ge \frac{1}{\Delta p} + \alpha' \Delta p$ (Veneziano, Gross) (\leftarrow nonlocality)

(Proposed app. to BH info: LPSTU)

Let's investigate ...

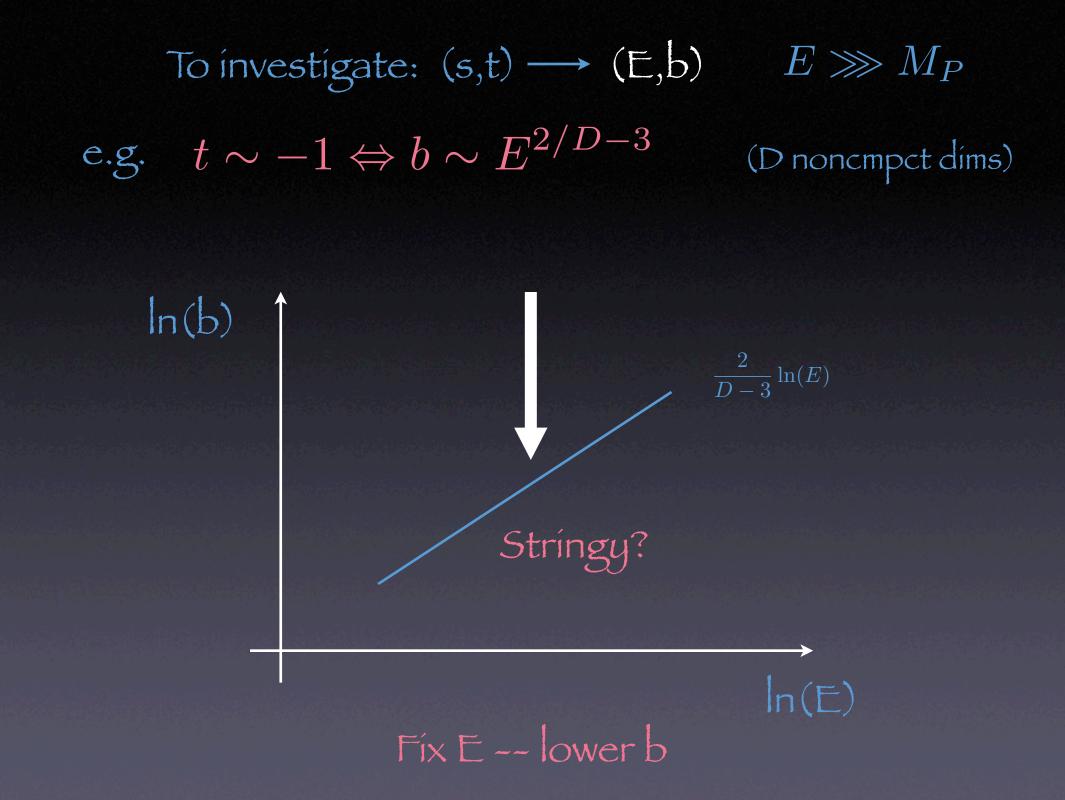
Begin w/tree-level amplitude: high E

$$\mathcal{A}_{0}^{\rm string}(s,t) \propto g_{s}^{2} \frac{\Gamma(-t/8)}{\Gamma(1+t/8)} s^{2+t/4} e^{2-t/4}$$

V5.

$$\mathcal{A}_0^{\mathrm{grav}}(s,t) \propto G_D \frac{s^2}{t}$$

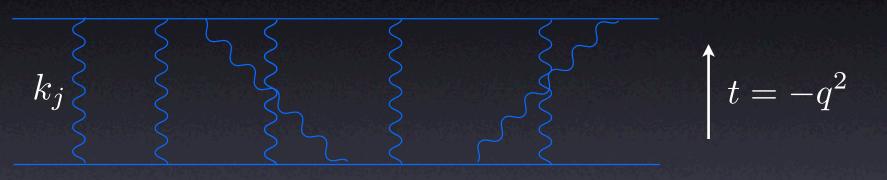
• No evidence for long strings; • But significant modifications for $\ t\sim -1$



To check, compare loops:

(Following Amati, Ciafaloni, Veneziano; Muzinich-Soldate; SBG, Gross, Maharana)

Ultrahigh-E: Eikonal



$$i\mathcal{A}_{N}^{\text{string}} = \frac{2s}{(N+1)!} \int \left[\prod_{j=1}^{N+1} \frac{d^{D-2}k_j}{(2\pi)^{D-2}} \frac{i\mathcal{A}_{0}^{\text{string}}(s, -k_j^2)}{2s} \right] (2\pi)^{D-2} \delta^{D-2} \left(\sum_{j} k_j - q_\perp \right)^{D-2} \delta^{D-2} \left(\sum_{j=1}^{N+1} \frac{d^{D-2}k_j}{(2\pi)^{D-2}} \frac{i\mathcal{A}_{0}^{\text{string}}(s, -k_j^2)}{2s} \right)^{D-2} \delta^{D-2} \left(\sum_{j=1}^{N+1} \frac{d^{D-2}k_j}{(2\pi)^{D-2}} \frac{i\mathcal{A}_{0}^{\text{strin$$

 $\prod_{j=1}^{N+1} \frac{E^{2-\alpha' k_j^2}}{k_j^2} \qquad \qquad 1) \ k_j \approx q/(N+1)$ $2) \ E^{-\alpha' q^2/(N+1)}$

Thus at large N, string corrections get smaller Which N dominates? Can sum eikonal series: $i\mathcal{A}_{eik}(s,t) = 2s \int d^{D-2}\mathbf{b}e^{-iq_{\perp}\cdot\mathbf{b}}(e^{i\chi(b)}-1)$ with $\chi(b) \sim G_D \frac{E^2}{bD-4}$

Dominant N: $N \sim \frac{G_D E^2}{b^{D-4}}$;

 \Leftrightarrow

At $t \sim -1$: $N \sim (G_D E^2)^{\frac{1}{D-3}}$

. Large loop order dominates.

Two Aichelburg-Sex shocks (ACV: checks)



But - can excite strings: "diffractive excitation" (ACV)

Indeed, unexcited (elastic) amplitude, near Schwarzschild radius:

So:

$$\mathcal{A}_{el} \sim \exp\left\{-E^{(D-4)/(D-3)}\right\} \qquad !!$$

?? No black hole?? Info carried away? (Venezíano, 2004)

But - intuition: string only "spread out" "after" collision?? However, string spreading is a notoriously fuzzy concept...

Where is the string?

Karlíner, Klebanov, Susskind: ít depends





"low resolution"

"high resolution"

So: need to check for process in question ...

A test:

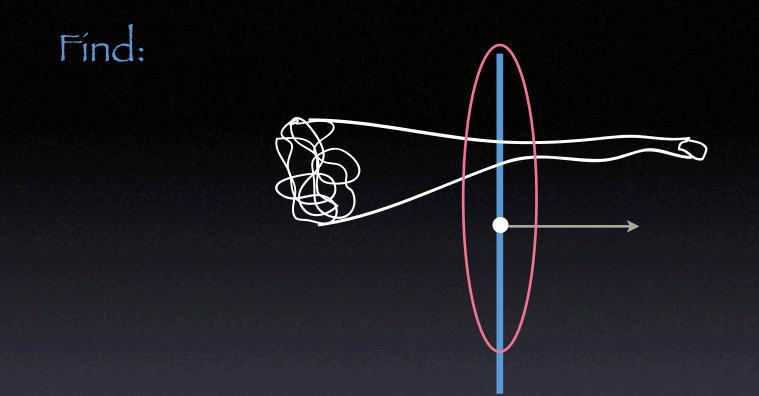


 $ds^{2} = -dudv + dx^{i}dx^{i} + \Phi(\rho)\delta(u)du^{2}$ $\Phi(\rho) = -8G\mu\ln\rho \quad , \quad D = 4$ $\Phi(\rho) = \frac{16\pi G\mu}{\Omega_{D-3}(D-4)\rho^{D-4}} \quad , \quad D > 4$

Scattering in a plane-wave metric: de Vega and Sanchez; Horowitz and Steif *Light cone quantization*

Compute for incoming unexcited string: $\langle \hat{X}^i_{\epsilon}(\tau,\sigma) \hat{X}^i_{\epsilon}(\tau,\sigma) \rangle$

Where $\hat{X}^i_\epsilon(\tau,\sigma)$ is deviation from CM of string, w/world sheet regulator ϵ



Indeed, origin of effect is "tidal string excitation" $(\Delta X)^2 \sim |\ln \epsilon| + \left[\frac{G_D E^2}{b^{D-2}}\tau\right]^2 |\ln \tau| \qquad \epsilon \ll \tau$

For small tau: inside trapped surface

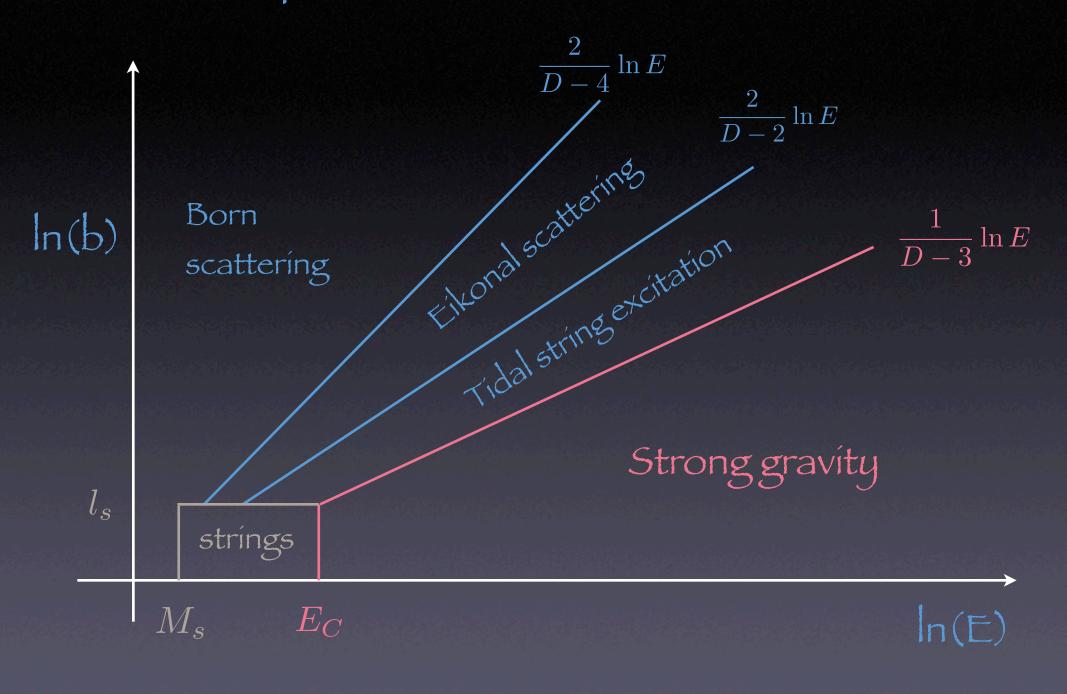
Thus:

String appears to behave ~locally during collision
Trapped surface (aka black hole) appears to form

What conclusions can we draw?

1. No evidence string effects prevent BH formation

Suggested "phase diagram:"



Moreover:

BH formation corresponds to breakdown of the gravitational loop expansion:



In what sense is string theory a complete theory of quantum gravity?

 $\mathcal{O}(R_S/b)$

2. No obvious precise notion of locality

Local QFT bounds

Strong gravity/ black hole regime:

 $\sigma_T \le c(\ln E)^{D-2}$
Froissart

 $\sigma_T(E) \sim [R_S(E)]^{D-2} \sim E^{\frac{D-2}{D-3}}$

 $|\mathcal{A}_{el}(s,t)| \ge e^{-f(\theta)E\ln E}$ Cerulus-Martín $\mathcal{A}_{el}(s,t) \sim e^{-S_{BH}}$ $\sim e^{-ER_S(E)} \sim e^{-E^{(D-2)/(D-3)}}$

3. Scattering appears dominated by strong gravitational effects; this suggests that any nonlocality would have its origin in gravitational dynamics (as opposed to, e.g., string extendeness)

side comment: perhaps unitarity is a deeper issue than renormalizability in quantum gravity (compare EW physics); BH info paradox suggests it can be respected at the price of locality 4. Suggested correspondence boundary where does GR+LQFT break down?

2 part Fock sp.:

 $\phi_{x,p}\phi_{y,q}|0
angle$

(mín uncertaínty wavepackets)

good description for $|x - y|^{D-3} > G|p + q|$ where $G \sim G_{Newton}$

> "the locality bound" (extends off shell?)

Other versions of the locality bound: Measurement limit: $\Delta t (\Delta x)^{D-3} \ge G\hbar$ N-particle: $\phi_{x_1,p_1}\cdots\phi_{x_N,p_N}|0\rangle$ not good for $\operatorname{Max}|x_i - x_j|^{D-3} < G|\sum P_i|$ de Sítter: see SBG and Marolf, arXív:0705.1178 Suggestion: perhaps these are special cases of a broader "nonlocality principle," stating that the nonperturbative physics that unitarizes gravity in domains where gravitational perturbation theory fails is nonlocal

Of course, we have other hints of nonlocality:

Possible indicators of nonlocality: hints from AdS/CFT & holographic beliefs (though don't yet fully address issues) formulation of approx. local ("proto-local") observables, w/limitations SBG, Marolf, and Hartle, hep-th/0512200 SBG and M. Gary, hep-th/0612191 conundrums of cosmology (eternal inflation/ landscape; Boltzmann brains, etc.) and likely breakdown of GR+LQFT Arkaní-Hamed et al, arXív:0704.1814 SBG hep-th/0703116; SBG and Marolf, arXiv:0705.1178

Returning to the original problem:

How is the information paradox resolved?

Logical possibilities:

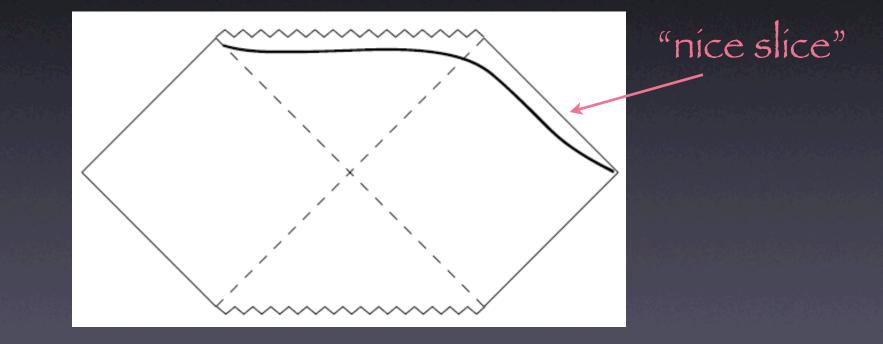
A. Compare classical - quantum, the H atom: classical physics doesn't break down at r_{atom} , it is replaced

B. Actual breakdown of semiclassical gravity

evidence for B...

Hawking's calculation (w/ updates):

$\Delta I = S \quad \leftarrow \rho \quad \leftarrow |\psi\rangle \quad \leftarrow |\psi\rangle_{NS}$

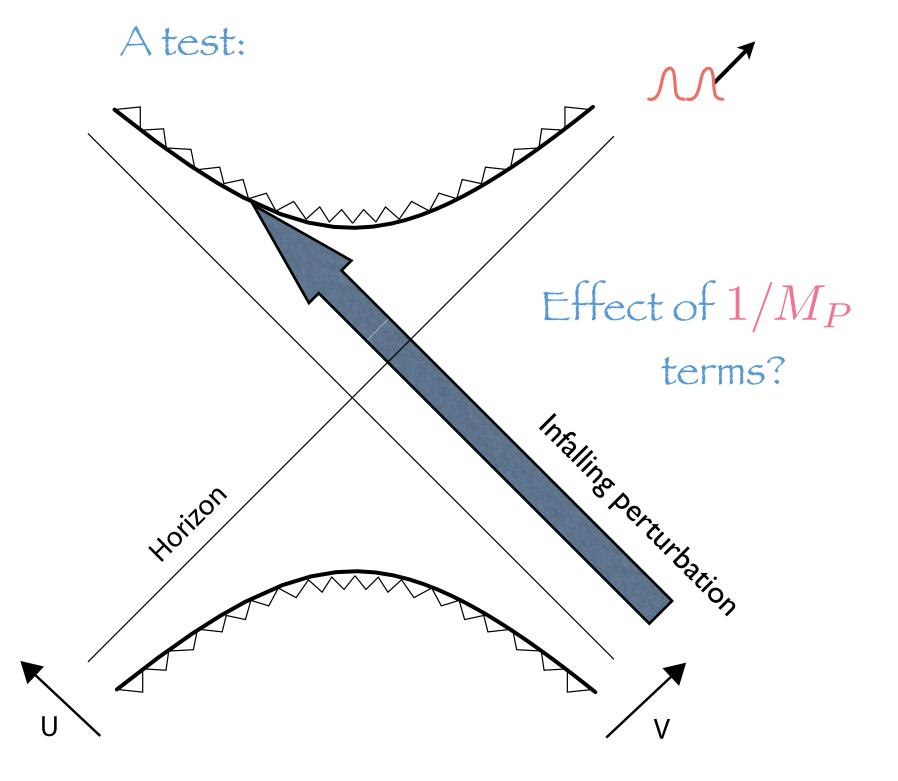


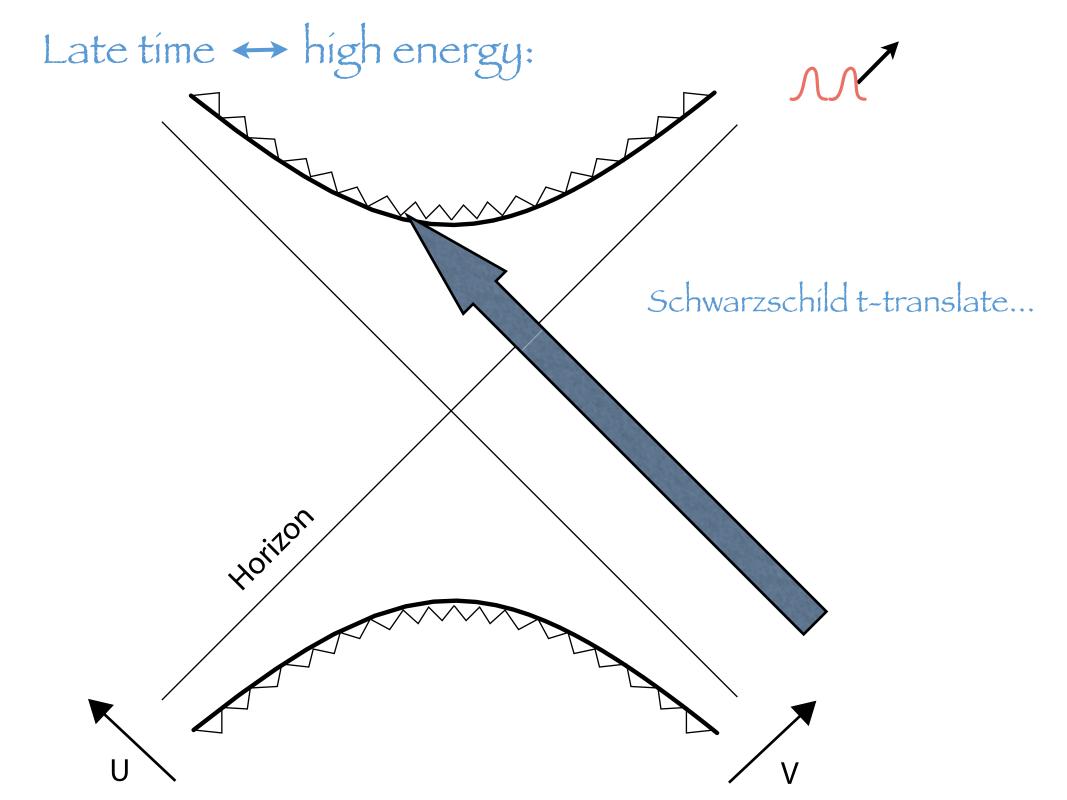
Can we justify this in a reliable approximation?

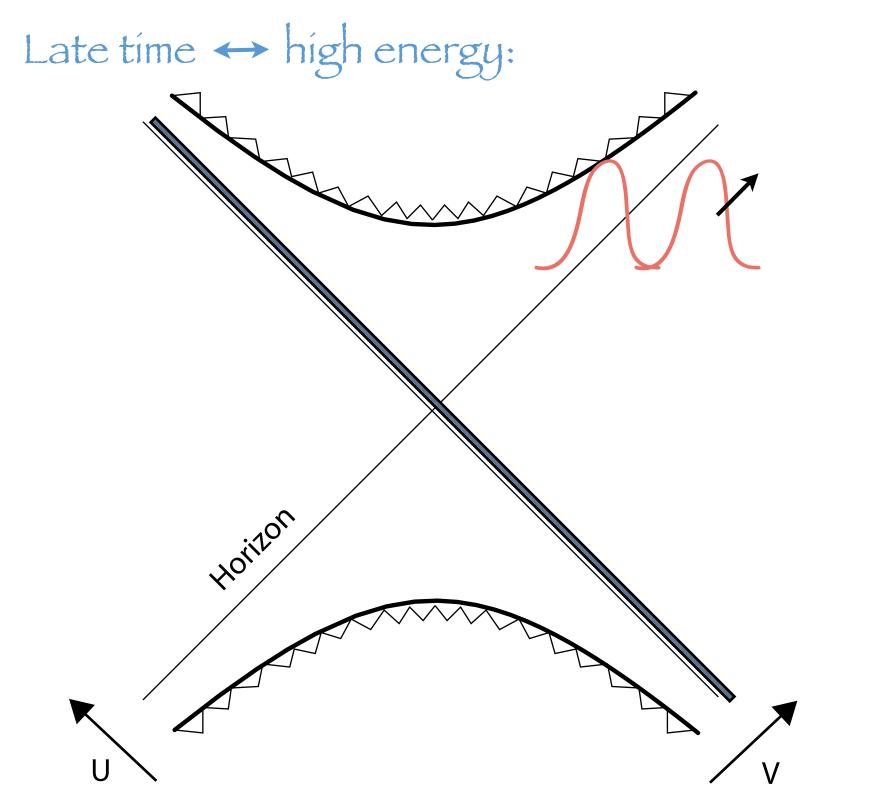
leading contribution - perturbative expansion in $1/M_P \label{eq:massed}$ (fix $M_P^{-2}M$)

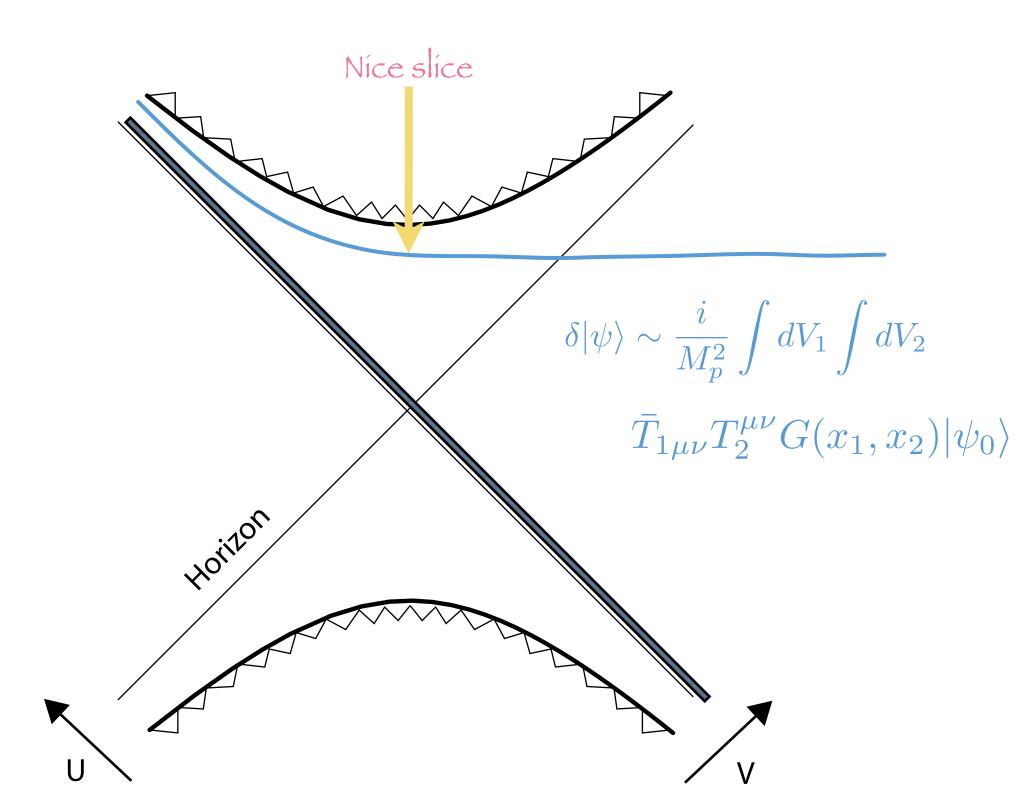
 \leftarrow QFT in semiclassical background (matter: ϕ)

 $|\Psi\rangle \sim \int_{\Psi_{in}} \mathcal{D}h\mathcal{D}\phi e^{iS}$ $g_{\mu\nu} = g_{\mu\nu}^0 + M_P^{-1} h_{\mu\nu}$ $S \sim \int d^4x \sqrt{-g_0} \Big\{ - (\nabla_0 \phi)^2 + h \Delta_L^0 h \Big\}$ semíclassical $+\frac{h}{M_P} \left[T^{\phi}_{\mu\nu} + (\nabla h)^2 \right] \cdots \right\} \qquad \frac{1}{M_P^n} \text{ terms}$ IMPORTANT?









w/ $1/M_P$ terms: estímate $\mathcal{O}(1)$ change in the nice slice state by the time $t \sim M^3$ Suggests breakdown of perturbation theory, at the required timescale (Page)

This leads to a proposal for how the information paradox is resolved:

1) Hawking's argument for information loss is not reliable; to accurately compute nice slice state need non-perturbative gravity (thus, no paradox)

2) The remaining information problem could be resolved if non-perturbative gravity has appropriate nonlocality Could supply more details, but another talk ...

similar considerations suggest that perturbative treatment of dS breaks down at time scale $t \sim R^3$ due to large fluctuations.

(possibly related to arguments of Arkani-Hamed et al)

Some conclusions:

- Pert. theory breaks down in HE scattering at sufficiently small b -- intrinsically gravitational
- No clear essential role for strings
- Possible intrinsically gravitational nonlocality?
- Pert thy apparently fails when computing ΔI
- Suggests: nonpert effects restore unitarity, at the price of locality?

What is needed: the nonperturbative, and quite possibly nonlocal, dynamics

(Is it string theory??)

Analogy to emergence of quantum mechanics, pre 1925

QMħ Hydrogen atom UV catastrophes Old quantization rules Uncertainty principle (Noble gases) Wave function Schrodinger eqn

? (NLM) \hbar, G Black hole Information paradox, ... Holographic princ; I=A/4 Nonlocality principle (locality bound, ...) (Extremal black holes)

What is this "nonlocal mechanics"?