

Kekulé Spirals in Twisted Bilayer Graphene

Steve Simon, Oxford University



Yves Kwan



→ Princeton

Glenn Wagner



→ Zurich

Nick Bultinck



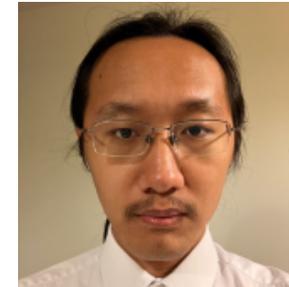
→ Ghent

Sid Parameswaran



Oxford

Ziwei Wang



Oxford



Tomohiro
Soejima

Berkeley



Michael
Zaletel

Berkeley



Erez
Berg

Weizmann



Yichen
Hu

Florida



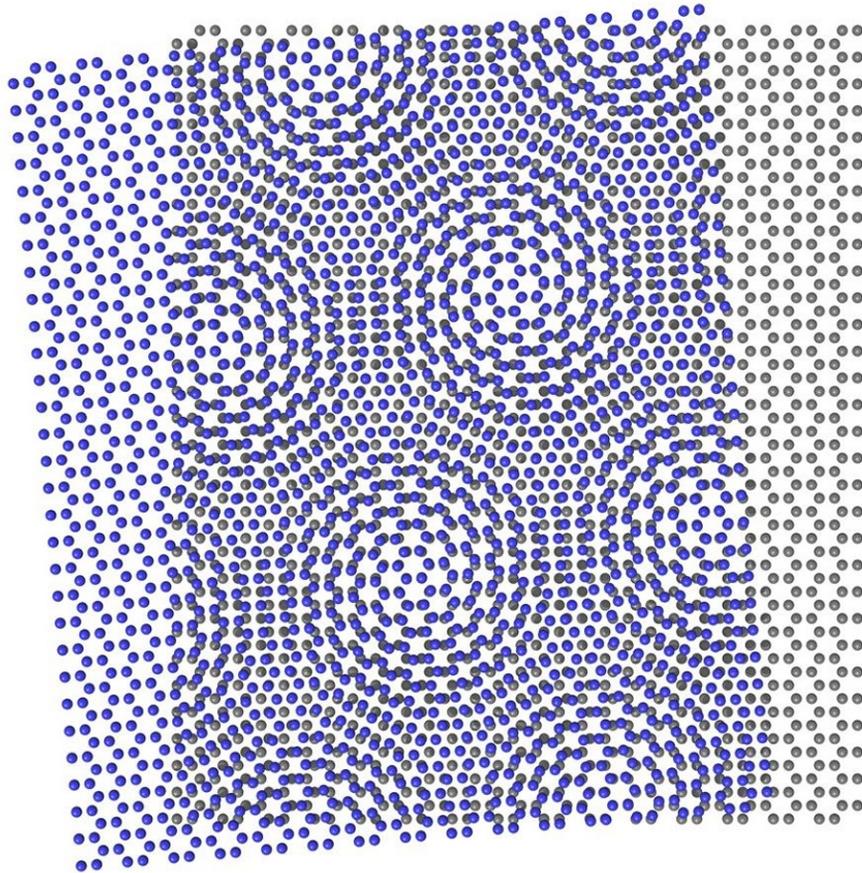
Nilotpal
Chakraborty

Dresden

[Kwan et al PRX 11, 041063 \(2021\)](#); PRX 12, 031020 (2022); PRB 110 (8), 085160 (2024);...
Wagner et al, PRL 128, 156401 (2022)...; Wang et al PRB 109, L201119 (2024)...

Do not capitalize moiré (unless it starts a sentence).

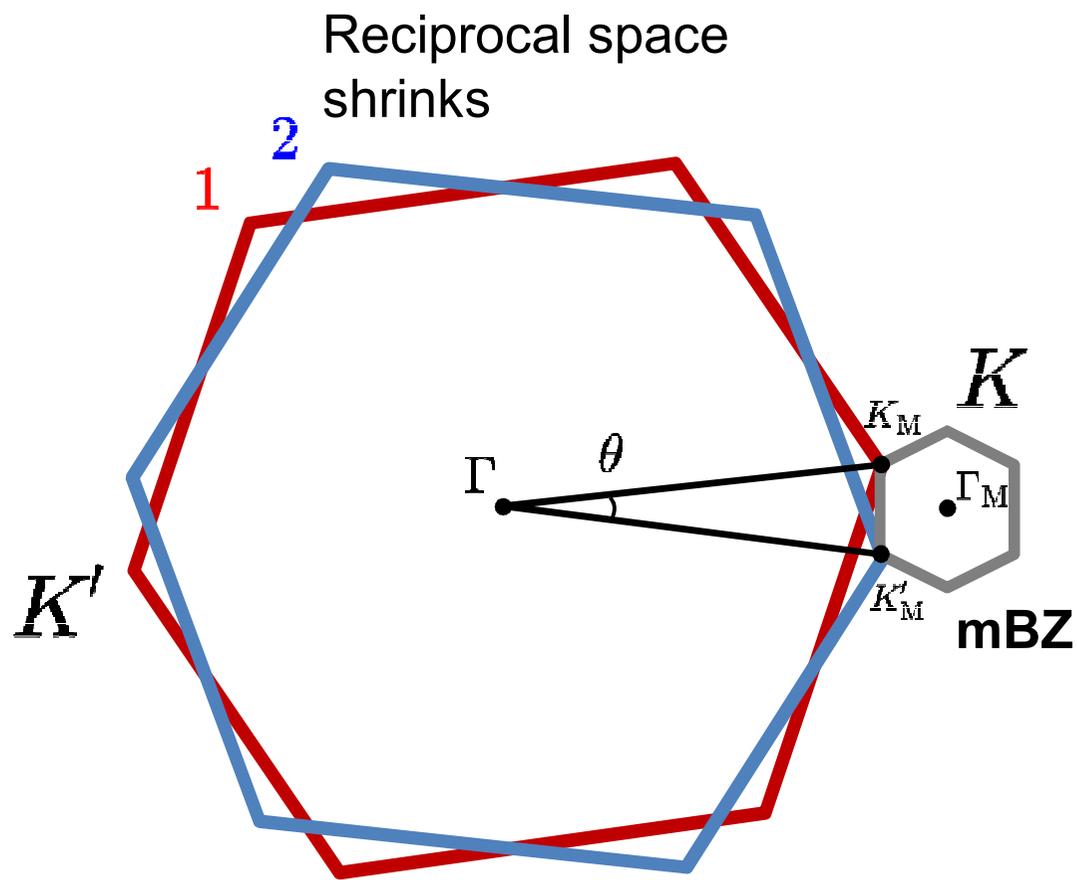
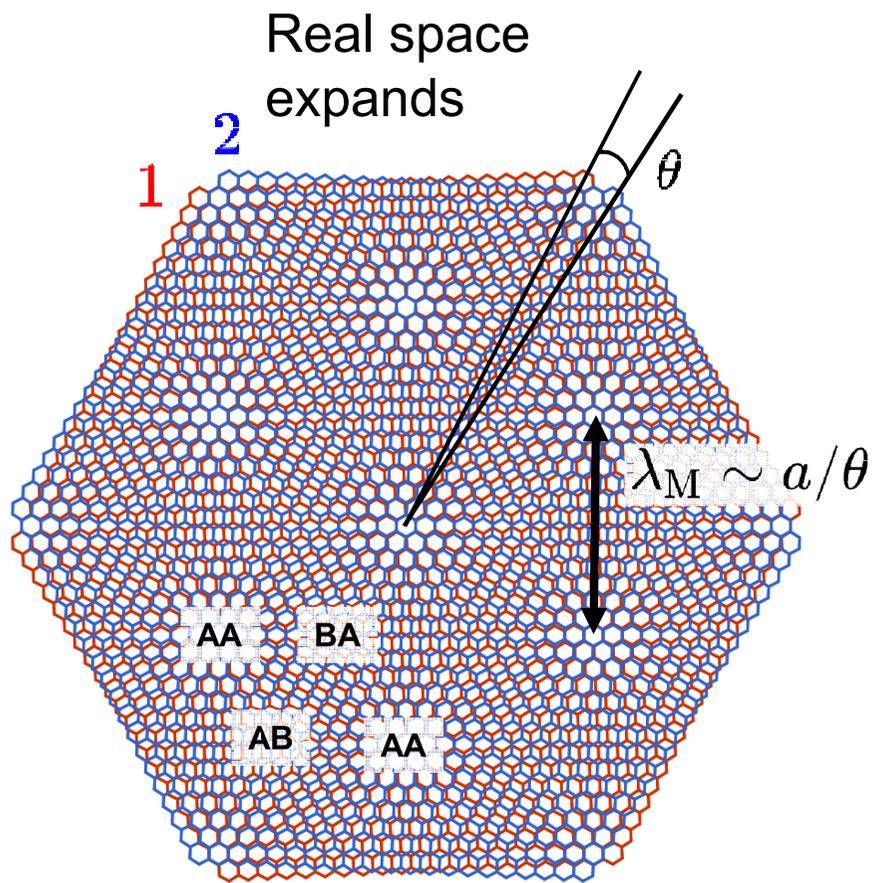
moiré: from *moirer*, from mohair, from *mukhayyar* or *khayyara* [chosen or best]



Moiré is the best!

Picture: NIST

moiré pattern

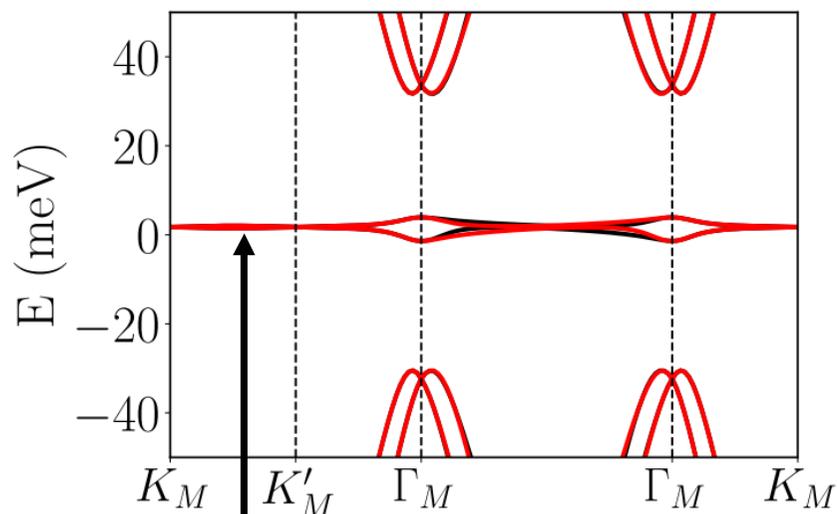


moiré pattern can generate flat bands:

Twisted Bilayer Graphene

Bistritzer & MacDonald, *PNAS* (2011)

FLAT BAND →
INTERACTIONS
DOMINATE



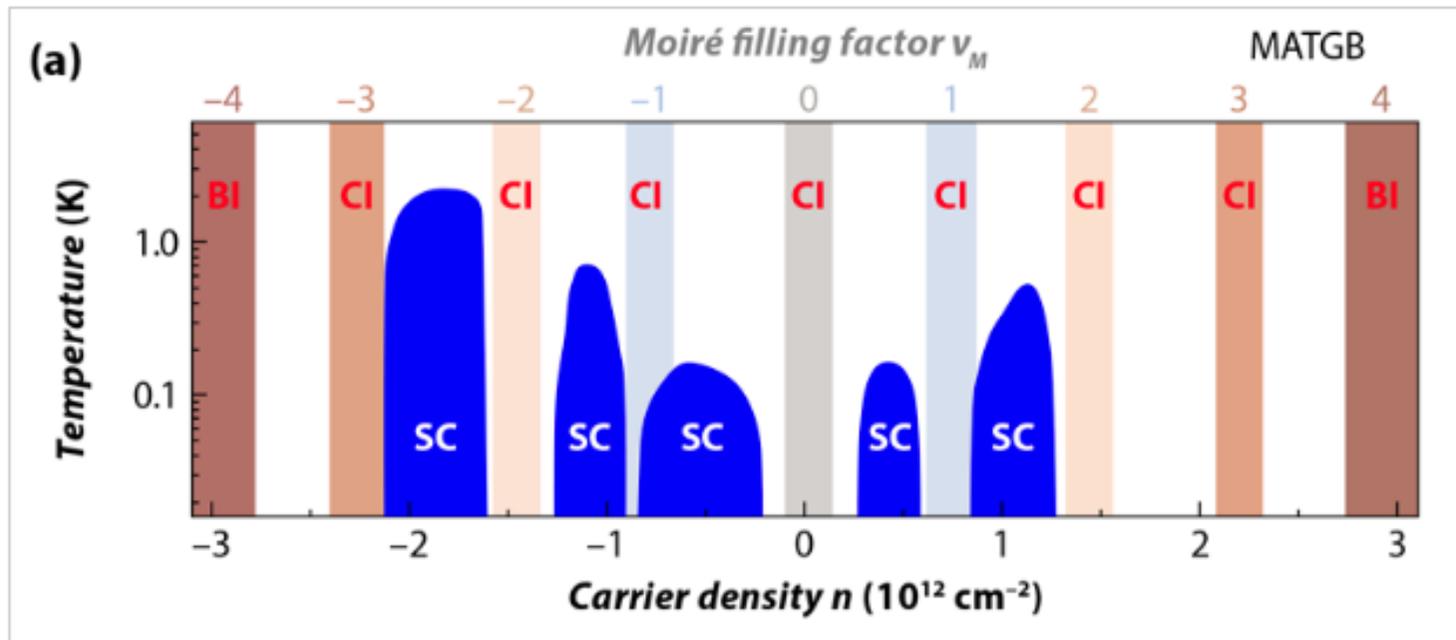
Magic Angle TBG

8 central bands
incl. spin and valley

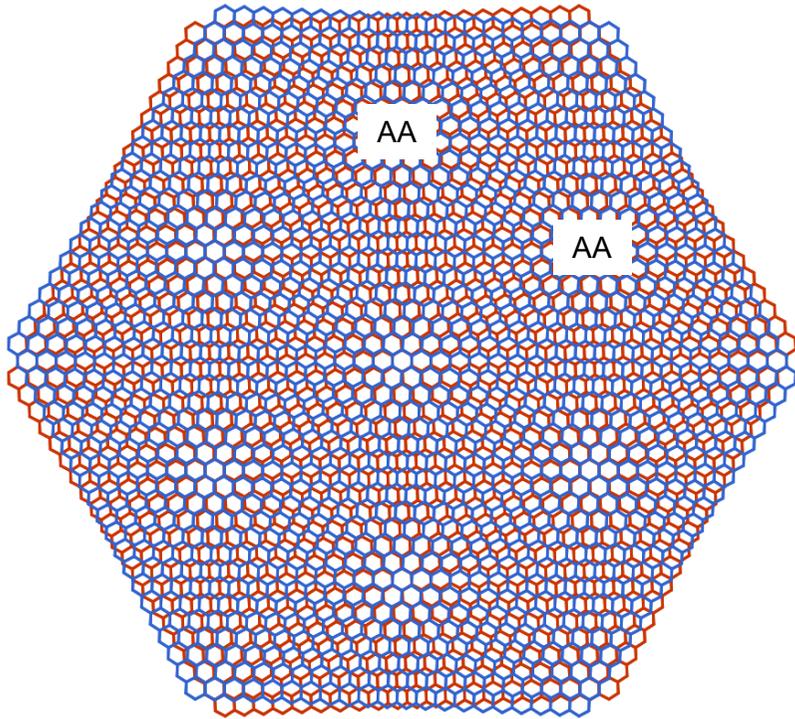
- focus on central bands

filling factor $\nu = -4, \dots, 0, \dots, 4$
(CNP)

Rough phase diagram (Stolen from MacDonald):

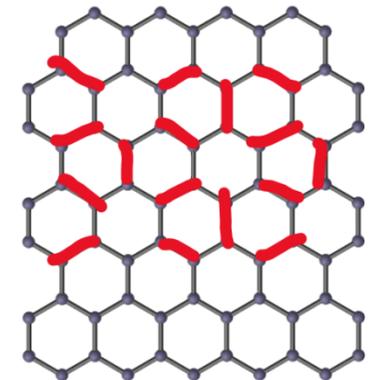
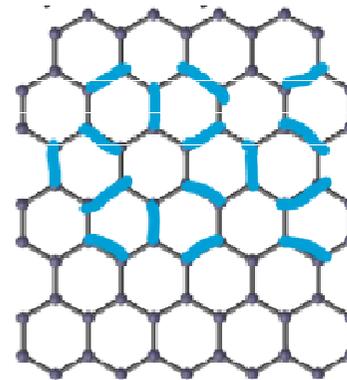
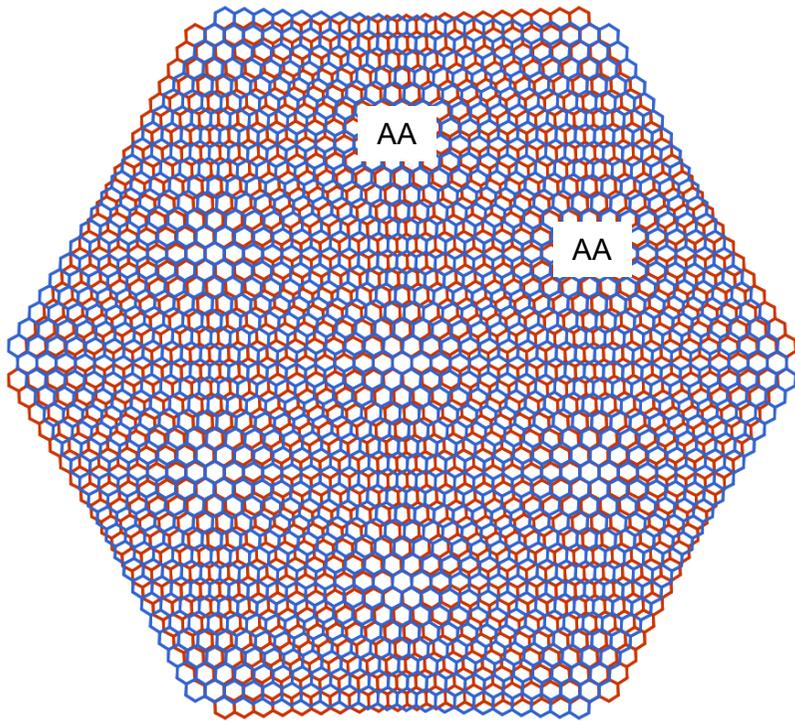


Q: What is the Correlated Insulator (Usually?)



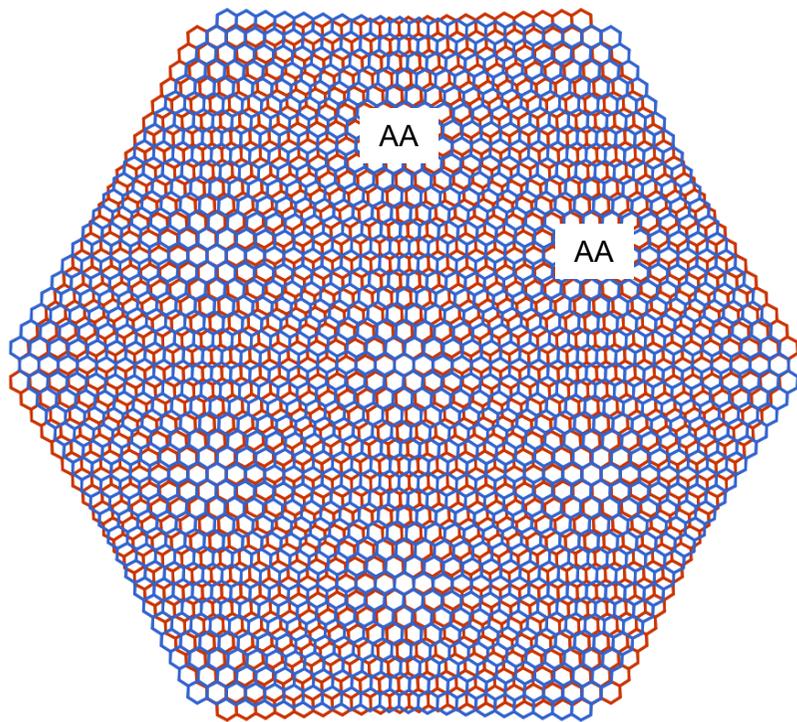
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A: Kekulé spiral!

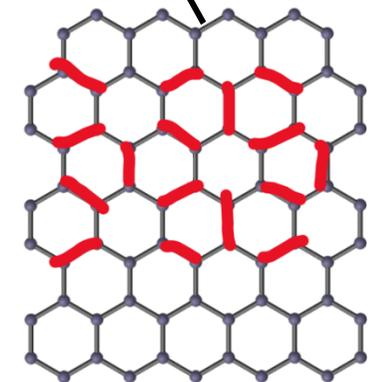
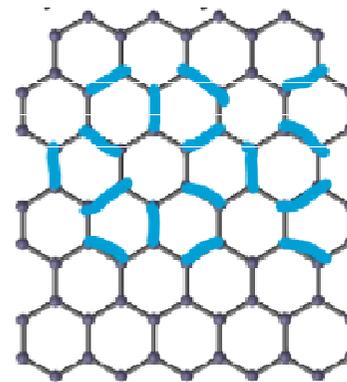
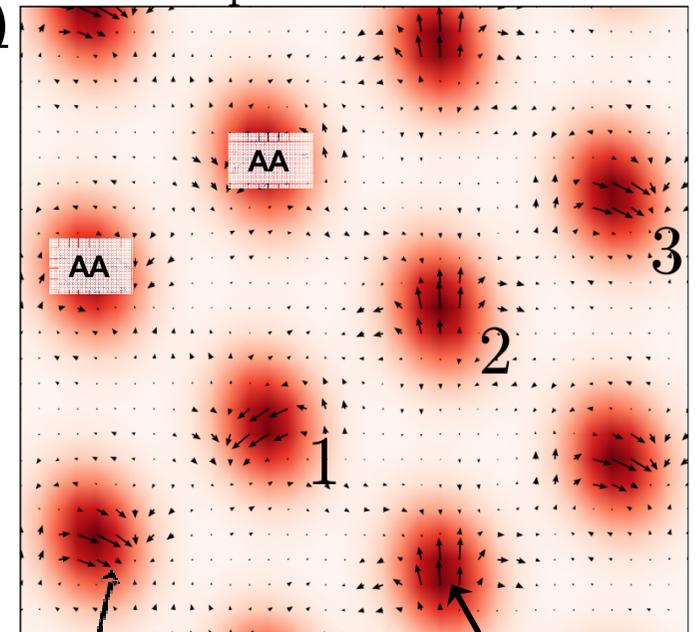


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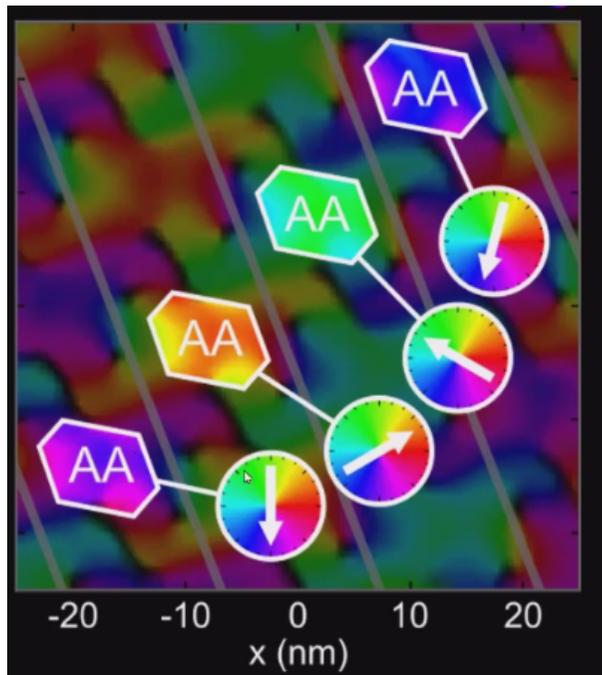
Superlattice scale



Q: What is the Correlated Insulator (Usually?)

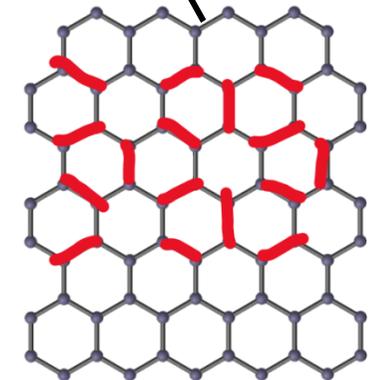
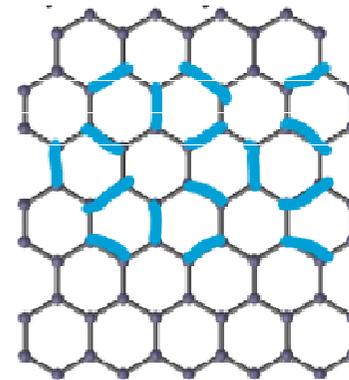
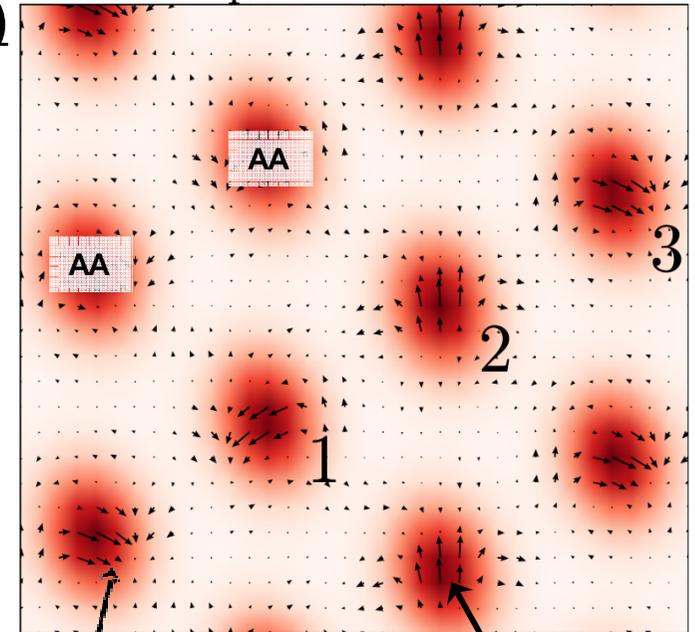
A: Kekulé spiral!

This is real:



(Screenshot from talk by Nuckolls/Yazdani)
Nature 620, 525 (2023) Nuckolls et al (Yazdani)

Superlattice scale

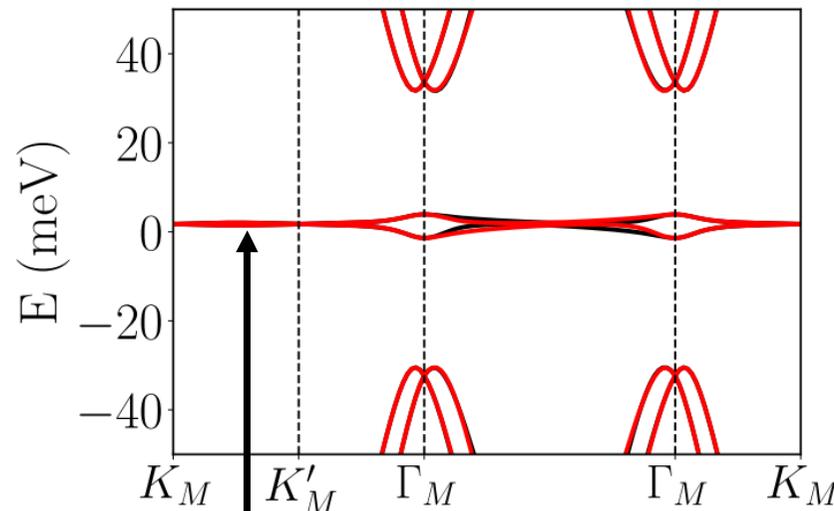


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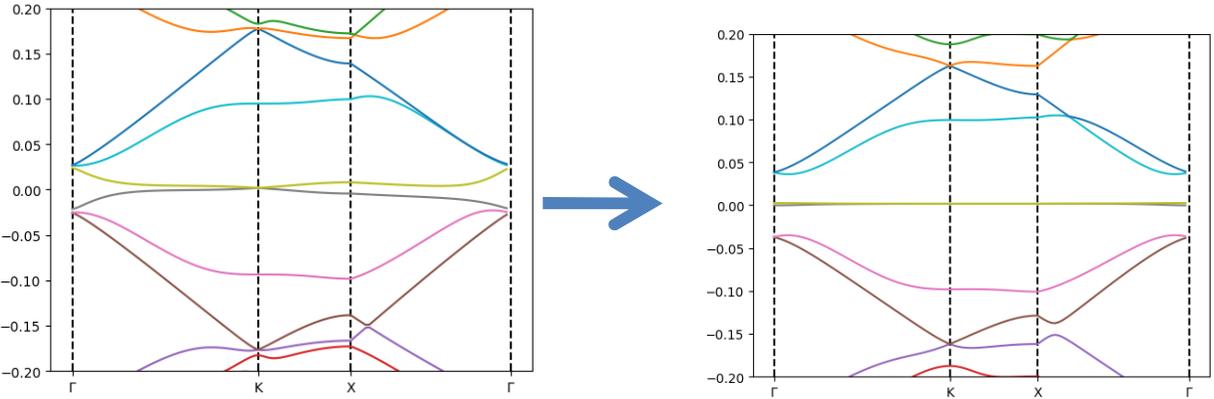
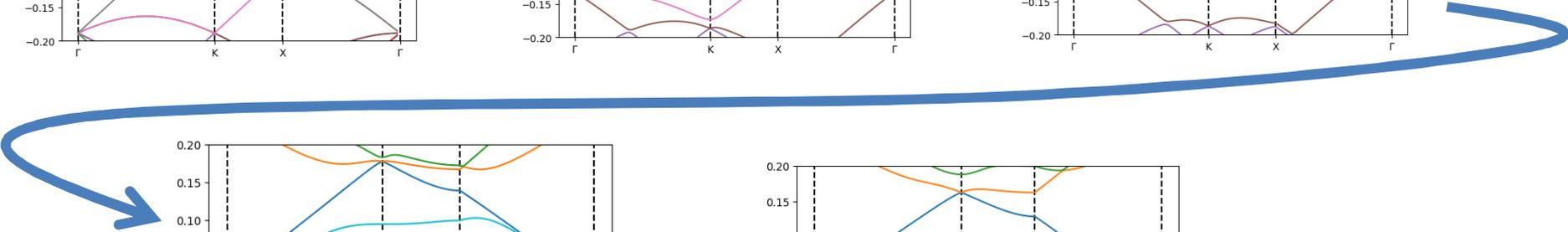
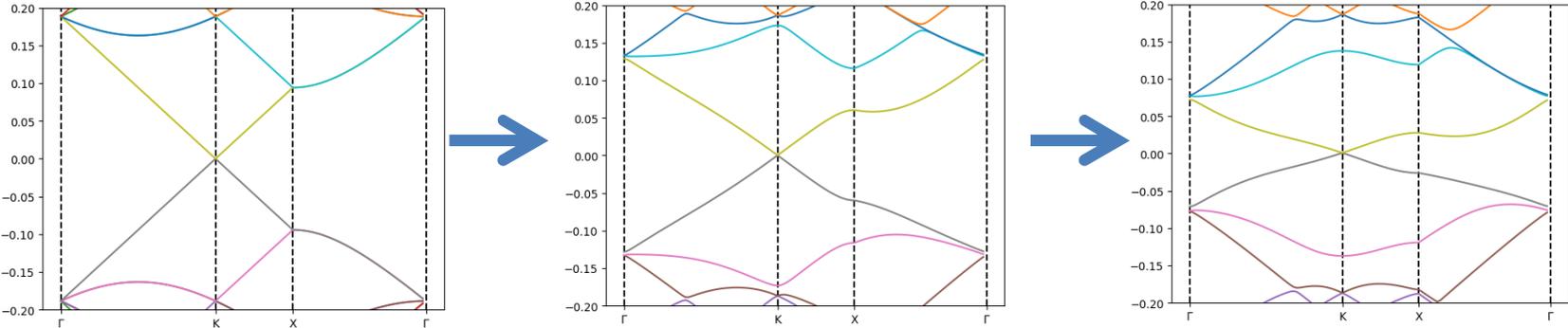
8 central bands
incl. spin and valley

- focus on central bands

filling factor $\nu = -4, \dots, 0, \dots, 4$
(CNP)

Slowly turning on the layer coupling:

No coupling

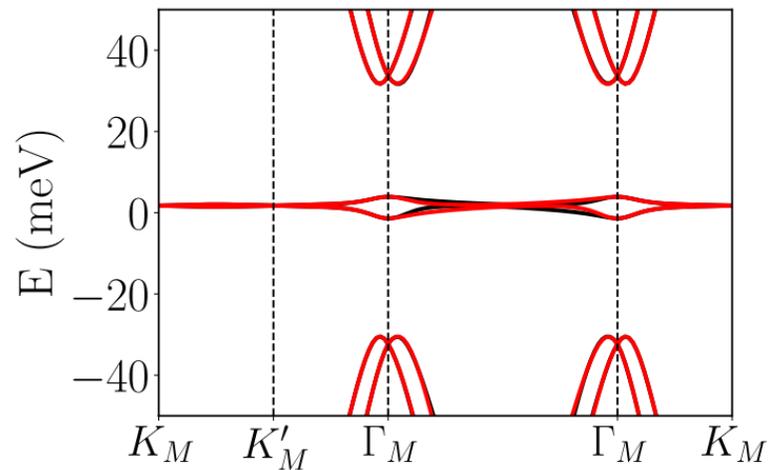


Full coupling

“Chiral” Limit of Bistritzer-MacDonald (BM) Model = TKV Model

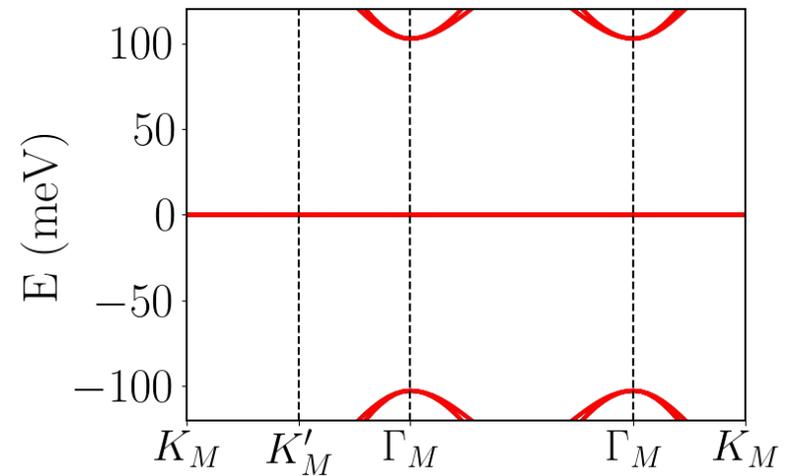
- chiral ratio $\eta = W_{AA}/W_{AB}$

$\eta = 0.8$



chiral limit

$\eta = 0$



Tarnopolsky, Kruchkov, Vishwanath, *PRL* (2019)

Chern Number:



Shiing-Shen Chern

Theorem:

A two-dimensional insulator has a quantized Hall conductivity

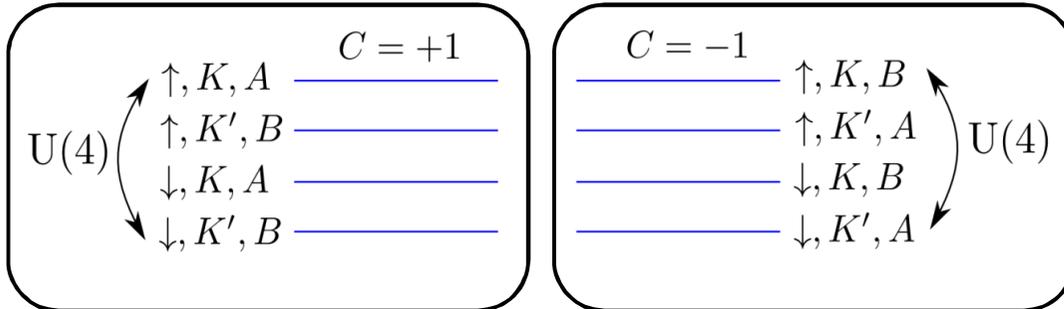
$$G_{xy} = C \frac{e^2}{h}$$

where C is the Chern number.

For non-interacting electrons:
(insulator=filled band and a gap)
 C is integer.

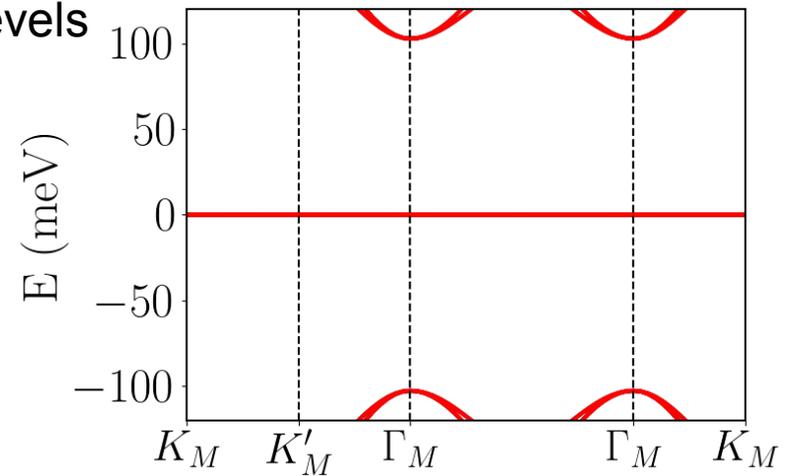
Chiral Limit (Strong Coupling)

Bultinck et al., *PRX* '20
 Bernevig et al., *PRB* '21
 Lian et al., *PRB* '21



Two 'Chern quartets' : like multicomponent Landau levels
Expect flavor ferromagnetism

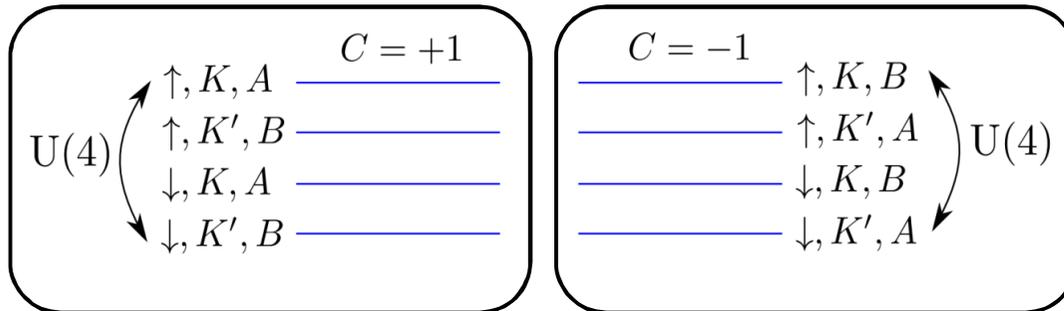
chiral limit
 $\eta = 0$



Tarnopolsky, Kruchkov, Vishwanath, *PRL* (2019)

Chiral Limit (Strong Coupling)

Bultinck et al., *PRX* '20
 Bernevig et al., *PRB* '21
 Lian et al., *PRB* '21



Two 'Chern quartets' : like multicomponent Landau levels

Expect flavor ferromagnetism = Slater determinant ground states

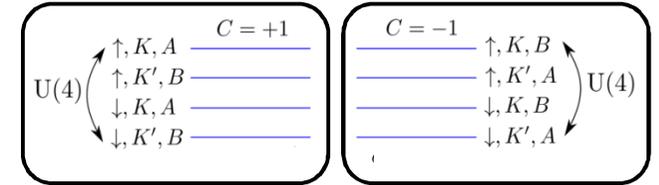
Fill any $4 + \nu$ out of 8 central Chern bands (generalized ferromagnets – can break T)

But can rotate within a Chern-Band, ex, fill $|KA \uparrow\rangle + e^{i\phi}|K'B \downarrow\rangle$
 tilting the pseudospin before filling the band

Superposing K and K' introduces wavevector $|K - K'|$ which triples graphene unit cell

Chiral Limit (Strong Coupling)

Bultinck et al., *PRX* '20
 Bernevig et al., *PRB* '21
 Lian et al., *PRB* '21

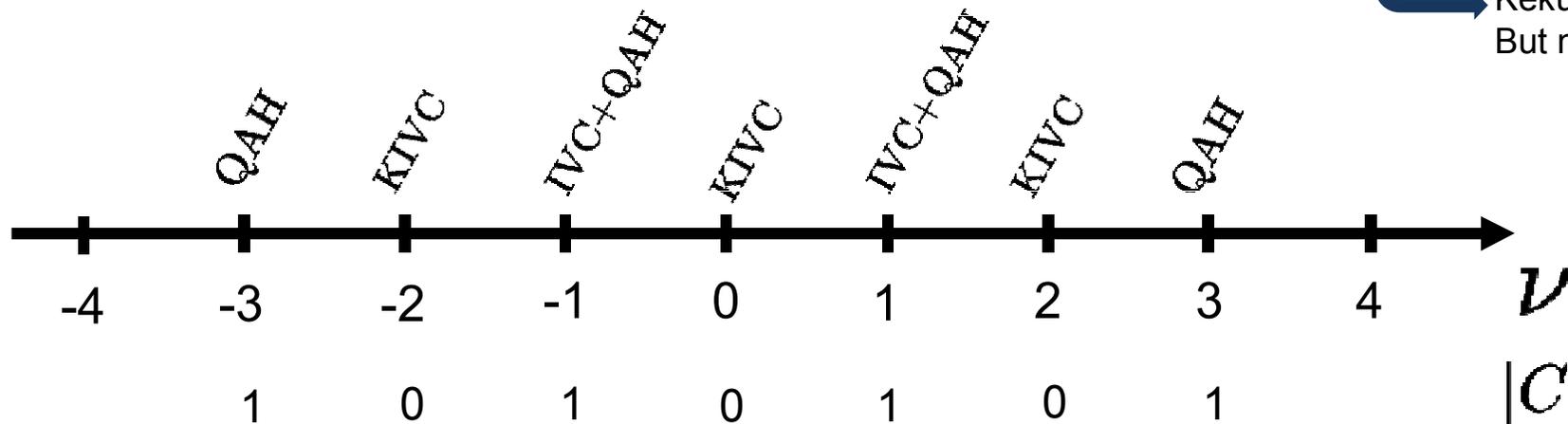


- single-particle terms and deviations from chiral limit treated perturbatively

select out subset of generalized ferromagnets
 i.e. 'strong-coupling' insulators

QAH = Quantum Anomalous Hall
 IVC = Intervalley Coherent
 KIVC = Kramers IVC

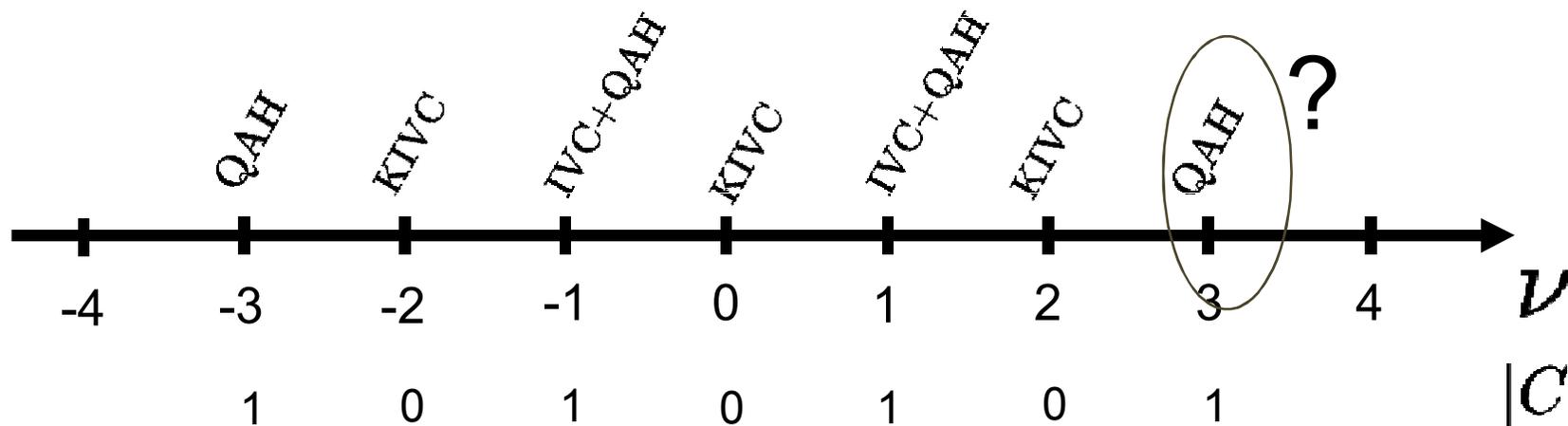
→ Kekule currents
 But not density



Bultinck et al., *PRX* '20
 Xie & MacDonald, *PRL* '20
 Cea & Guinea, *PRB* '20
 Zhang et al., *PRB* '20
 Lian et al., *PRB* '21
 Xie et al., *PRB* '21
 Potasz et al., *PRL* '21
 Parker et al., *PRL* '21
 Kwan et al., *PRX* '21
 ...

Strong Coupling & Experiments

- ✓ correlated insulators at integer ν as flavour ferromagnetism
- ✗ experimentally CNP is often semimetallic, but large-gap robust insulators in strong-coupling
- ✗ experimentally metallic/weak anomalies at $\nu = \pm 1$, but gapped insulators in strong-coupling
- ✗ observation of ‘Chern-odd’ insulators, e.g. $C=0$ at $\nu = 3$ Pierce et al., Nat Phys, '21



Beyond Strong Coupling

“Perturbations”

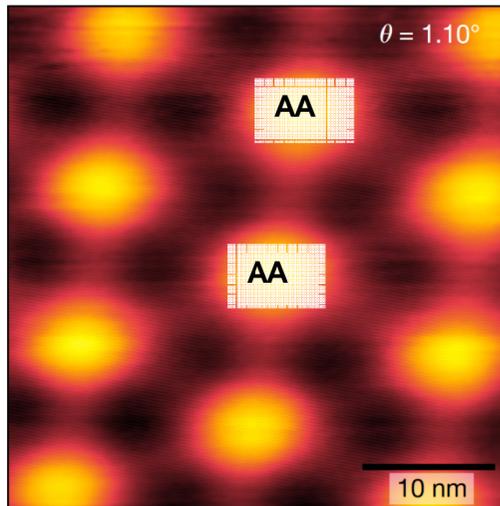
- Not exactly flat band
- p-h sym breaking terms
- substrate potential
- Strain 

Treat these more seriously from the start?

Strain

- uniaxial heterostrain ϵ

Bi, Yuan, and Fu PRB 2019



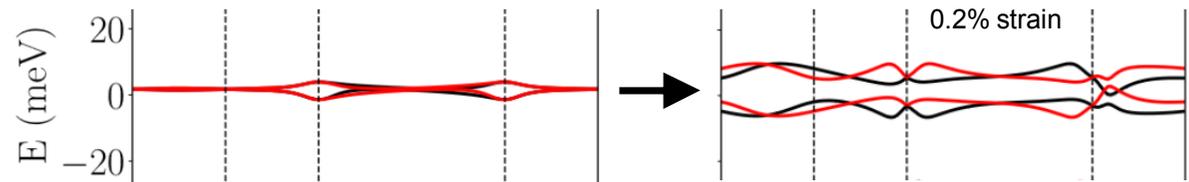
strain is ubiquitous in experiment

typically $\epsilon \simeq 0.1 - 0.7\%$

Xie et al., *Nature* '19

Choi et al., *Nat. Phys.* '19

Kerelsky et al., *Nature* '19



- theoretically known to degrade strong-coupling KIVC gap at charge neutrality (many expts see no gap at $\nu=0$)

Parker et al., *PRL* '21

Results

Kwan et al., *PRX* 2021

substrate potential
↑
strain →

IKS: Incommensurate Kekulé Spiral

QAH: quantized anomalous Hall state

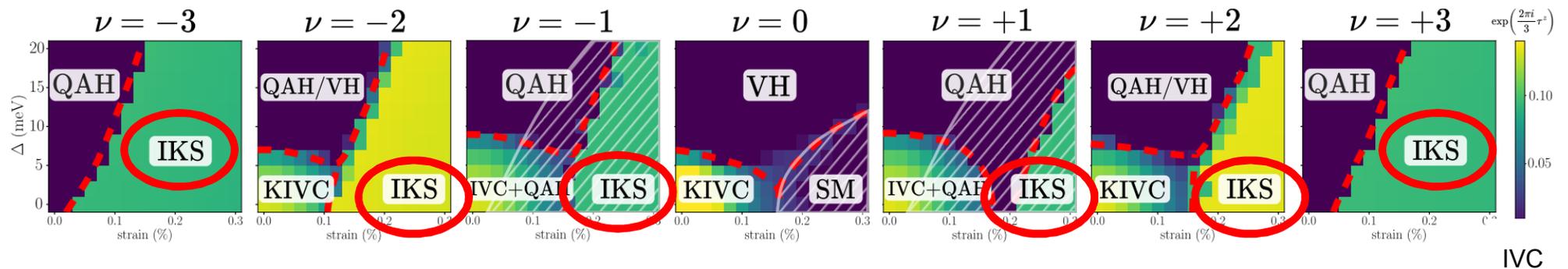
KIVC: Kramers intervalley coherent state

VH: valley Hall state

IVC: intervalley coherence

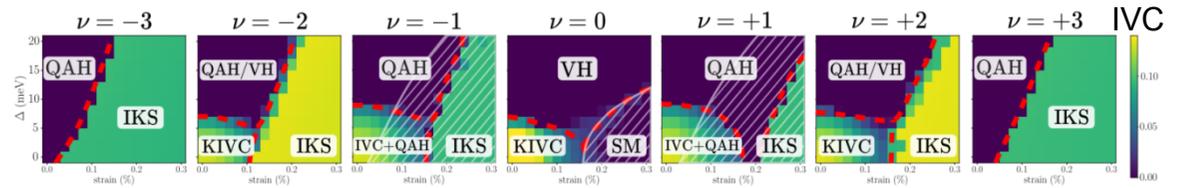
SM: semimetal

Hatched: no HF gap



- self-consistent Hartree-Fock with realistic parameters
- check for completely general translational symmetry-breaking
- **IKS dominates all non-zero integers with strain**

What is IKS? Symmetries



For any integer ν :

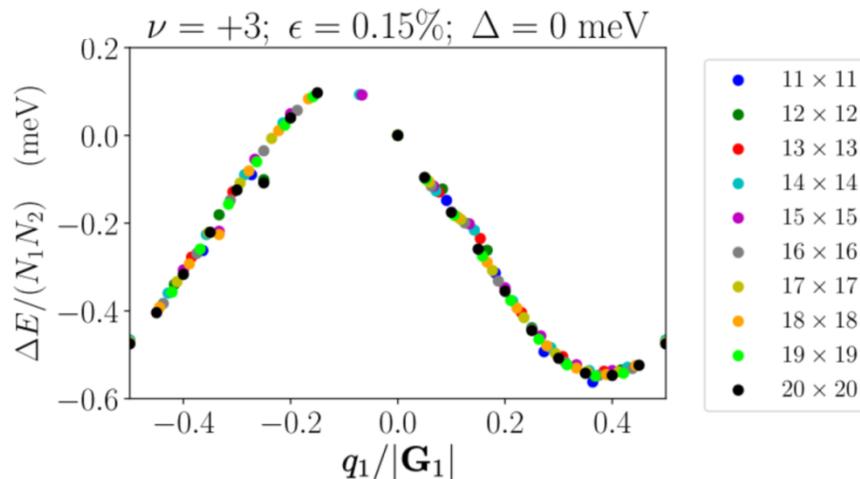
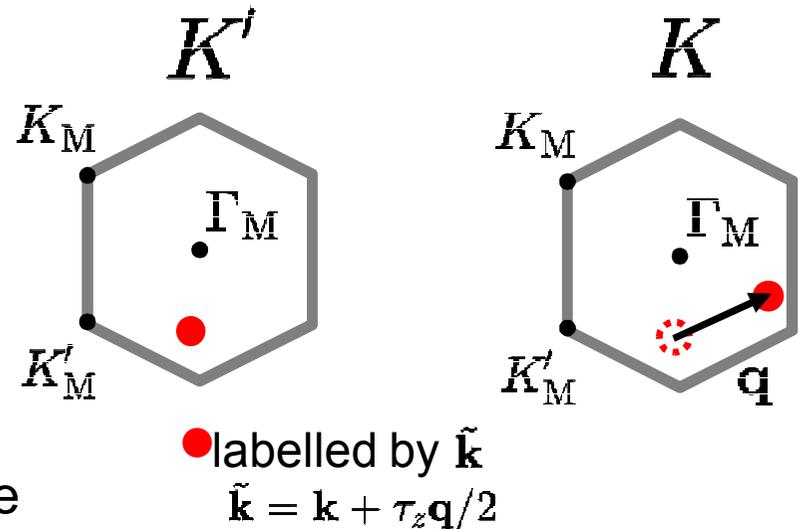
- Superposes K and K' (i.e., Kekule distortion, breaks $U(1)_\nu$)
- breaks translation symmetry $\hat{T}_{\mathbf{a}_i}$ \mathbf{q} is wavevector of the spiral
- preserves modified translation symmetry $\hat{T}'_{\mathbf{a}_i} = \hat{T}_{\mathbf{a}_i} e^{i\mathbf{a}_i \cdot \mathbf{q} \tau_z / 2}$
 - combination of moiré translation and valley rotation

What is IKS? \mathbf{q}

- valleys hybridize at finite wavevector \mathbf{q} :

$$u_{\mathbf{k}}|K, \mathbf{k} - \mathbf{q}/2\rangle + v_{\mathbf{k}}|K', \mathbf{k} + \mathbf{q}/2\rangle$$

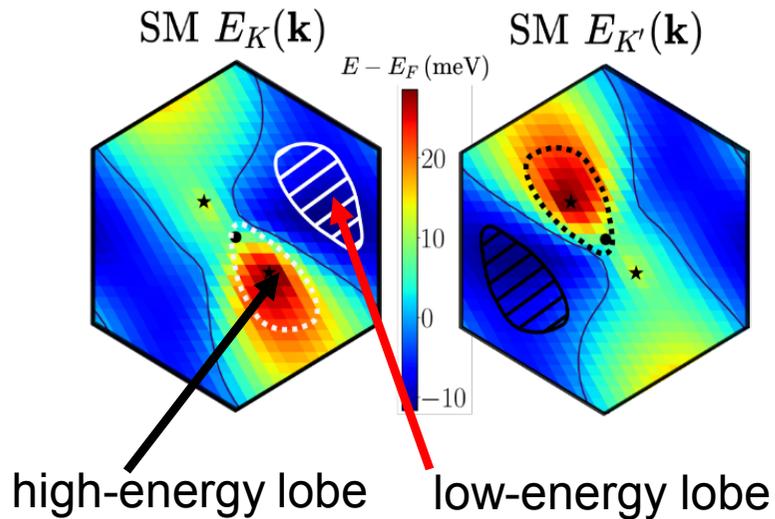
- \mathbf{q} is very soft, and generally incommensurate



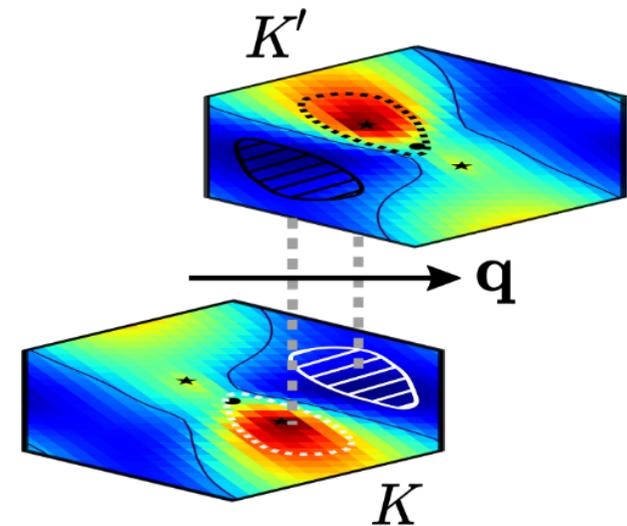
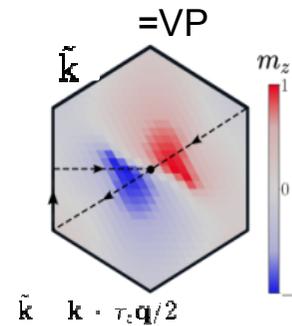
IKS: Incommensurate Kekulé Spiral

Physical mechanism for IKS

consider HF bands of symmetric metal



- Align high-energy lobe in one valley with low-energy lobe in other valley to find \mathbf{q} . IVC everywhere else.



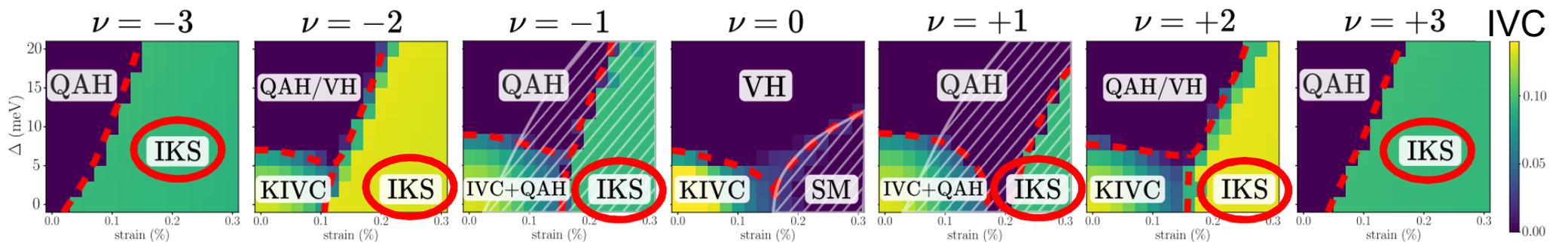
Strain increases dispersion \rightarrow Strong Coupling Ferromagnetism Fails!

Connections of IKS/HF to experiment

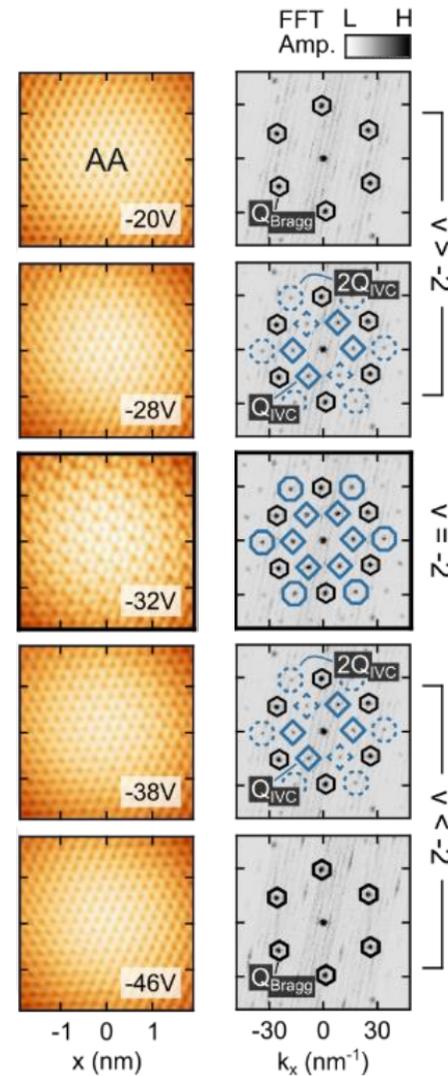
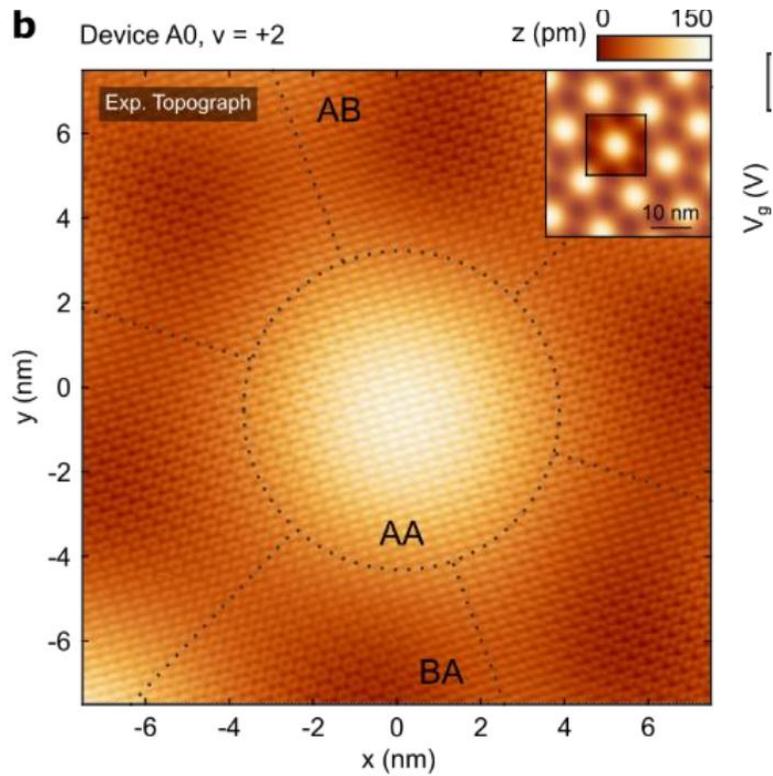
- ✓ strain significantly degrades gap and induces semimetallic behaviour at CNP
- ✓ IKS at $\nu = \pm 1$ is a 'near-insulator'
- ✓ gapped IKS at $\nu = \pm 2$ is spin-unpolarized Yankowitz et al., *Science* '19
- ✓ gapped IKS at $\nu = \pm 3$ has $C=0$ Pierce et al., *Nat Phys*, '21
- ✓ IKS emerges at strain ratios well within experimental limits

✓ KEKULÉ SPIRAL OBSERVED IN STM!

Nature 620, 525 (2023) Nuckolls et al (Yazdani)



Kwan et al., *PRX* '21

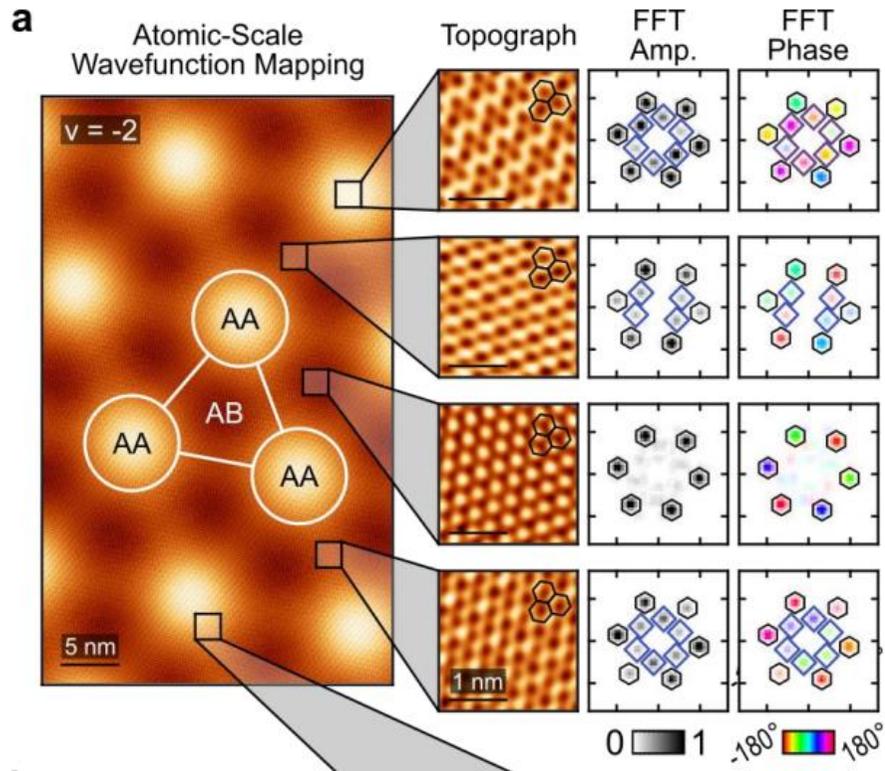


Also Kekulé distortion

Kekulé distortion =
Intervalley Coherence,
.....**BUT NOT KIVC!**

Hong et al *PRL* 22
Calugaru et al *PRL* 22

“Quantum textures of the many-body wavefunctions in magic-angle graphene”



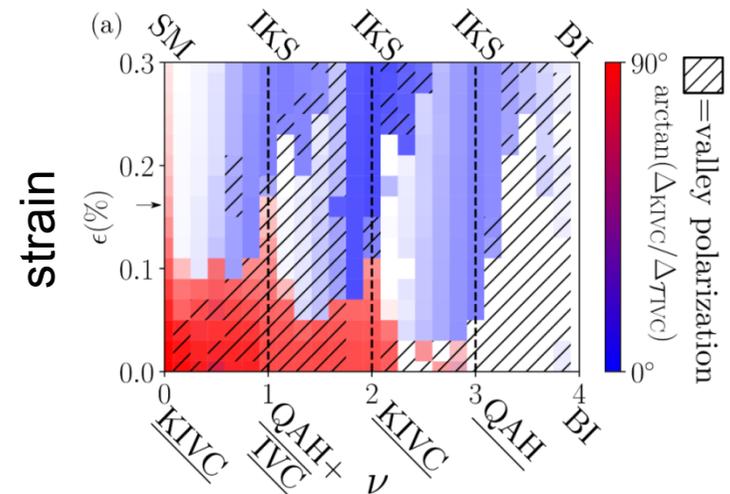
Kekule Bragg Peaks seen almost everywhere, but pattern changes

Analysis by phases of FFT peaks

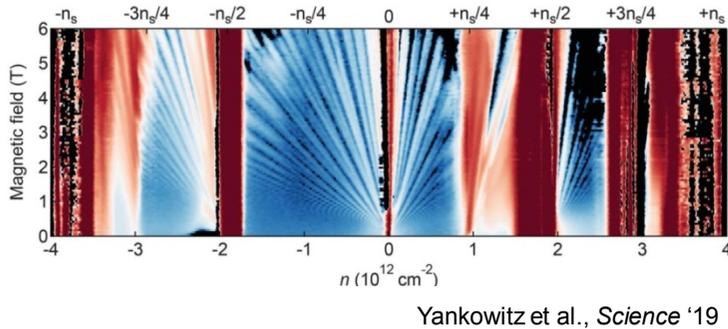
Non-integer fillings

Wagner et al., *PRL* '22

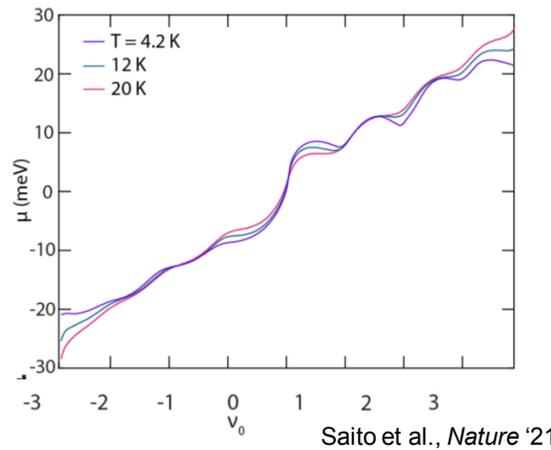
- IKS persists for a large range of dopings at finite strain



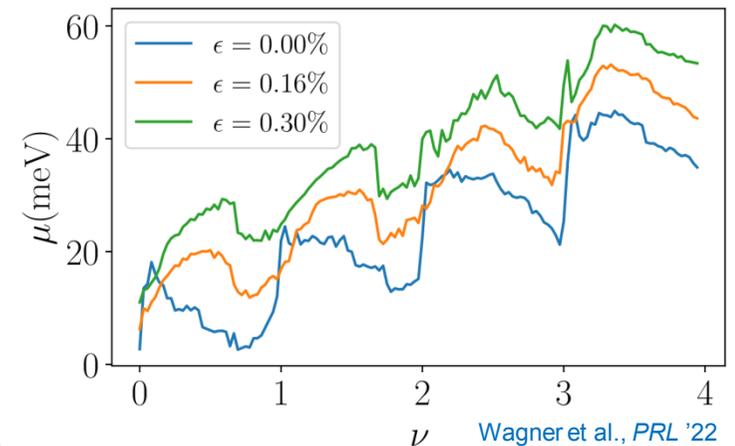
- ✓ Landau fan degeneracy at finite strain is consistent with experiments
4,-,2,1 degeneracy at fillings 0,1,2,3
- ✓ 'cascade' transitions consistent with experiment Zondiner et al., *Nature* '20



Yankowitz et al., *Science* '19

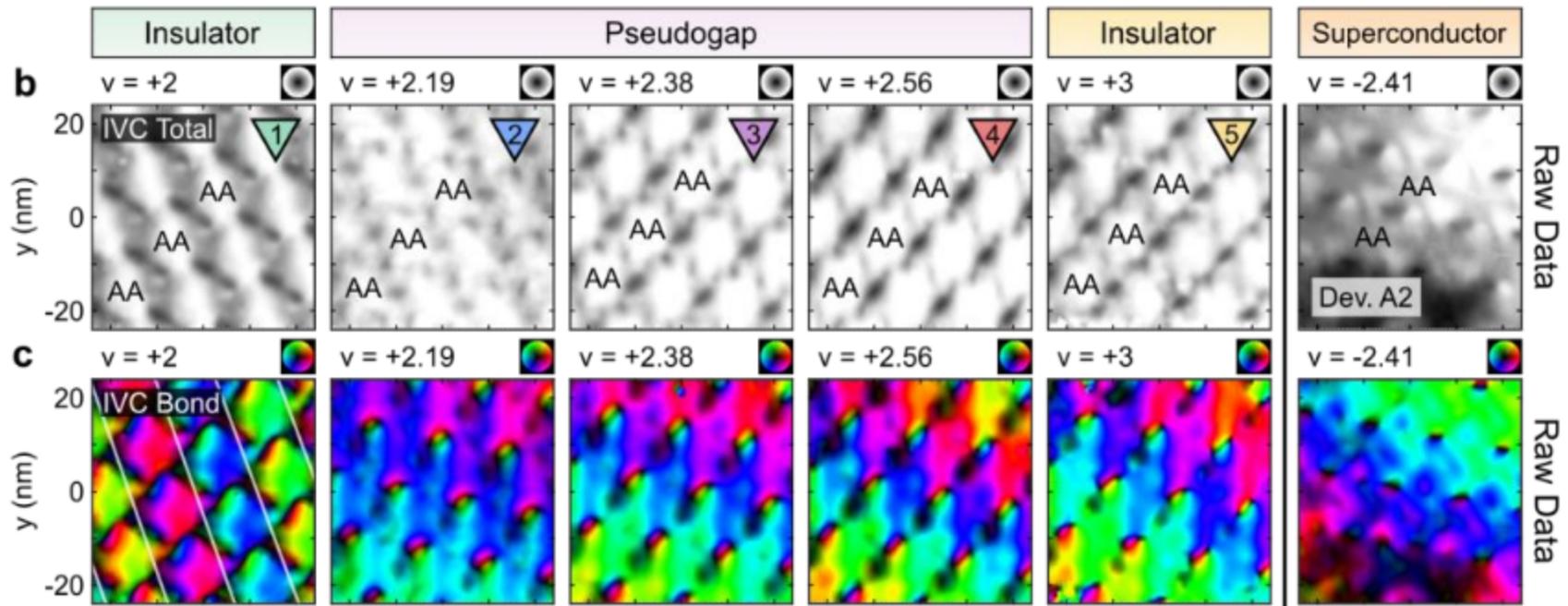


Saito et al., *Nature* '21



Wagner et al., *PRL* '22

Non-integer fillings

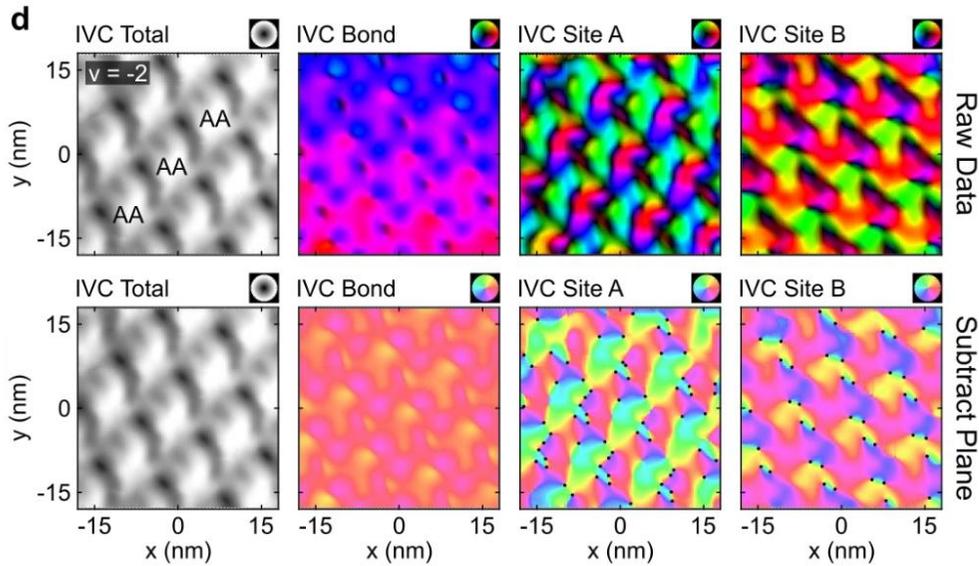


IKS-ish?

What about highly unstrained samples?

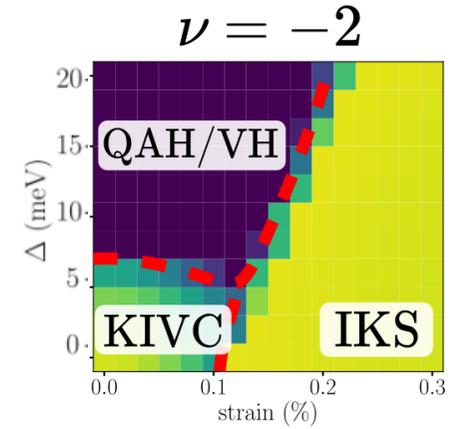
Hartree-Fock says KIVC

--- should show no Kekulé density pattern

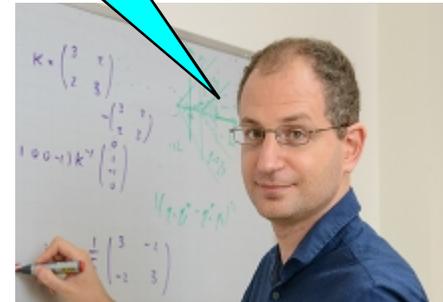


Nature 620, 525 (2023) Nuckolls et al (Yazdani)

But data sees Kekulé density pattern ($q=0$)



Phonons!



Electron-phonon coupling and competing Kekulé orders in twisted bilayer graphene

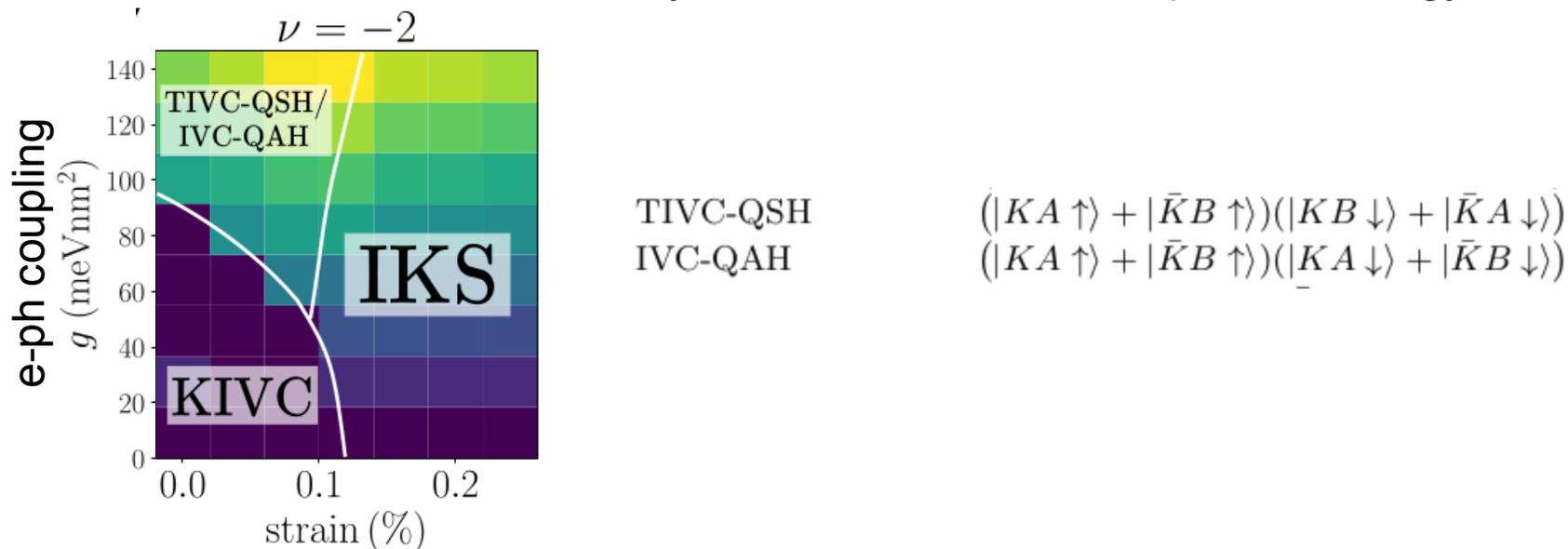
Yves H. Kwan,¹ Glenn Wagner,² Nick Bultinck,^{3,4} Steven H. Simon,³ Erez Berg,⁵ and S.A. Parameswaran³

Phys. Rev. B **110**, 085160 (2024)

$$\hat{H} = \epsilon a^\dagger a + \gamma a + \gamma^* a^\dagger \quad \longrightarrow \quad \hat{H} = \epsilon b^\dagger b - \frac{|\gamma|^2}{\epsilon}$$

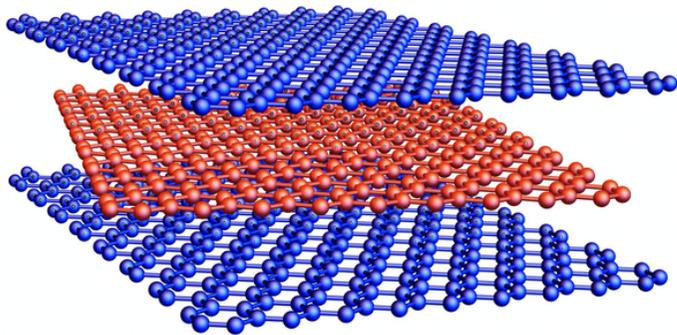
$$b^\dagger = a^\dagger + \gamma/\epsilon$$

An electron state with density modulation lowers total phonon energy.



“Imaging inter-valley coherent order in magic-angle twisted *trilayer* graphene”

Kim et al (Nadj-Perge) Nature 693 942, 2023

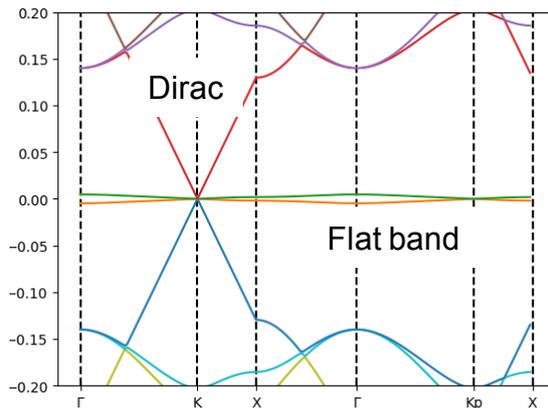


Twisted trilayer graphene: 1st and 3rd layers aligned

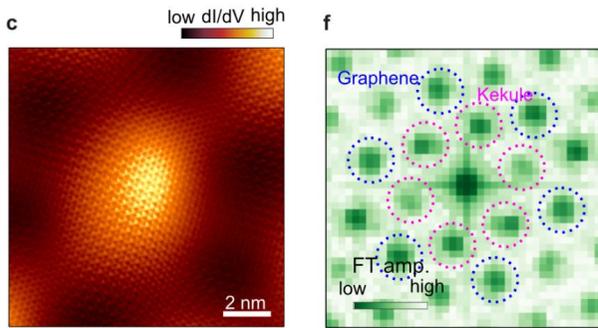
Noninteracting spectrum = Twisted Bi-Layer graphene + One Dirac Cone left over

symmetric combination
of layers 1 and 3 couples
to layer 2.

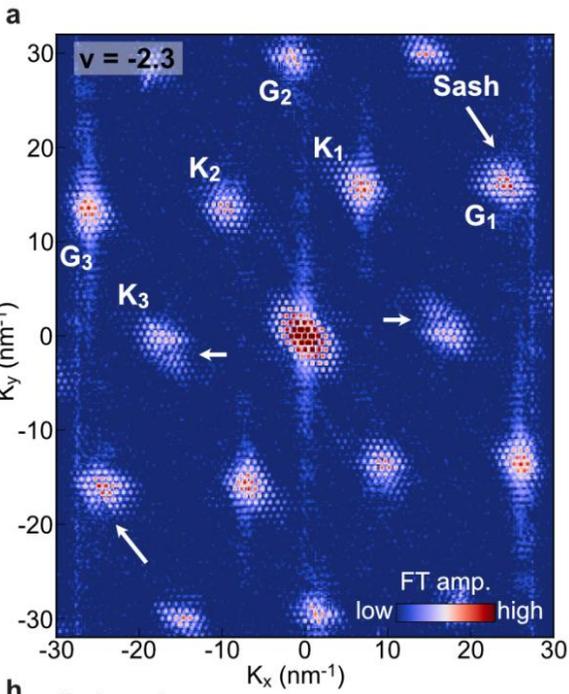
antisymmetric combination
of layers 1 and 3 decouples



How different is this from bilayer?

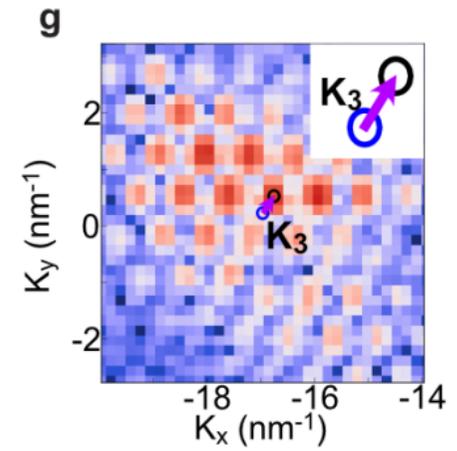
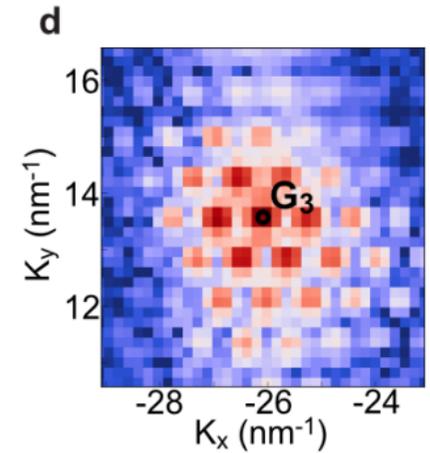


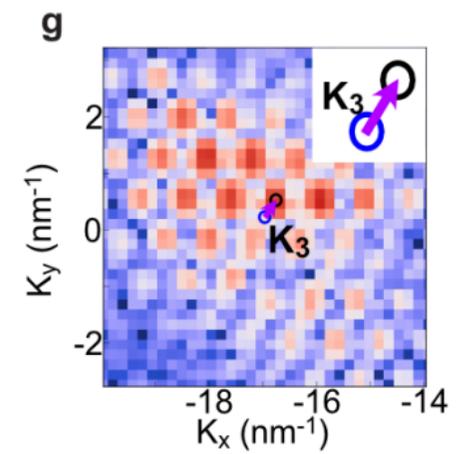
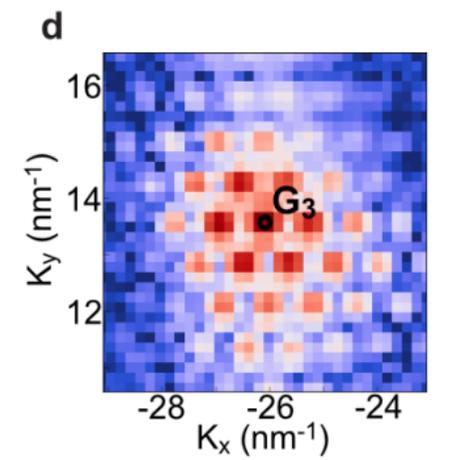
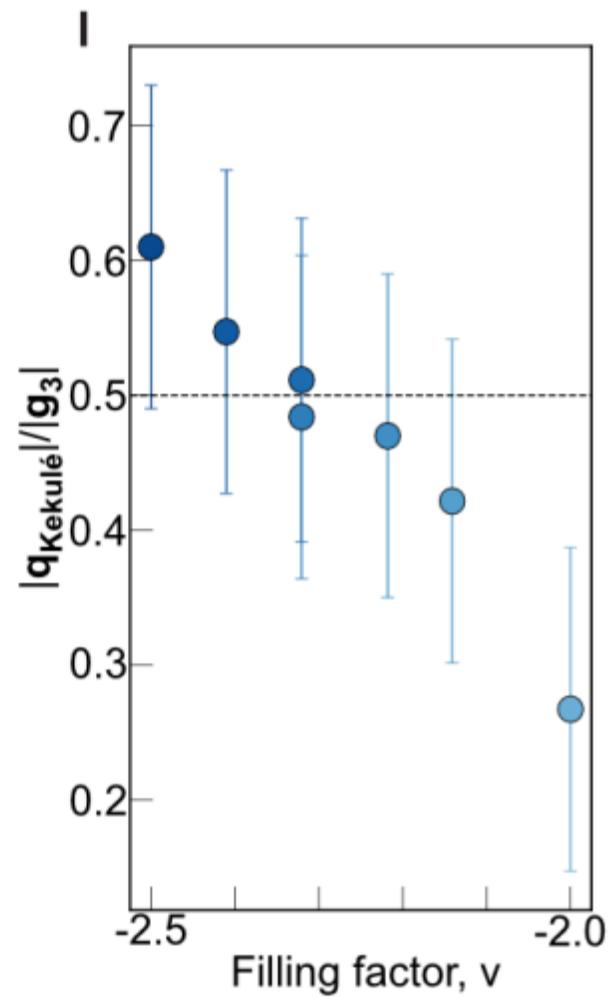
Looking for an additional wavevector q_{IKS}



Additional periodicity From Moire

Displacement of pattern at G compared to pattern at K gives another wavevector!





Z. Wang et al: Arxiv 2310.16094:

Kekule spirals and charge transfer cascades in twisted symmetric trilayer graphene

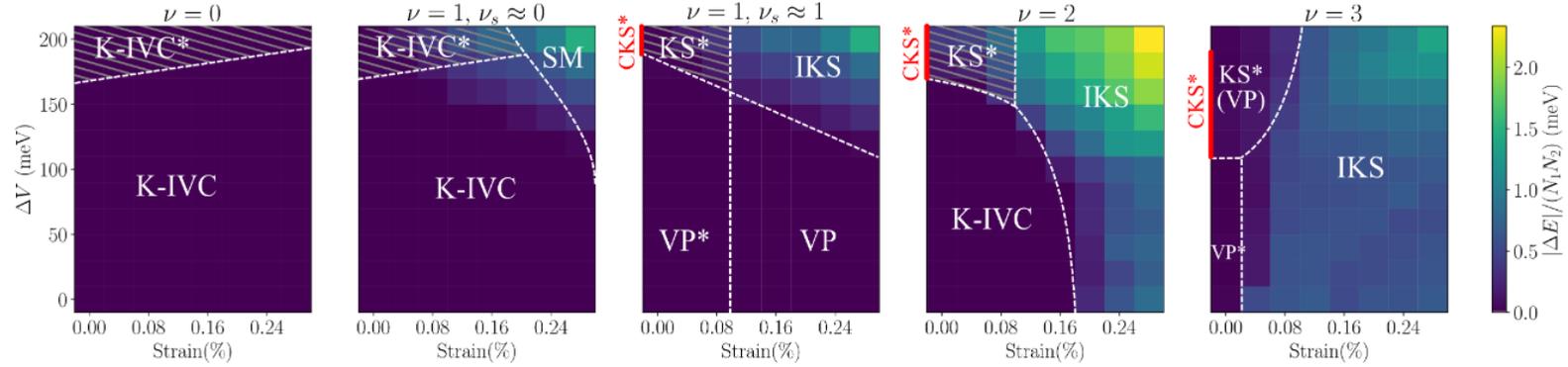


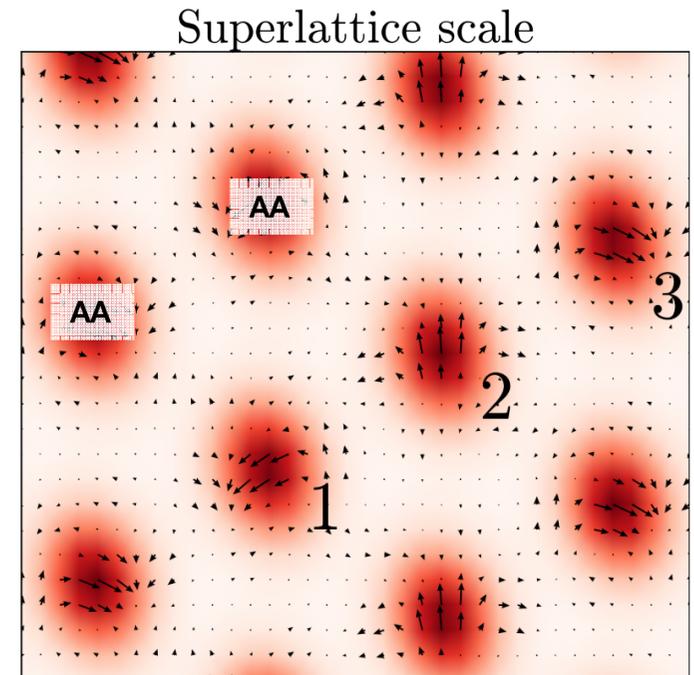
FIG. 2. Phase diagrams as a function of strain and interlayer potential at integer fillings ν . We have performed 10×10 self-consistent Hartree-Fock calculations, scanning over all possible Kekulé spiral vectors \mathbf{q} . ΔE is the energy difference between the $\mathbf{q} = 0$ state with the overall energy minimum. The acronyms are: K-IVC for Kramers-intervalley coherence, VP for valley polarization, (I)KS for (incommensurate) Kekulé spiral, SM for (compensated) semi-metal. Asterisk (*) denotes states that break C_2T -symmetry. The KS*(VP) region in $\nu = 3$ consists of Kekulé spiral states with possible finite valley polarization. The red lines on the zero-strain axis indicate *commensurate* Kekulé spirals. The shaded regions have a finite charge gap. For $\nu = 1$, we have shown the phases and charge gaps of the two spin sectors separately, where ν_s denotes the filling of each spin sector. As ΔE cannot be separately defined for each spin sector, the data are simply duplicated in the two plots of $\nu = 1$.

Summary:

IKS: Incommensurate Kekulé Spiral

It is there in “typical” samples

Explains a lot.... how much more will it explain?



Kekulé Spirals in Twisted Bilayer Graphene

Thank you
for listening!

Steve Simon
Oxford University

Yves Kwan



→ Princeton

Glenn Wagner



→ Zurich

Nick Bultinck



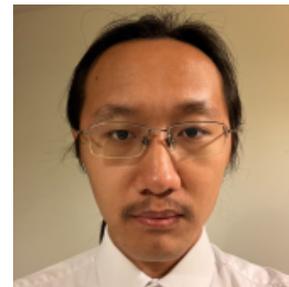
→ Ghent

Sid Parameswaran



Oxford

Ziwei Wang



Oxford



[Kwan et al PRX 11, 041063 \(2021\)](#); arXiv:2303.13602; PRX 12, 031020 (2022); ...
Wagner et al, PRL **128**, 156401 (2022). Wang et al, arxiv 2310.16094