

Quantum simulations with superconducting qubits

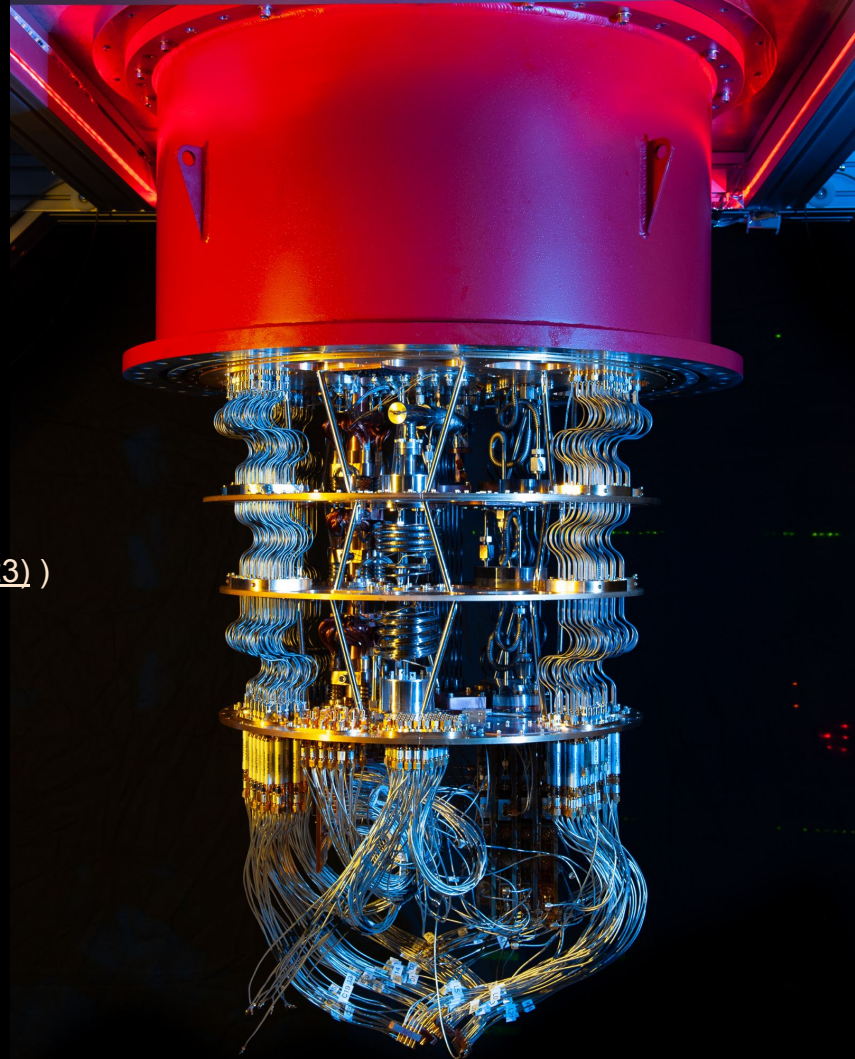
pedramr@google.com, @PedramRoushan



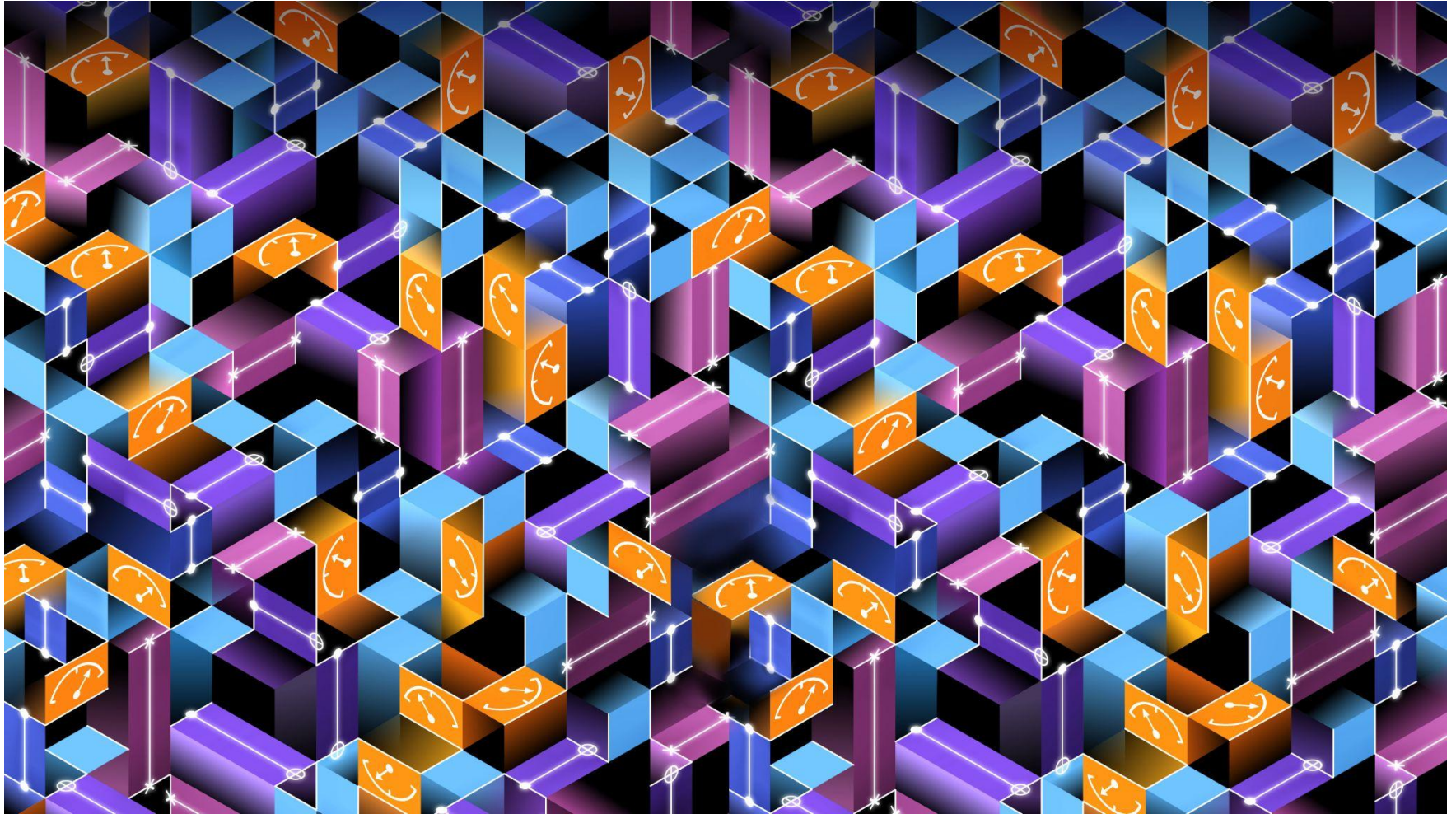
Google Quantum AI,

Measurement-induced Phases of quantum matter
Measurement-induced entanglement and teleportation (Nature 622 (2023))

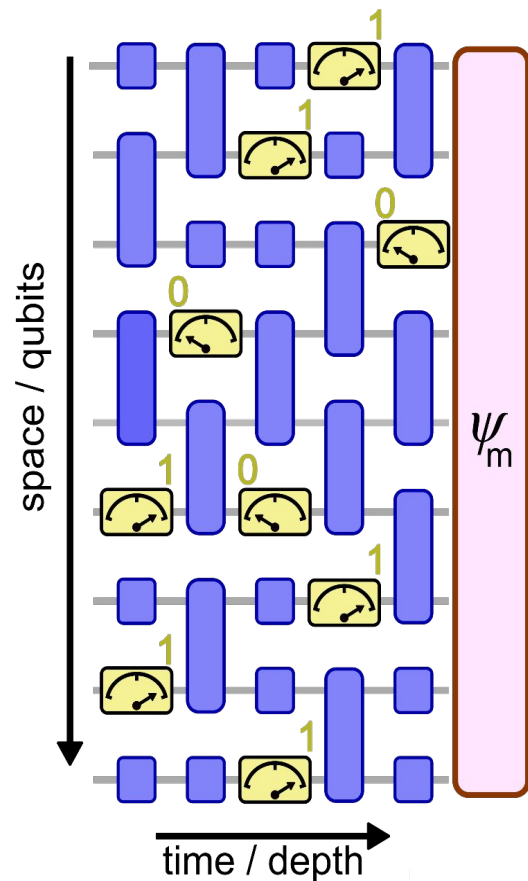
Quantum connection , Stockholm 2025



Quantum information phases in space-time



Measurement-induced entanglement and teleportation



Measurement is key in many protocols

A genuine NISQ problem
Can noise destabilize phases



Jesse Hoke
(Stanford)



Elliott Rosenberg
(Cornell U)



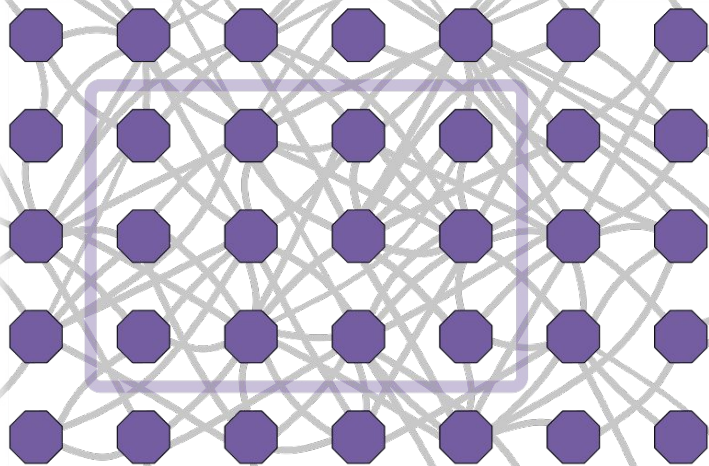
Vedika Khemani
Stanford



Matteo Ippoliti
Stanford

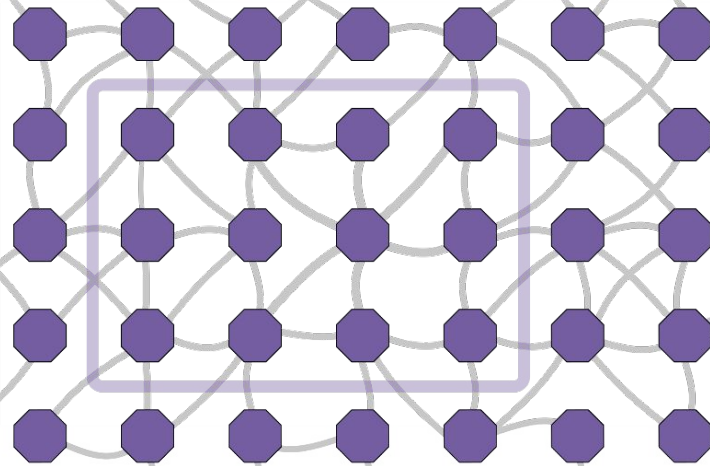
Common entanglement structures

Highly-entangled state



e.g. thermalizing state, $S \sim l^d$

Low-entangled state

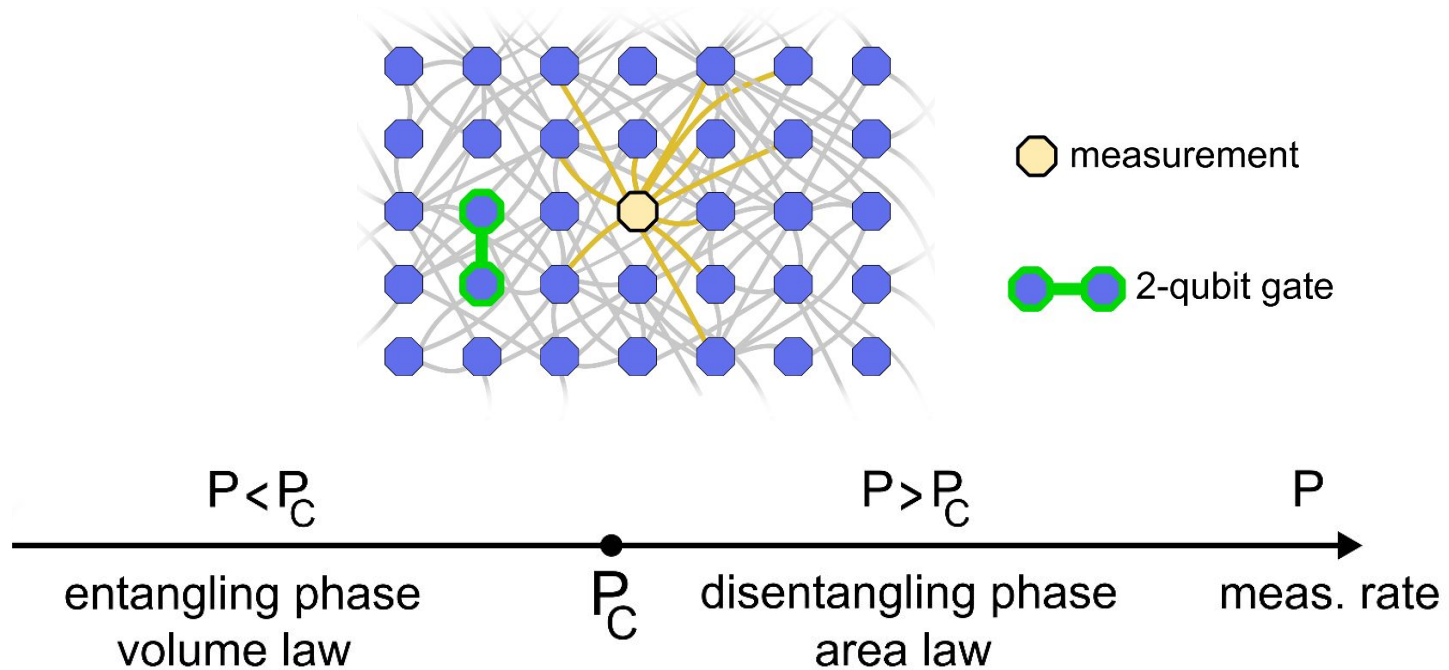


e.g. MBL, $S \sim l^{d-1}$

What happens if we add mid-circuit measurements ?

Physics folklore of circa 2017: Entangled states are very fragile, entanglement would be no match for measurement.

Phase transition in quantum information



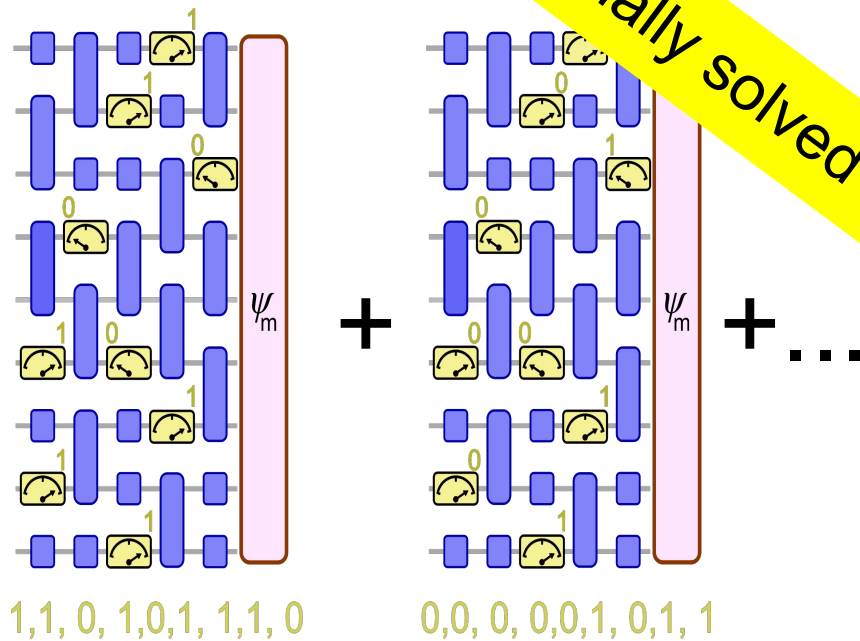
“Vandalized resistor grid”

Cao, Tilloy, De Luca, SciPost Physics (2019)

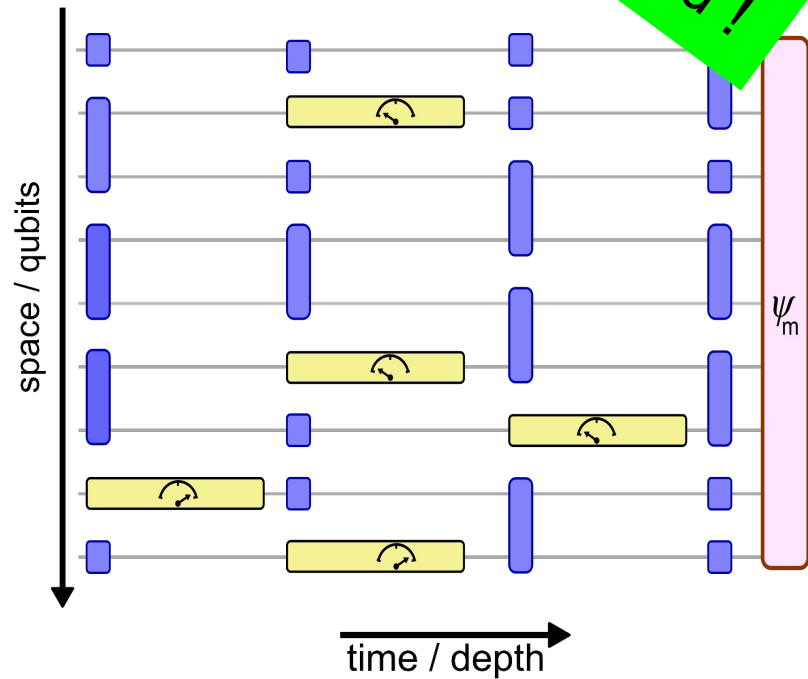
Skinner, Ruhman, Nahum, PRX (2019), Li, Chen, Fisher, PRB (2018), Chan, Nandkishore, Pretko, Smith, PRB (2019)

Challenges in studying monitored circuits

1. Post-selection: staying on a specific trajectory

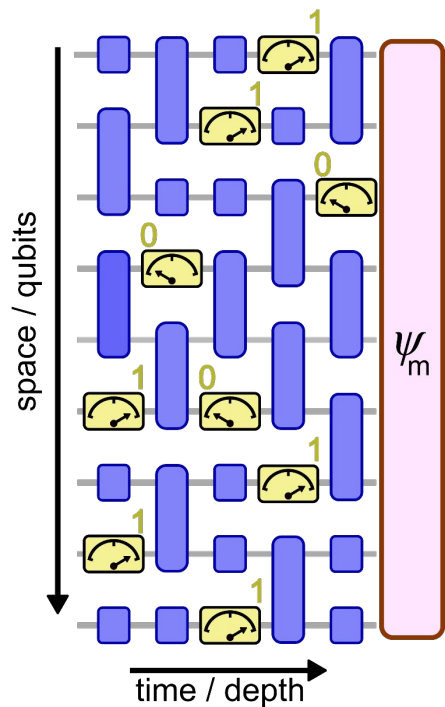


2. Mid-circuit measurement

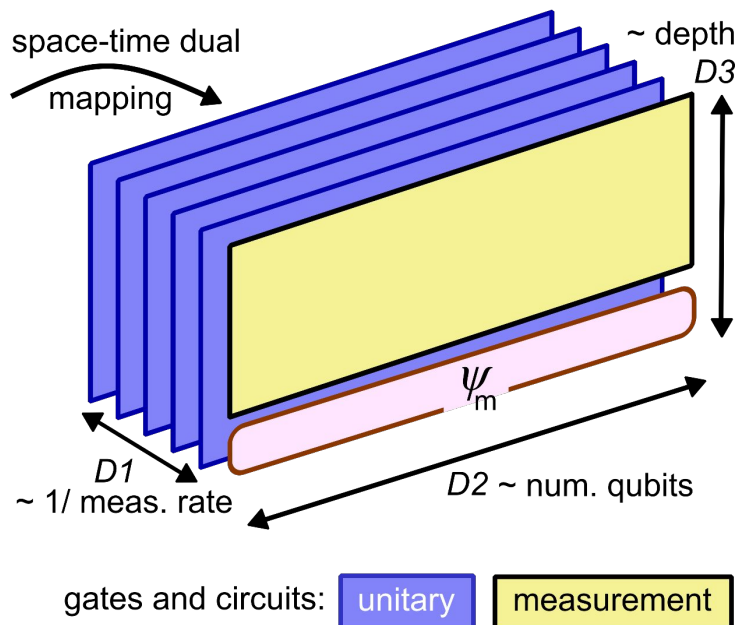


$$\sum_m \langle \psi_m | \hat{C} | \psi_m \rangle = \sum_m \text{Tr} \left(|\psi_m\rangle \langle \psi_m| \hat{C} \right) = \text{Tr} \sum_m |\psi_m\rangle \langle \psi_m| \hat{C} = \text{Tr}(\rho_{\text{ave}} \hat{C})$$

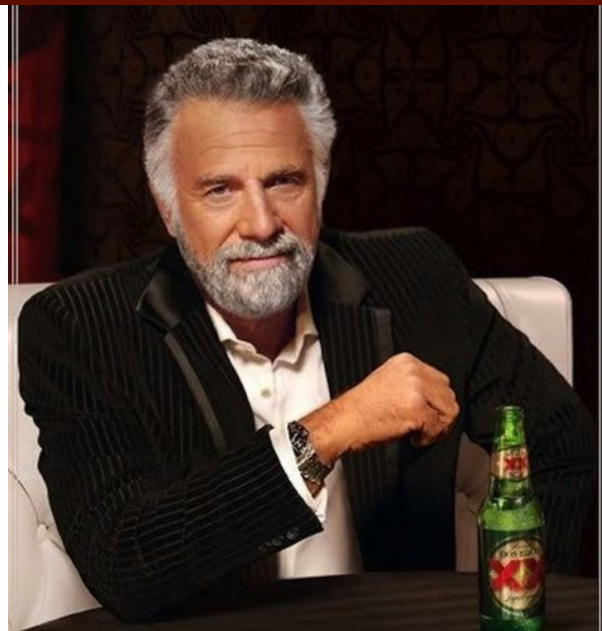
circuit for monitored evolution



equivalent dual circuit



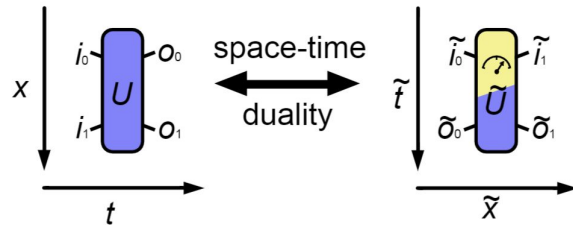
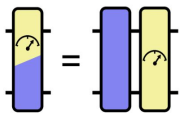
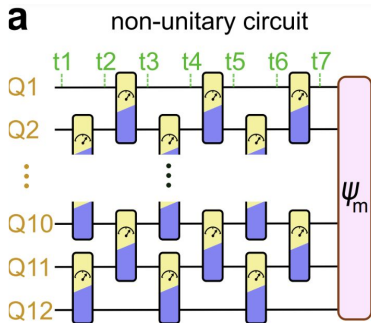
I usually do not study monitored circuits



But when I do, I use space-time duality

Absence of causality : “arrow of time” loses its unique role
 → network of quantum information in space-time

Implementation of space-time duality in 1D



$$U = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \theta & -i \sin \theta & 0 \\ 0 & -i \sin \theta & \cos \theta & 0 \\ 0 & 0 & 0 & e^{-2i\theta} \end{pmatrix} \mapsto \tilde{U} = \begin{pmatrix} 1 & 0 & 0 & \cos \theta \\ 0 & 0 & -i \sin \theta & 0 \\ 0 & -i \sin \theta & 0 & 0 \\ \cos \theta & 0 & 0 & e^{-2i\theta} \end{pmatrix}$$

$\tilde{U} = 2VH$ where V is unitary and H is Hermitian

$$H = \frac{1}{2} \sqrt{1 + 3 \cos^2(\theta)} |\psi_\theta\rangle\langle\psi_\theta| + \frac{1}{2} |\sin(\theta)| (I - |\psi_\theta\rangle\langle\psi_\theta|)$$

$$|\psi_\theta\rangle \equiv \frac{1}{\sqrt{2}} (e^{-i\theta} |00\rangle + e^{i\theta} |11\rangle)$$

$$V' = \begin{pmatrix} 0.6487 - 0.3993i & 0.4935 + 0.4198i \\ 0.4935 + 0.4198i & 0.2901 - 0.7043i \end{pmatrix} \quad (\theta = \pi/10)$$

Implementation of space-time duality in 1D

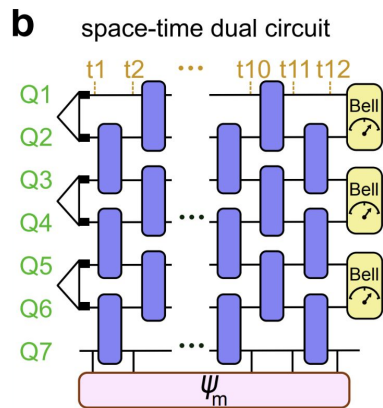
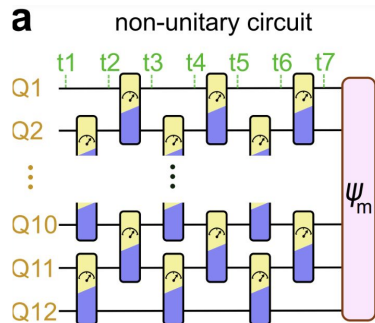
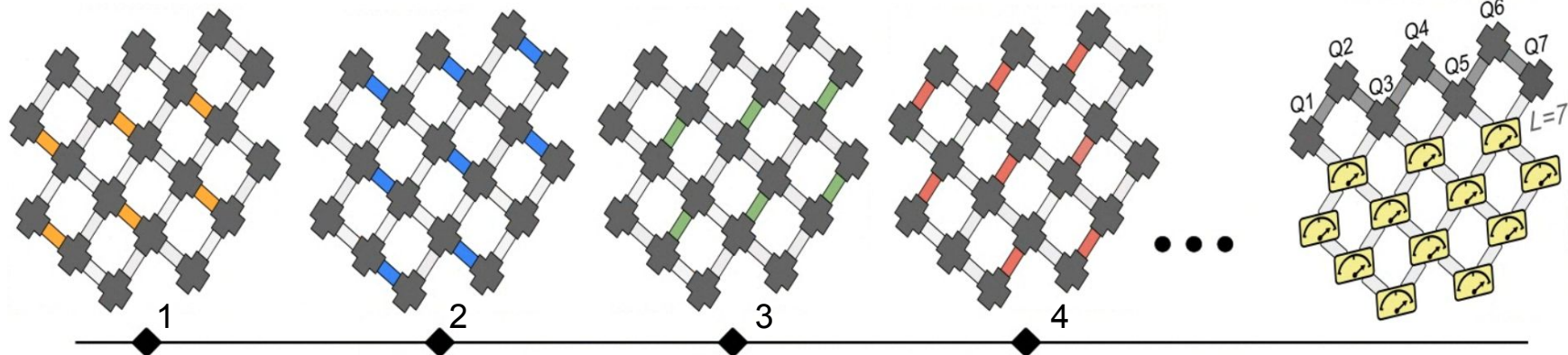


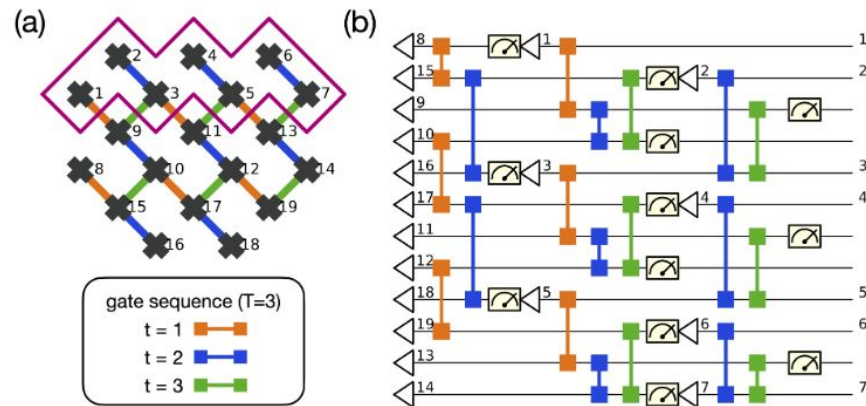
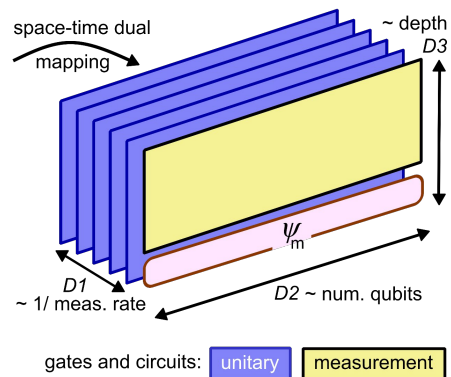
Diagram illustrating the space-time duality transformation. The transformation maps a gate with a meter to a sequence of two gates (blue and yellow). This is then transformed via "space-time dual" into a single blue gate, which is then shown to be equivalent to a $f\text{Sim}(\theta, \phi)$ gate followed by two Z^n gates.

$$f\text{Sim}(\theta, \phi) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos\theta & i\sin\theta & 0 \\ 0 & i\sin\theta & \cos\theta & 0 \\ 0 & 0 & 0 & e^{i\phi} \end{pmatrix}$$

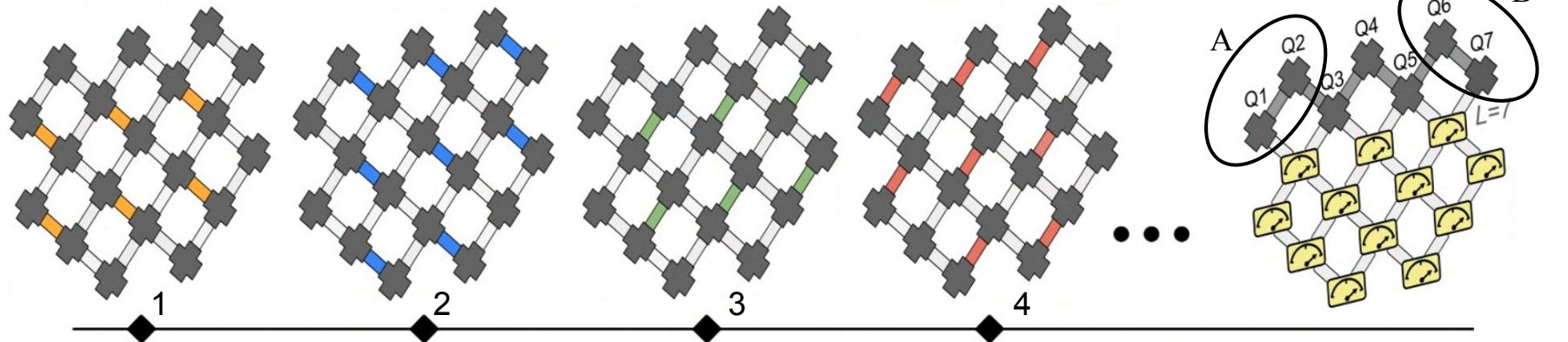
1D entanglement phases from 2D shallow quantum circuits



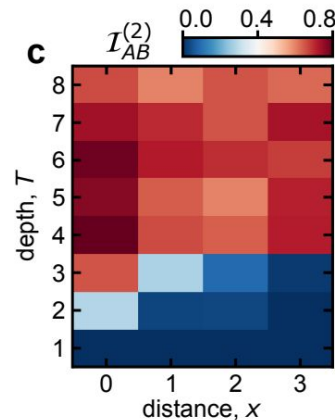
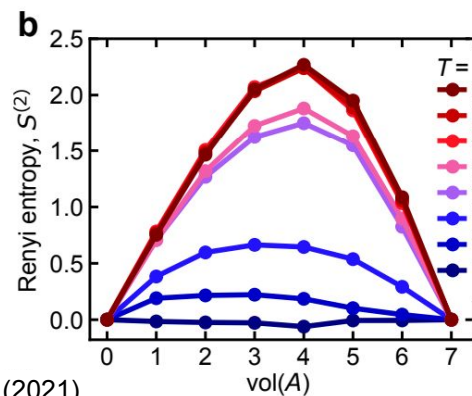
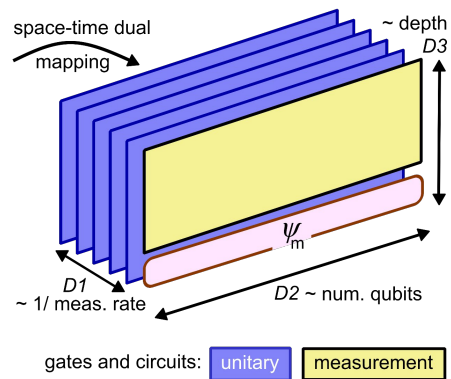
1D monitored circuits can be mapped onto shallow circuits in 2D, with final measurements on all but a 1D subsystem



1D entanglement phases from 2D shallow quantum circuits

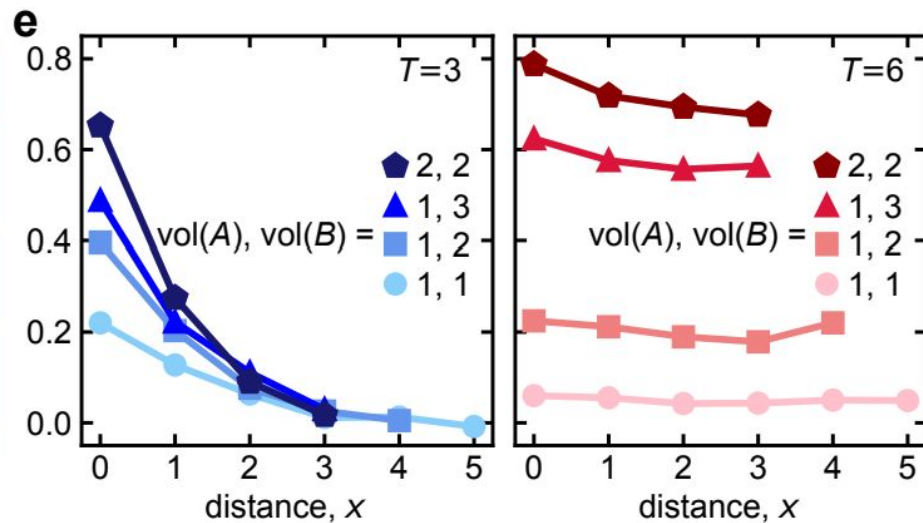
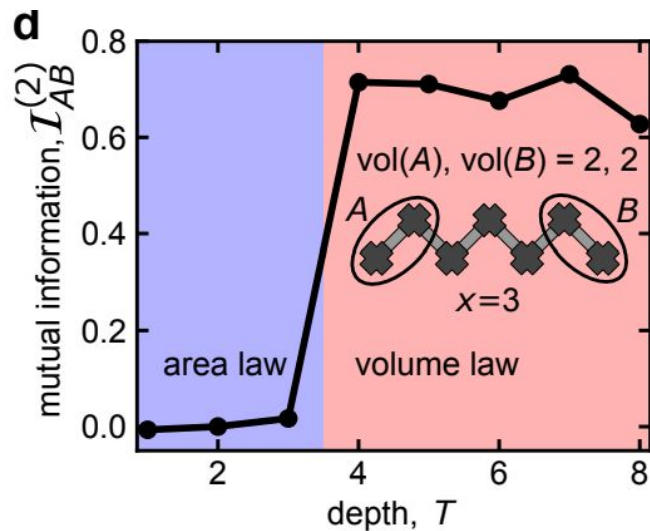
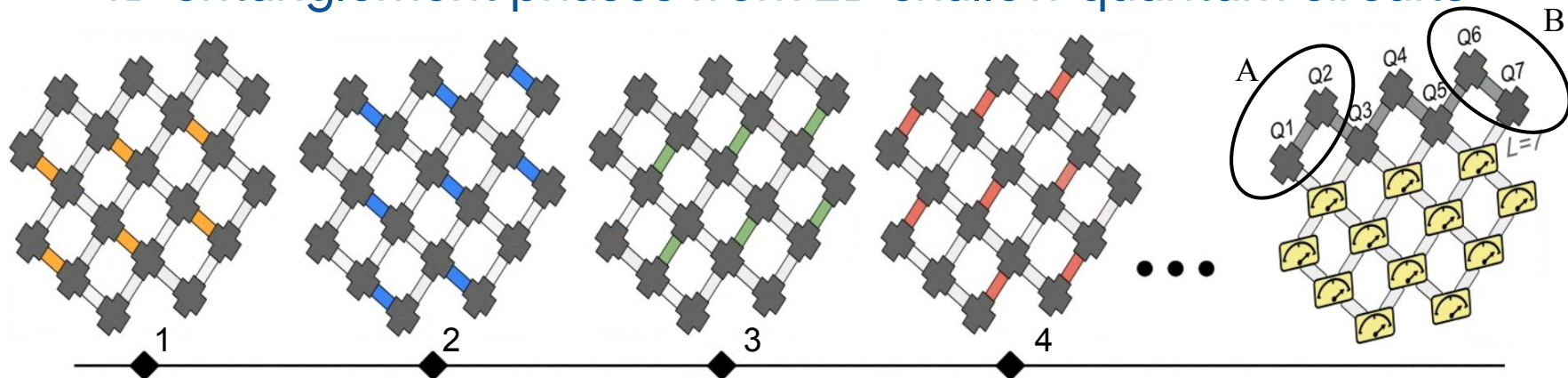


$$\mathcal{I}_{AB}^{(2)} = S_A^{(2)} + S_B^{(2)} - S_{AB}^{(2)}$$

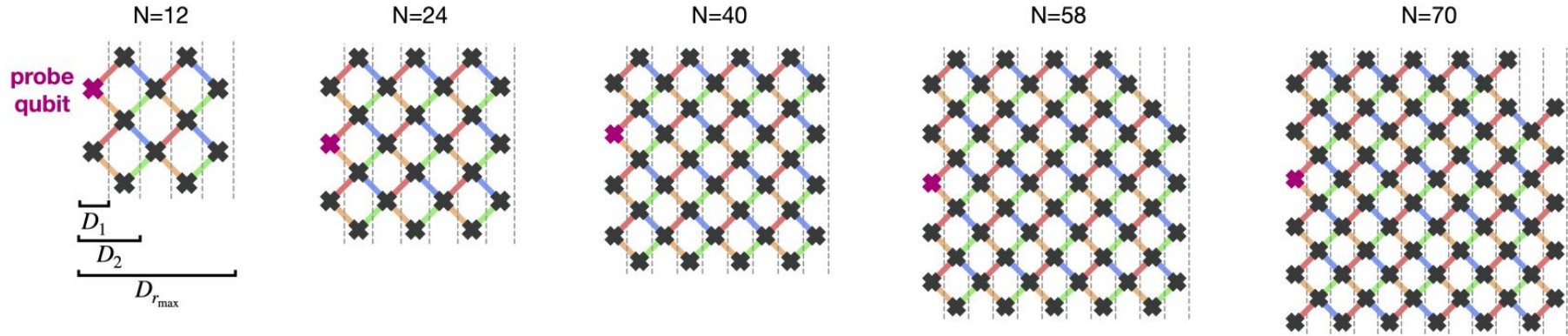


Theory work: Bao *et al.*, arXiv:2110.06963 (2021)

1D entanglement phases from 2D shallow quantum circuits

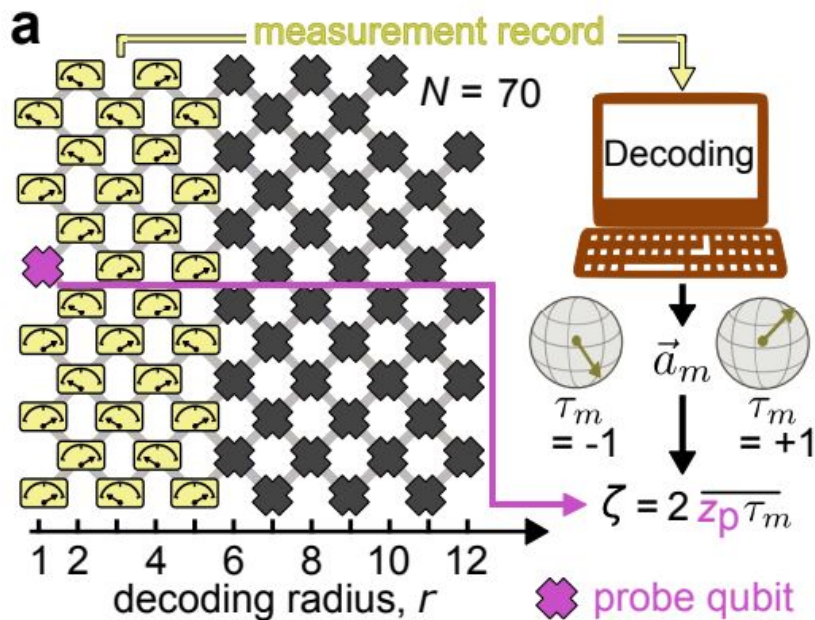
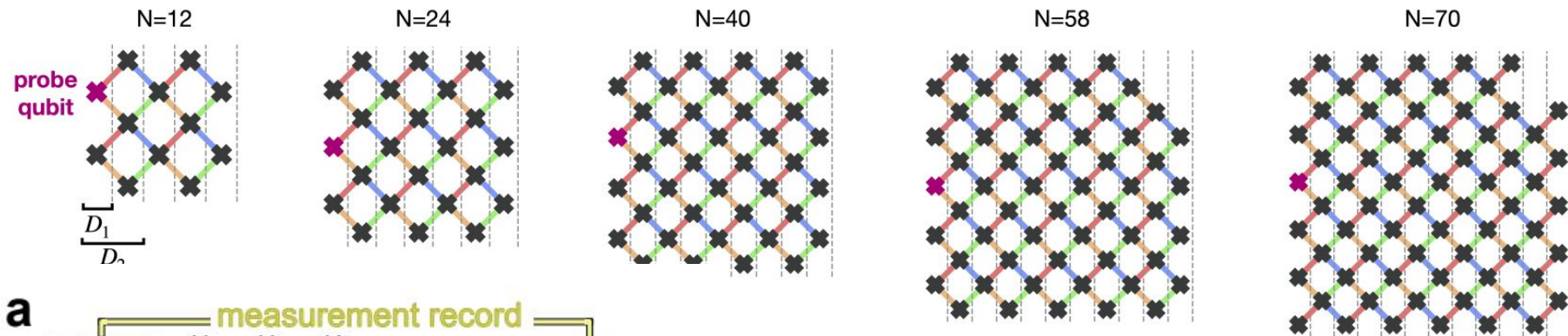


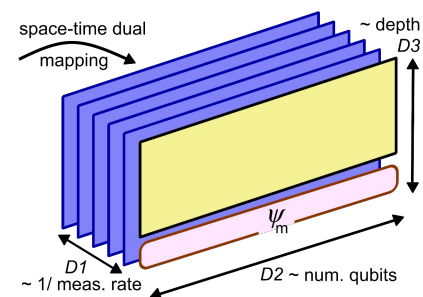
Decoding to overcome the post-selection challenge



In the entangling phase, an initially mixed state (exponential in the system size) purified.
In the disentangling phase, initially mixed states are easily purified.

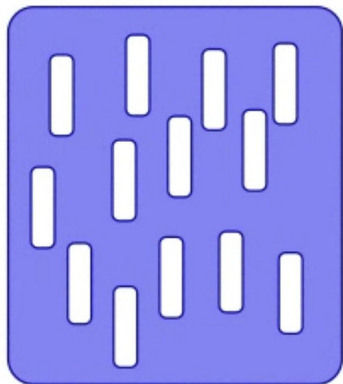
Decoding to overcome the post-selection challenge



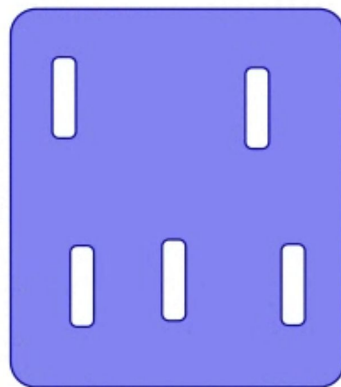


Decoding of local order parameter

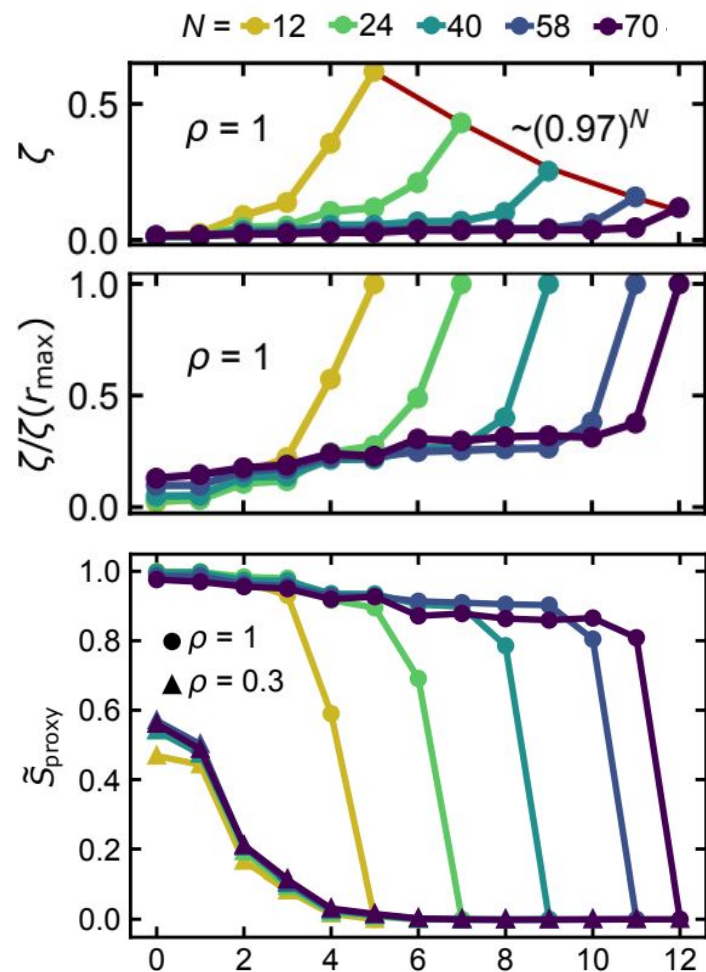
Cross-correlator $\rightarrow$
 “probe” qubit entanglement $\rightarrow$
 proxy for the system entanglement



(weak unitary)
 $\rho \rightarrow 0$

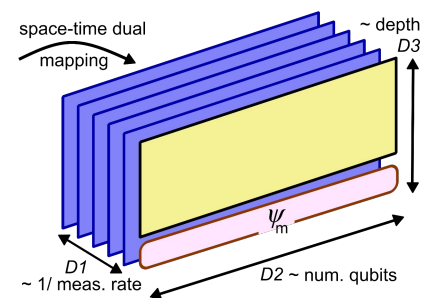


(strong unitary)
 $\rho \rightarrow 1$

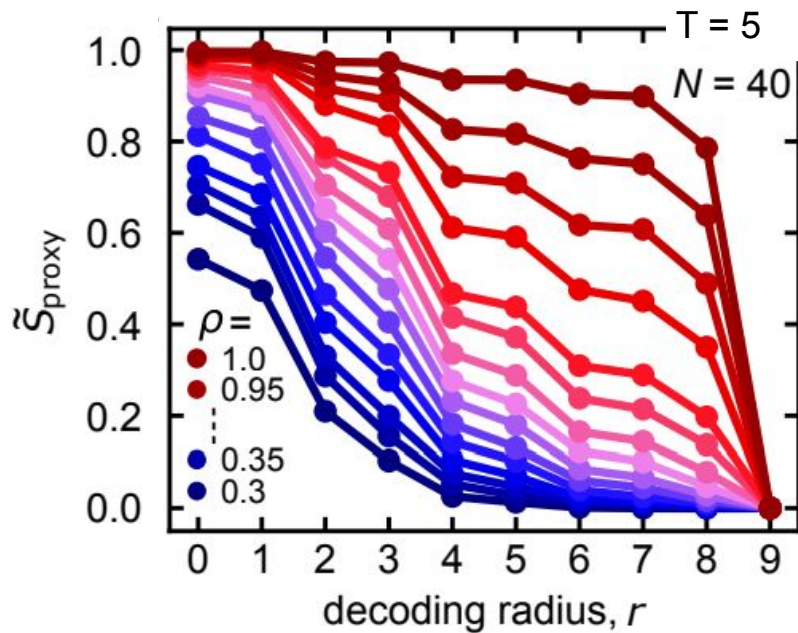


$$S_{\text{proxy}} = -\log_2[(1 + \zeta^2)/2]$$

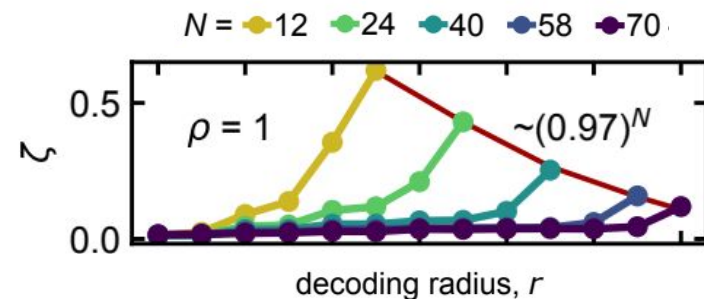
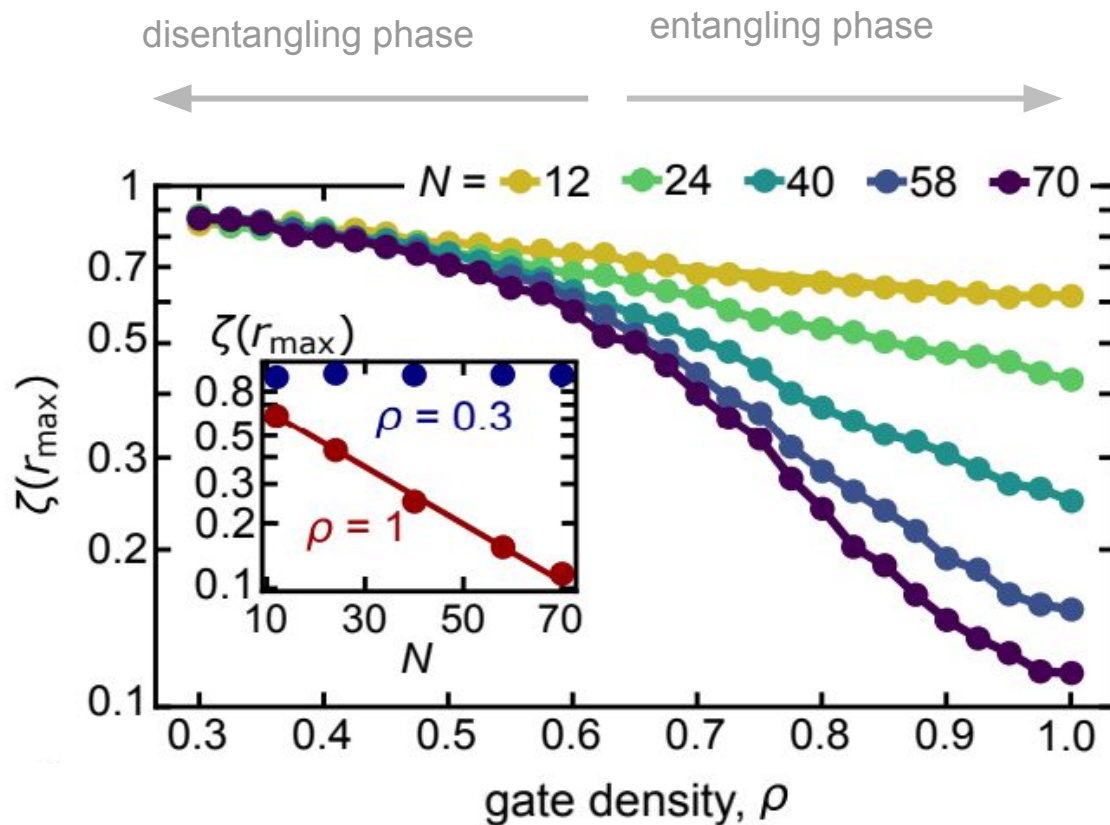
Crossover between entanglement structures



$$S_{\text{proxy}} = -\log_2[(1 + \zeta^2)/2]$$



Noise as a probe of the entanglement structure



Coherence $\rightarrow$ upper limit on qubit array sizes of about 12×12