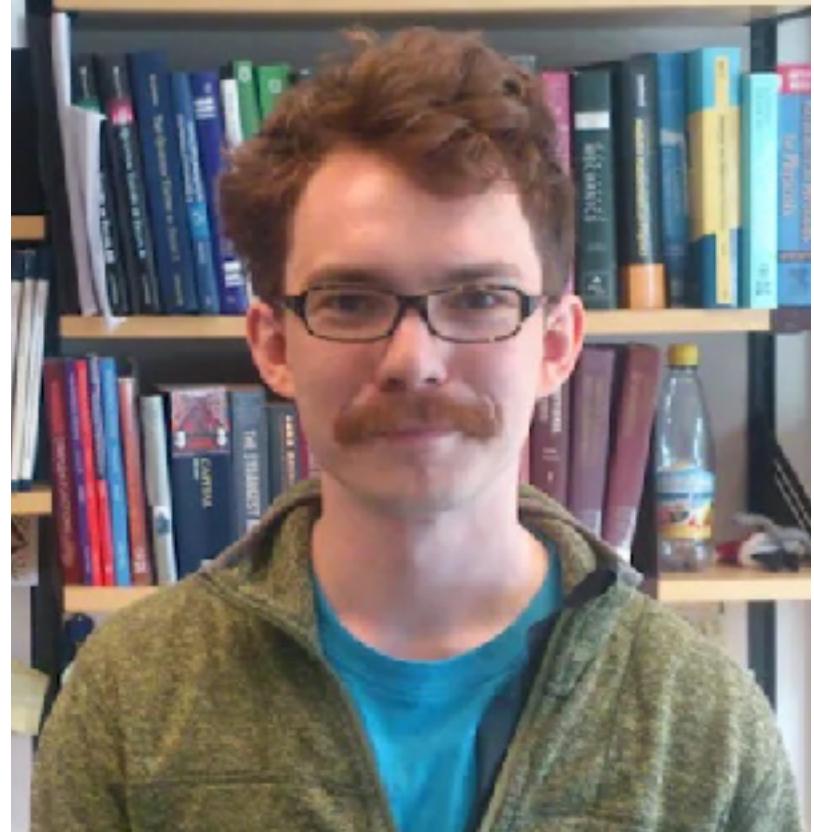




Local-Information Time Evolution

Claudia Artiaco

KTH Royal Institute of Technology, Stockholm, Sweden



Thomas Klein Kvorning
KTH



David Aceituno
KTH



Christoph Fleckenstein
KTH → ParityQC, Innsbruck



Loïc Herviou
LPMMC, Grenoble



Jens H. Bárðarson
KTH

- ★ T. Klein Kvorning, L. Herviou, J. H. Bárðarson, **SciPost Phys. 13 (4), 080 (2022)**
- ★ C. Artiaco, C. Fleckenstein, D. Aceituno, T. Klein Kvorning, J. H. Bárðarson, **PRX Quantum 5 (2), 020352 (2024)**
- ★ K. Harkins, C. Fleckenstein, N. D'Souza, P. M. Schindler, D. Marchiori, C. Artiaco, Q. Reynard-Feytis, U. Basumallick, W. Beatrez, A. Pillai, M. Hagn, A. Nayak, S. Breuer, X. Lv, M. McAllister, P. Reshetikhin, E. Druga, M. Bukov, A. Ajoy, **Sci. Adv., 11 (13), eadn 9021 (2025)**
- ★ C. Artiaco, T. Klein Kvorning, D. Aceituno, L. Herviou, J. H. Bárðarson, **Phys. Rev. Lett. 134 (19), 190401 (2025)**
- ★ N. P. Bauer, B. Trauzettel, T. Klein Kvorning, J. H. Bárðarson, C. Artiaco, **arXiv:2505.00537**

“Entanglement is [...] the characteristic trait of quantum mechanics”

Erwin Schrödinger

“Information is physical”

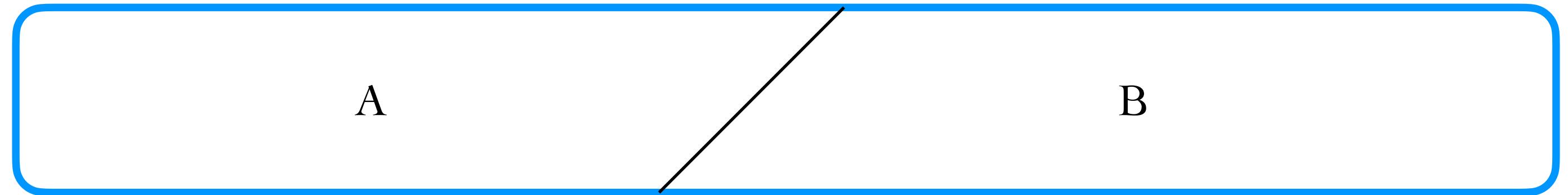
Rolf Landauer

Entanglement poses significant challenges in simulating
the out-of-equilibrium dynamics of interacting quantum systems

What is entanglement? And how is it related to information?

$$\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$$

Composite system



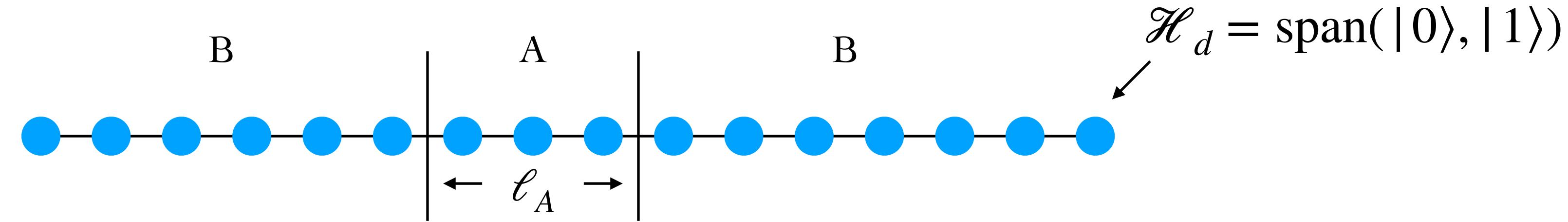
$$\left. \begin{array}{l} \rho \text{ acts on } \mathcal{H} \\ \rho^2 = \rho \end{array} \right\} \quad \rho \neq \rho_A \otimes \rho_B \text{ in any basis} \implies \rho \text{ has entanglement between A and B}$$

Information of state of subsystem A is contained in $\rho_A = \text{Tr}_B(\rho)$ via $\langle O_A \rangle = \text{Tr}_A(\rho_A O_A)$

If ρ is entangled, then all information in the state can **not** be accessed locally in A and B!

Entanglement is nonlocal information

The von Neumann information



If $\rho_A^2 = \rho_A$ then we have access to $\ln(2^{\ell_A})$ bits of information

If $\rho_A \propto \mathbf{1}_A$ then we have access to 0 bits of information

In between, the accessible information is the von Neumann (Shannon) information

$$I(\rho_A) = \ln 2^{\ell_A} - S(\rho_A)$$

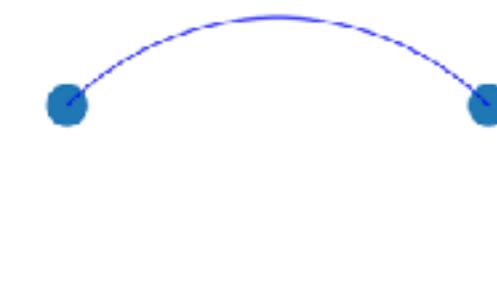
$$S = -\text{Tr}_A(\rho_A \ln \rho_A)$$

Q: What is the distribution of information in a given state?

Example: information in a singlet

$$|\psi\rangle = \frac{1}{\sqrt{2}} (| \uparrow \downarrow \rangle - | \downarrow \uparrow \rangle)$$

$$\rho = \frac{1}{4} \mathbf{1}_4$$



There is **no** information in the single sites!

$$\rho_A = \frac{1}{2} \mathbf{1}_2 = \rho_B$$

$$S_A = \ln 2 = S_B$$

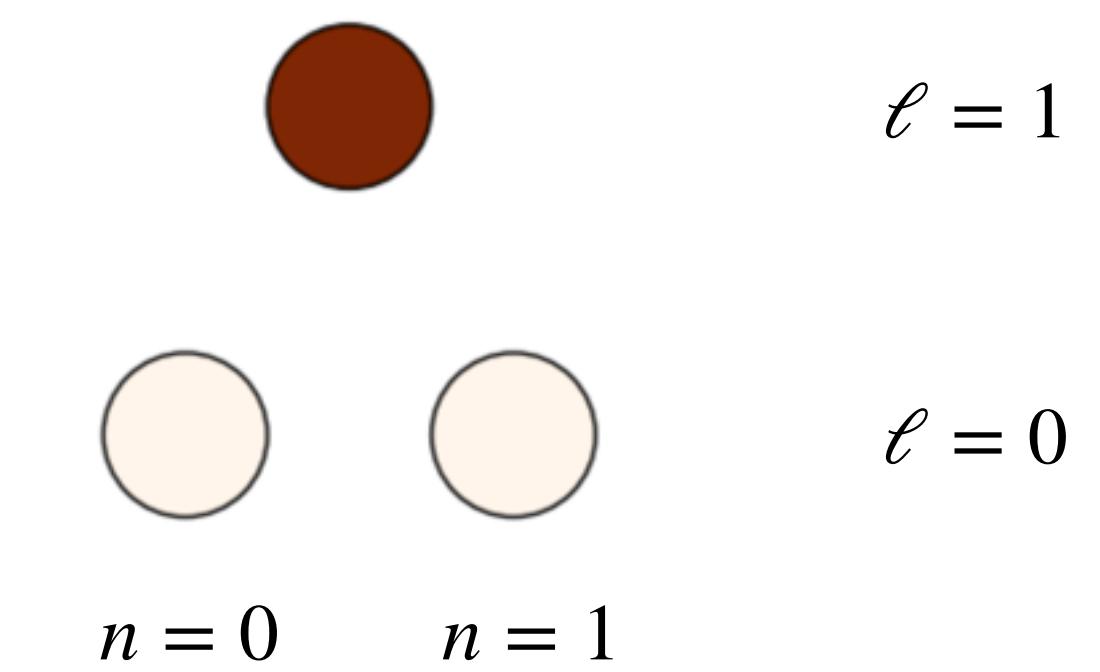
$$I_A = 0 = I_B$$

Information on single site says nothing about information on two sites

$$\rho_A = \frac{1}{2} \mathbf{1}_2 = \rho_B$$

$$S_A = \ln 2 = S_B$$

$$I_A = 0 = I_B$$



All information is on the two site state

$$S_{A \cup B} = 0$$

$$I_{A \cup B} = \ln 4$$

Fully mixed two site state has no information

$$S_{A \cup B} = \ln 4$$

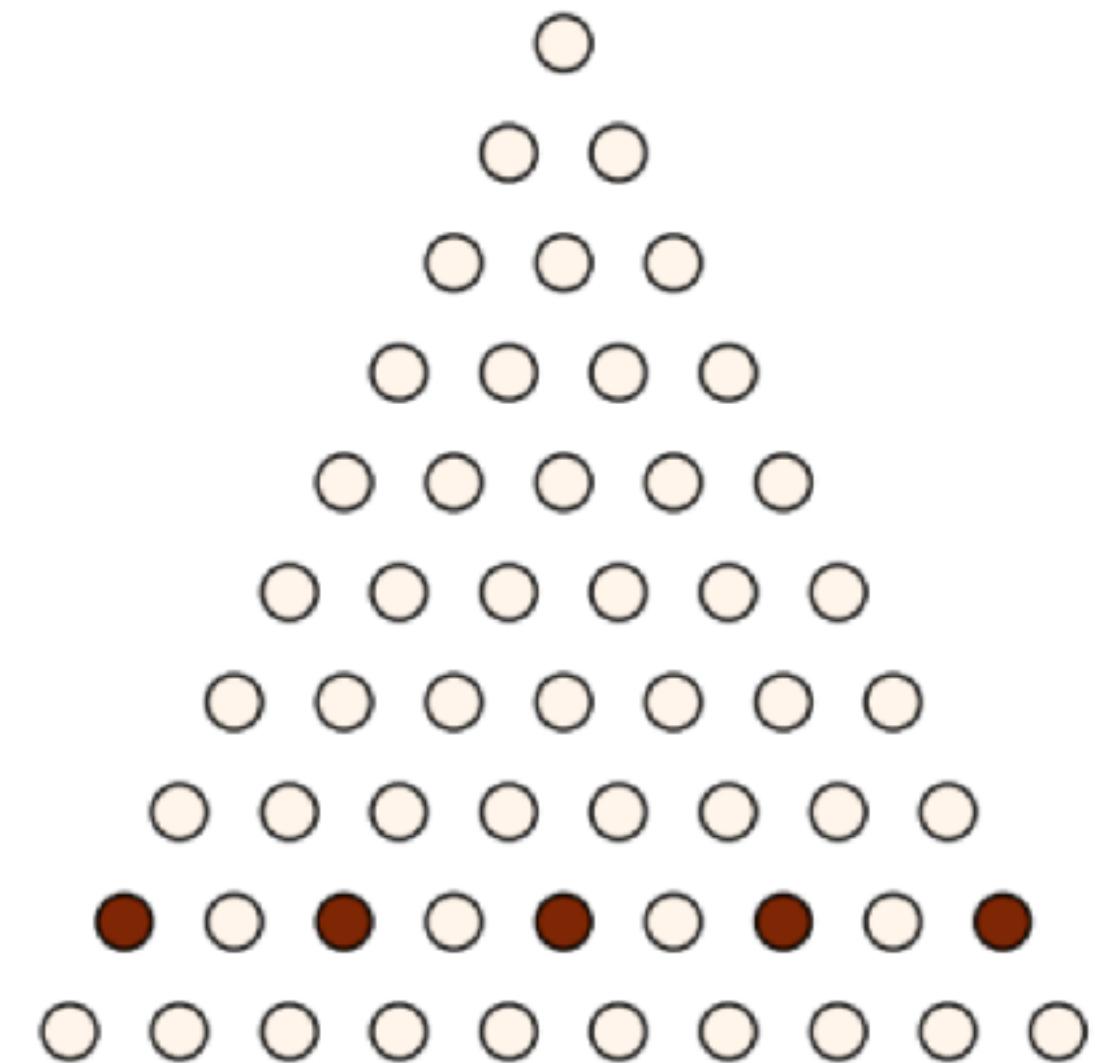
$$I_{A \cup B} = 0$$

We define **local information** as

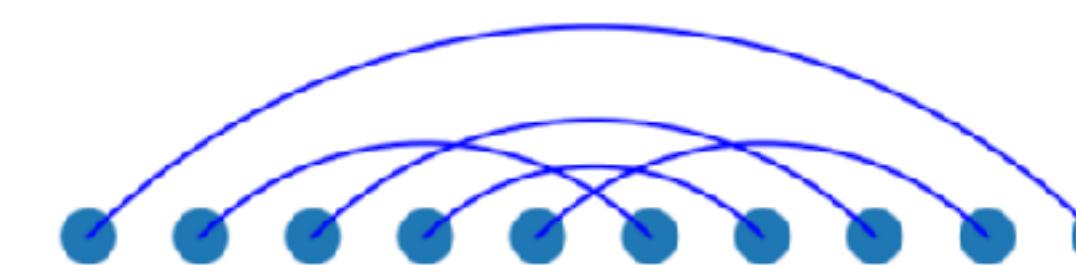
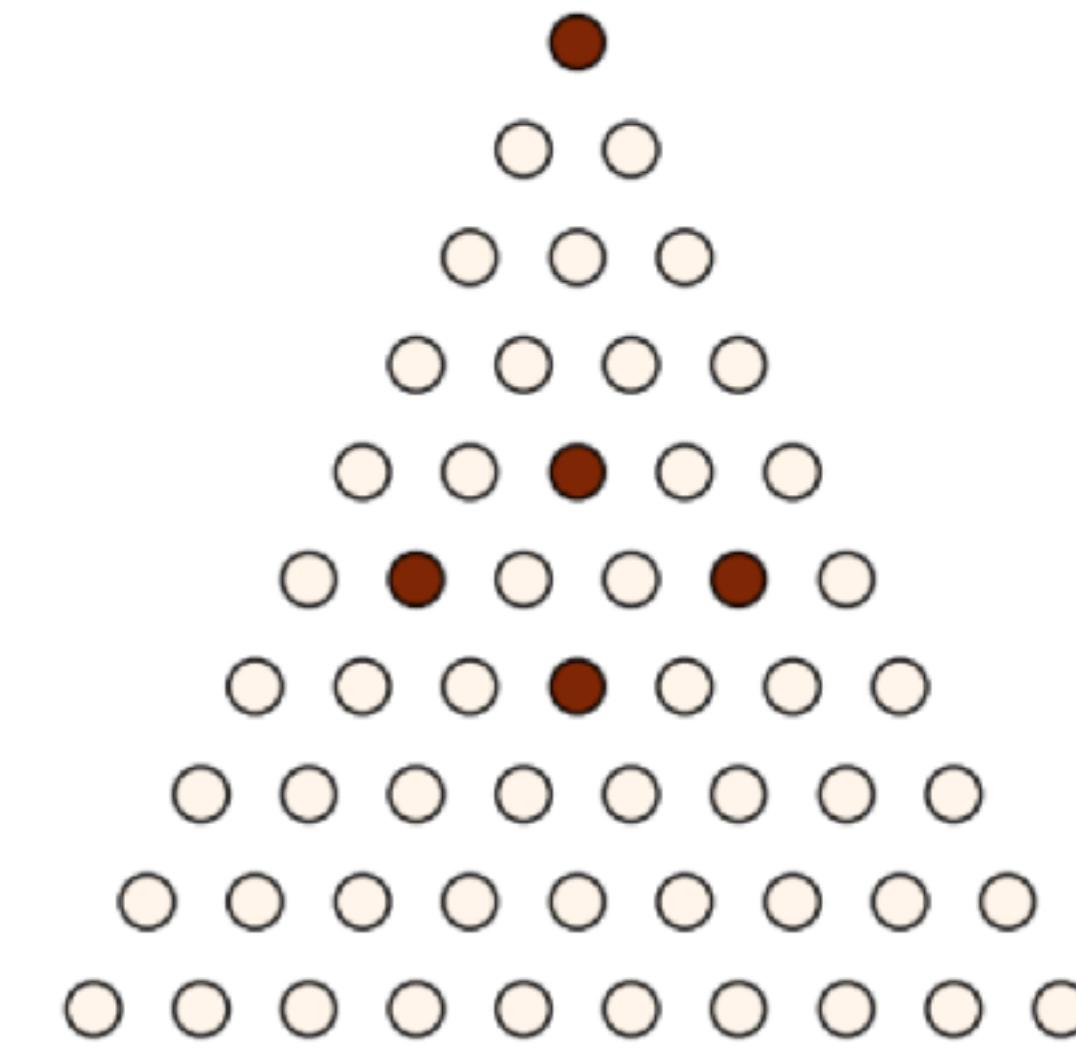
$$i_n^\ell = I_n^\ell - I_{n-1/2}^{\ell-1} - I_{n+1/2}^{\ell-1} + I_n^{\ell-2}$$

Local information within the information lattice

Product state of neighboring singlets



Product state of singlets

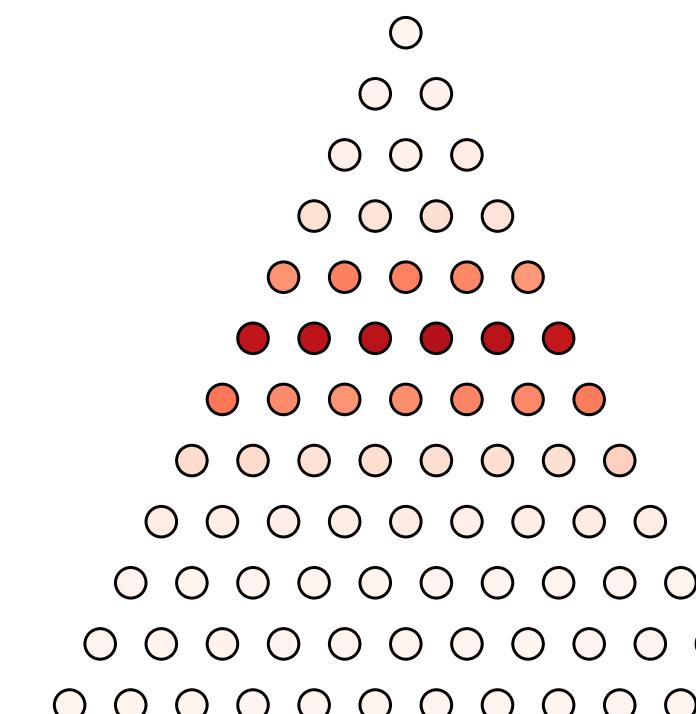
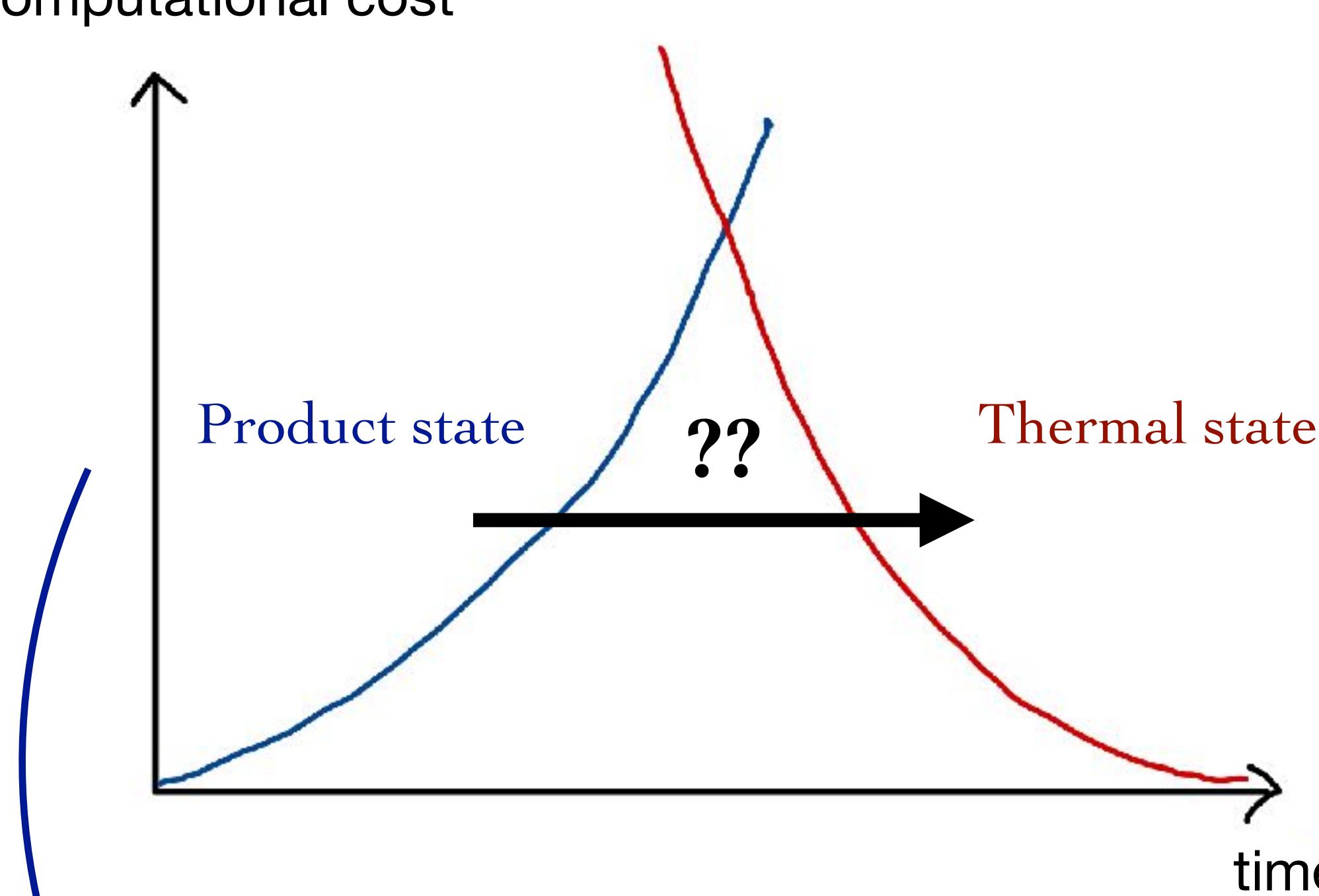


- i_n^ℓ decomposes the total information in a region centered at n according to the corresponding scales
- i_n^ℓ is local: a local unitary on nearest neighbors can only move it from ℓ to $\ell \pm 1$

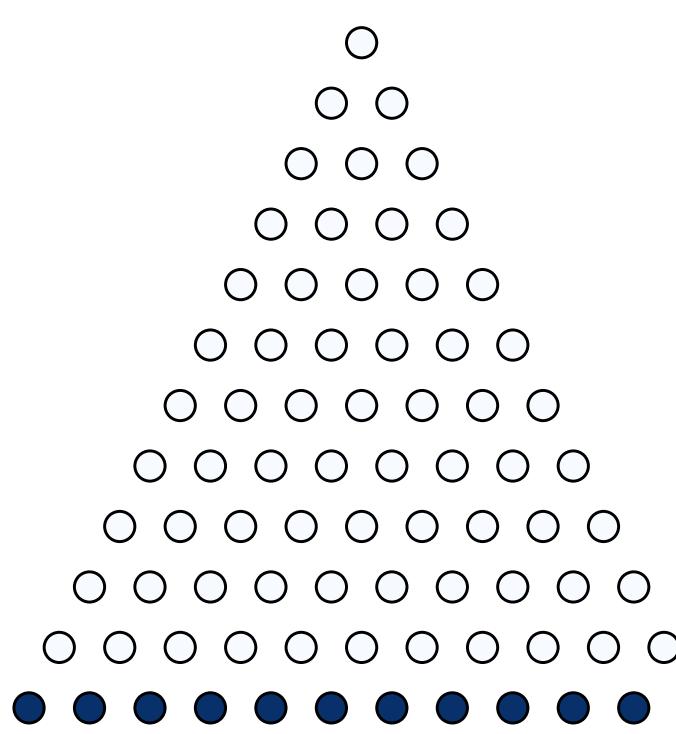
Local information for the efficient time evolution of interacting quantum systems

Q: How to simulate the dynamics of thermalizing interacting quantum systems?

Computational cost



New approach: Local information time evolution (LITE)



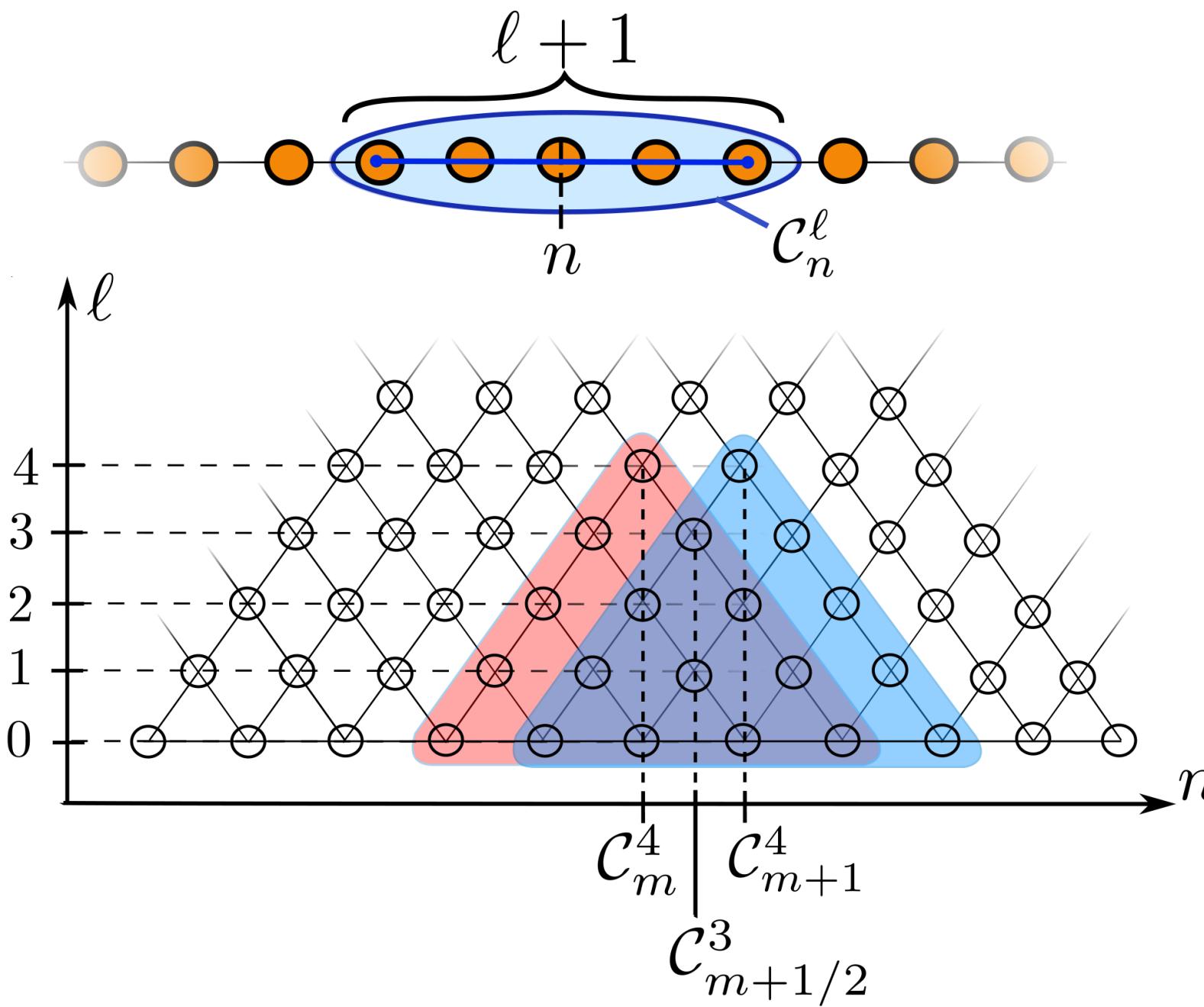
Efficient simulation of the dynamics
of local observables for interacting local Hamiltonians

[T.Klein Kvorning, L.Herviou, J.H.Bárðarson, SciPost Physics 2022]

[C.Artiaco, C.Fleckenstein, D.Aceituno Chávez, T.Klein Kvorning, J.H.Bárðarson, PRX Quantum 2024]

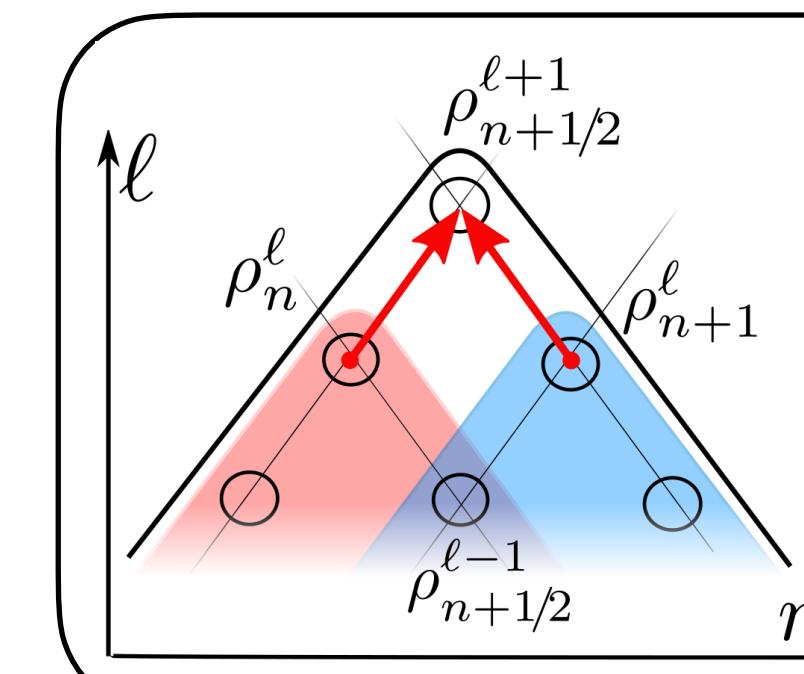
Time evolution on subsystems

We decompose the system in smaller subsystems...



... to solve the von Neumann equation in parallel

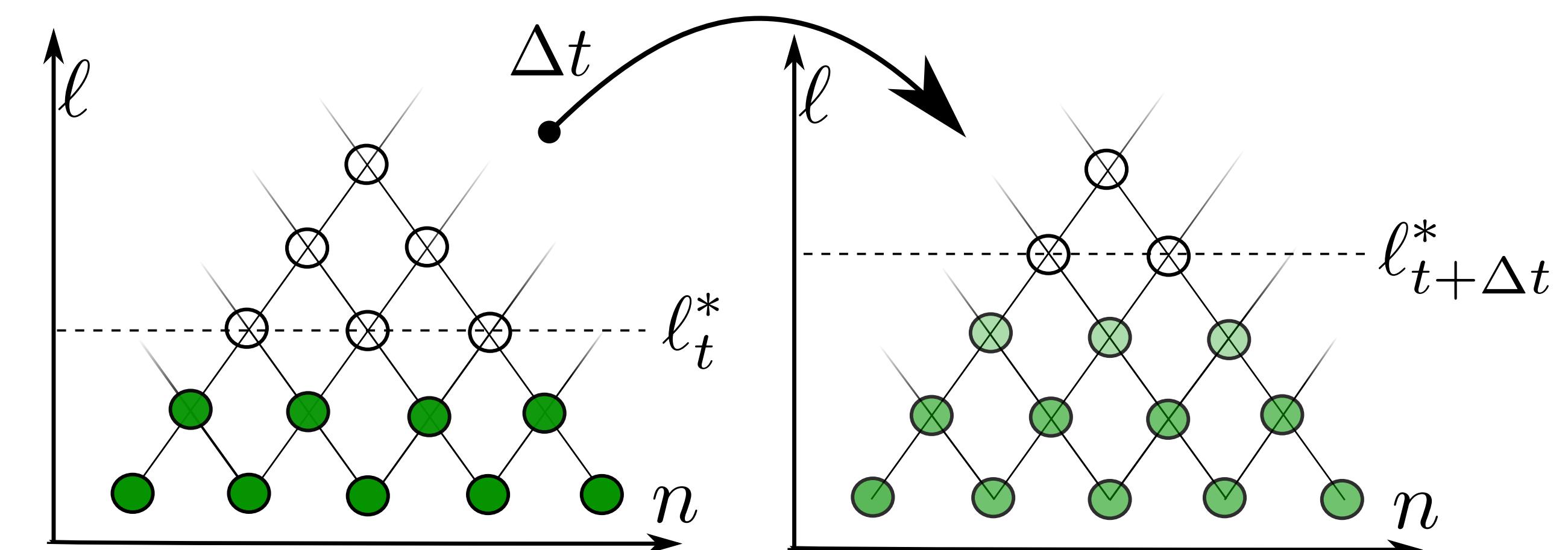
$$\begin{aligned}\partial_t \rho_n^\ell &= -i [H_n^\ell, \rho_n^\ell] \\ &\quad - i\text{Tr}_L^r ([H_{n-r/2}^{\ell+r} - H_n^\ell, \rho_{n-r/2}^{\ell-r}]) \\ &\quad - i\text{Tr}_R^r ([H_{n+r/2}^{\ell+r} - H_n^\ell, \rho_{n+r/2}^{\ell+r}])\end{aligned}$$



$r = 1$: If there is **no information on scale $\ell + 1$** , we can construct $\rho^{\ell+1}$ via a **Petz recovery map**

$$\rho_{n+1/2}^{\ell+1} = \exp [\ln(\rho_n^\ell) + \ln(\rho_{n+1}^\ell) - \ln(\rho_{n+1/2}^{\ell-1})]$$

closing the von Neumann equation



But ℓ^* (on which local information is 0) increases with time (growth of entanglement)!

This gives rise to a non-closed hierarchy of equations!

Petz recovery map at fixed scale fails

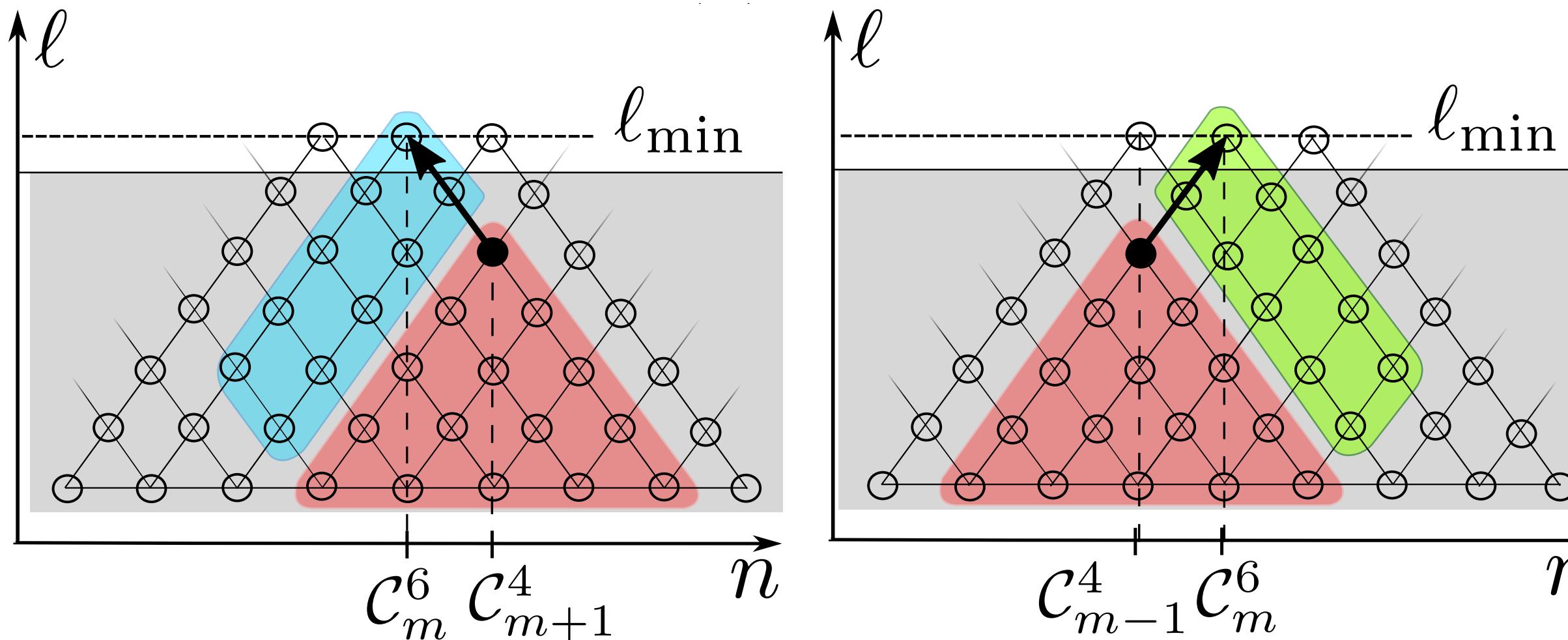
Minimize information on a large scale ℓ_{\min} keeping density matrices and information flow on smaller scales fixed

$$\bar{\rho}_n^\ell = \rho_n^\ell + \chi_n^\ell$$

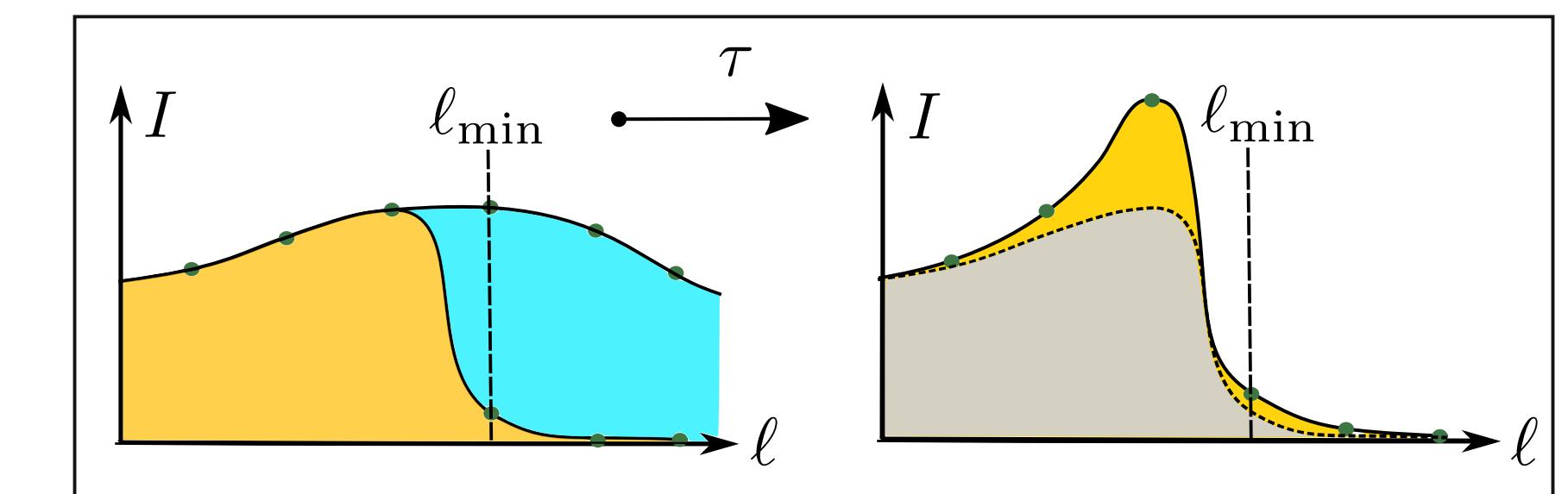
$$\begin{aligned} \text{I } & \left\{ \begin{array}{l} \text{Tr}_R^1(\chi_n^\ell) = 0 \\ \text{Tr}_L^1(\chi_n^\ell) = 0 \end{array} \right\} \\ \text{II } & \left\{ \begin{array}{l} J_R^r(\chi_n^\ell) = 0 \\ J_L^r(\chi_n^\ell) = 0 \end{array} \right\} \end{aligned}$$

Defines subspace to obtain $\xi = \mathbf{P}\chi$

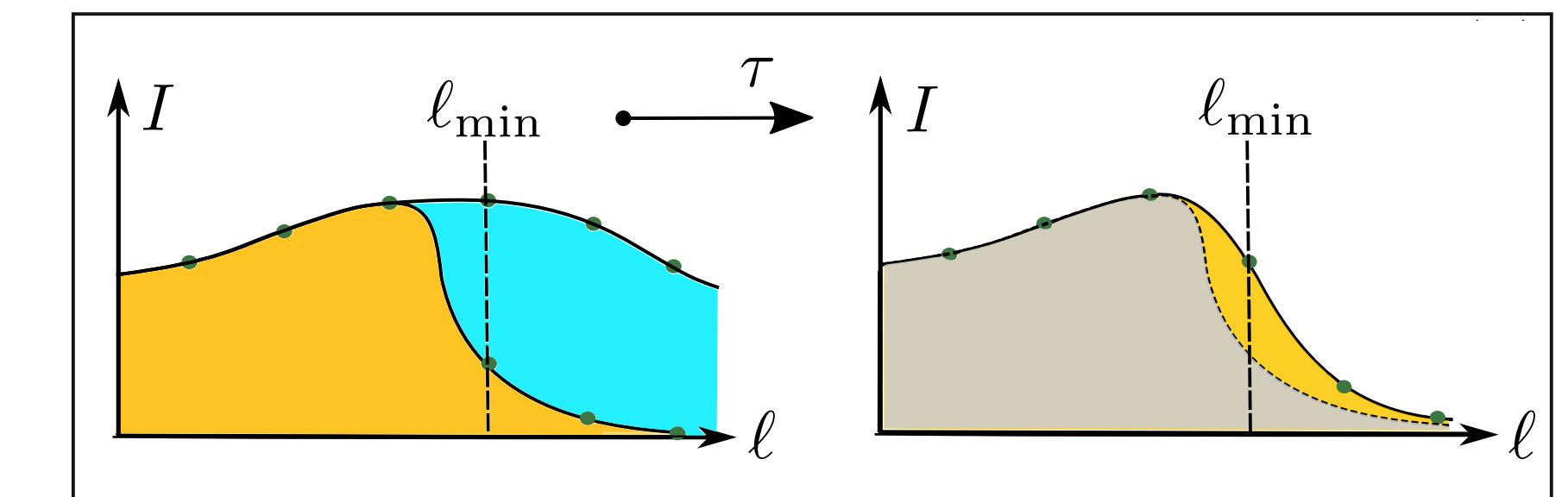
Maximize $S(\rho + \xi) = S(\rho) + \text{Tr}(\mathbf{P} \nabla_\rho S \chi) + \frac{1}{2} \text{Tr}(\chi \mathbf{P} \mathcal{H}_\rho \mathbf{P} \chi) + \mathcal{O}(\xi^3)$



Minimization under constraints I



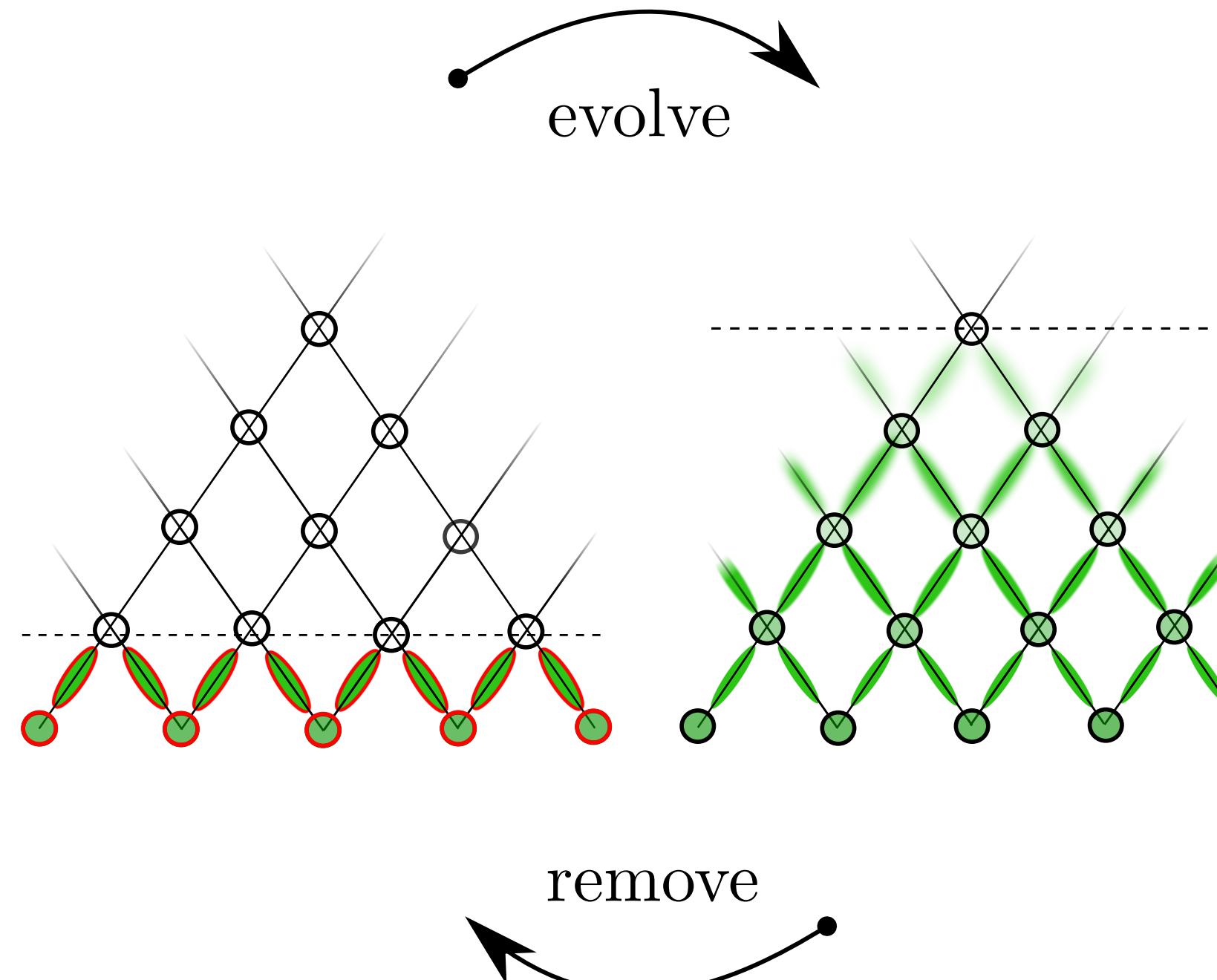
Minimization under constraints I + II



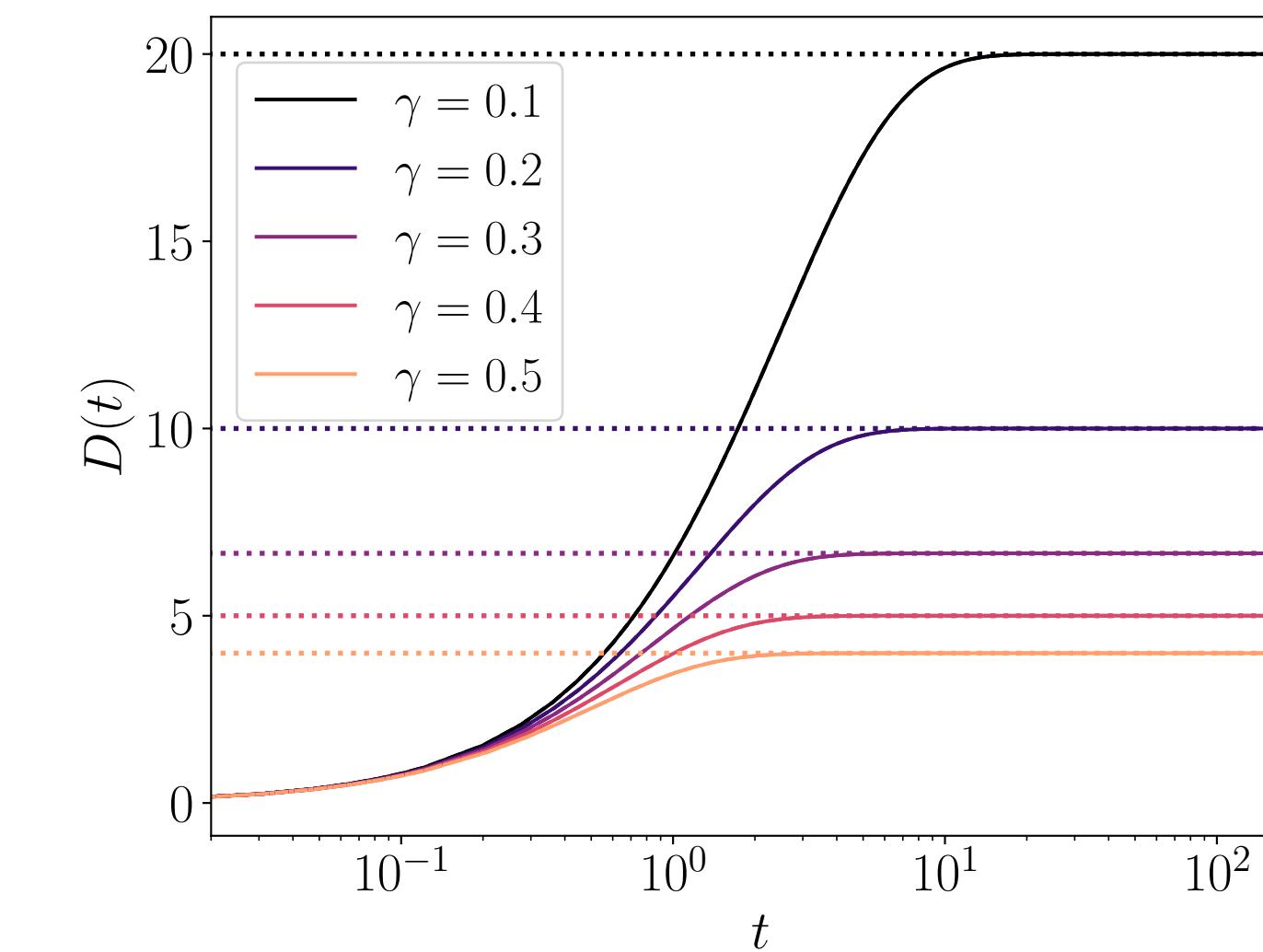
We get reliable information current
evolving information to larger scale $\ell_{\max} > \ell_{\min}$

Magnetization & energy diffusion with LITE

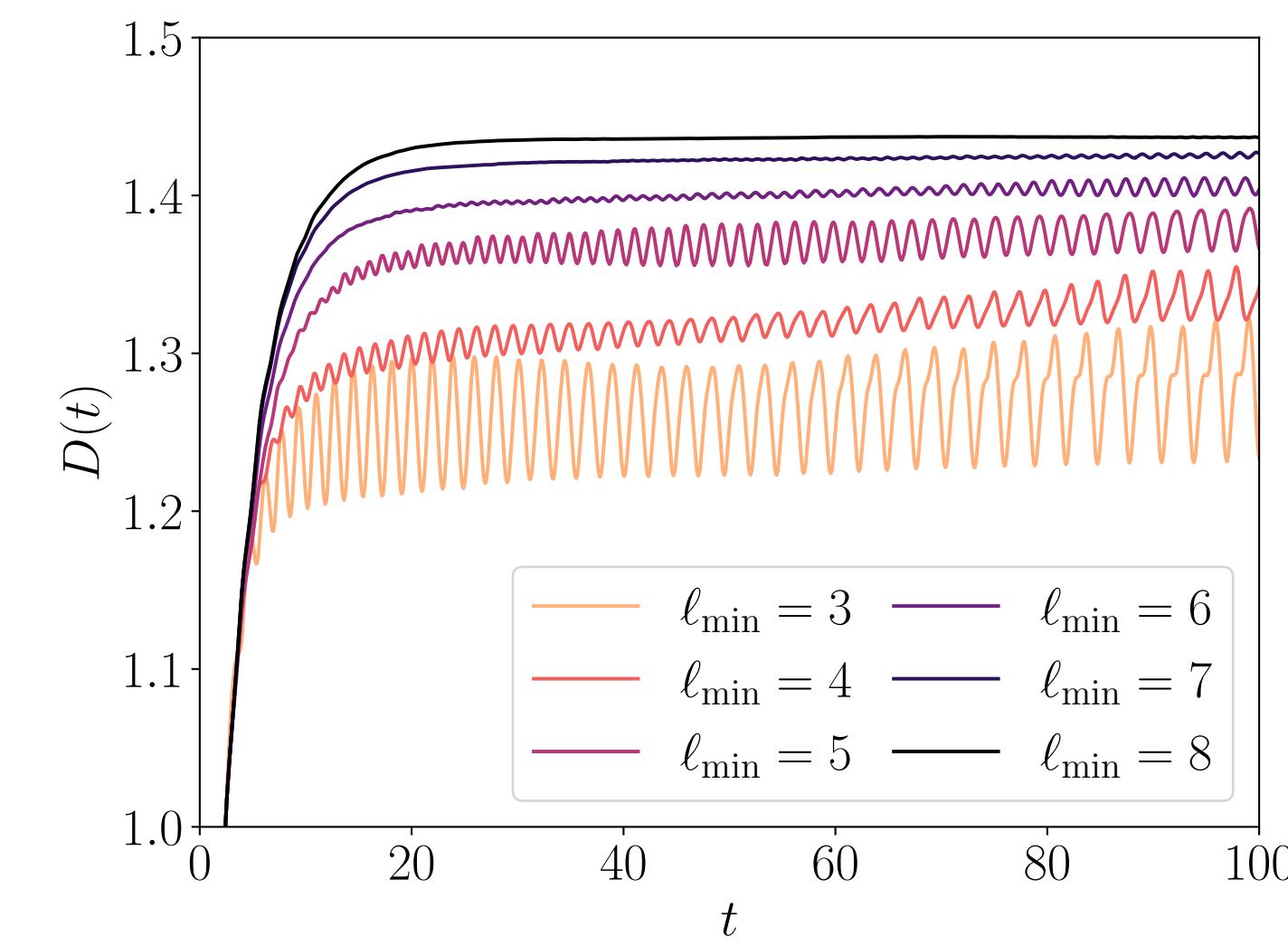
Remove information at large scales
preserving information currents



Open XX chain with dephasing



Mixed field Ising model

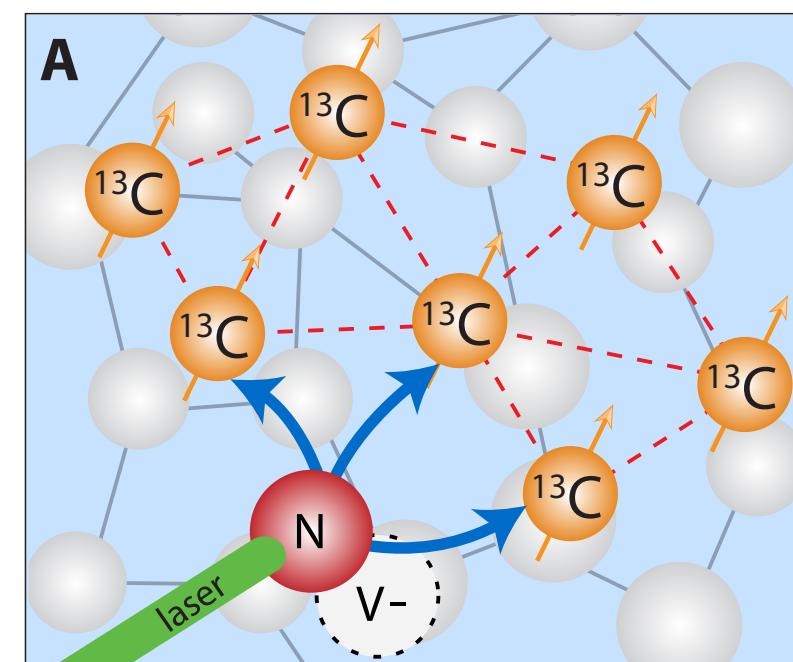


[C.Artiaco, C.Fleckenstein, D.Aceituno Chávez, T.Klein Kvorning, J.H.Bárðarson, PRX Quantum 2024]

[K.Harkins, C.Fleckenstein, N.D'Souza, P.M.Schindler, D.Marchiori, C.Artiaco, Q.Reynard-Feytis, U.Basumallick, W.Beatrez, A.Pillai, M.Hagn, A.Nayak, S.Breuer, X.Lv, M.McAllister, P.Reshetikhin, E.Drusa, M. Bukov, A.Ajoy, Science Advances 2025]

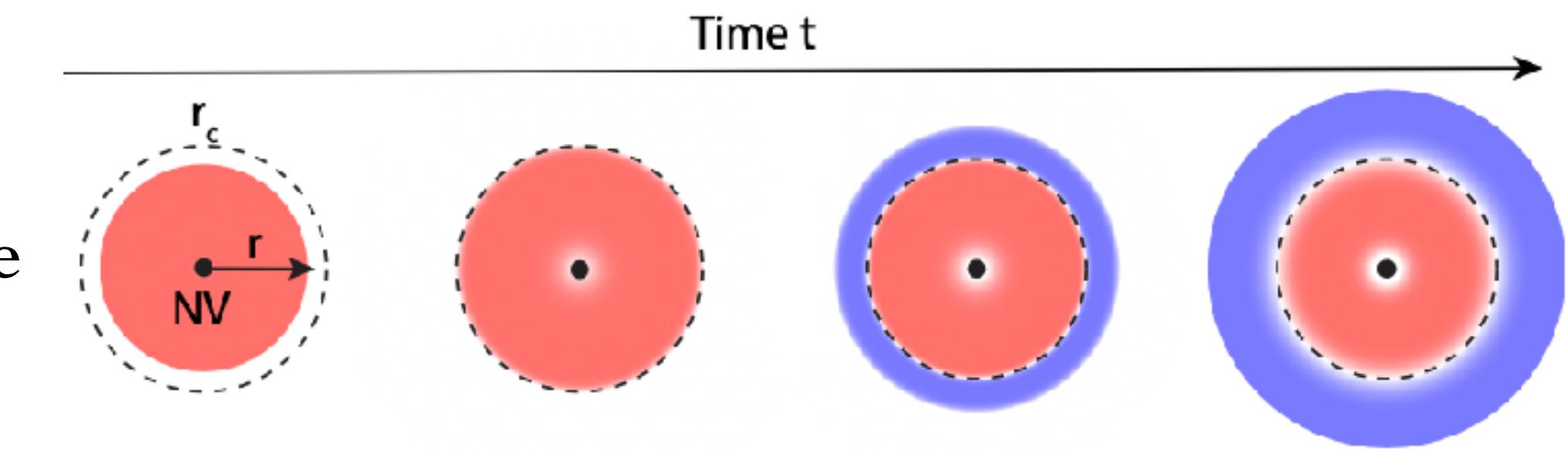
Nanoscale spin transport with LITE

Magnetization transport
in interacting Carbon-13 nuclear spins
surrounding a Nitrogen-Vacancy center
in diamond

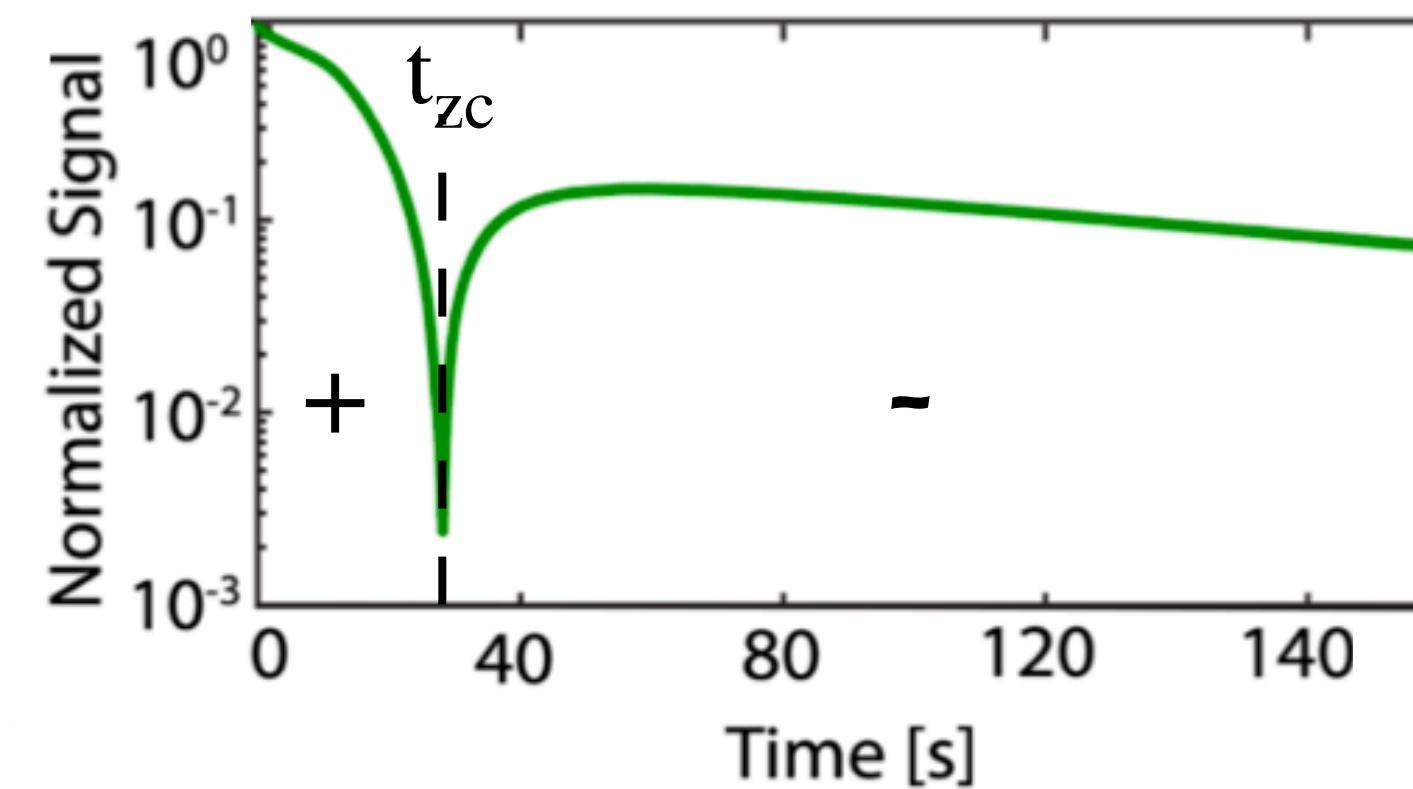


LITE is the state-of-the-art method for transport
→ reproduced the experimental results and gave access
to the spatial distribution of magnetization

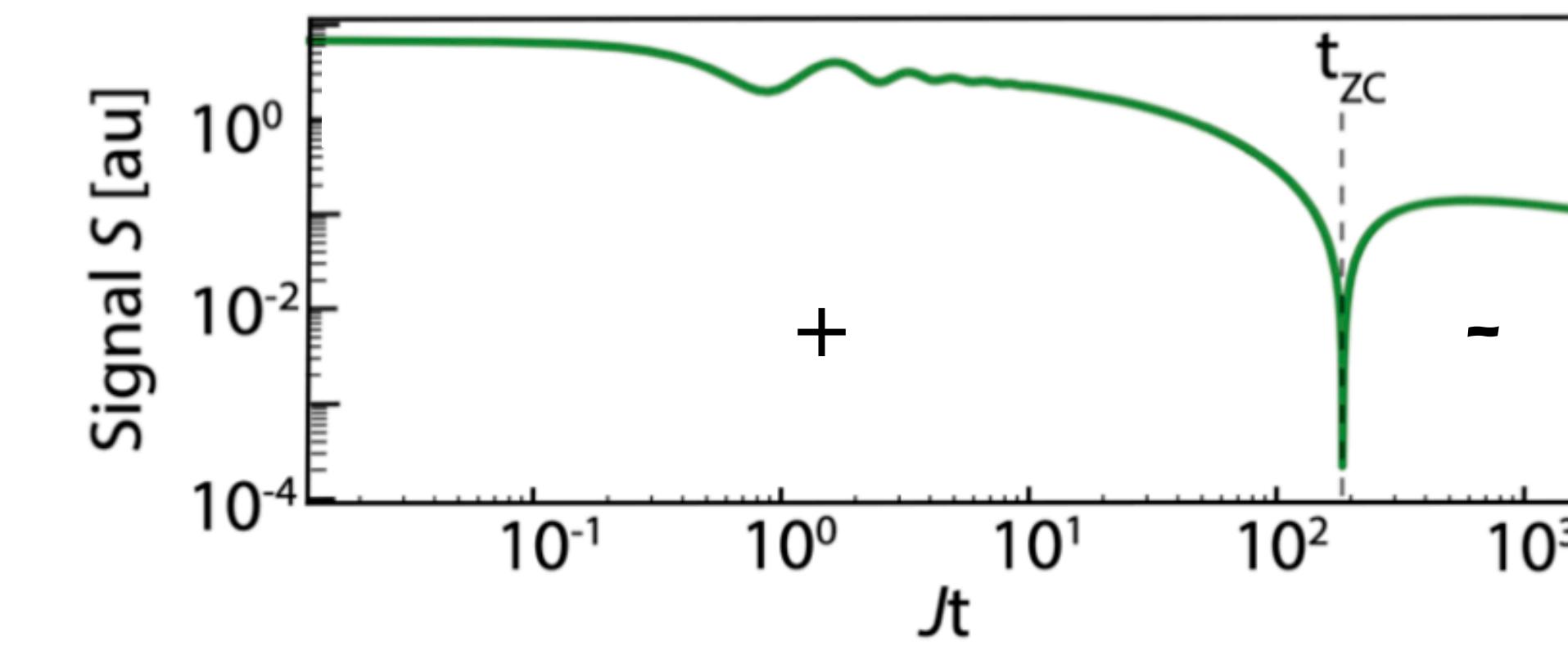
Stable spin texture



Experiment



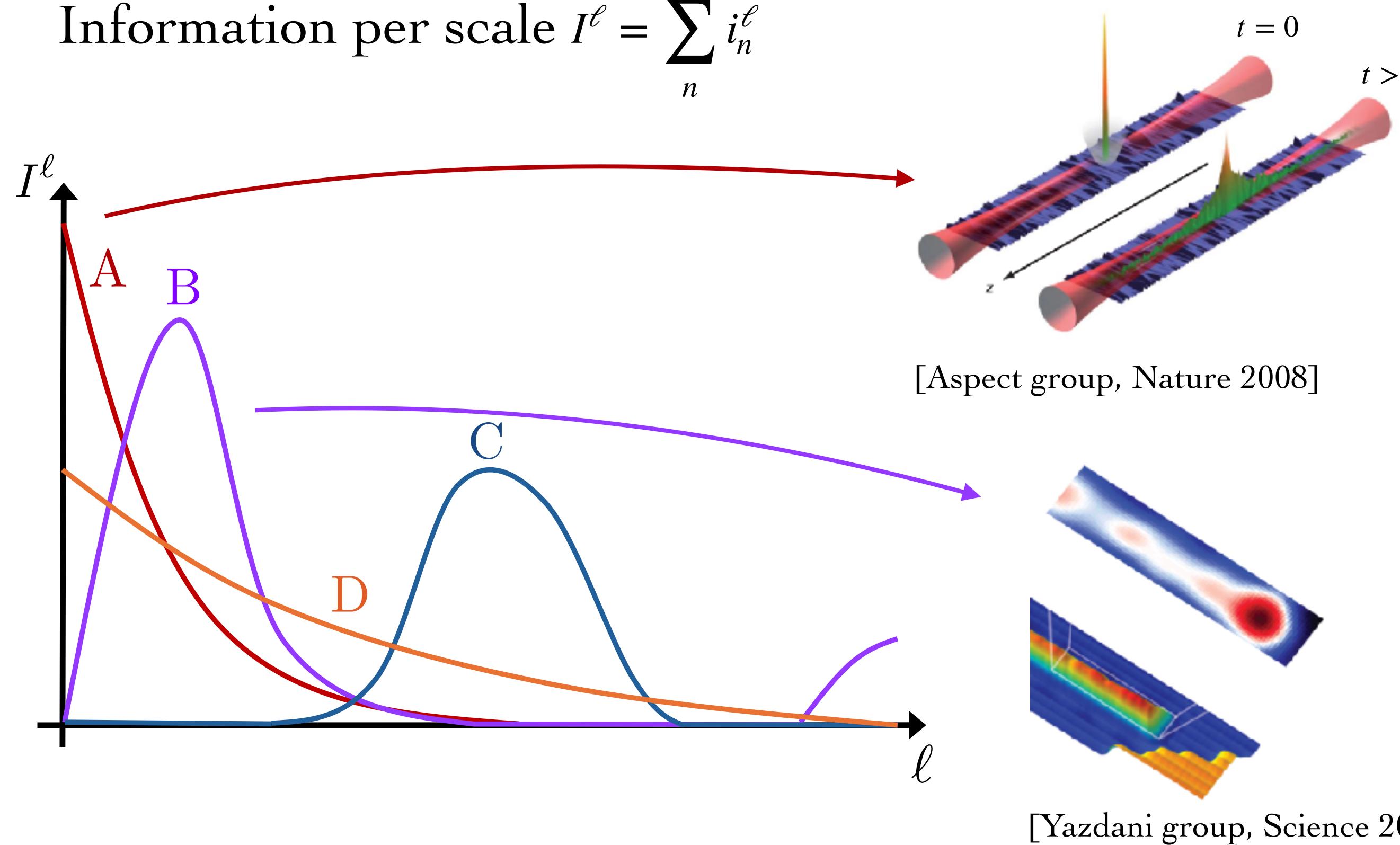
Numerics with LITE



Characterizing states through local information

The information lattice characterizes quantum states
through the distribution of local information

$$\text{Information per scale } I^\ell = \sum_n i_n^\ell$$



(A) Localized state

(B) Localized state with topological edge modes

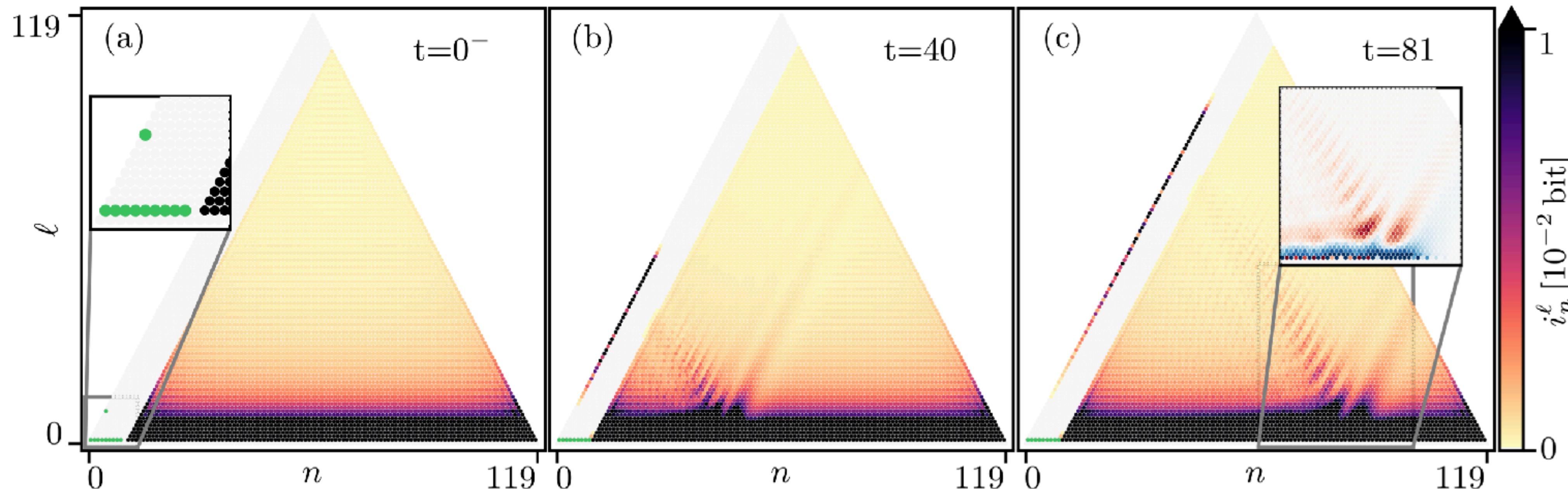
(C) Thermal state

(D) Critical state

Local information flow in quantum quench dynamics

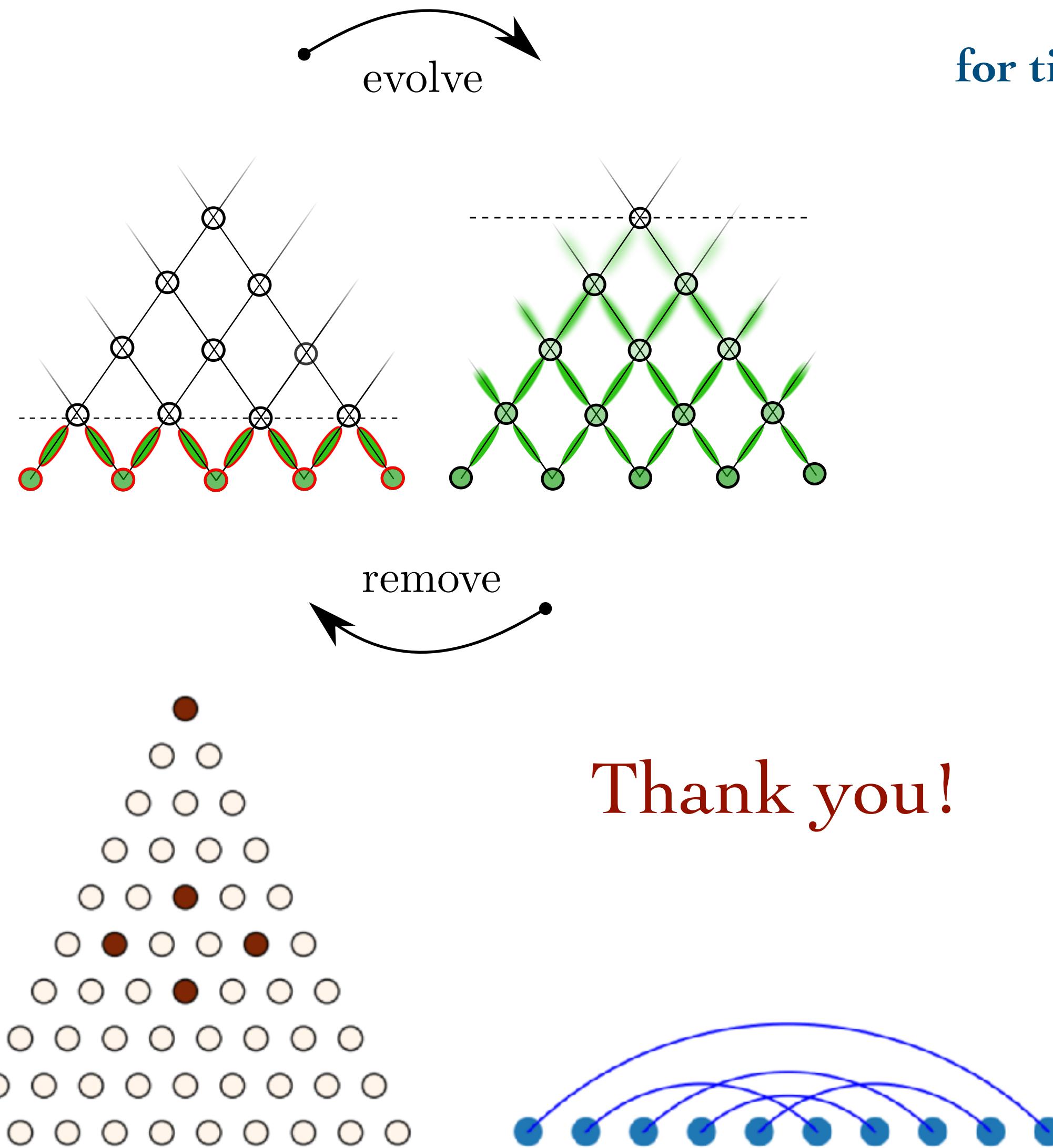
The local information flow characterizes quantum quench dynamics

Example: Topological Kitaev chain coupled to a tight-binding chain



It explains previously observed fractional signatures of entanglement entropy

Conclusions



**LITE is an efficient method
for time evolution of large many-body quantum systems up to large times**

LITE applies to generic local Hamiltonians:

- Disordered and localized systems
- Subdiffusive, diffusive, superdiffusive, and ballistic hydrodynamic behaviors
- Open quantum system governed by the Lindblad equation
- Driven systems
- Higher dimensional systems

The code will soon be publicly available and fully documented

**The information lattices a universal framework
for characterizing quantum states and quantum dynamics**