

Zeros of Complexified Survival Amplitude at Different Time-Scales

**Long-Range Interactions and Dynamics
in Complex Quantum Systems**



NORDITA



Jakub Novotný
Faculty
of Mathematics and Physics
Charles University

Outline

Survival Amplitude

Motivation

Distribution of the zeros in the complex plane

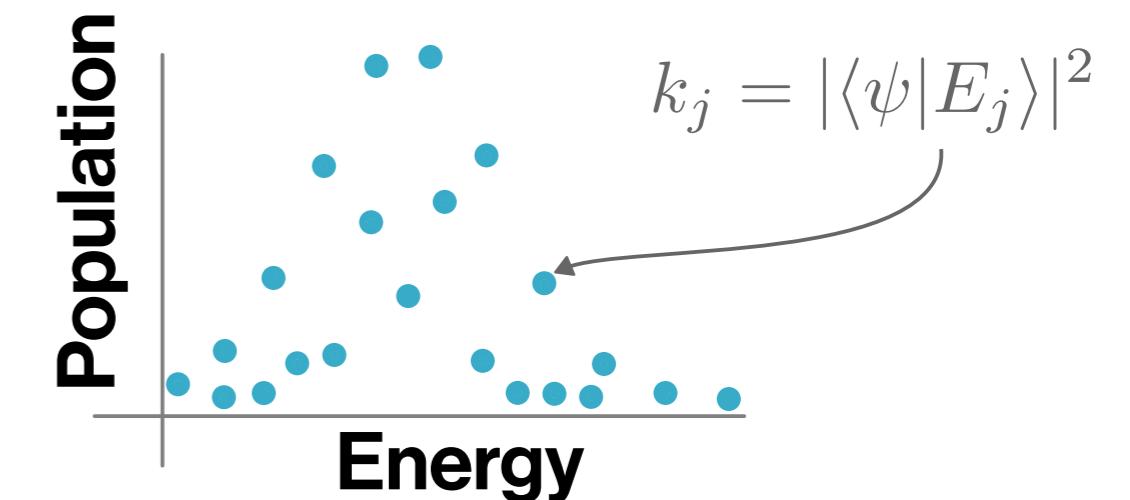
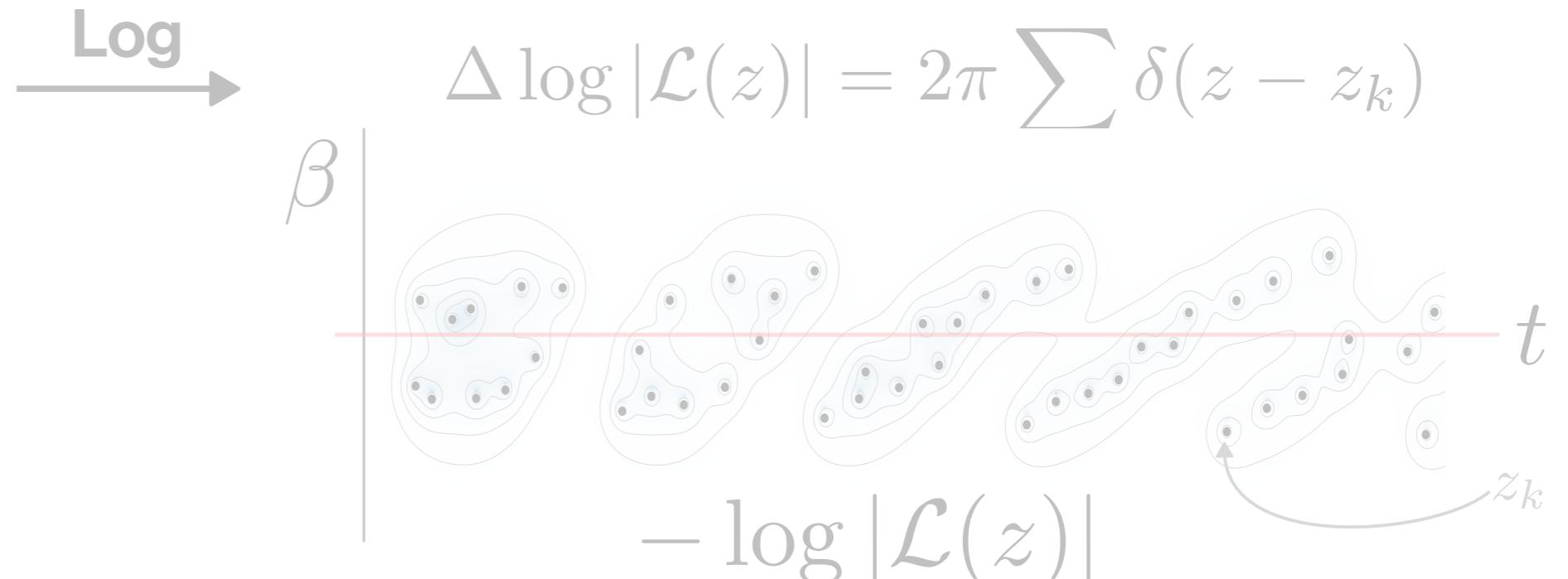
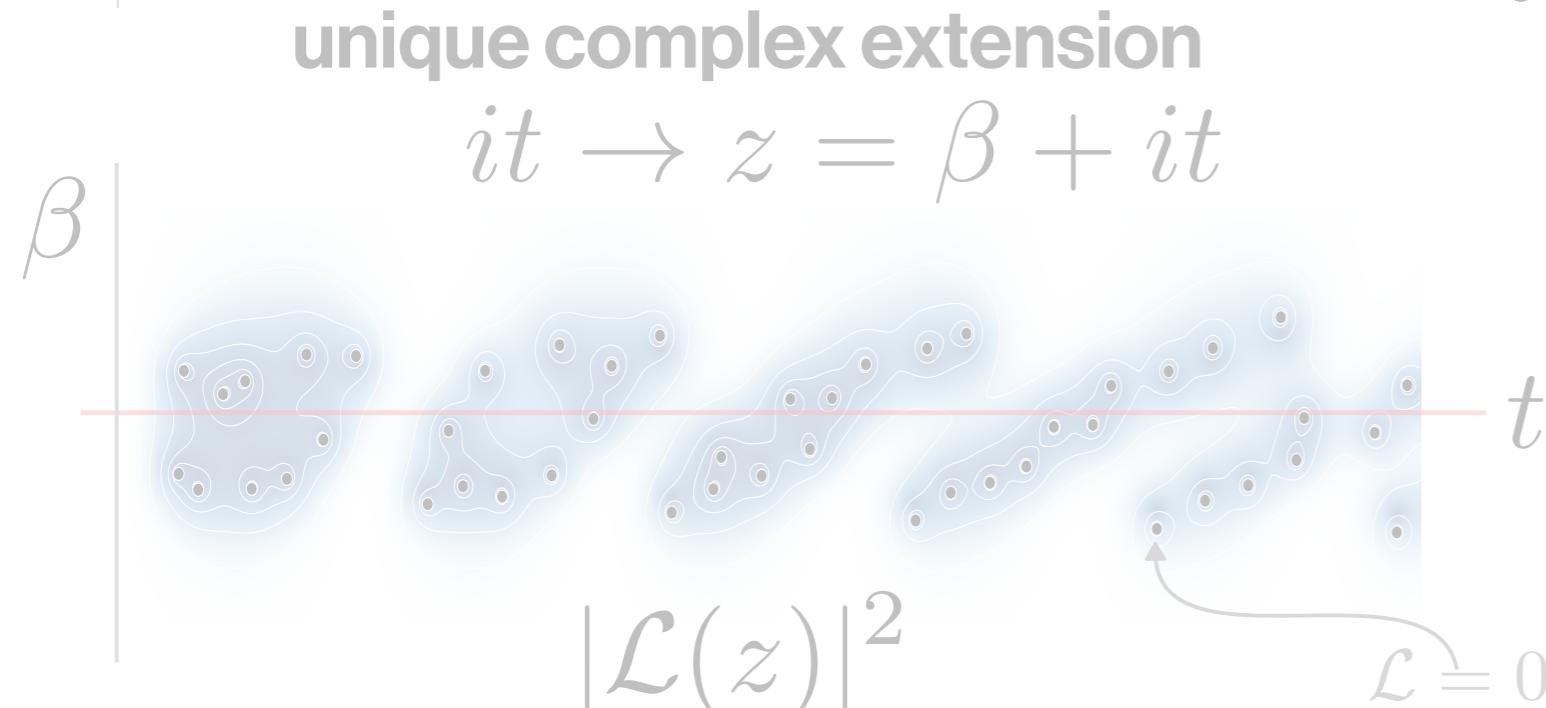
Long times

Short times

Remarks and Quenches

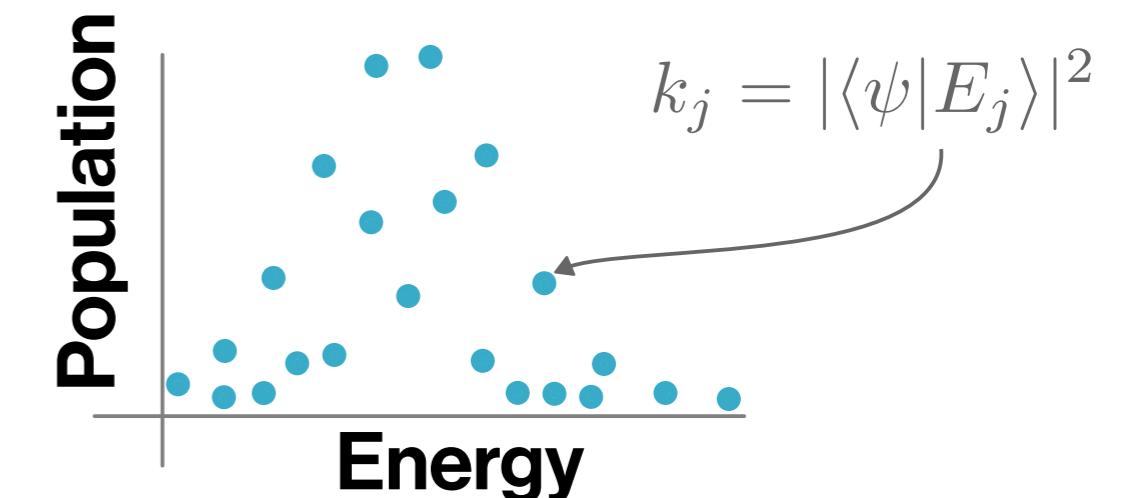
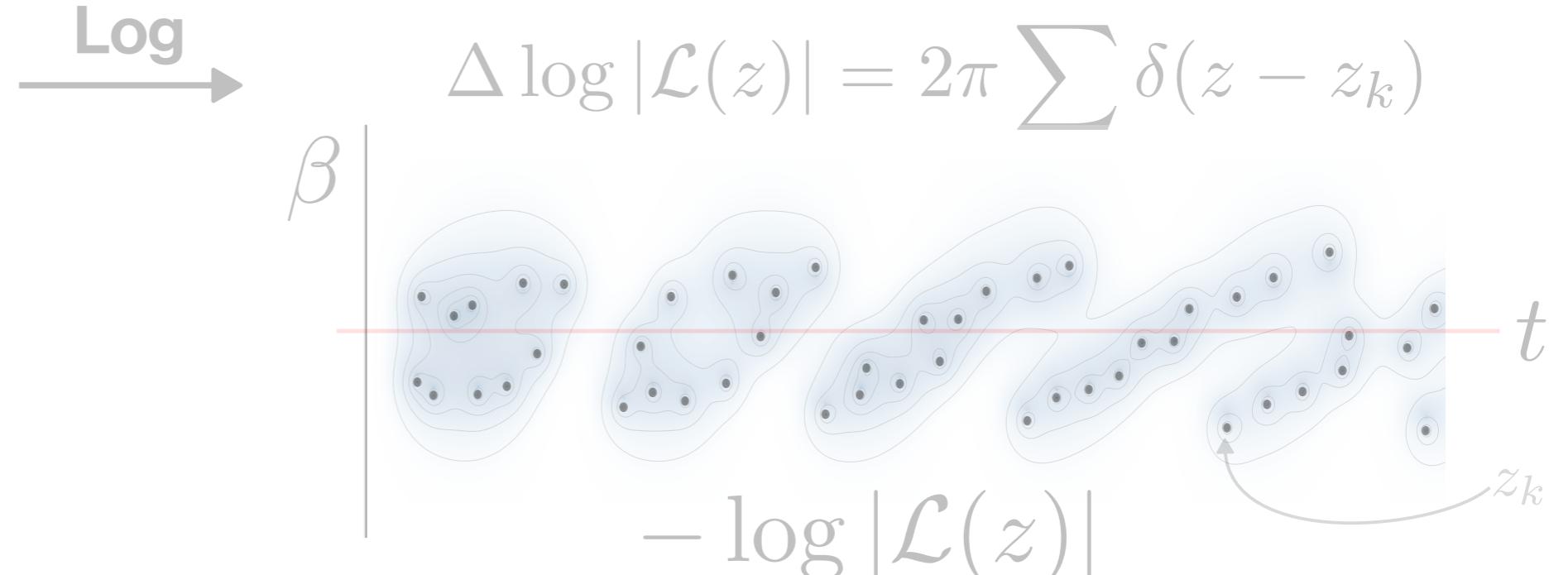
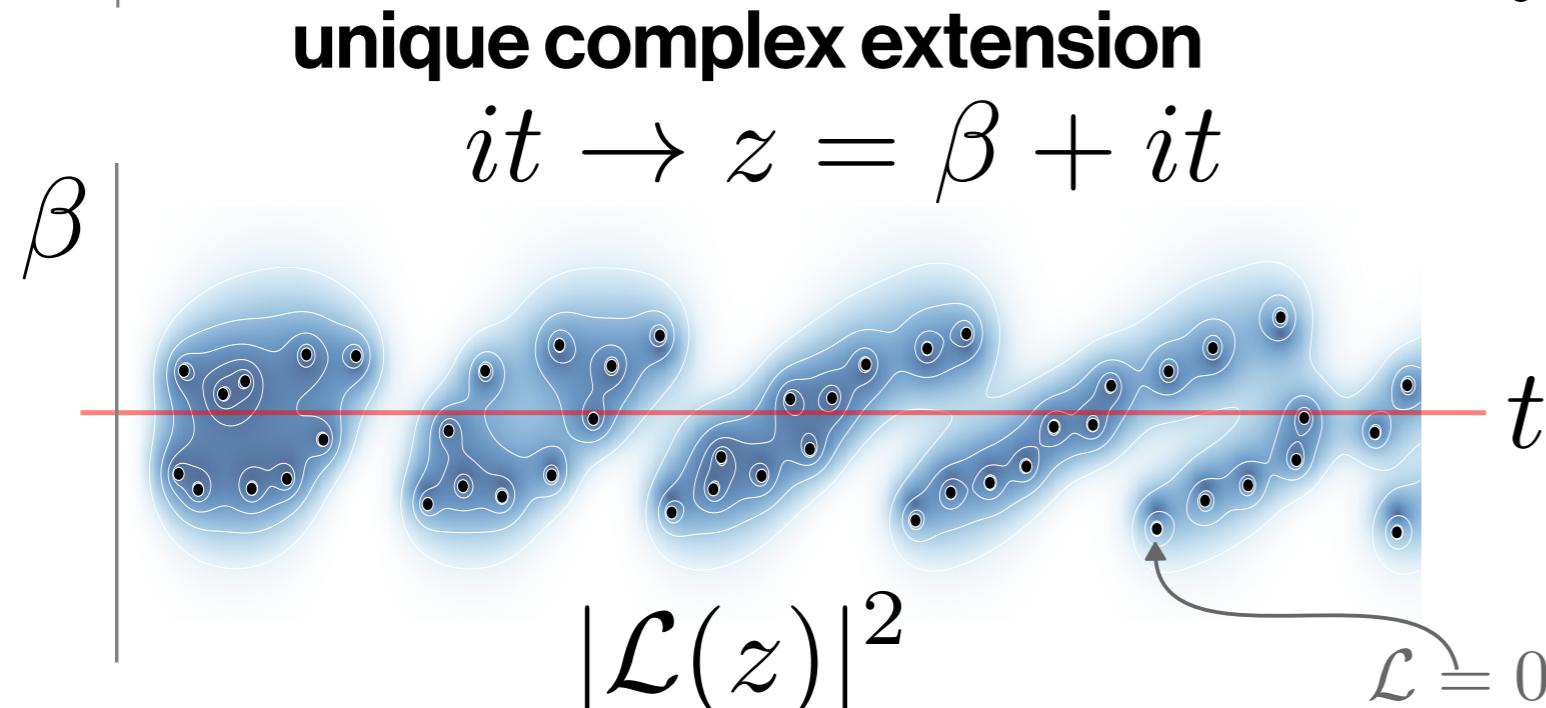
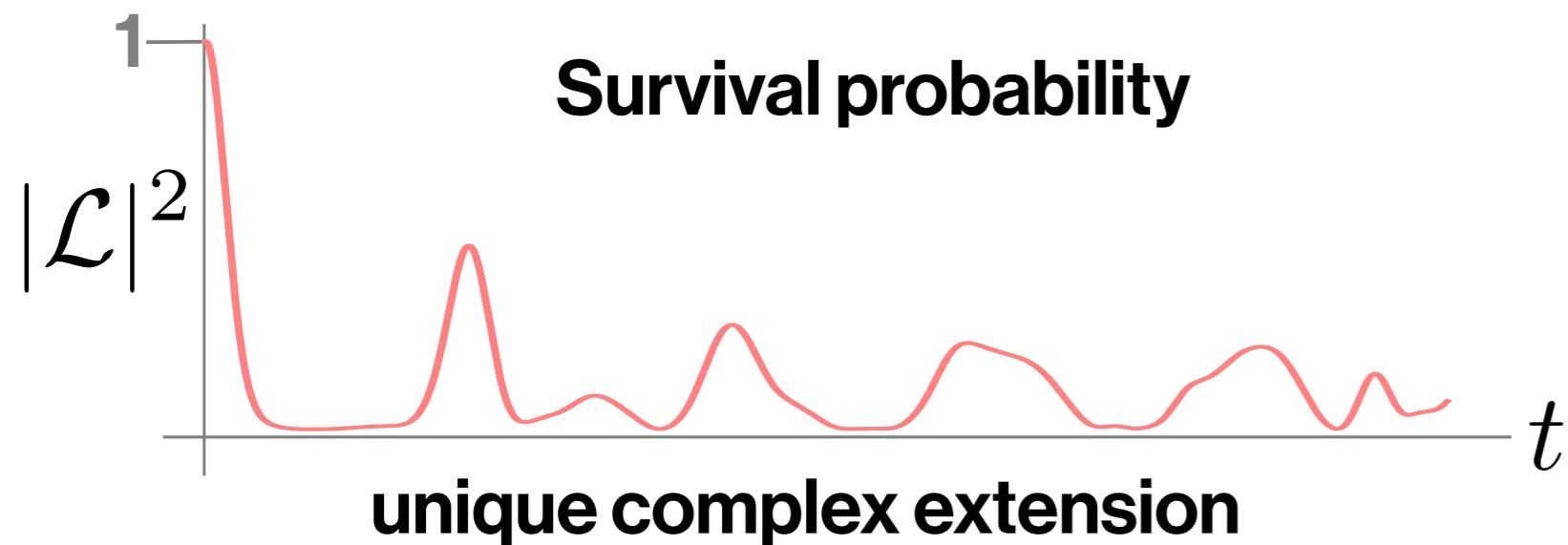
Survival amplitude

$$\mathcal{L}(t) = \langle \psi | \hat{U}(t) | \psi \rangle = \sum_j k_j e^{-iE_j t}$$



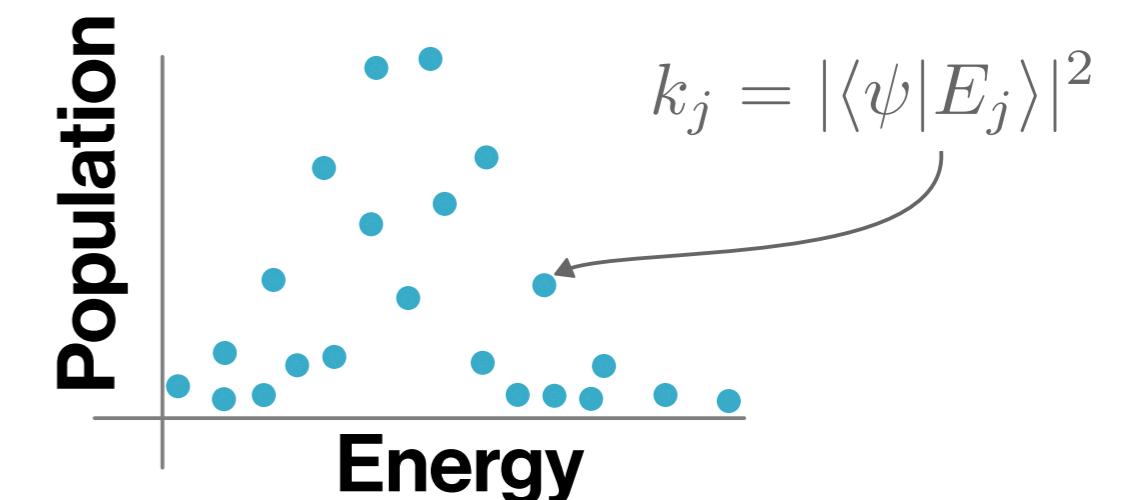
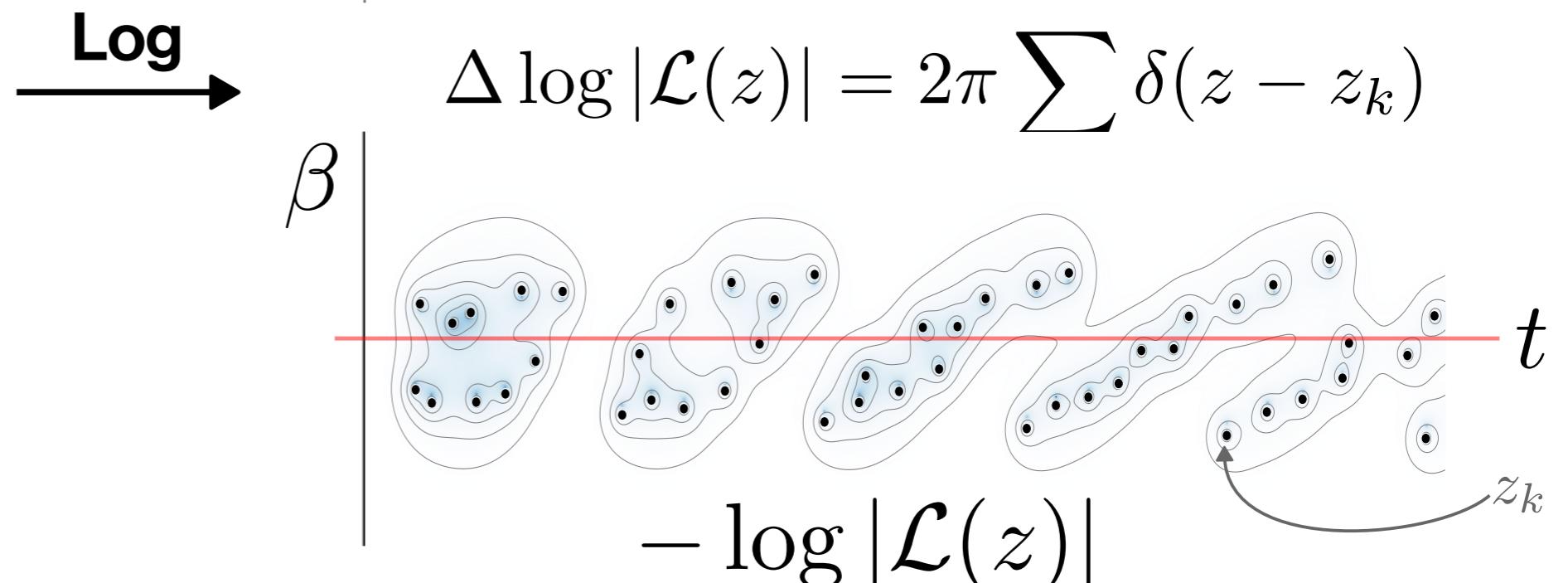
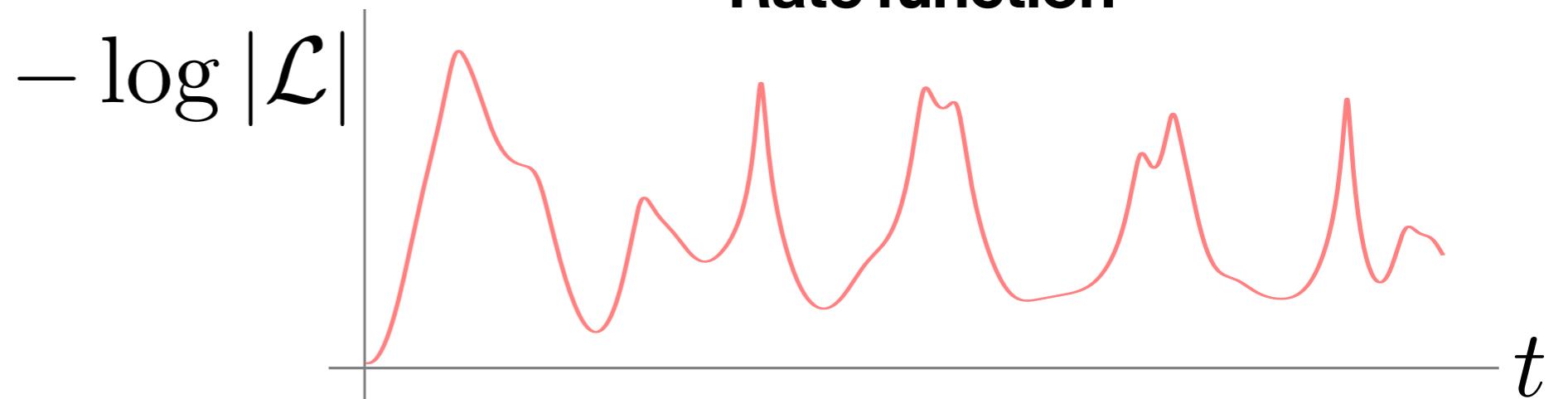
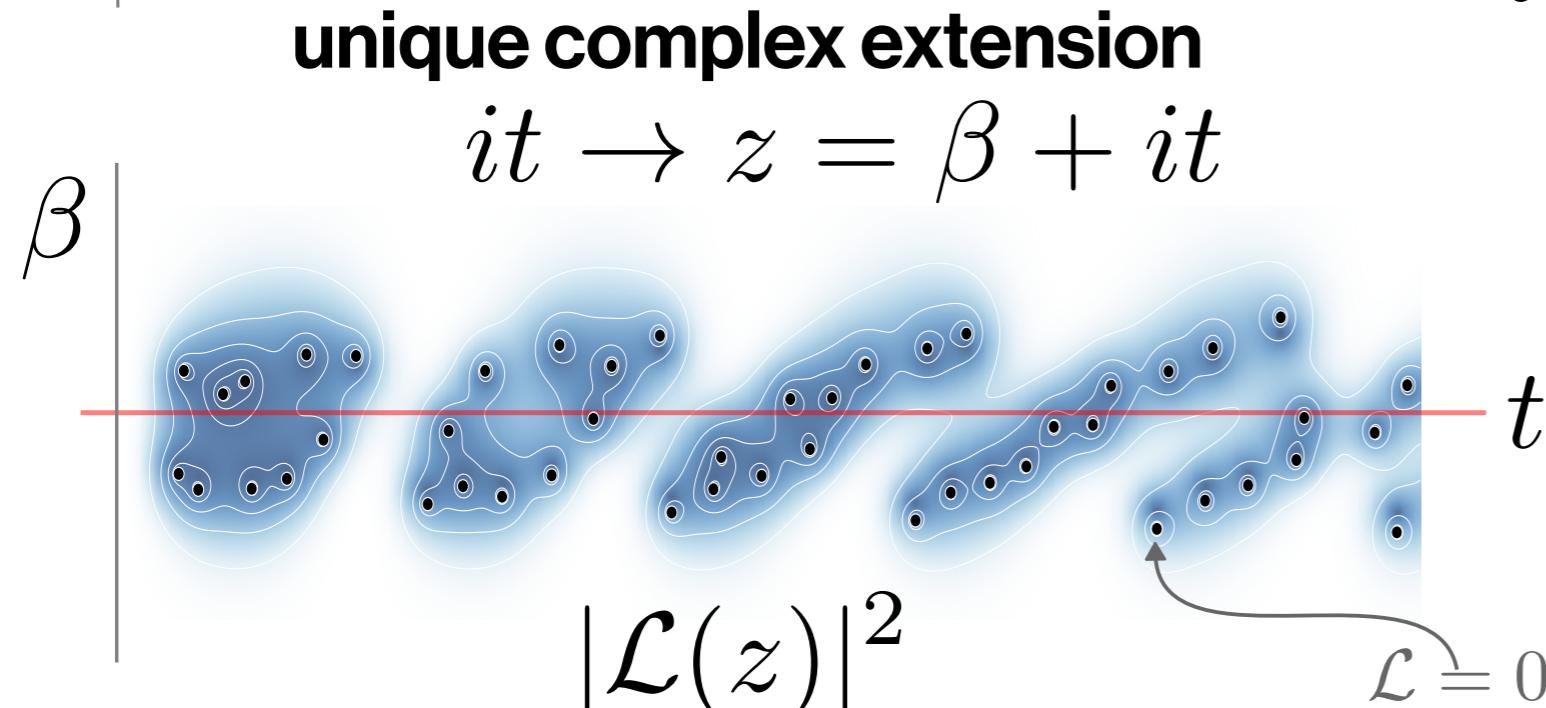
Survival amplitude

$$\mathcal{L}(t) = \langle \psi | \hat{U}(t) | \psi \rangle = \sum_j k_j e^{-iE_j t}$$



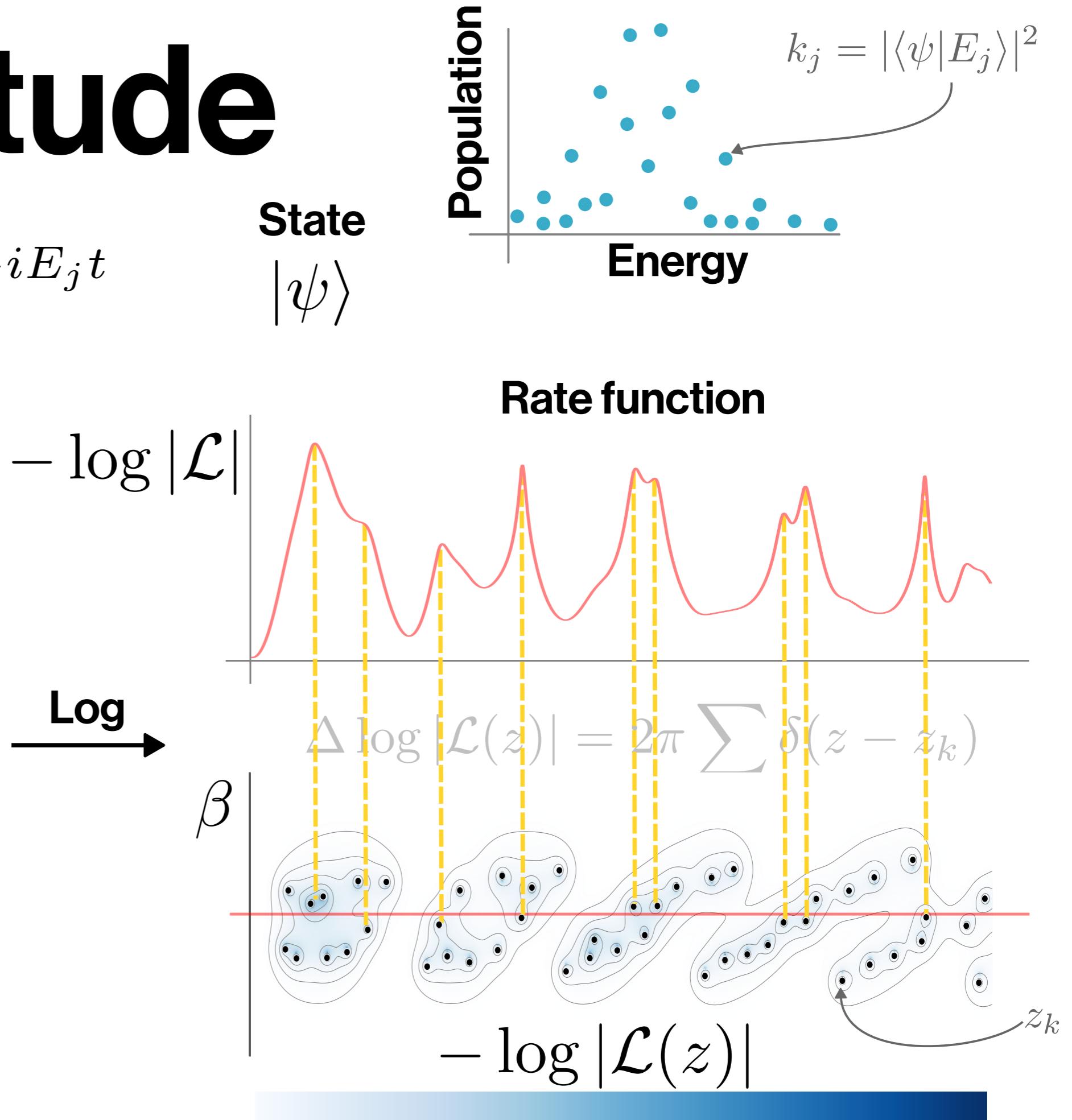
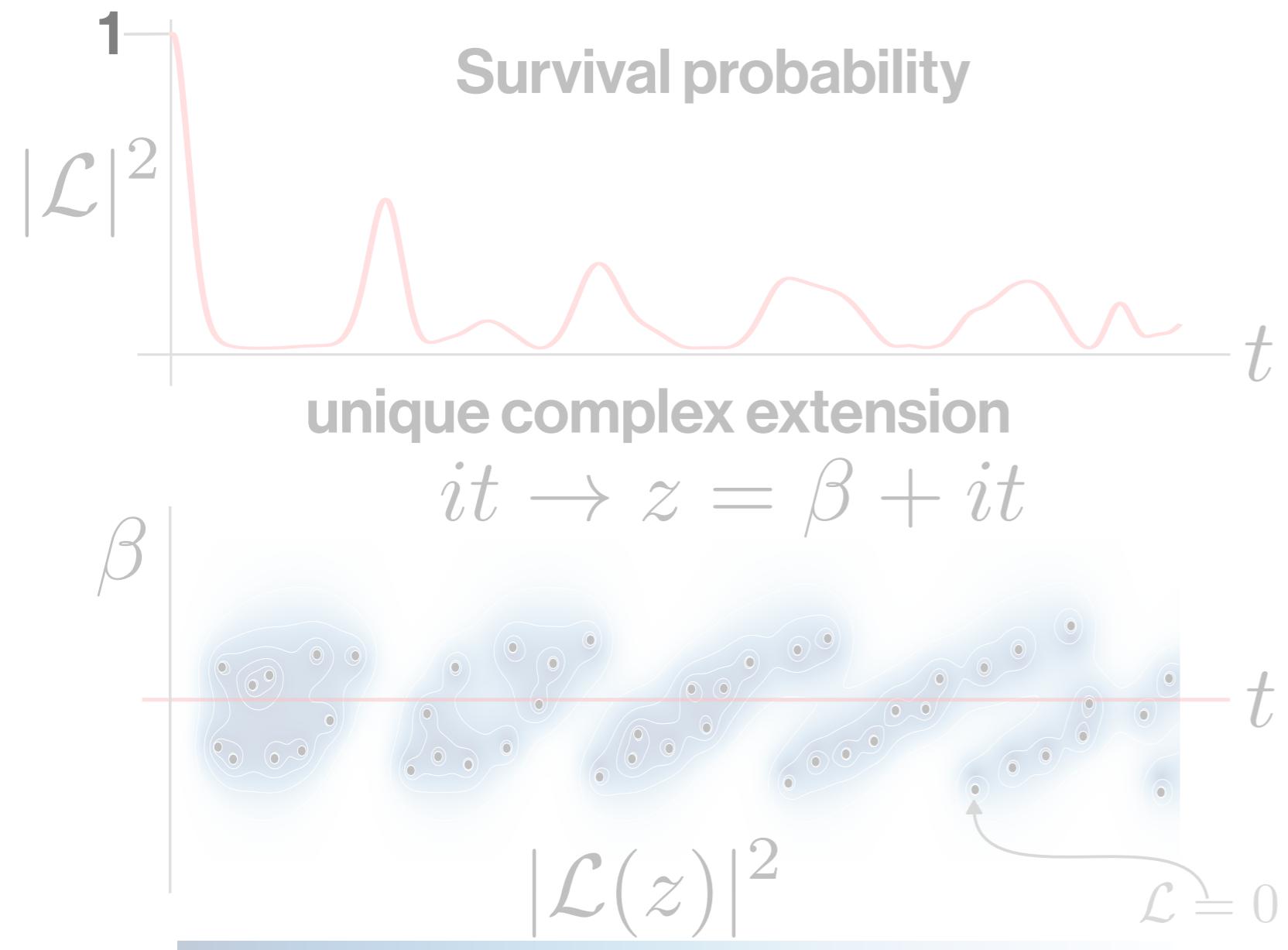
Survival amplitude

$$\mathcal{L}(t) = \langle \psi | \hat{U}(t) | \psi \rangle = \sum_j k_j e^{-iE_j t}$$



Survival amplitude

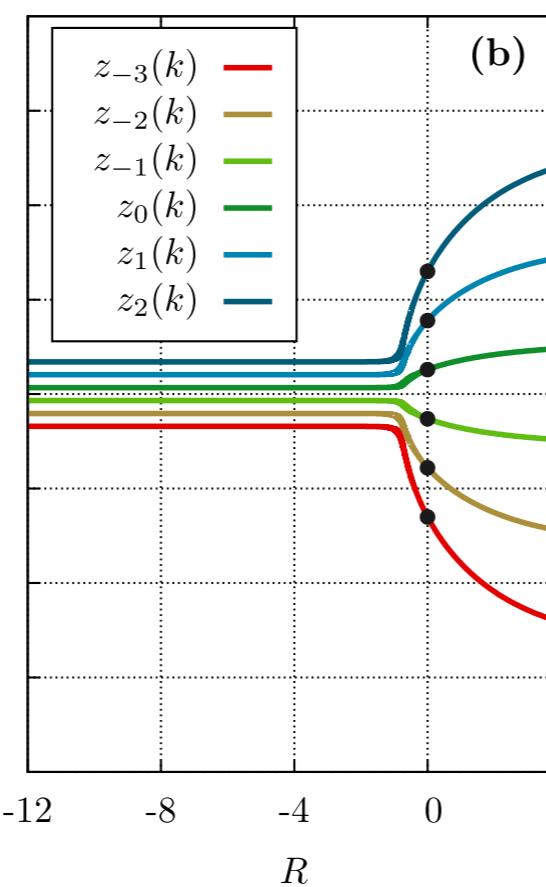
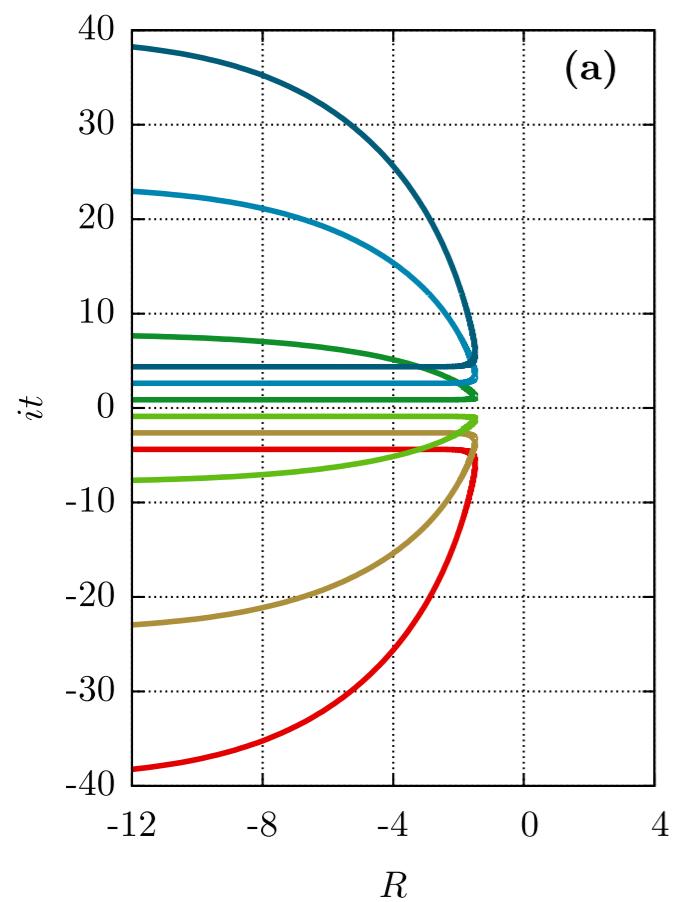
$$\mathcal{L}(t) = \langle \psi | \hat{U}(t) | \psi \rangle = \sum_j k_j e^{-iE_j t}$$



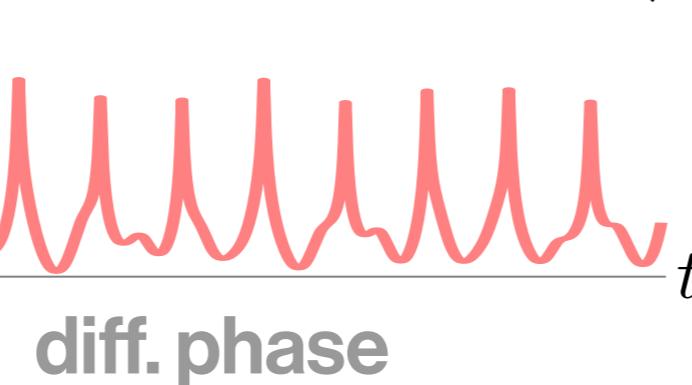
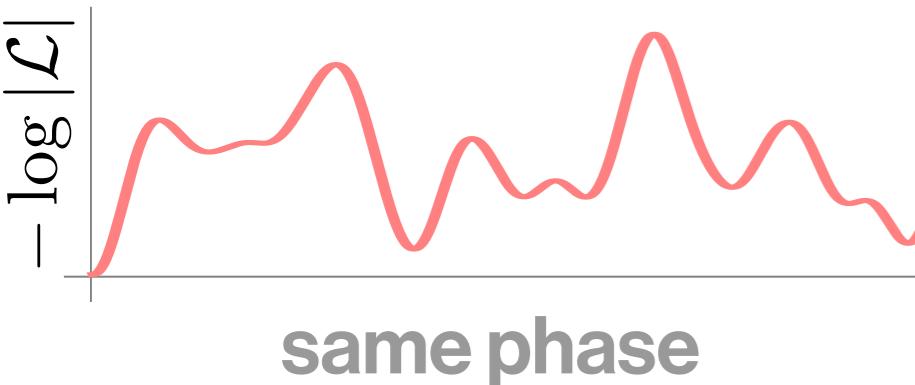
Heyl et al.
PRL 110, 135704 (2013)

Quench
same phase

$$\hat{H} = -J \sum_{i=1} \hat{\sigma}_i^z \hat{\sigma}_{i+1}^z + h \sum_{i=1} \hat{\sigma}_i^x$$

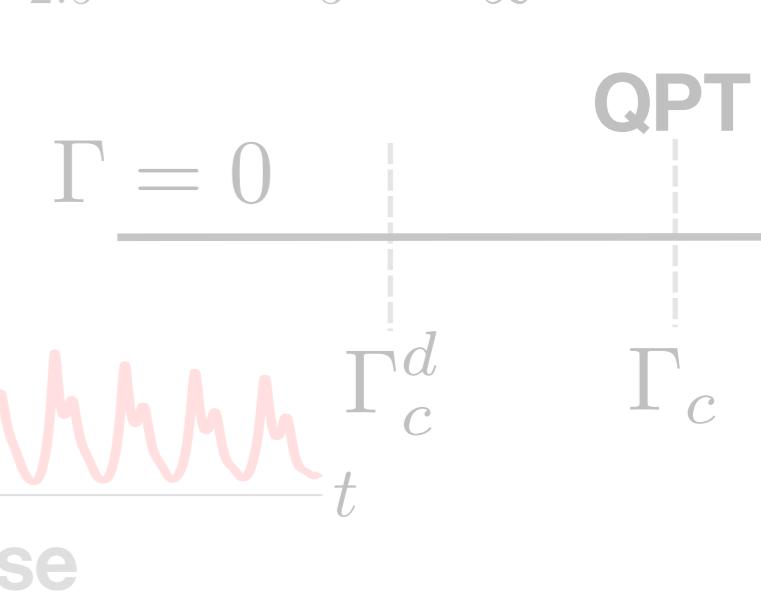
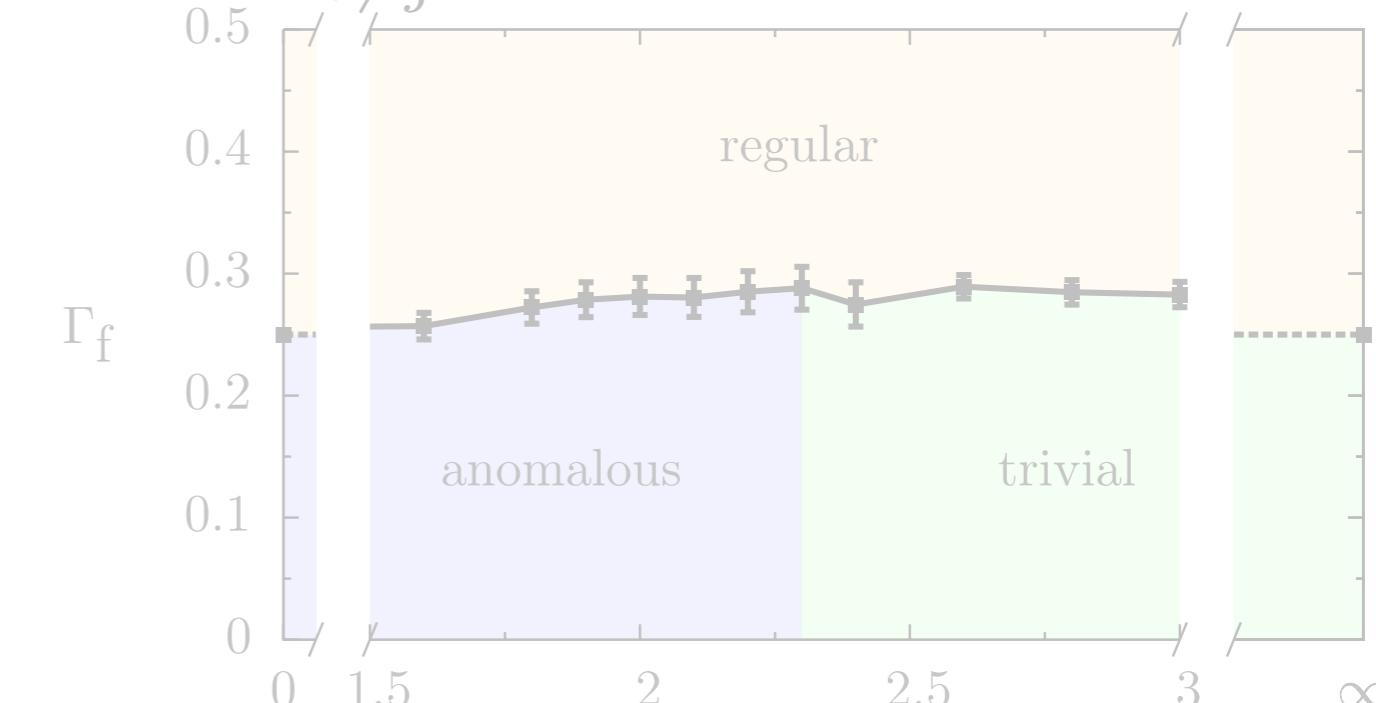


$N \rightarrow \infty$
DQPT (-II)



Homrighausen et al.
PRB 96, 104436 (2017)

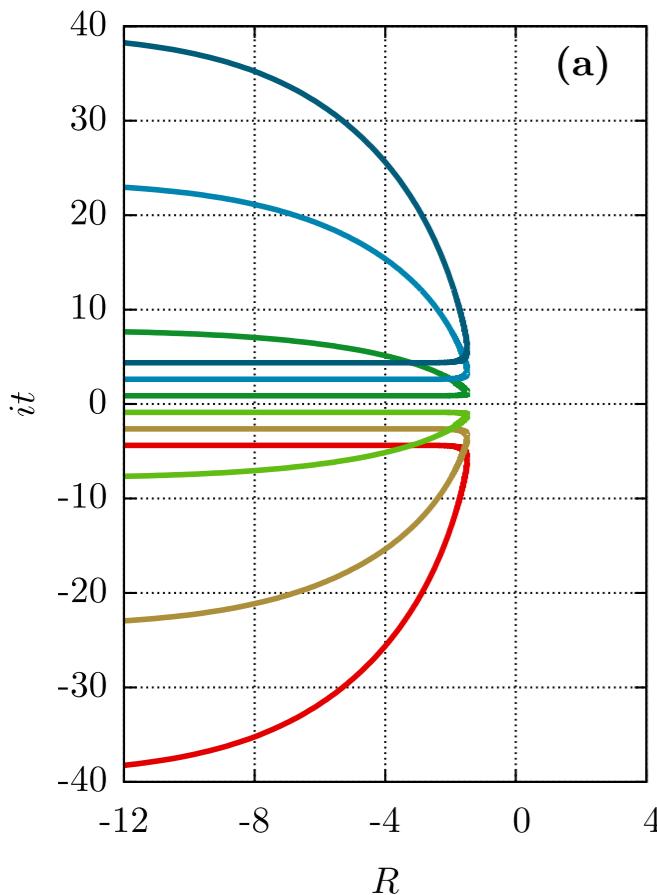
$$\hat{H} = -\frac{J}{2N} \sum_{i \neq j} \frac{1}{|i-j^\alpha|} \hat{S}_i^z \hat{S}_{i+1}^z + \Gamma \sum_{i=1} \hat{S}_i^x$$



Why zeros?

Heyl et al.
PRL 110, 135704 (2013)

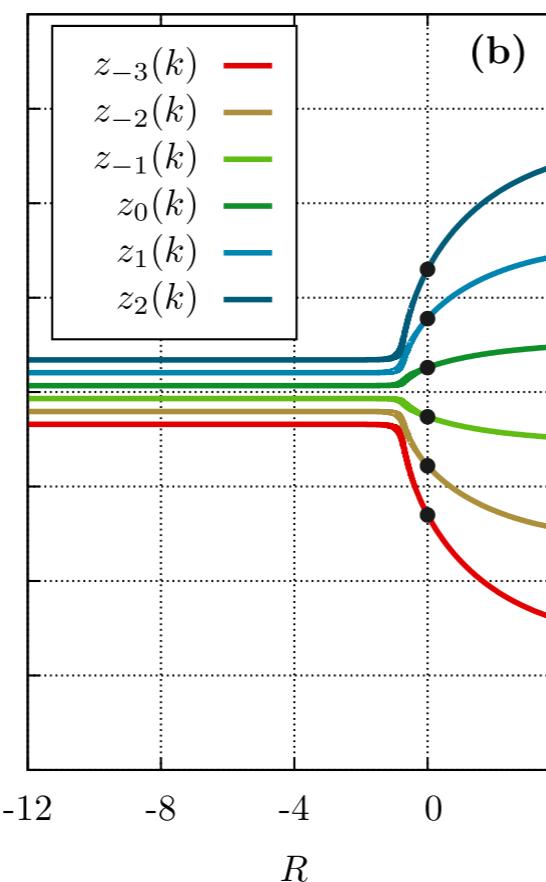
Quench
same phase



same phase

$$\hat{H} = -J \sum_{i=1} \hat{\sigma}_i^z \hat{\sigma}_{i+1}^z + h \sum_{i=1} \hat{\sigma}_i^x$$

QPT
diff. phase



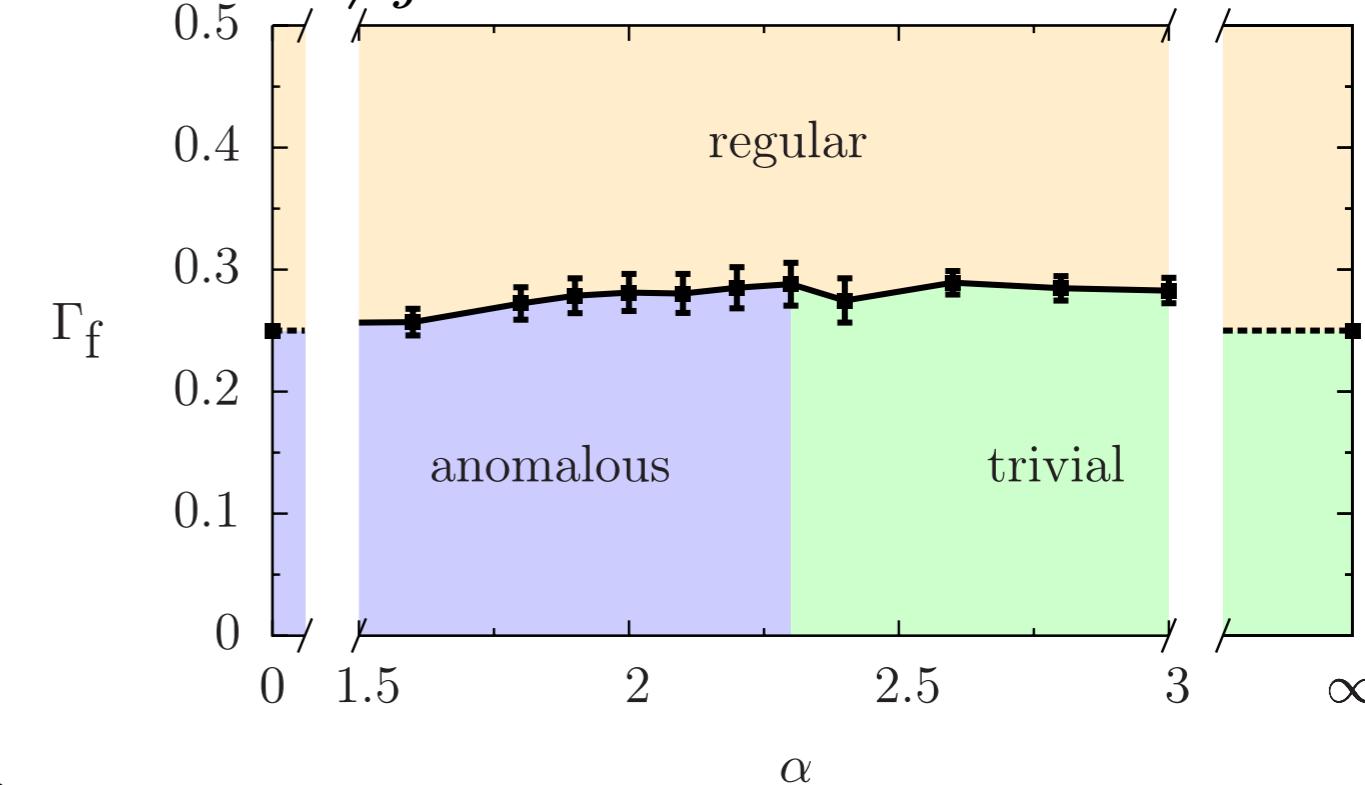
diff. phase

$N \rightarrow \infty$
DQPT (-II)

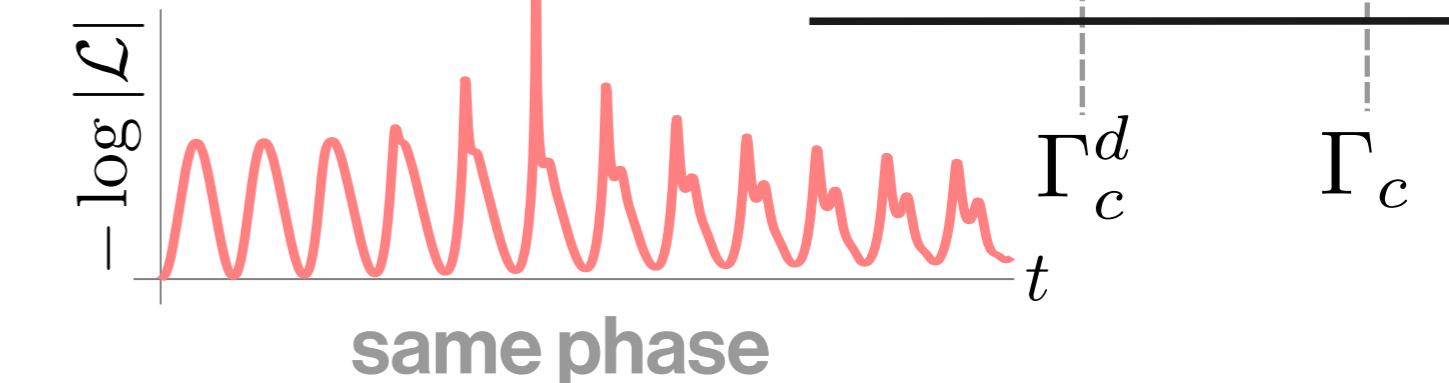
Why zeros?

Homrighausen et al.
PRB 96, 104436 (2017)

$$\hat{H} = -\frac{J}{2N} \sum_{i \neq j} \frac{1}{|i-j^\alpha|} \hat{S}_i^z \hat{S}_{i+1}^z + \Gamma \sum_{i=1} \hat{S}_i^x$$



QPT



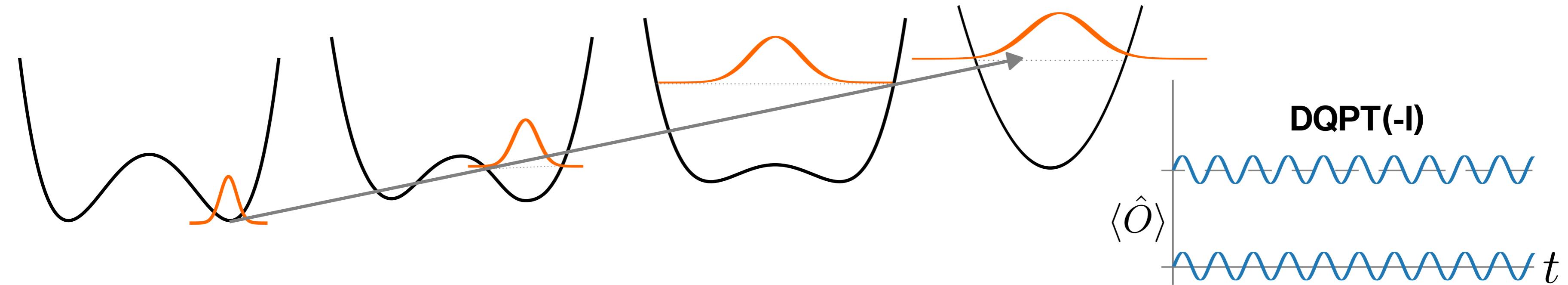
same phase

Γ_c^d

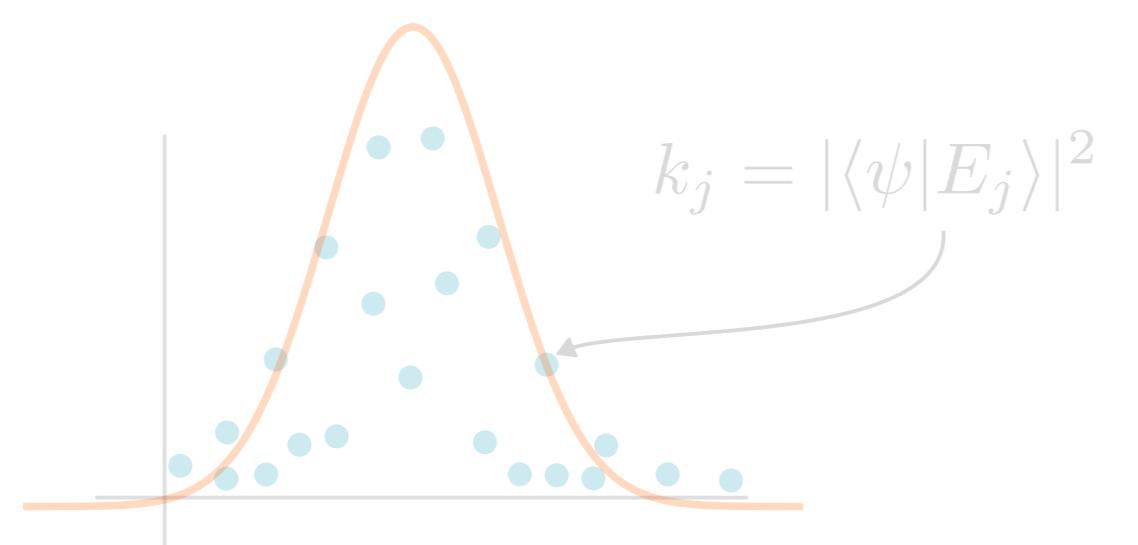
Γ_c

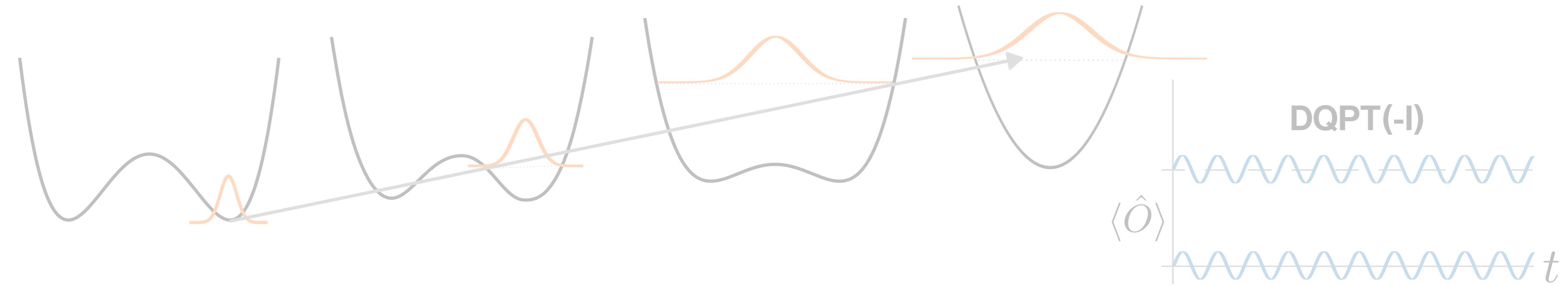
Γ_c^d

Γ_c

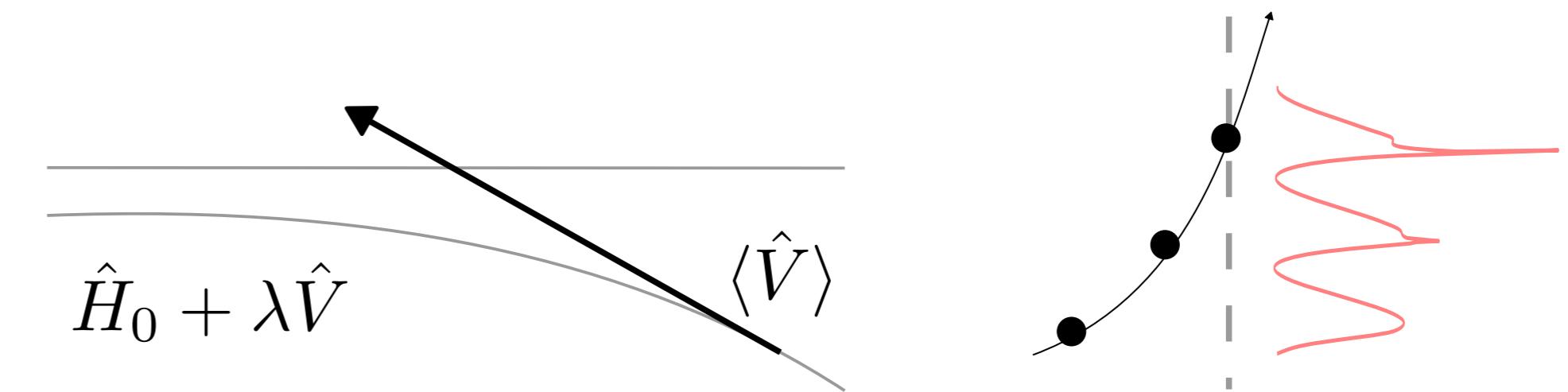
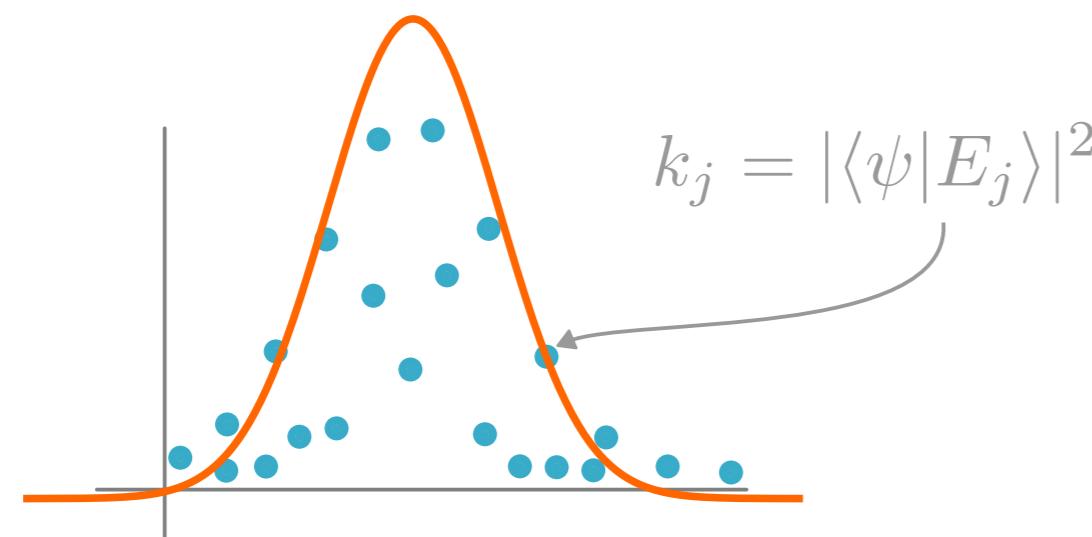


What determines the distribution of the zeros?





What determines the distribution of the zeros?



Zeros of holomorphic functions

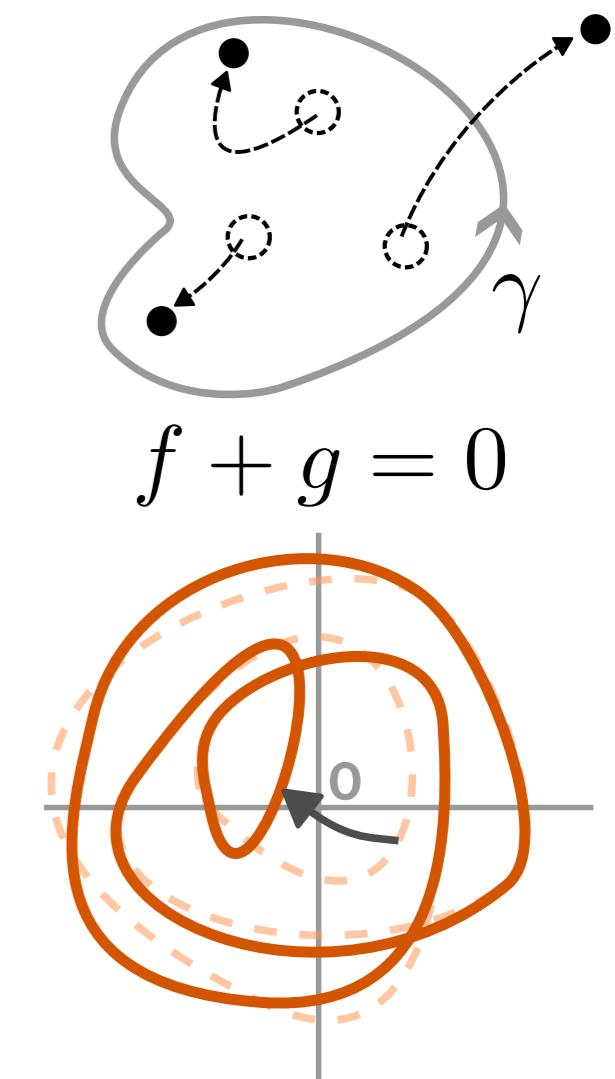
Zeros cannot be created or annihilated by a hol. perturbation

$$\# \text{ zeros of } f(z) \text{ in } \text{# zeros of } f(z) \text{ in }$$
$$= W[f \circ \gamma] = W[f \circ \gamma] = W[f \circ \gamma]$$

Num. of zeros is determined by the dominant part

$$\# \text{ zeros in } f(z) = \# \text{ zeros in } f(z) + g(z) \quad \text{if } g < f \text{ on } \gamma$$

Rouché's theorem



Zeros of holomorphic functions

Zeros cannot be created or annihilated by a hol. perturbation

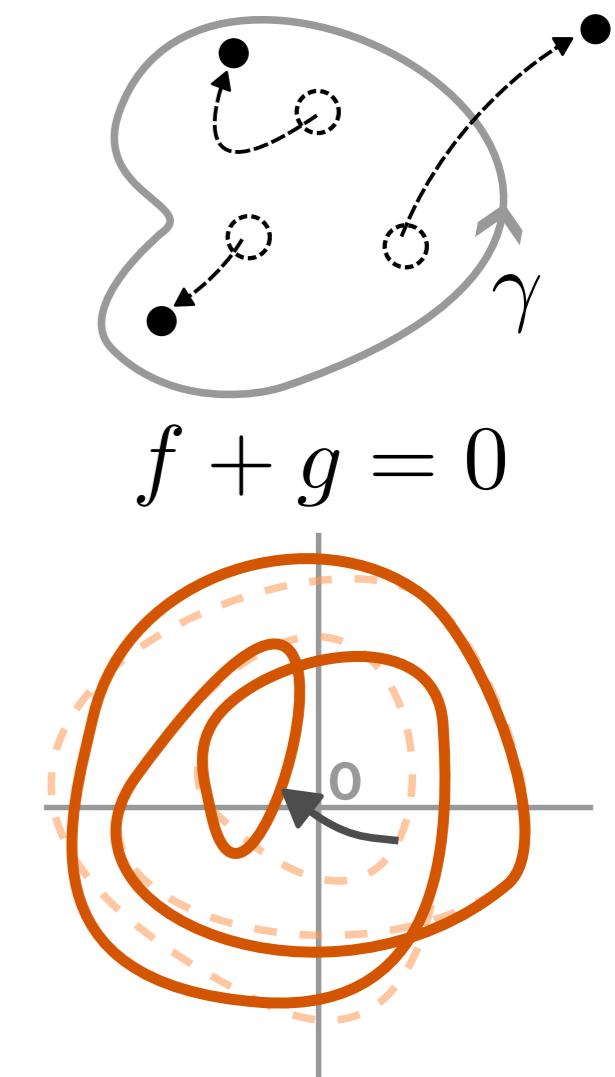
$$\# \text{ zeros of } f(z) \text{ in } \text{# zeros of } f(z) \text{ in }$$
$$= W[f \circ \gamma] = W[f \circ \gamma] = W[f \circ \gamma]$$

winding number

Num. of zeros is determined by the dominant part

$$\# \text{ zeros in } f(z) = \# \text{ zeros in } f(z) + g(z) \quad \text{if } g < f \text{ on } \gamma$$

Rouché's theorem

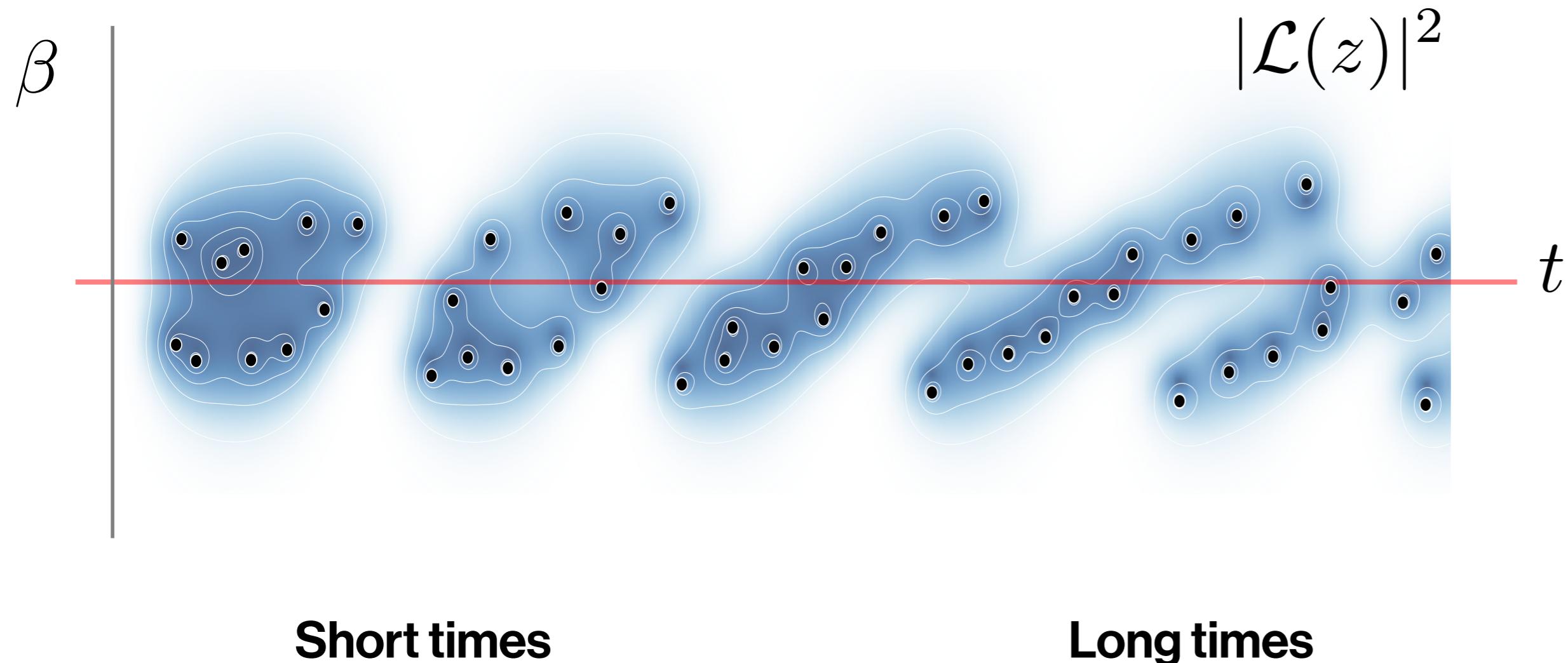


Different time scales

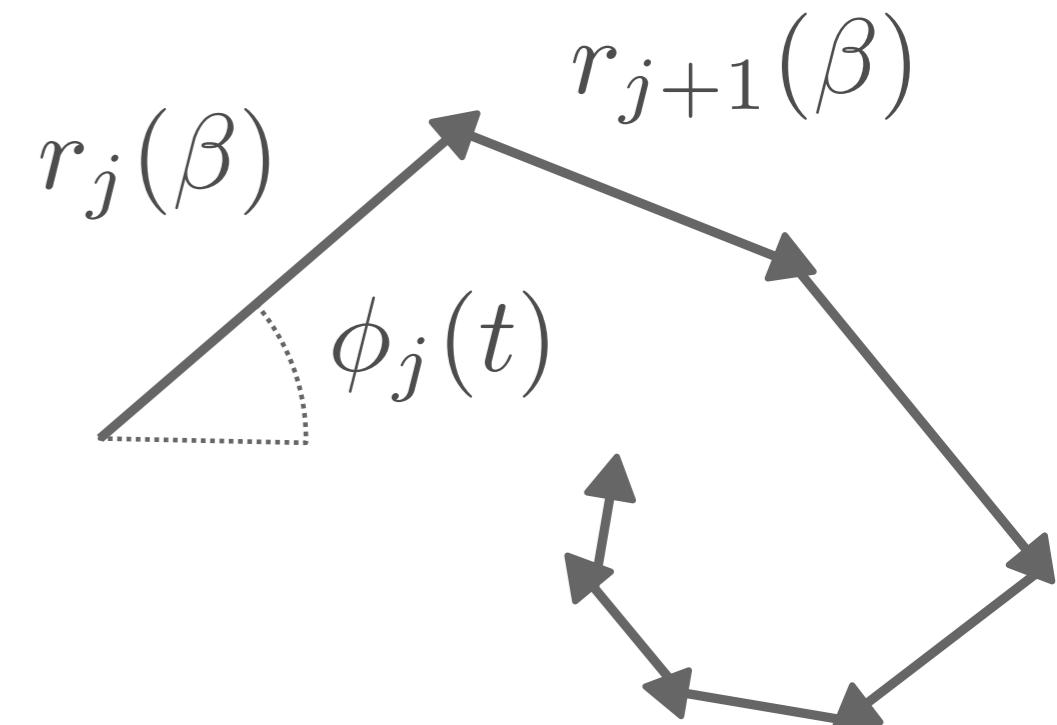
Level spacing
 $\Delta_i = \Delta + \varepsilon_i$



Dephasing



$$\mathcal{L}(\beta + it) = \sum_j k_j e^{-\beta E_j} e^{-i E_j t}$$
$$r_j(\beta)$$

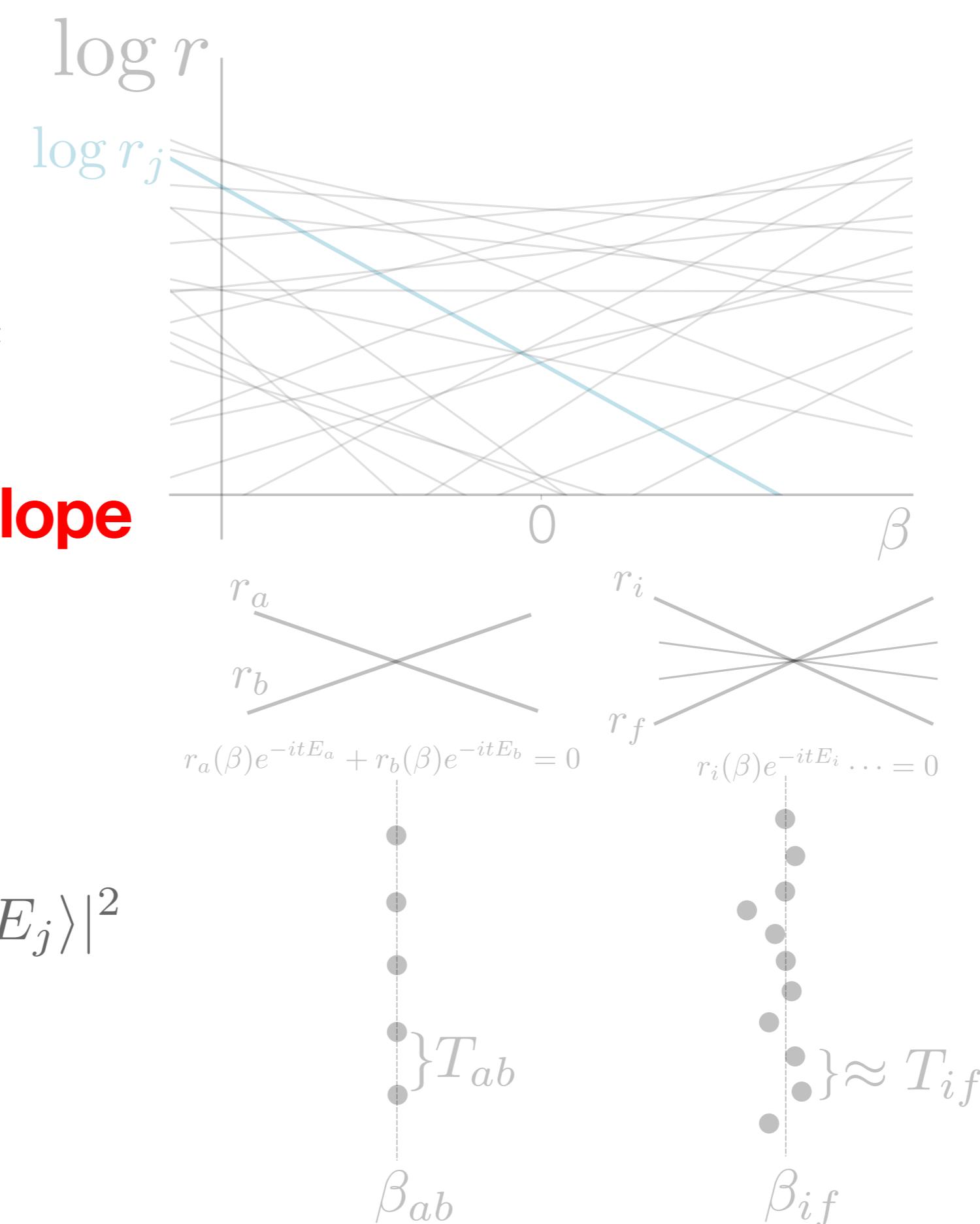
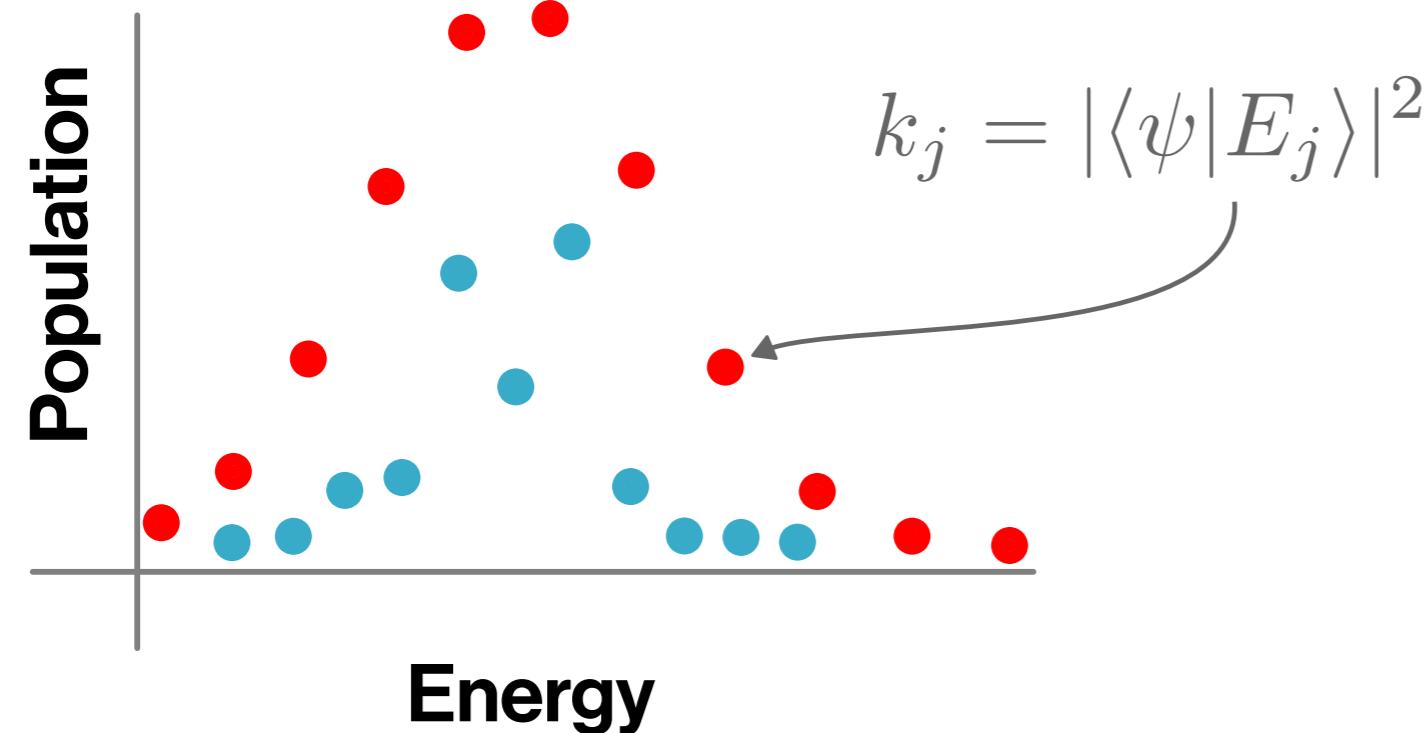


Long times

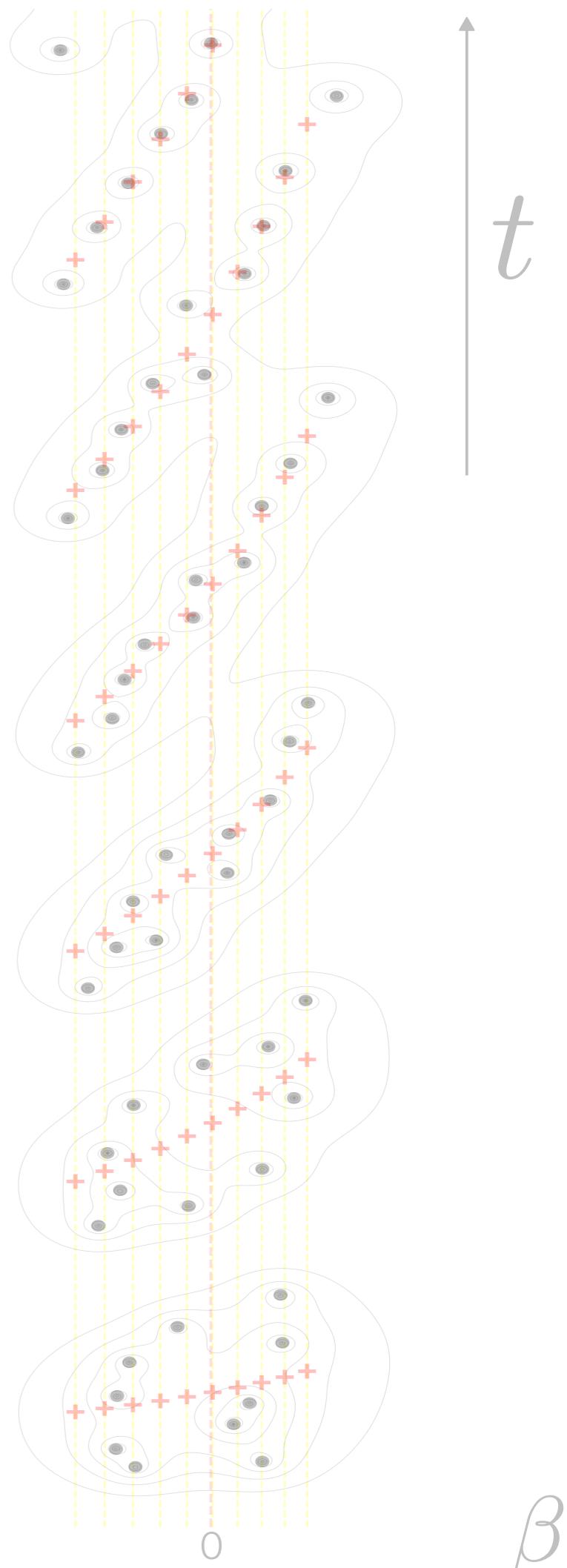
$$\mathcal{L}(\beta + it) = \sum_j r_j(\beta) e^{-iE_j t}$$

The dominant part = **The Envelope**

$$k_a \left(\frac{k_a}{k_b} \right)^{-\frac{\Delta_{ac}}{\Delta_{ab}}} \geq k_c, \quad \forall k_c, E_c$$



$$z_{ab}^*(n) = \frac{1}{\Delta_{ab}} \left(\log \left(\frac{k_a}{k_b} \right) + i\pi (2n+1) \right)$$

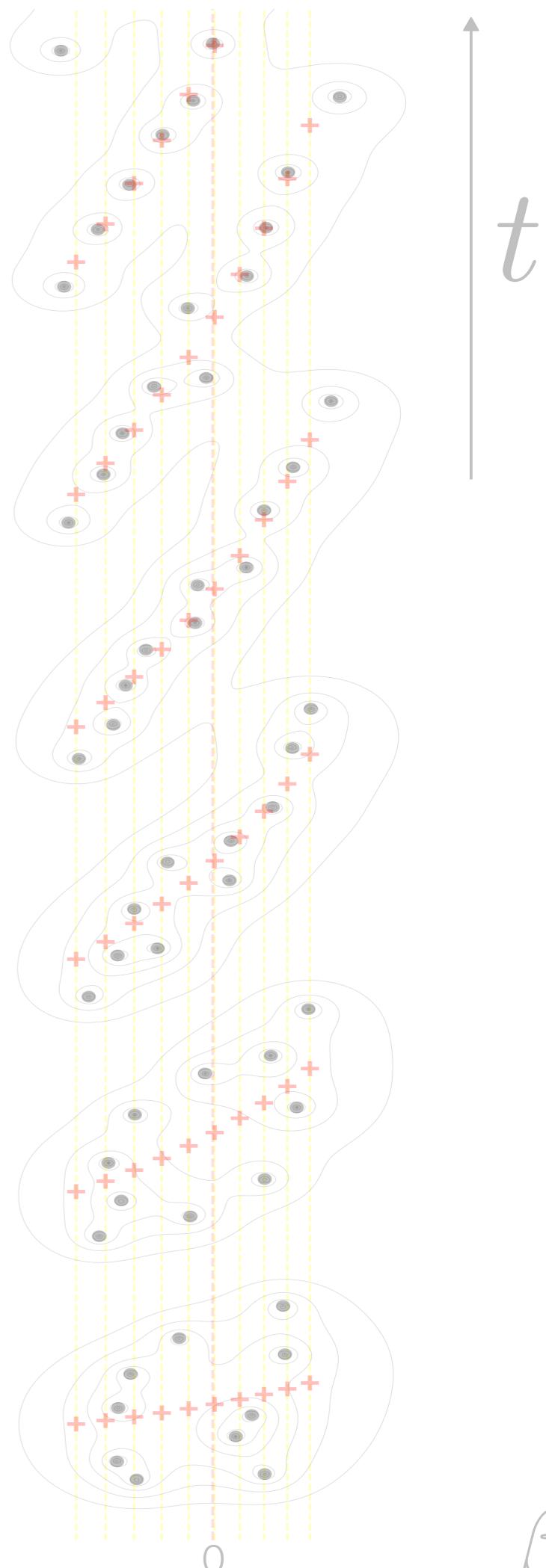
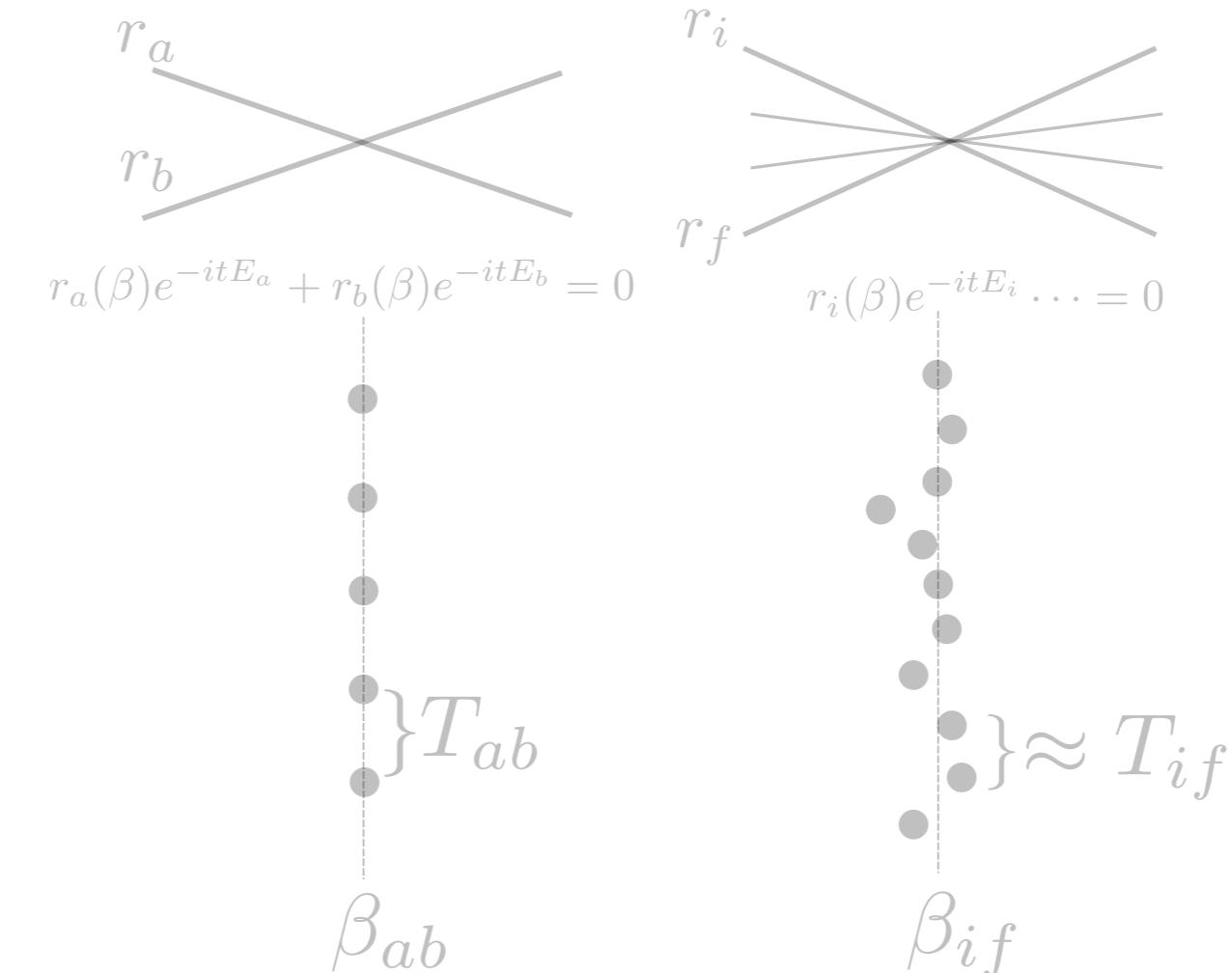
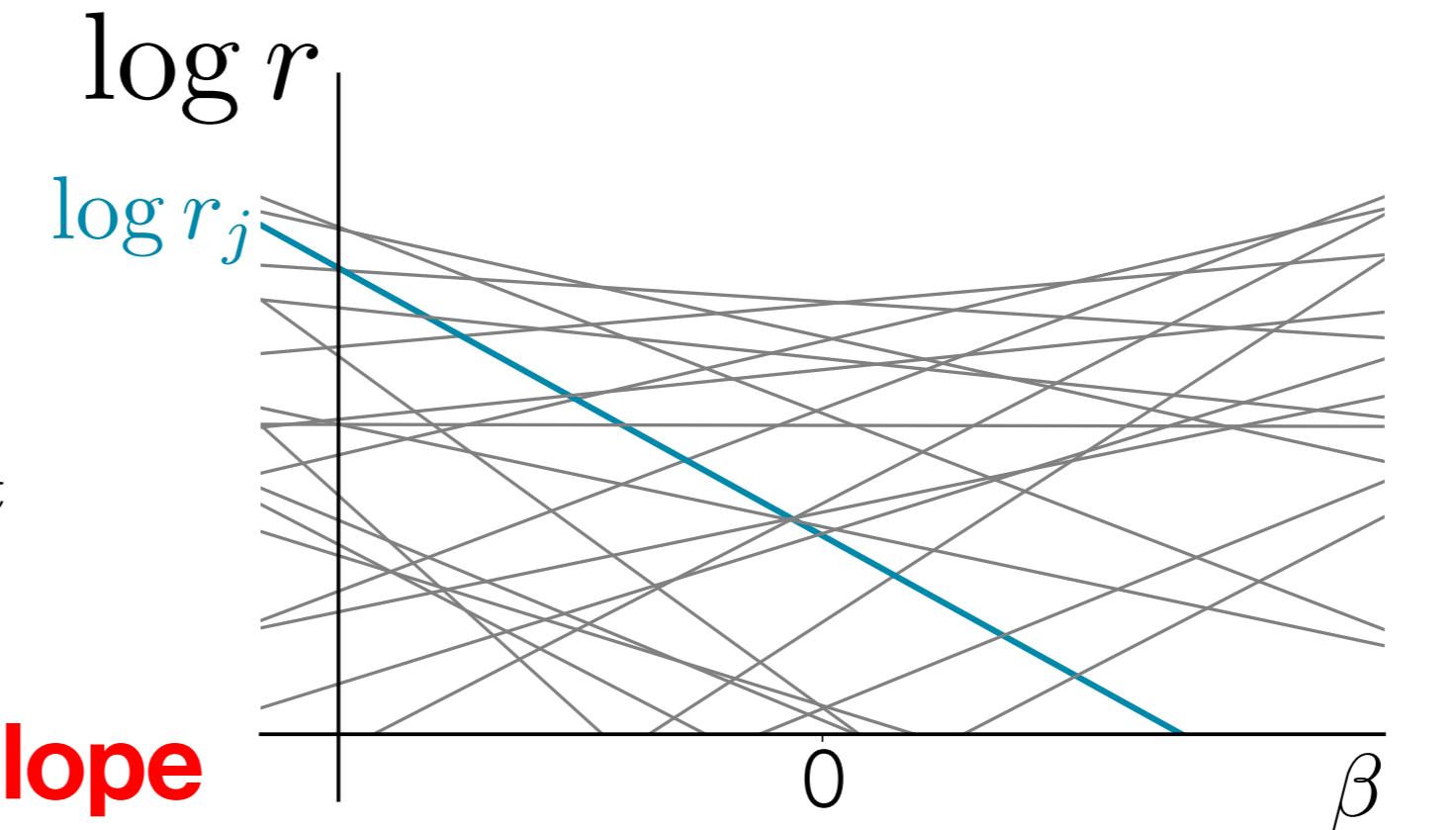
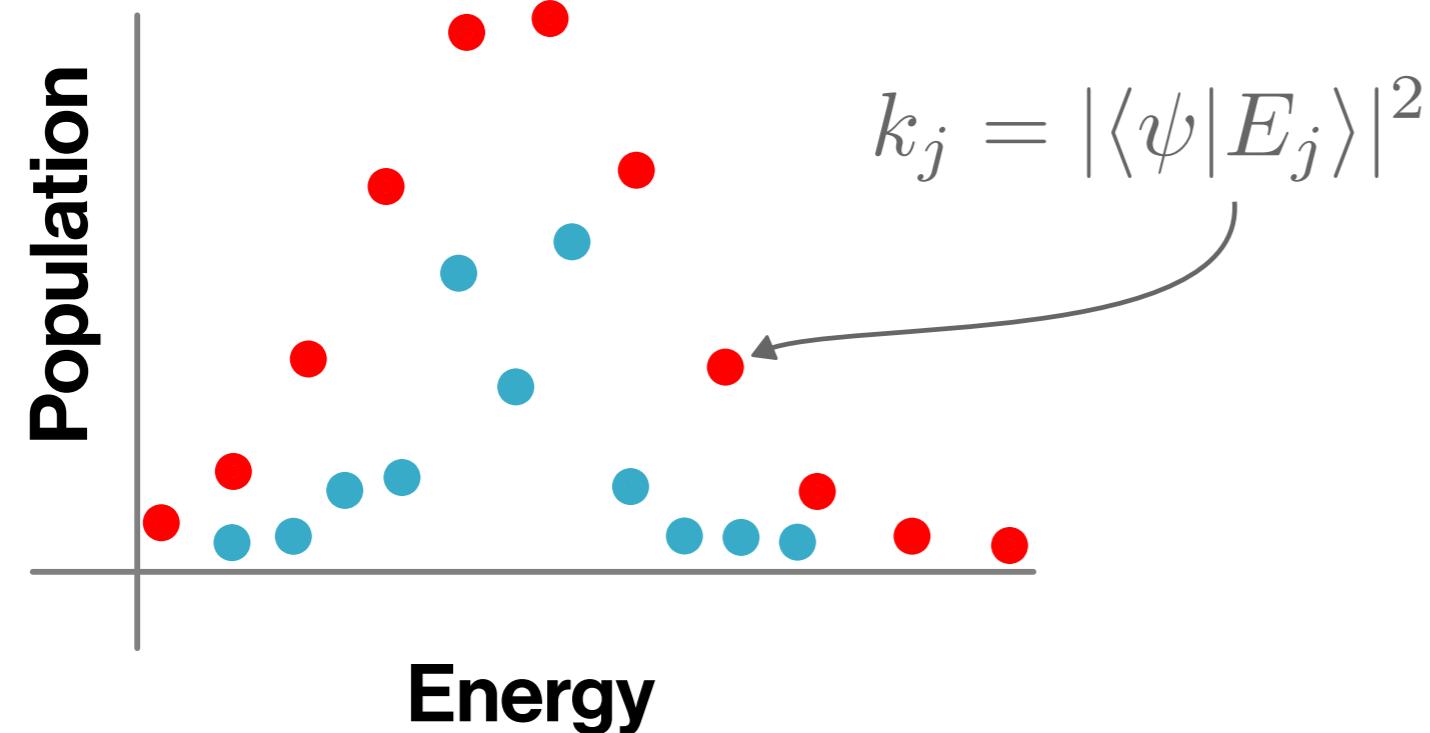


Long times

$$\mathcal{L}(\beta + it) = \sum_j r_j(\beta) e^{-iE_j t}$$

The dominant part = **The Envelope**

$$k_a \left(\frac{k_a}{k_b} \right)^{-\frac{\Delta_{ac}}{\Delta_{ab}}} \geq k_c, \quad \forall k_c, E_c$$

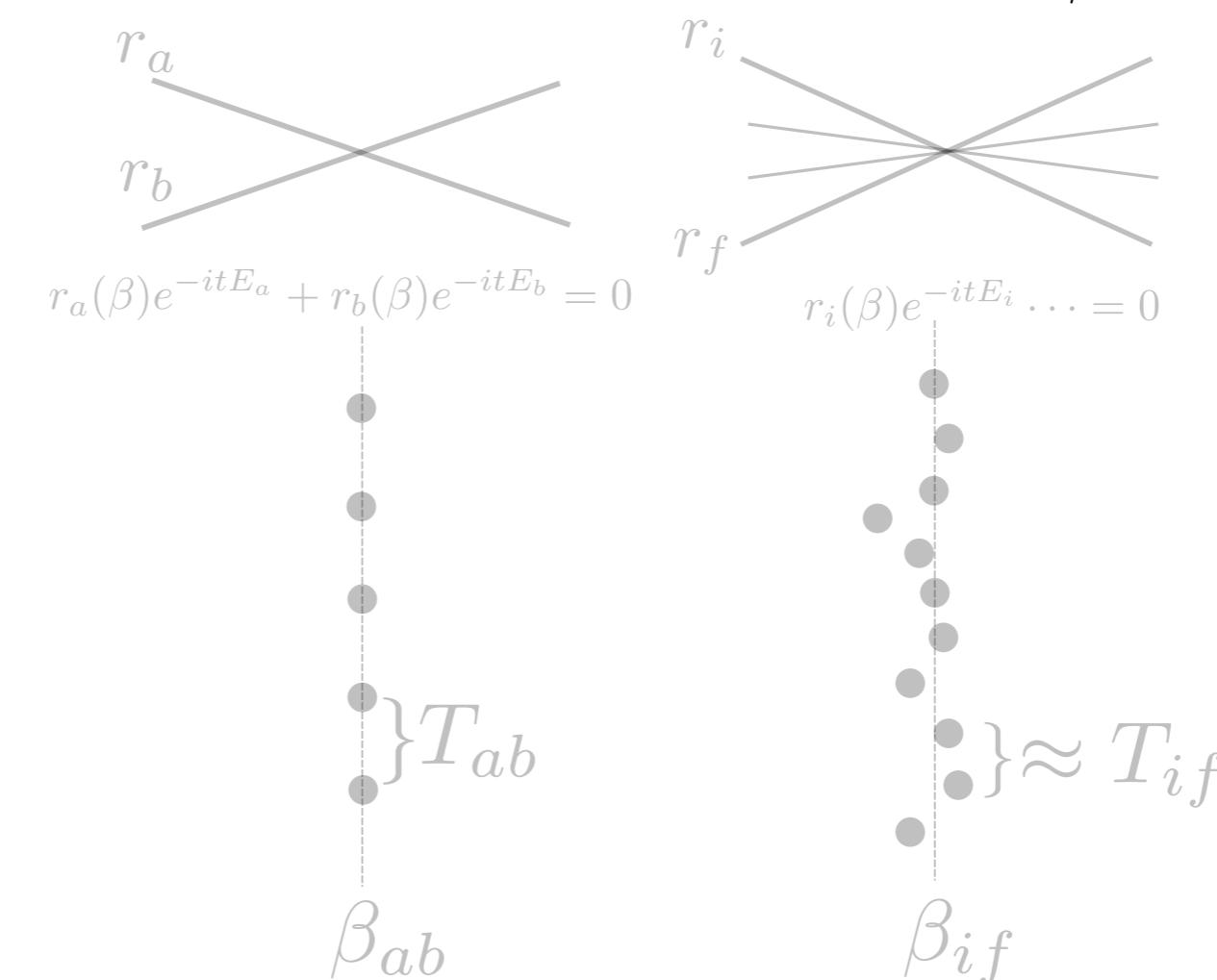
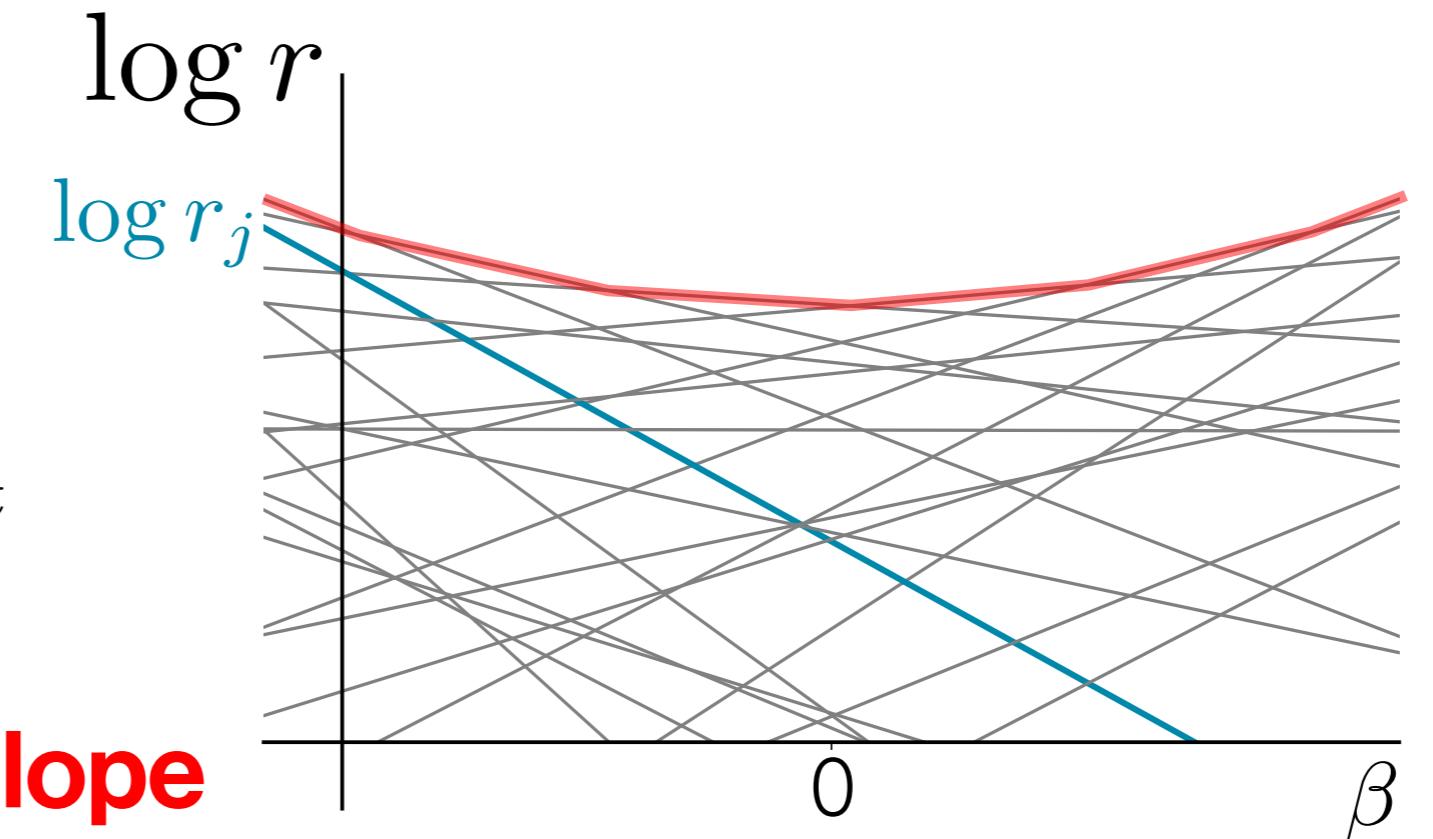
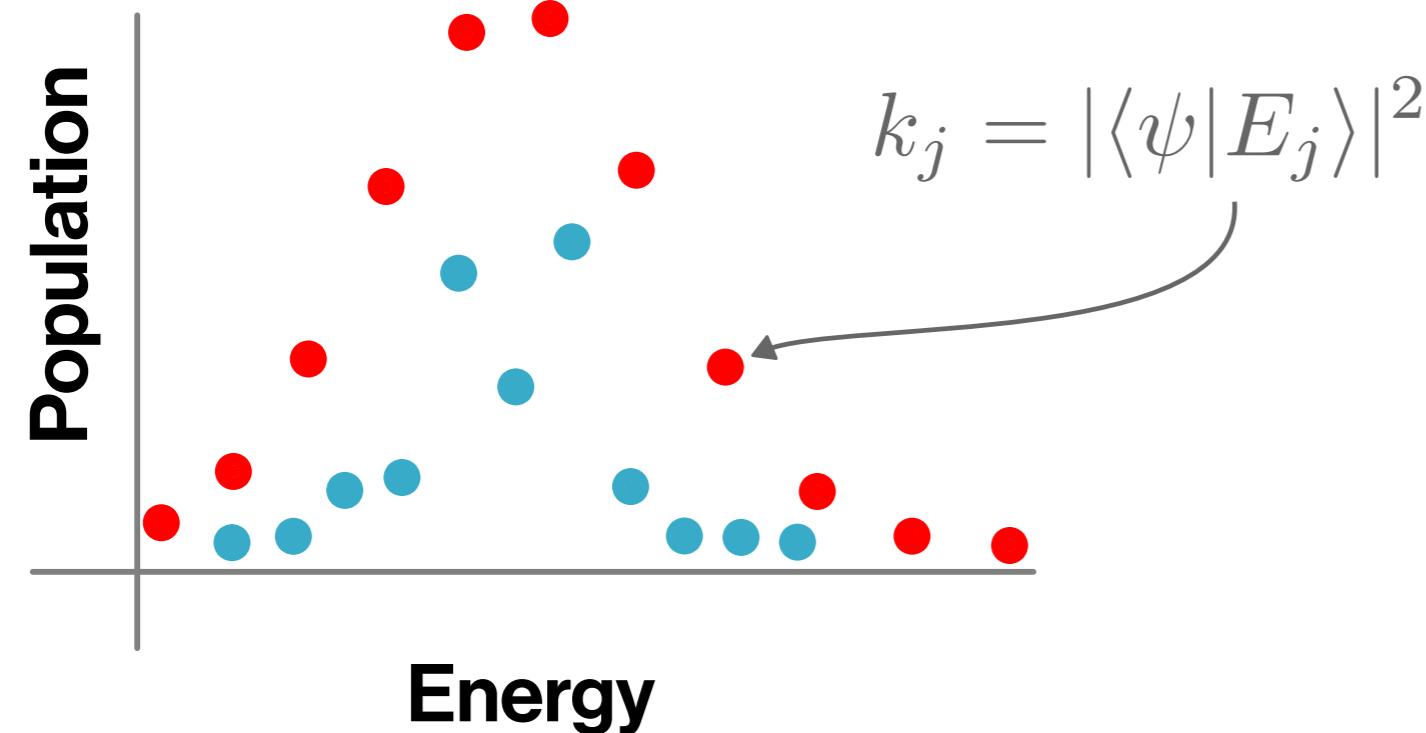


Long times

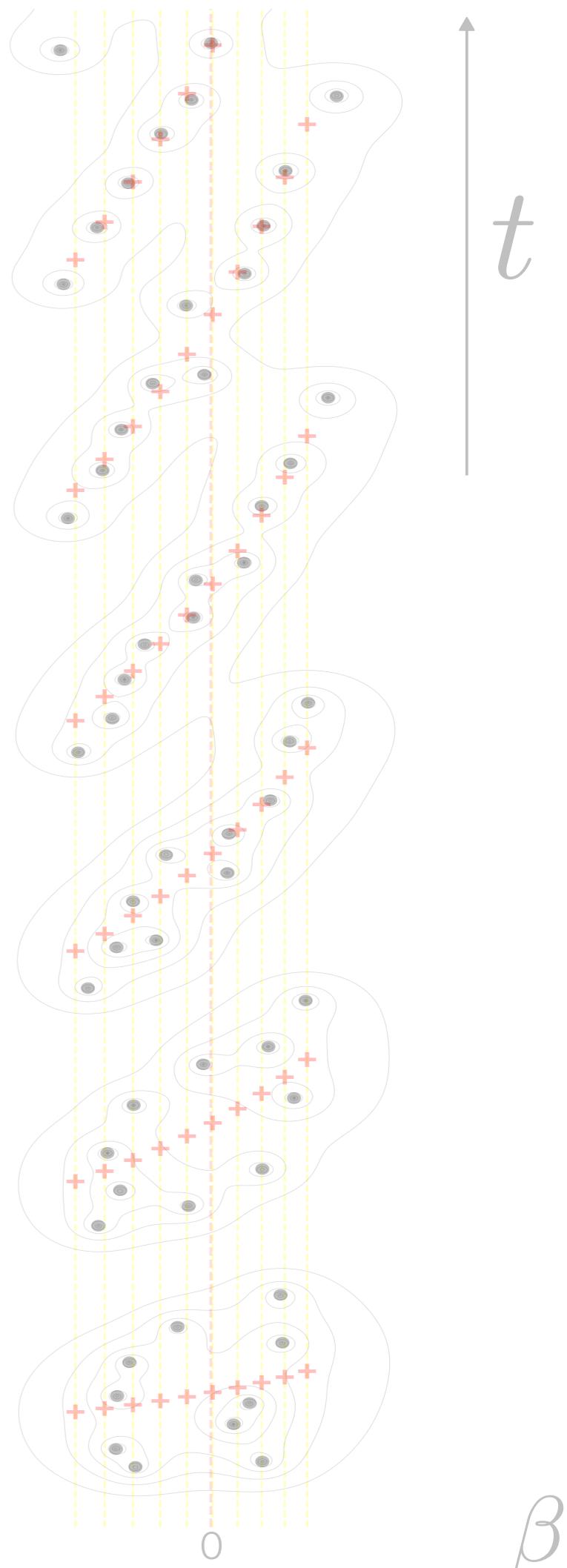
$$\mathcal{L}(\beta + it) = \sum_j r_j(\beta) e^{-iE_j t}$$

The dominant part = **The Envelope**

$$k_a \left(\frac{k_a}{k_b} \right)^{-\frac{\Delta_{ac}}{\Delta_{ab}}} \geq k_c, \quad \forall k_c, E_c$$



$$z_{ab}^*(n) = \frac{1}{\Delta_{ab}} \left(\log \left(\frac{k_a}{k_b} \right) + i\pi (2n+1) \right)$$

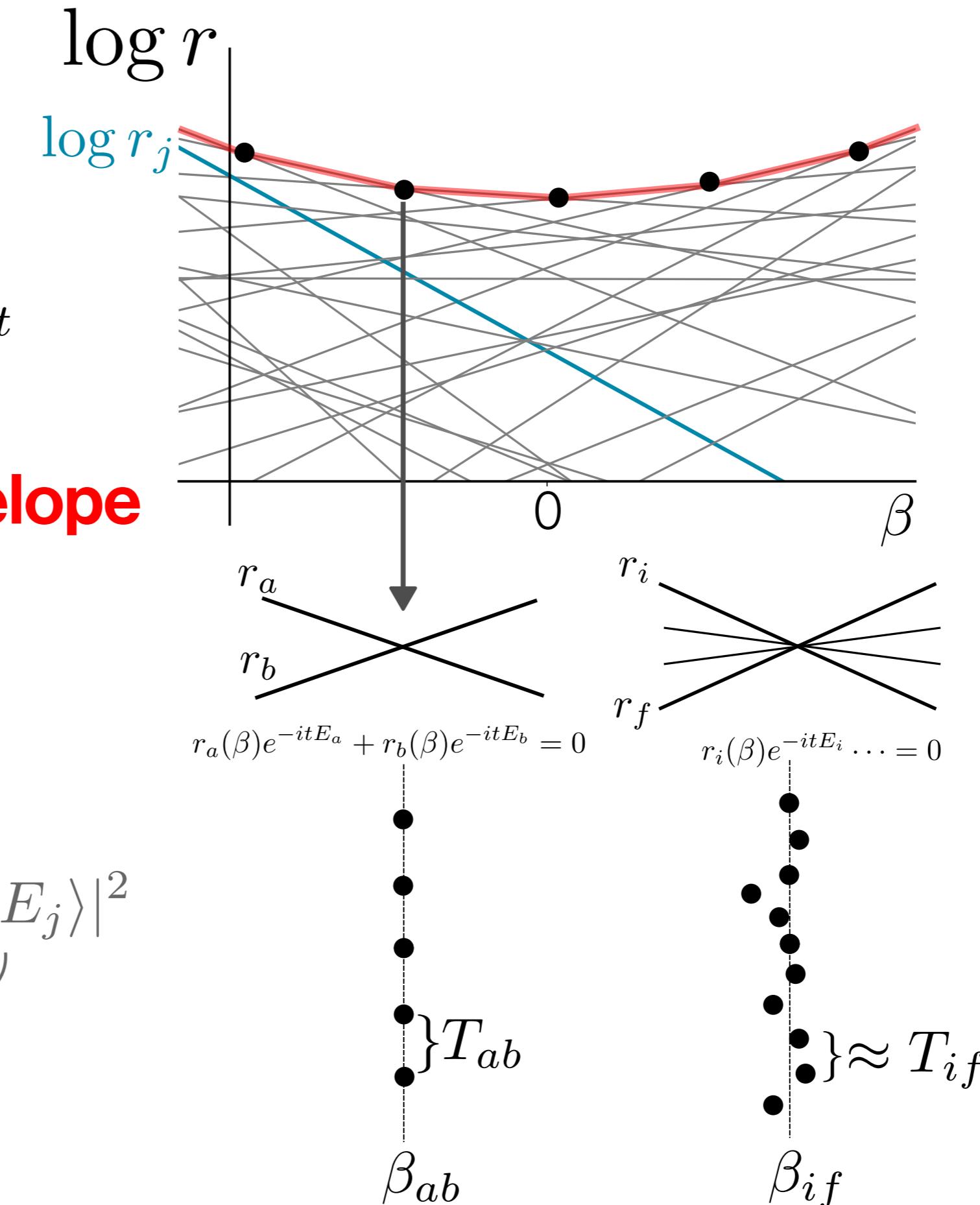
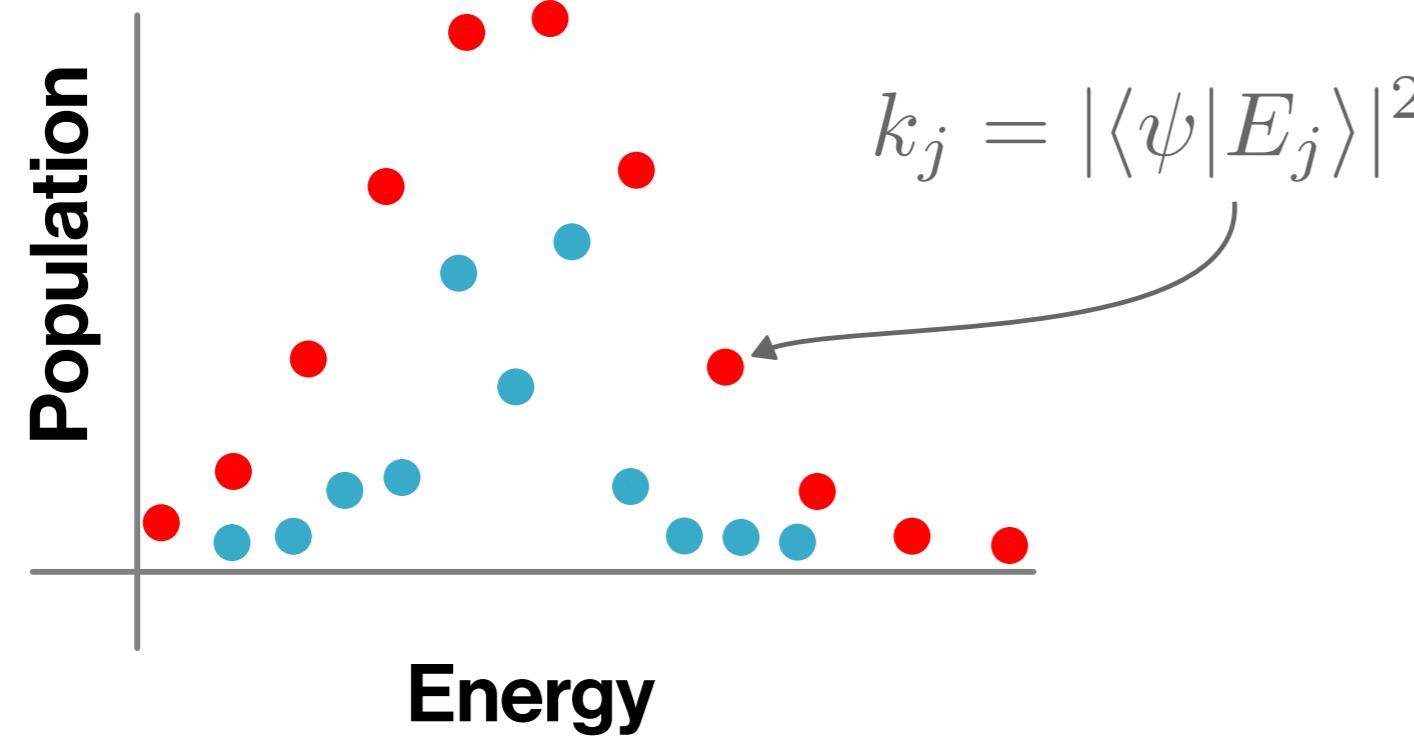


Long times

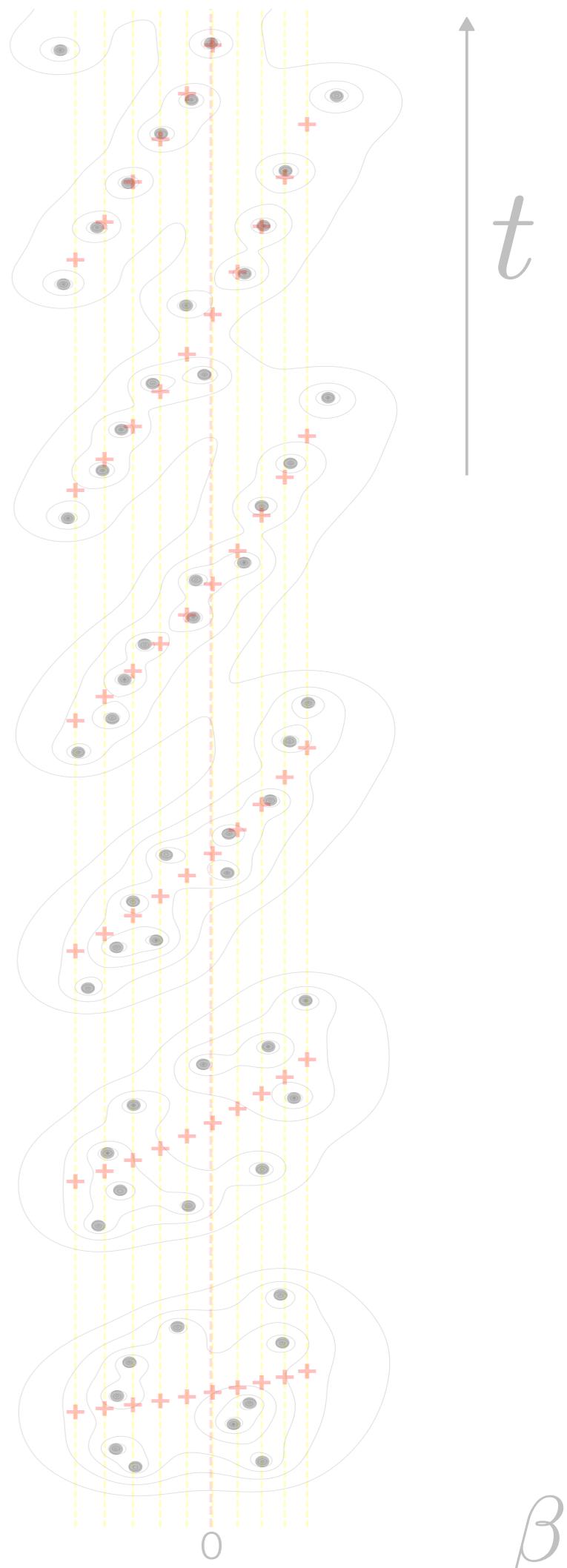
$$\mathcal{L}(\beta + it) = \sum_j r_j(\beta) e^{-iE_j t}$$

The dominant part = **The Envelope**

$$k_a \left(\frac{k_a}{k_b} \right)^{-\frac{\Delta_{ac}}{\Delta_{ab}}} \geq k_c, \quad \forall k_c, E_c$$



$$z_{ab}^*(n) = \frac{1}{\Delta_{ab}} \left(\log \left(\frac{k_a}{k_b} \right) + i\pi(2n+1) \right)$$

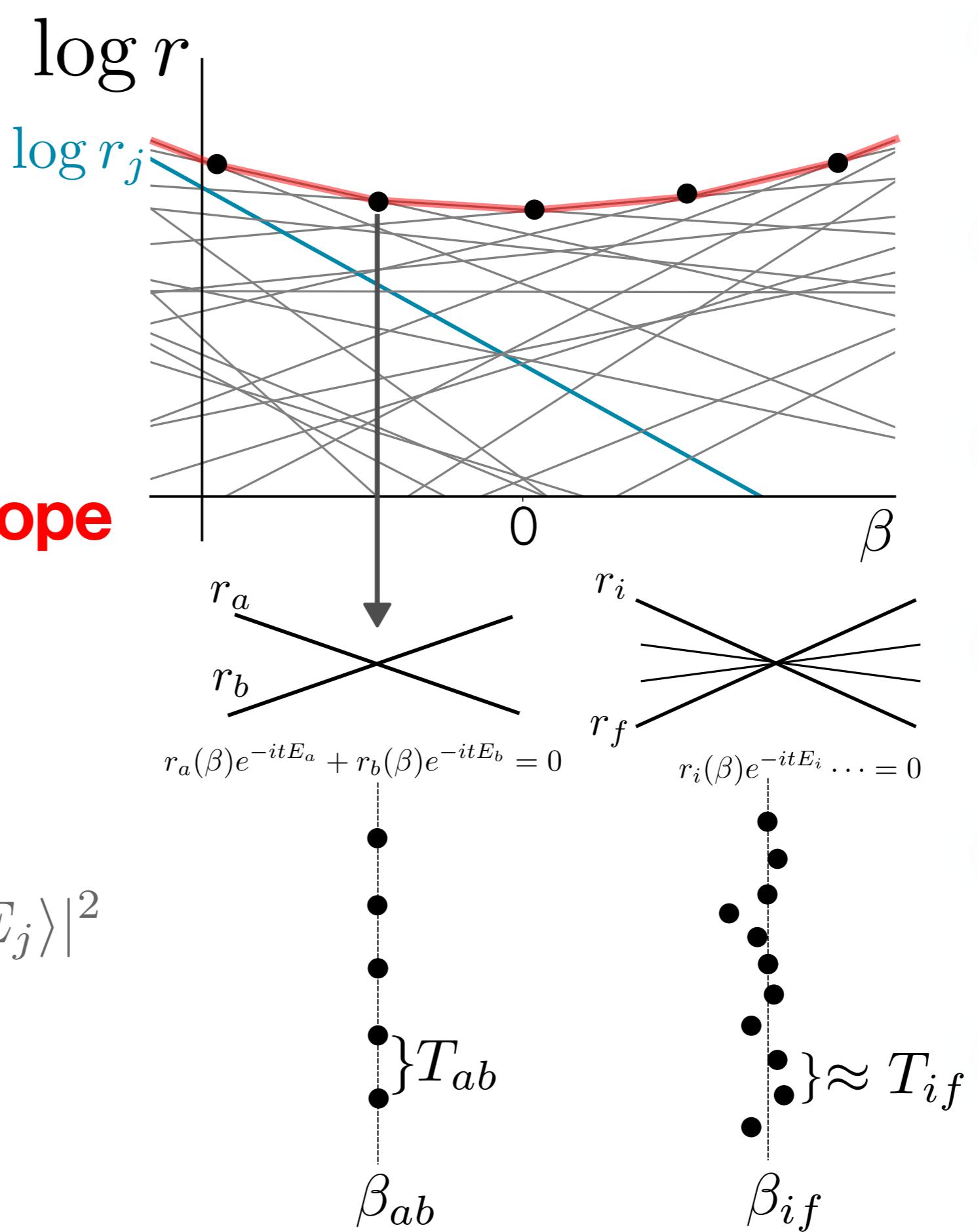
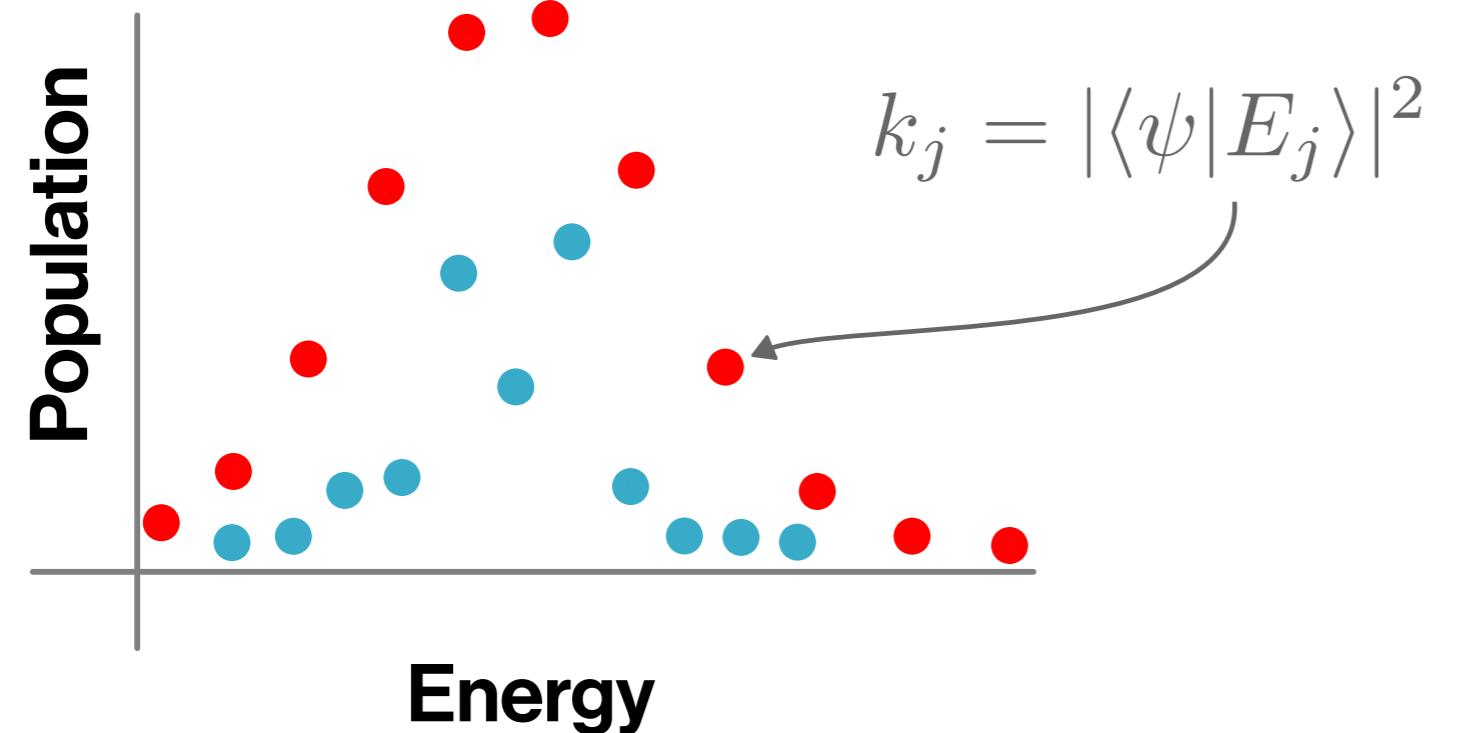


Long times

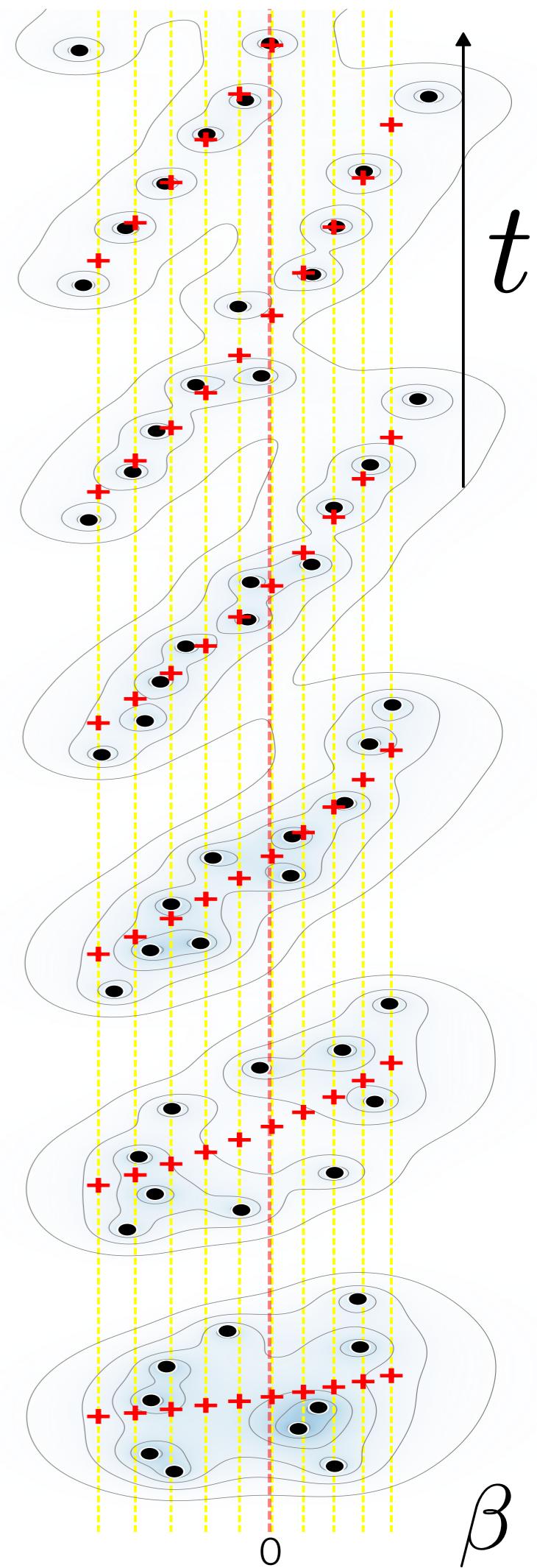
$$\mathcal{L}(\beta + it) = \sum_j r_j(\beta) e^{-iE_j t}$$

The dominant part = **The Envelope**

$$k_a \left(\frac{k_a}{k_b} \right)^{-\frac{\Delta_{ac}}{\Delta_{ab}}} \geq k_c, \quad \forall k_c, E_c$$



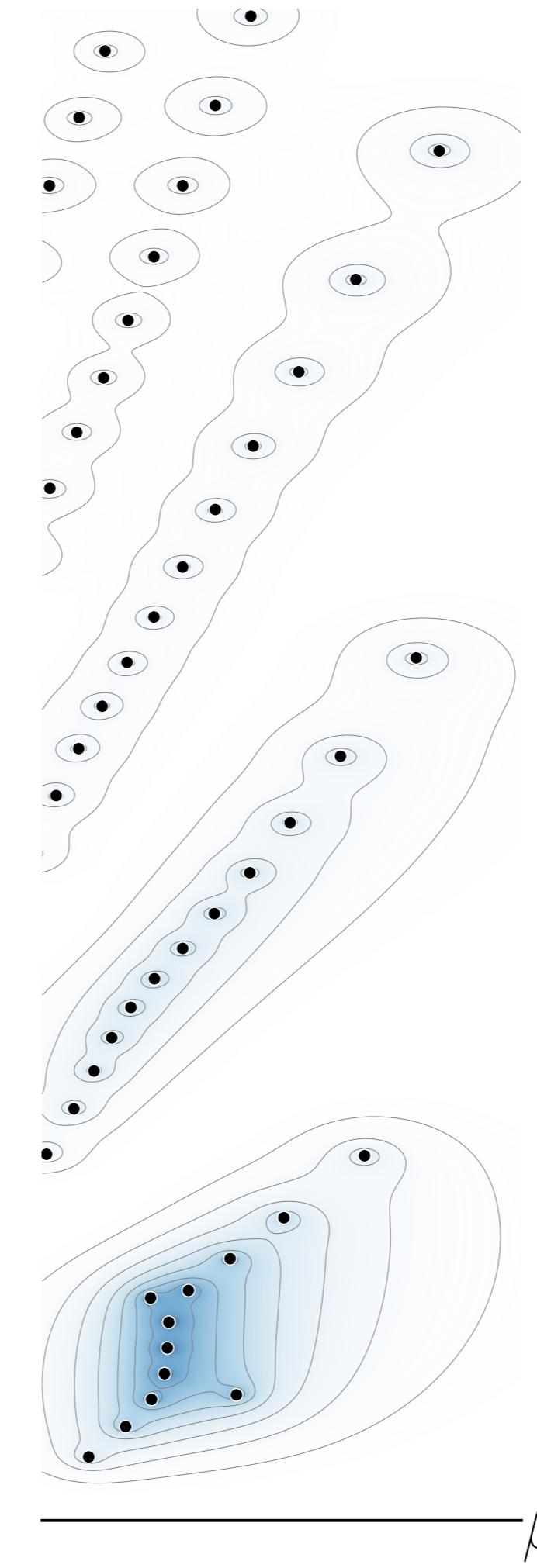
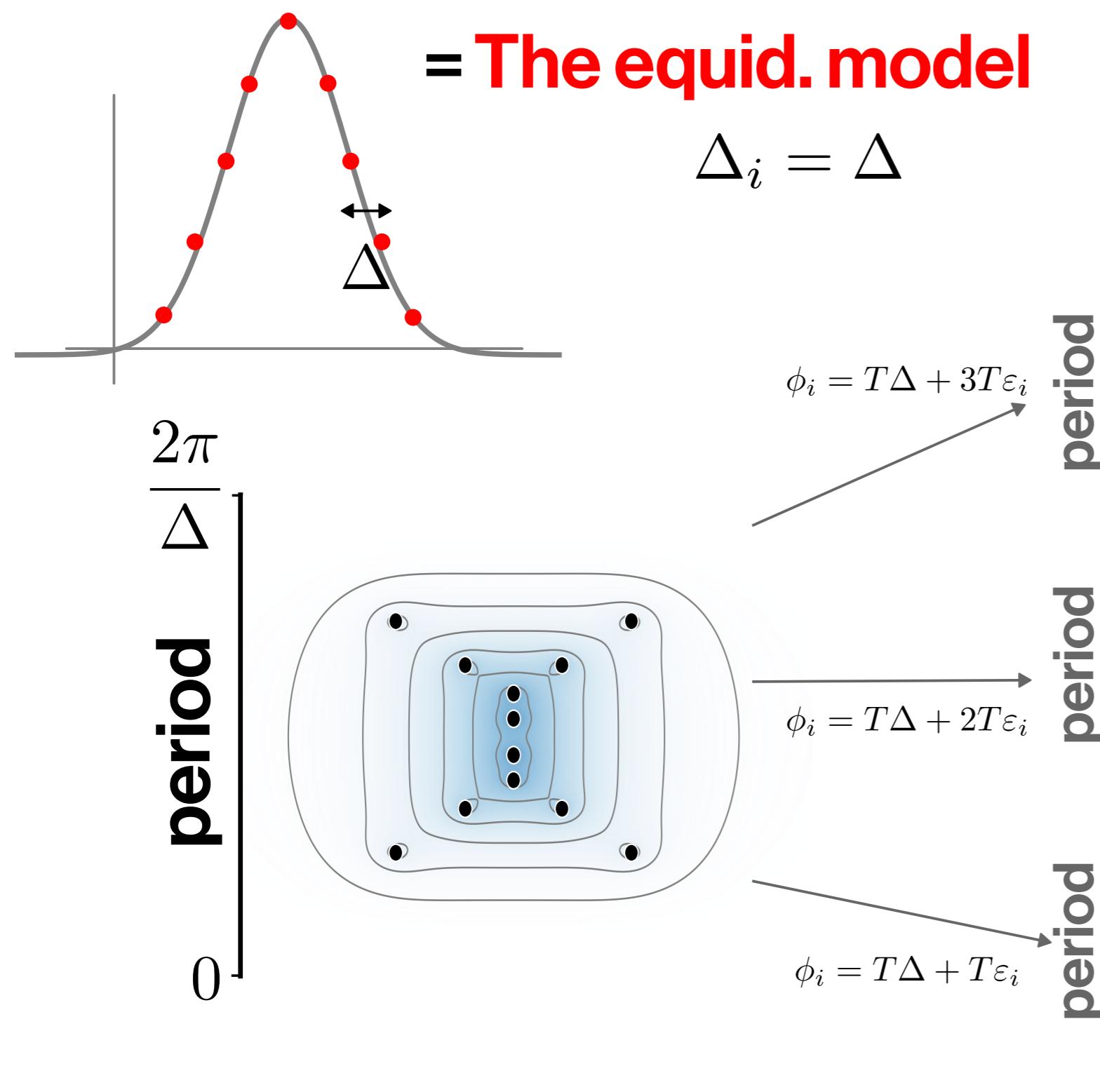
$$z_{ab}^*(n) = \frac{1}{\Delta_{ab}} \left(\log \left(\frac{k_a}{k_b} \right) + i\pi(2n+1) \right)$$



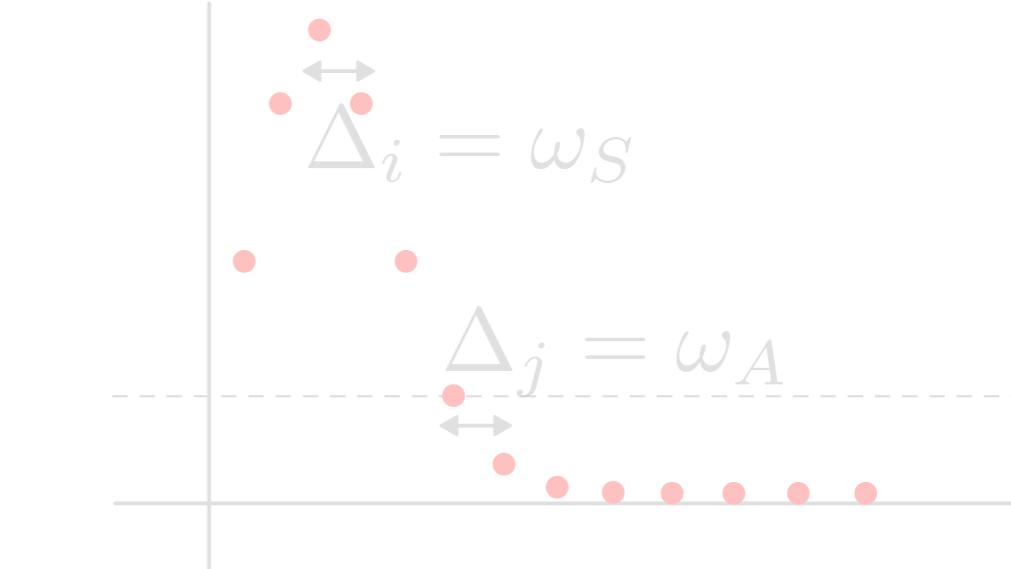
Short times

The dominant part

= The equid. model



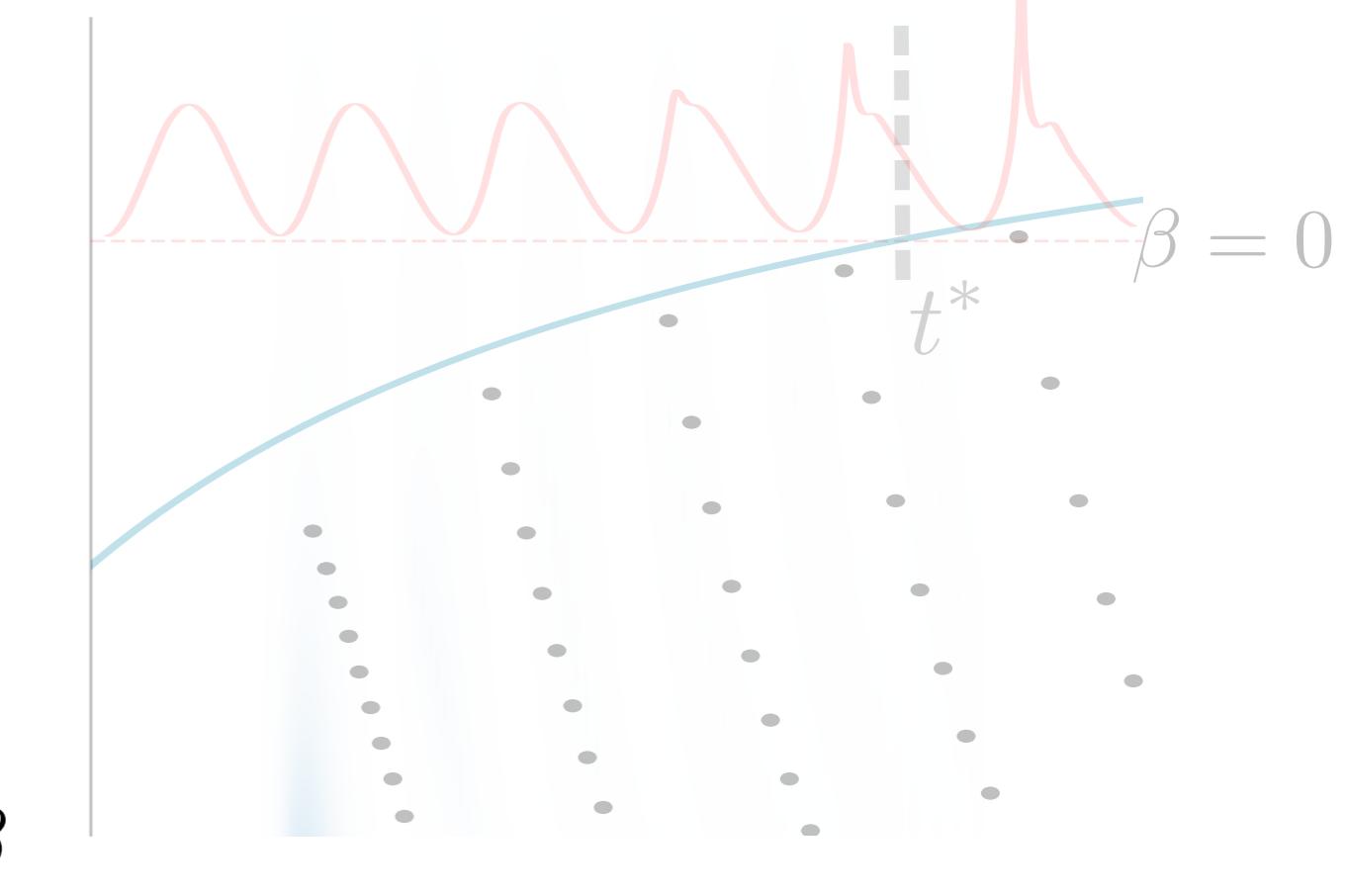
Asymmetry → running zeros



Arrival of the First zero

$$t^*(\beta) \approx t_0 + \frac{\beta - \beta_0}{|\omega_S(\beta) - \omega_A(\beta)|}$$

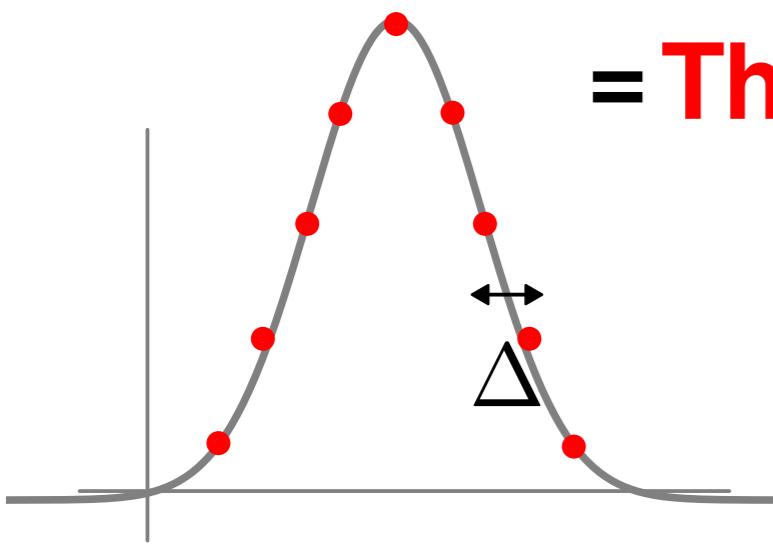
Lipkin model $\hat{H} = -\frac{J}{N}\hat{S}_x^2 + h\hat{S}_z$



Short times

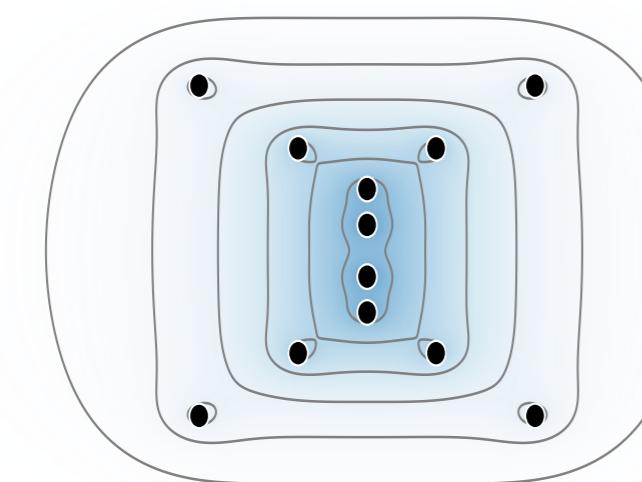
The dominant part

= The equid. model



$\frac{2\pi}{\Delta}$

period



$$\phi_i = T\Delta + 3T\varepsilon_i$$

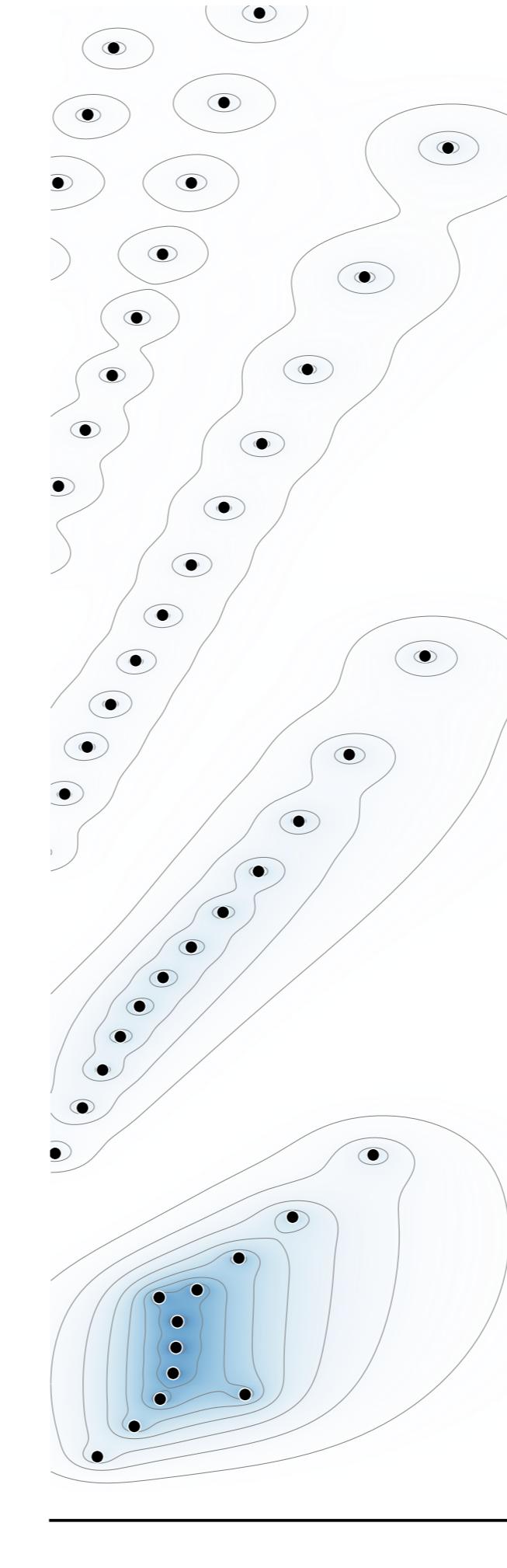
$$\phi_i = T\Delta + 2T\varepsilon_i$$

$$\phi_i = T\Delta + T\varepsilon_i$$

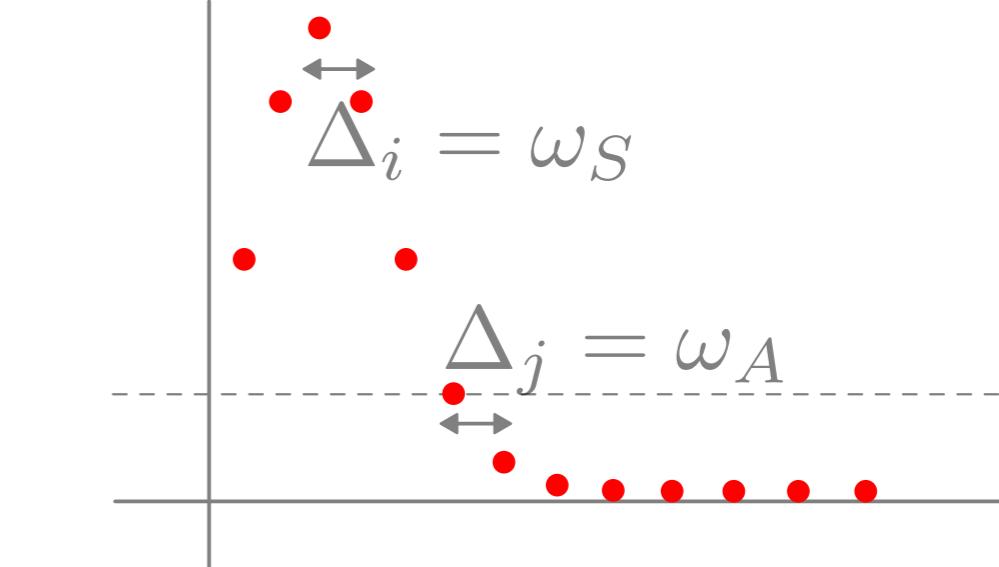
period

period

period



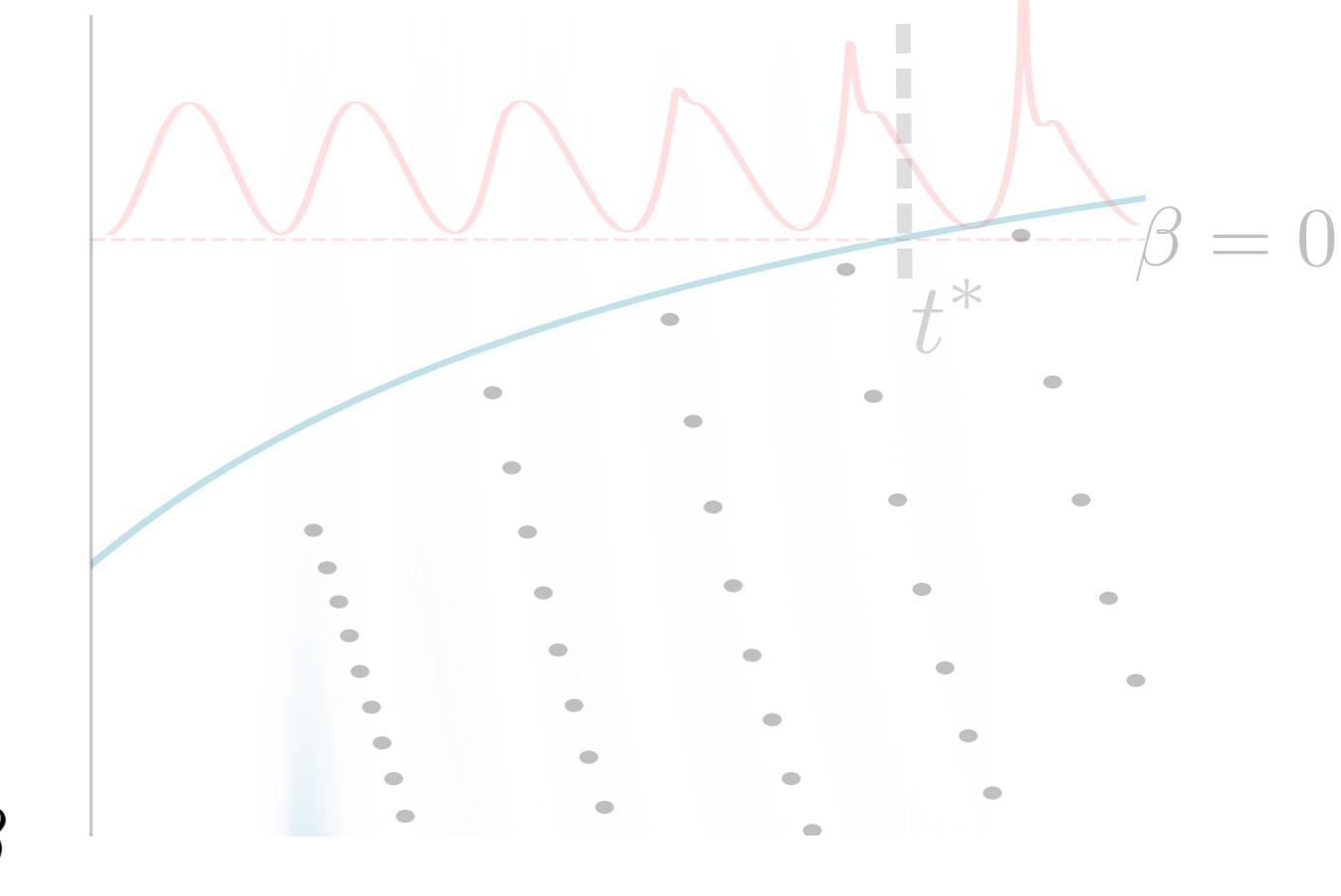
Asymmetry → running zeros



Arrival of the First zero

$$t^*(\beta) \approx t_0 + \frac{\beta - \beta_0}{|\omega_S(\beta) - \omega_A(\beta)|}$$

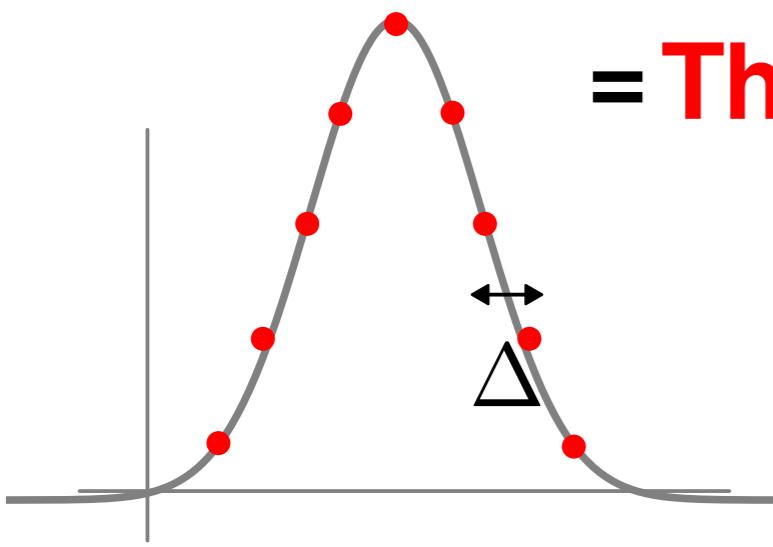
Lipkin model $\hat{H} = -\frac{J}{N}\hat{S}_x^2 + h\hat{S}_z$



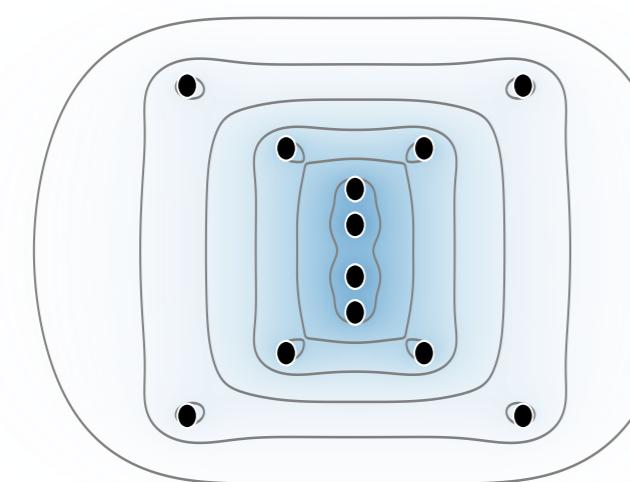
Short times

The dominant part

= The equid. model



$$\frac{2\pi}{\Delta}$$



period

$$\Delta_i = \Delta$$

$$\phi_i = T\Delta + 3T\varepsilon_i$$

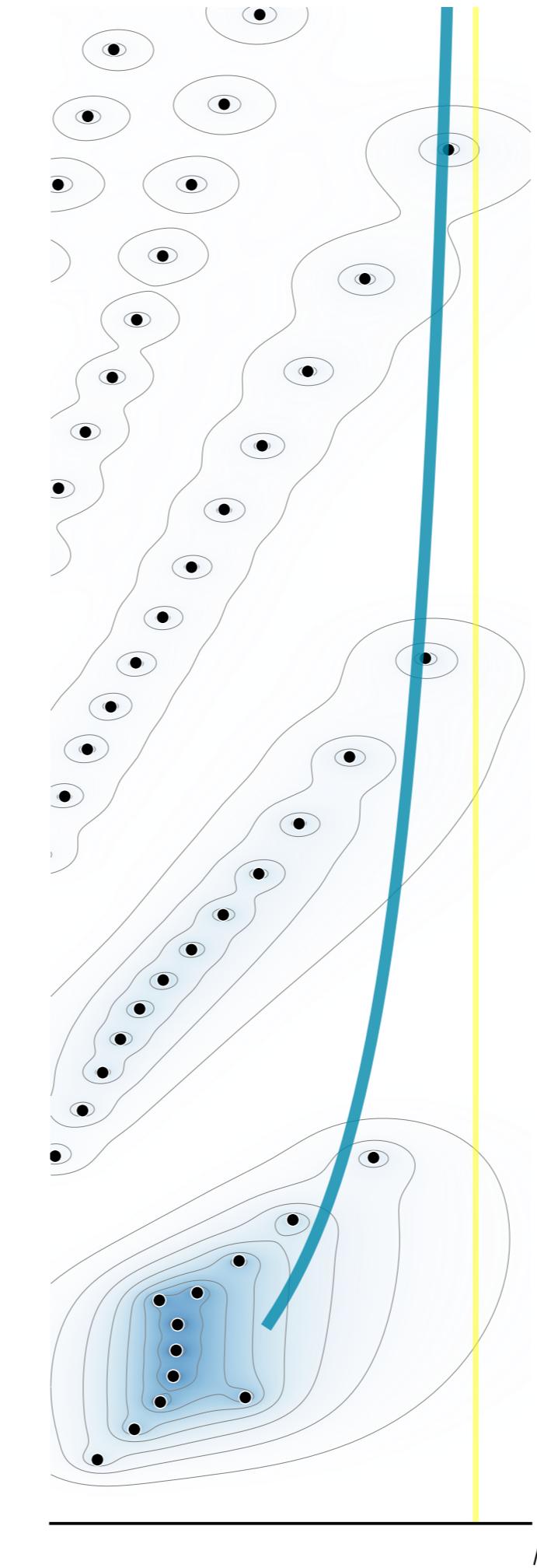
$$\phi_i = T\Delta + 2T\varepsilon_i$$

$$\phi_i = T\Delta + T\varepsilon_i$$

period

period

period



Asymmetry → running zeros

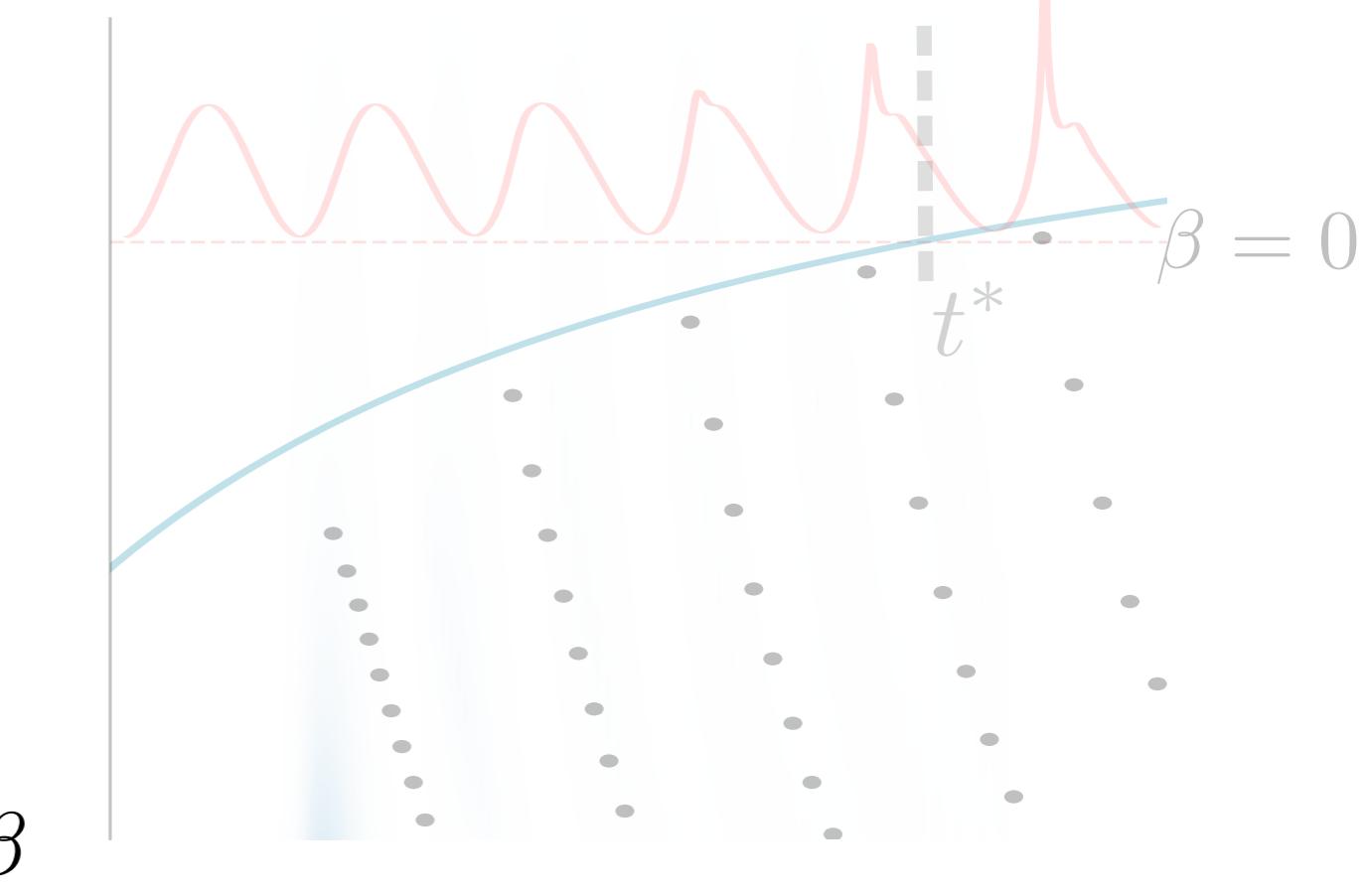
$$\Delta_i = \omega_S$$

$$\Delta_j = \omega_A$$

Arrival of the First zero

$$t^*(\beta) \approx t_0 + \frac{\beta - \beta_0}{|\omega_S(\beta) - \omega_A(\beta)|}$$

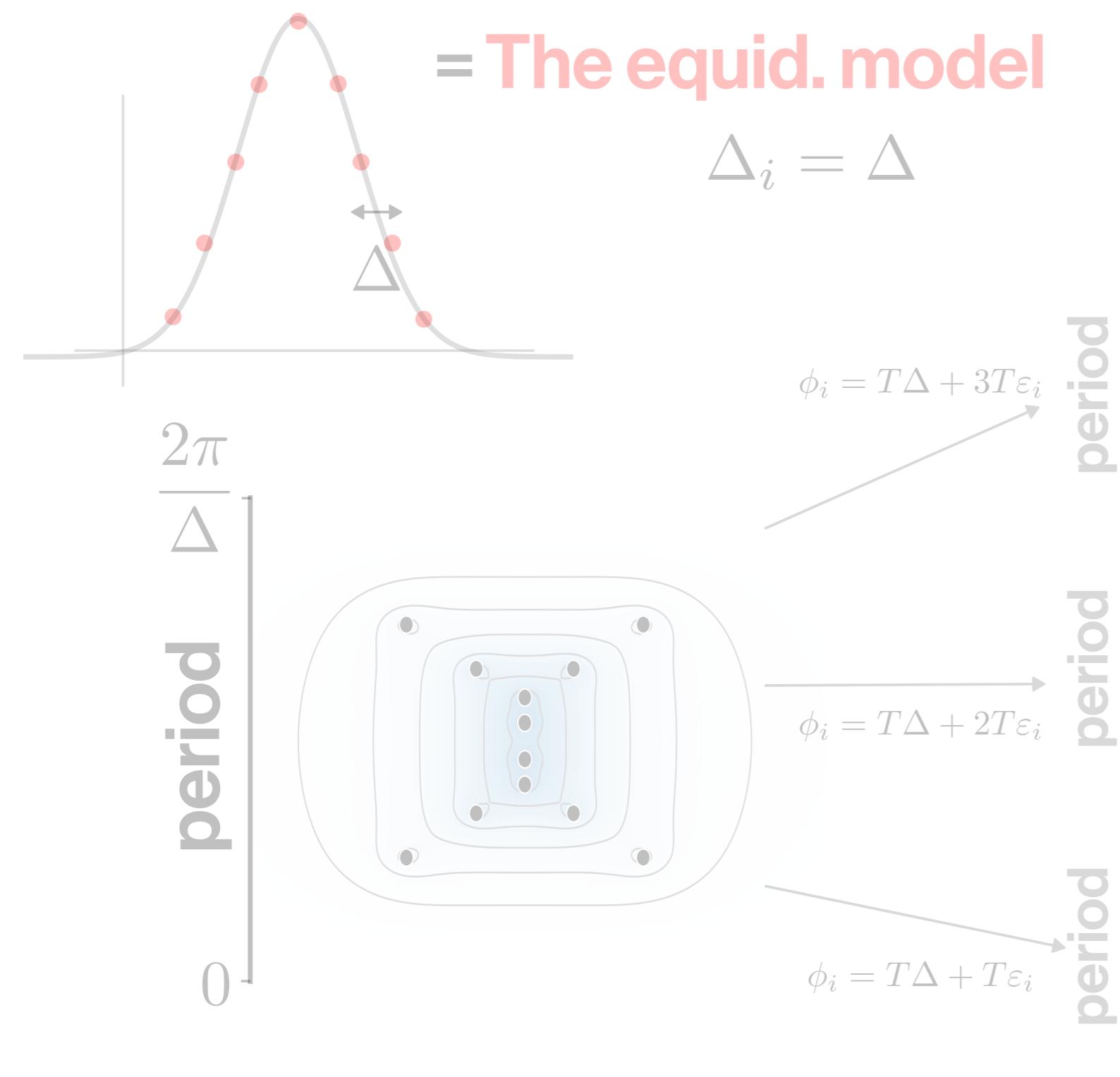
Lipkin model $\hat{H} = -\frac{J}{N}\hat{S}_x^2 + h\hat{S}_z$



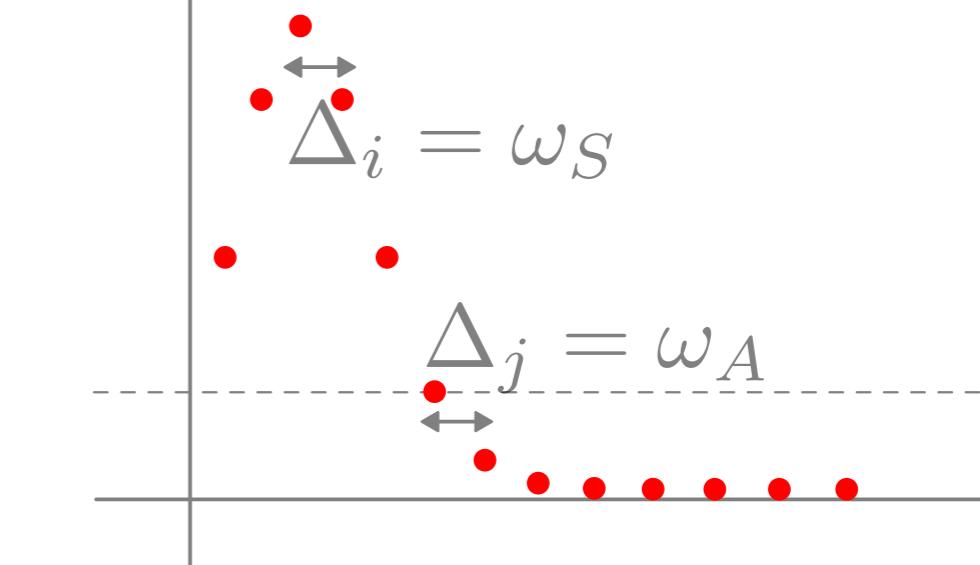
Short times

The dominant part

= The equid. model



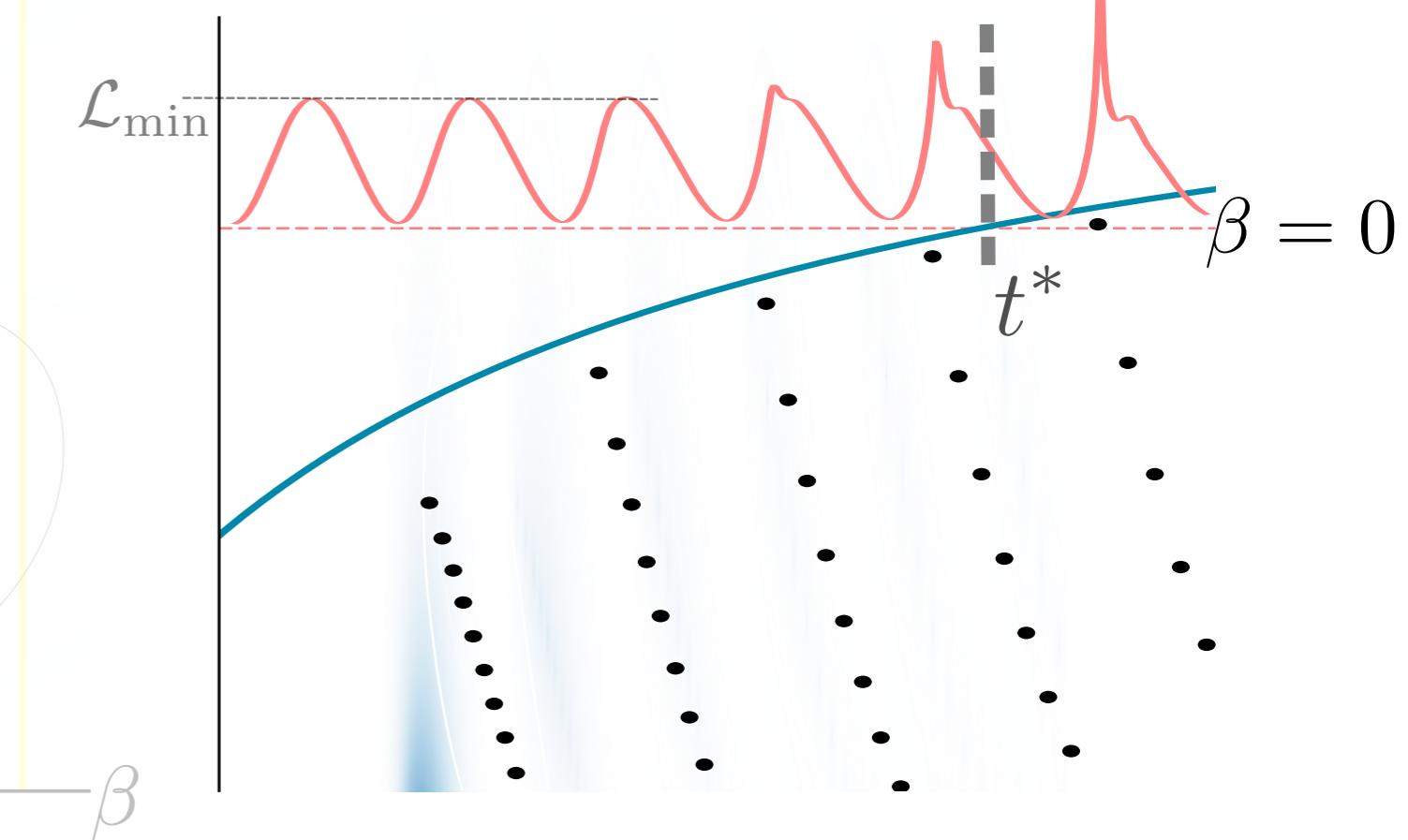
Asymmetry \rightarrow running zeros



Arrival of the First zero

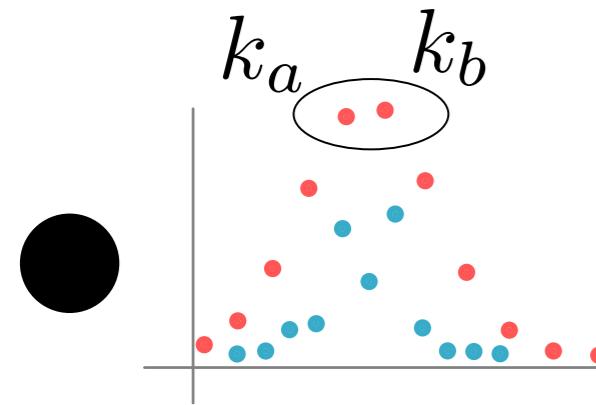
$$t^*(\beta) \approx t_0 + \frac{\beta - \beta_0}{|\omega_S(\beta) - \omega_A(\beta)|}$$

Lipkin model $\hat{H} = -\frac{J}{N}\hat{S}_x^2 + h\hat{S}_z$

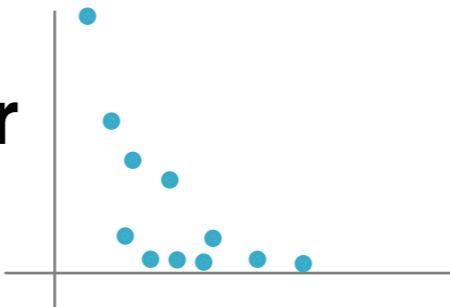


Quenches

- $\max\{k_j\} > \frac{1}{2}$



no zeros near the time axis

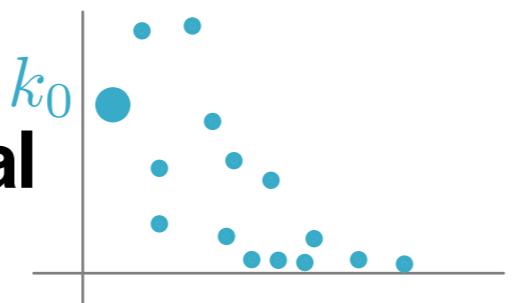


maximum determines the closest zeros to time axis

$$T_{ab} \quad \beta_{ab}$$

- $k_0 > \mathcal{L}_{\min}$

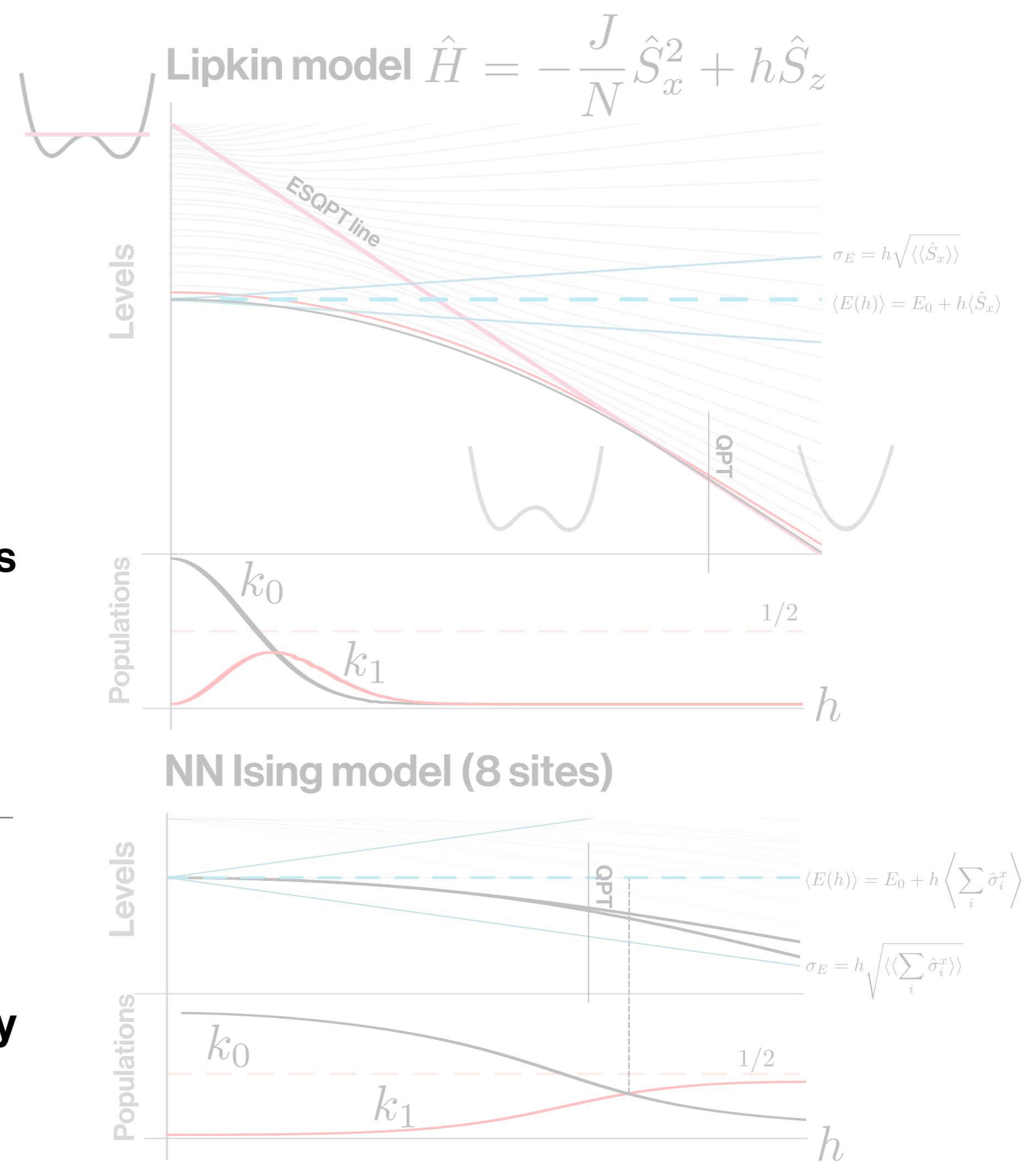
late time arrival



Ground state quench

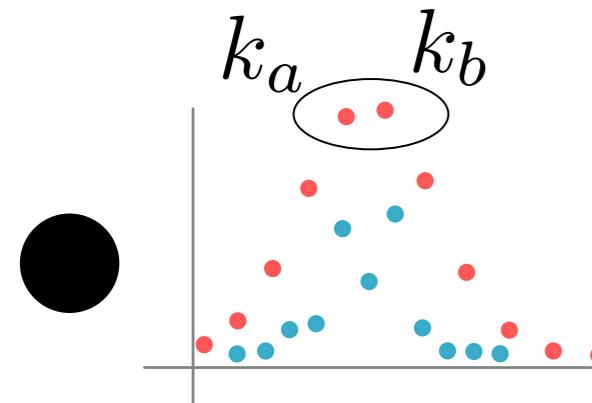
"zeros close to the time axis"

↔ "two (or more) highly populated states"

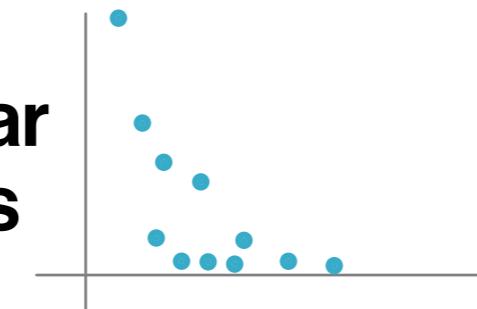


Quenches

- $\max\{k_j\} > \frac{1}{2}$



no zeros near the time axis

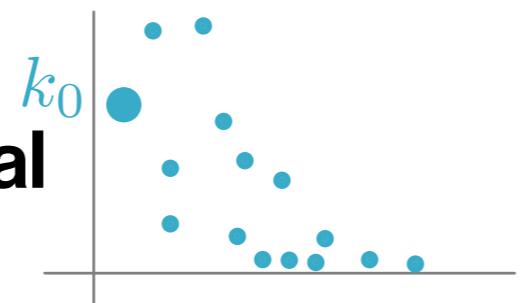


maximum determines the closest zeros to time axis

$$T_{ab} \quad \beta_{ab}$$

- $k_0 > \mathcal{L}_{\min}$

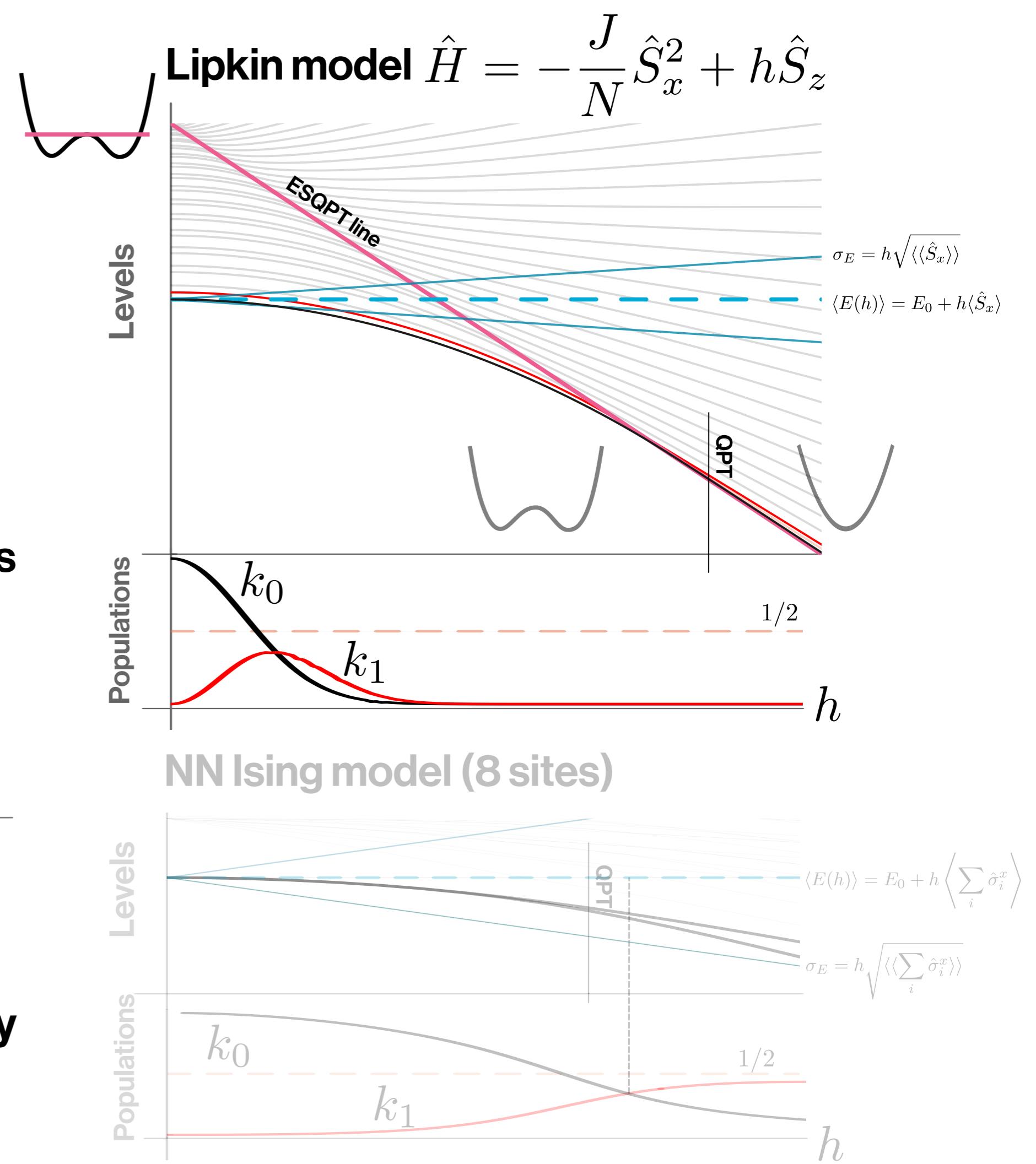
late time arrival



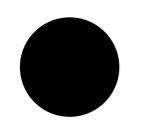
Ground state quench

"zeros close to the time axis"

"two (or more) highly populated states"

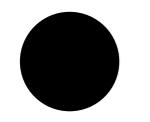
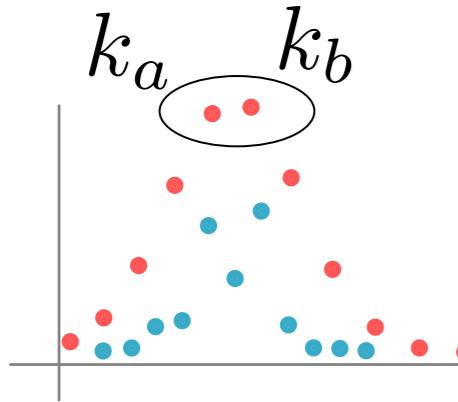


Quenches

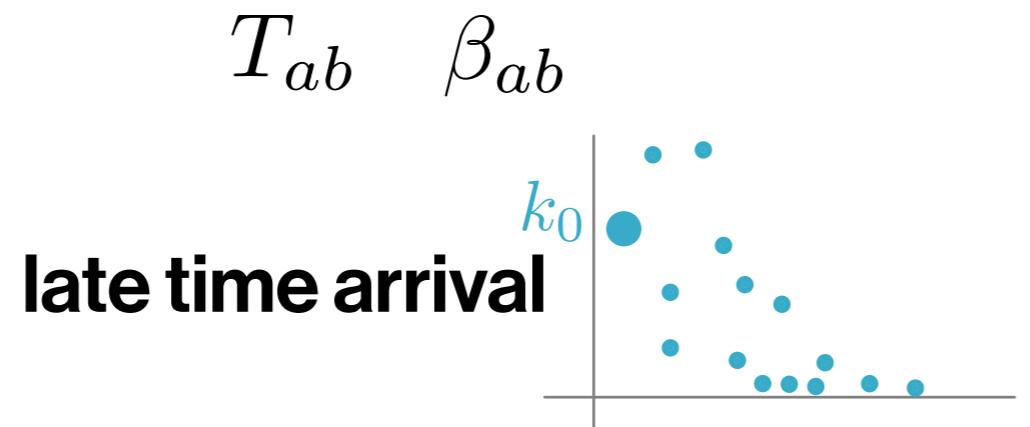


$$\max\{k_j\} > \frac{1}{2}$$

no zeros near the time axis



$$k_0 > \mathcal{L}_{\min}$$



Ground state quench

"zeros close to the time axis"

↔ "two (or more) highly populated states"

Cejnar et al.
JPA (2021)

Lipkin model $\hat{H} = -\frac{J}{N}\hat{S}_x^2 + h\hat{S}_z$

Levels

Populations

Levels

Populations

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

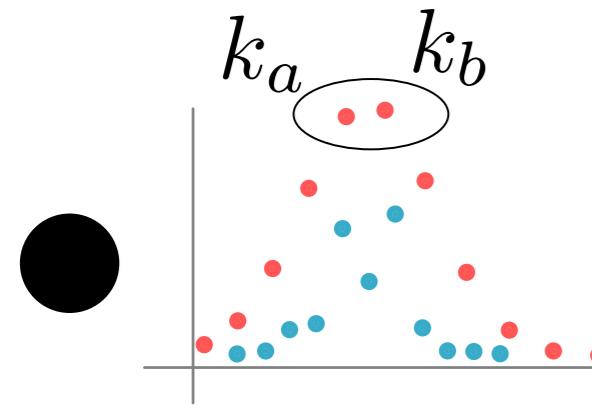
$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

$$\sigma_E = h\sqrt{\langle\langle \hat{S}_x \rangle\rangle}$$

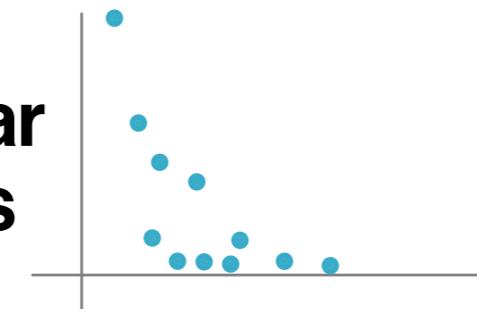
$$\langle E(h) \rangle = E_0 + h\langle \hat{S}_x \rangle$$

Quenches

- $\max\{k_j\} > \frac{1}{2}$



no zeros near the time axis

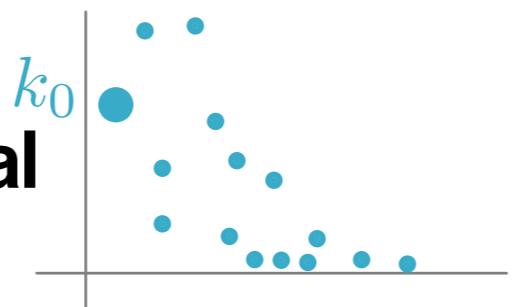


maximum determines the closest zeros to time axis

- $k_0 > \mathcal{L}_{\min}$

$$T_{ab} \quad \beta_{ab}$$

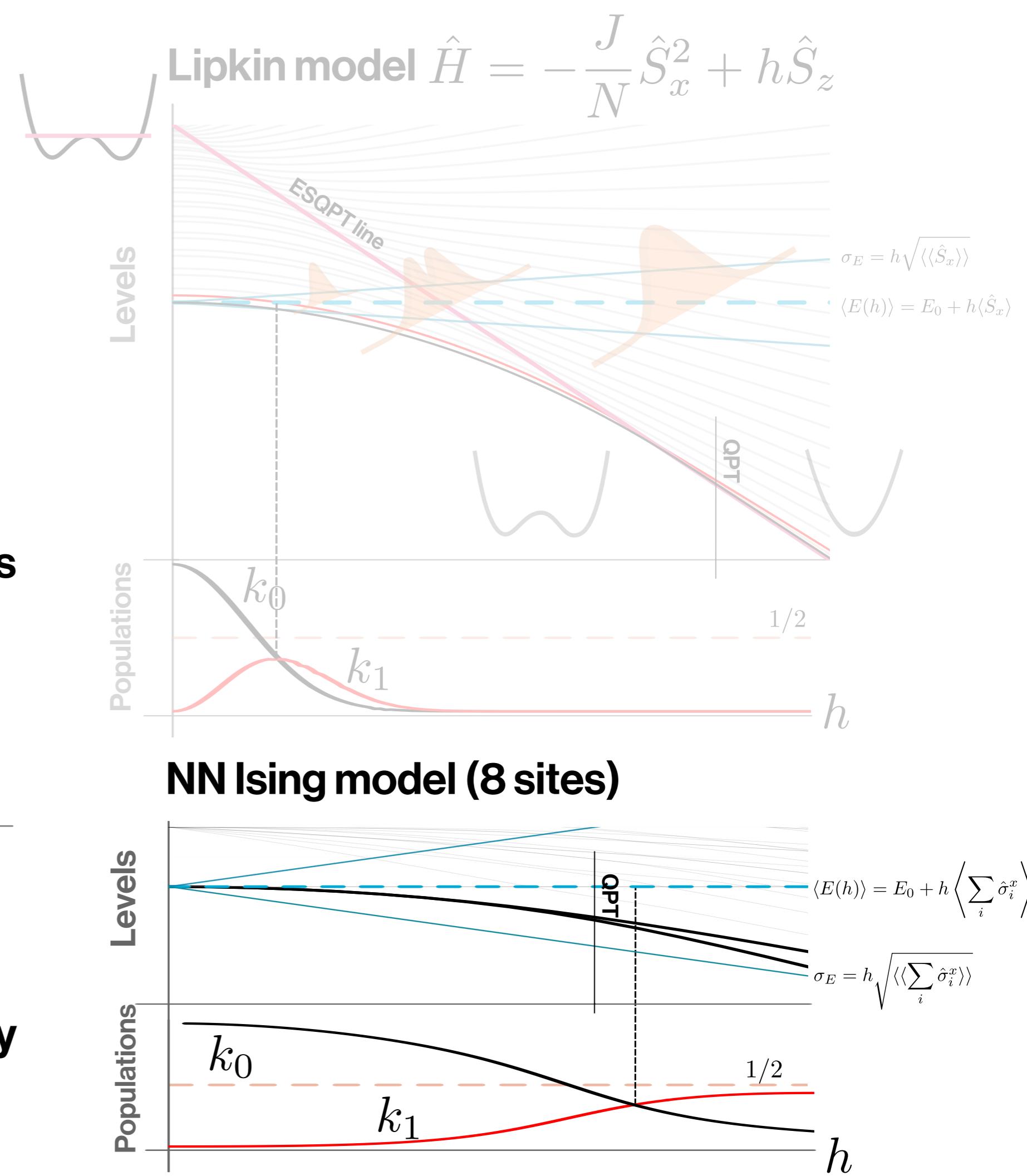
late time arrival



Ground state quench

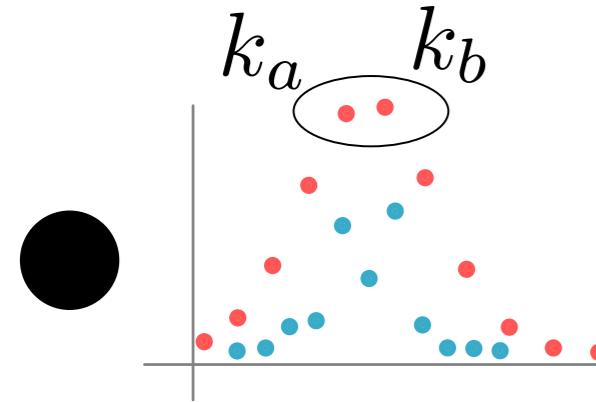
"zeros close to the time axis"

"two (or more) highly populated states"

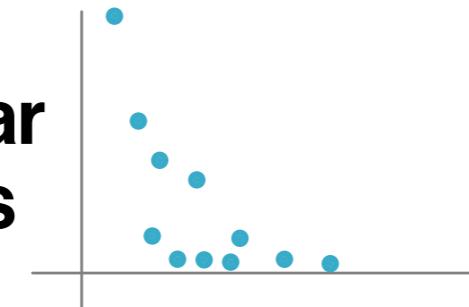


Quenches

- $\max\{k_j\} > \frac{1}{2}$



no zeros near the time axis

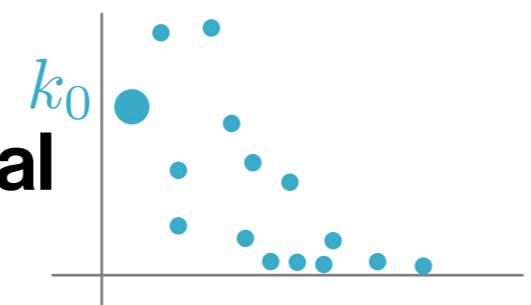


maximum determines the closest zeros to time axis

- $k_0 > \mathcal{L}_{\min}$

$$T_{ab} \quad \beta_{ab}$$

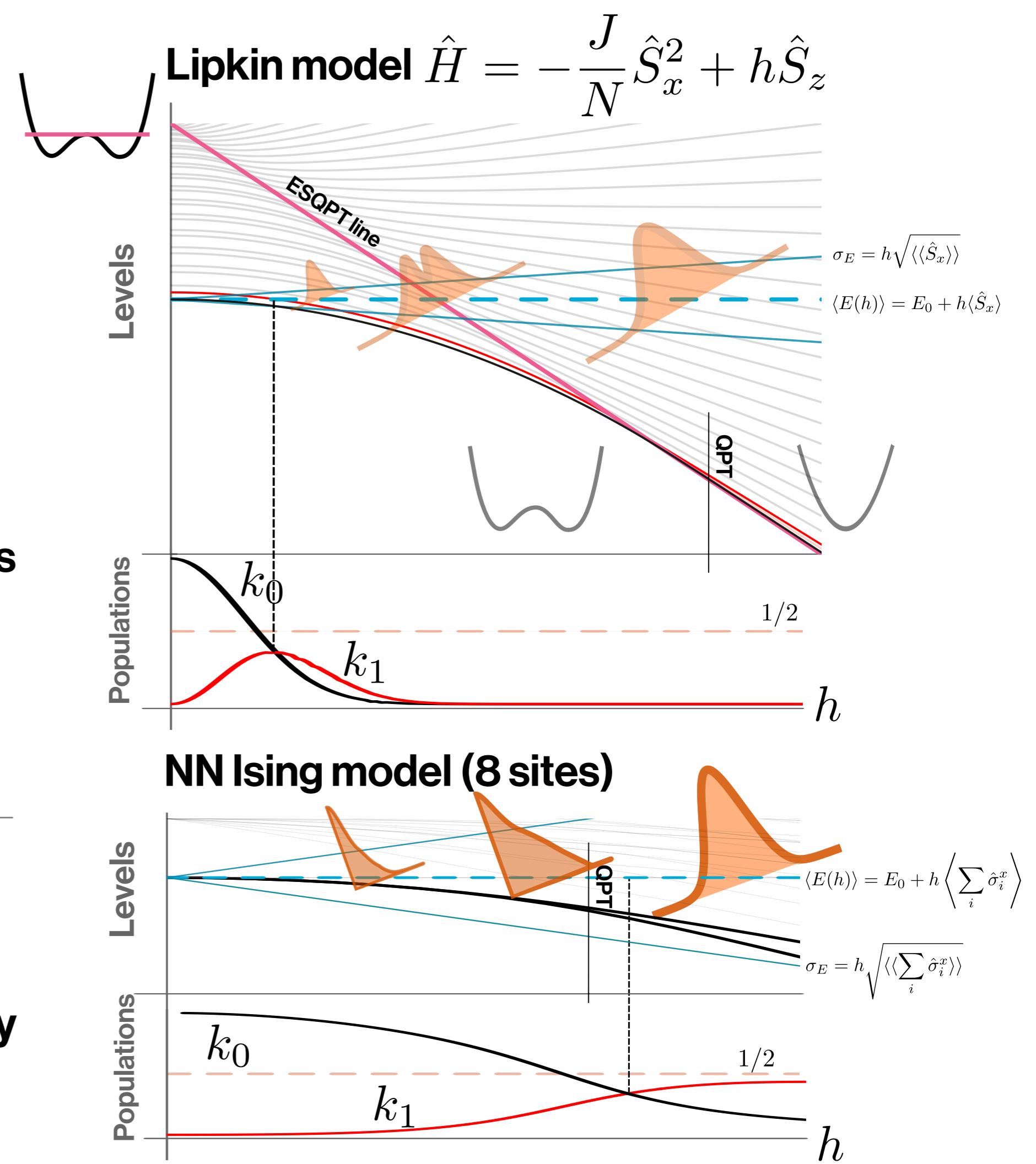
late time arrival



Ground state quench

"zeros close to the time axis"

"two (or more) highly populated states"



Conclusions

Zeros comes from the dominant parts of the initial state
- the envelope

Zeros form chain-like structures around β_{ab} with (quasi-) periods T_{ab}

Asymmetry of the initial state may delay the first zero

"zeros close
to the time axis" \longleftrightarrow "two (or more) highly
populated states"

**Thank you
for
your attention**

The work was supported by the The Charles University Grant Agency
(grant no. 215323)

