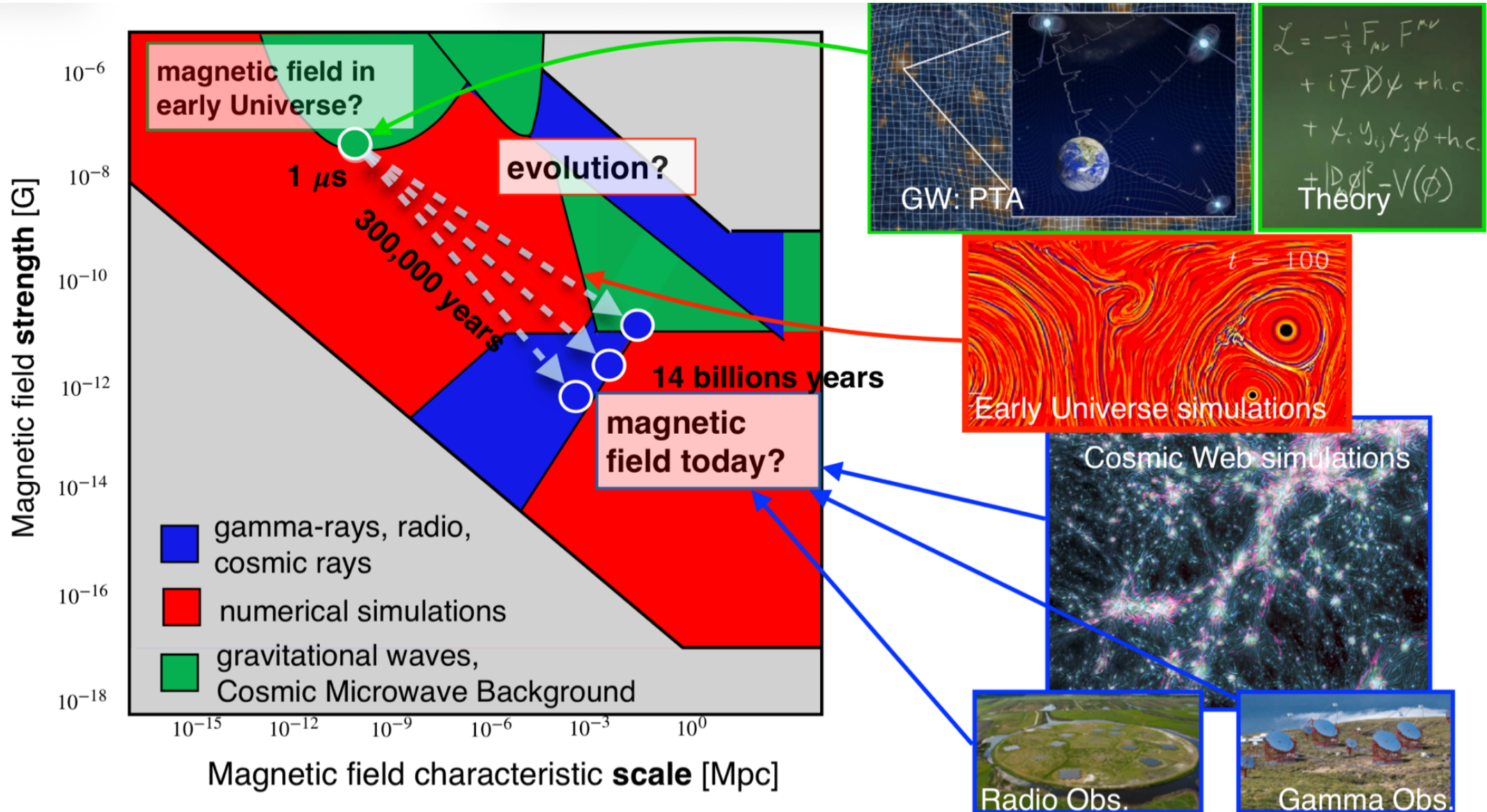


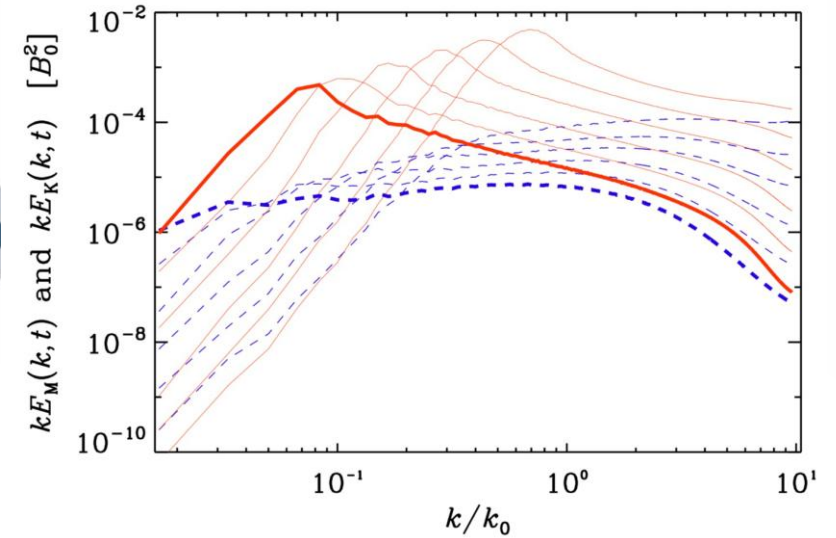
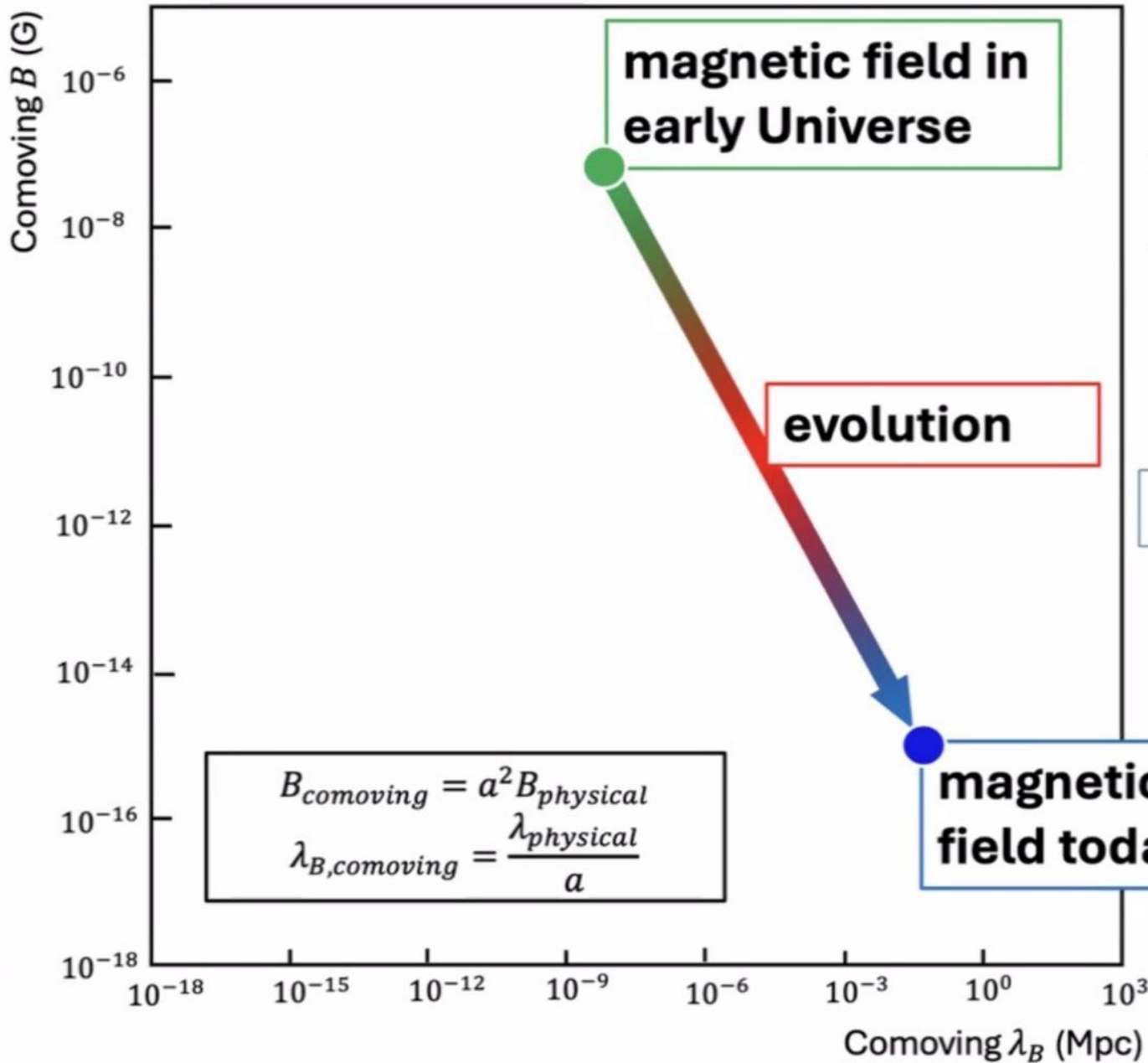
Magnetic fields as probes of the first microseconds

Axel Brandenburg
& Oksana Iarygina



Magnetic field: inversely cascading

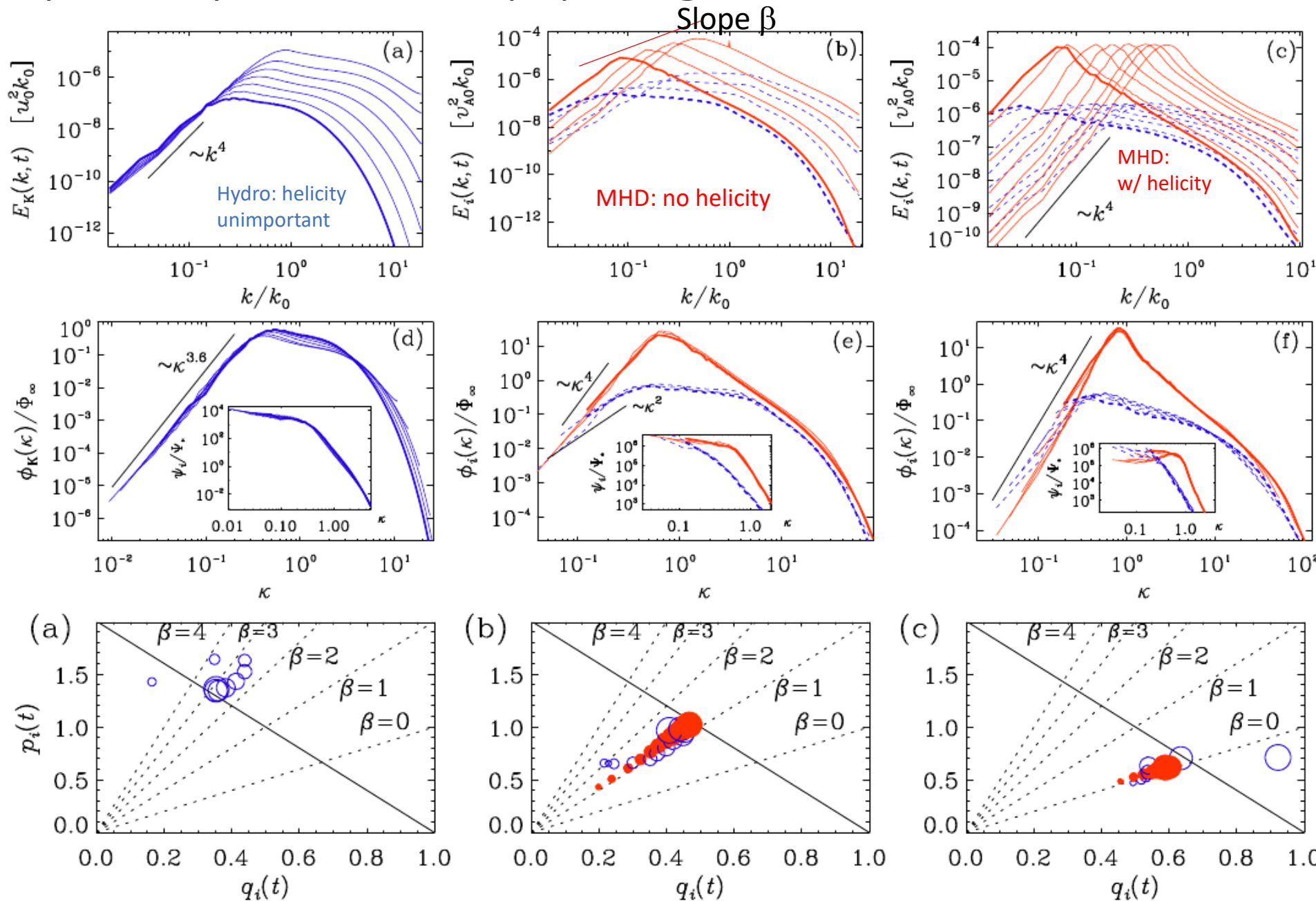
Present-day relic magnetic field parameters are related to the initial field characteristics. They potentially provide information on the mechanism of production of the field in the early Universe.



- Turbulent decay
 - Magnetic field drives gas motions
 - The velocity stirs up the field further
- But: length scale can increase
 - Energy at large scales (small k) can increase!
 - \rightarrow Inverse cascade

Collapsed spectra and pq diagrams

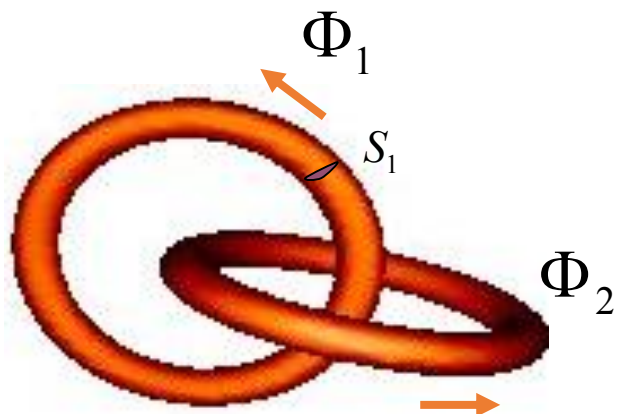
$$-p_i(t) = d \ln \mathcal{E}_i / d \ln t, \quad q_i(t) = d \ln \xi_i / d \ln t,$$



Explanations
for slope β
Exponents p, q
(Hosking &
Schekochihin
2021+2023)

Inverse cascading
is unique to
magnetic decay

Magnetic helicity



$$H = \int_V \mathbf{A} \cdot \mathbf{B} dV$$

$$\mathbf{B} = \nabla \times \mathbf{A}$$

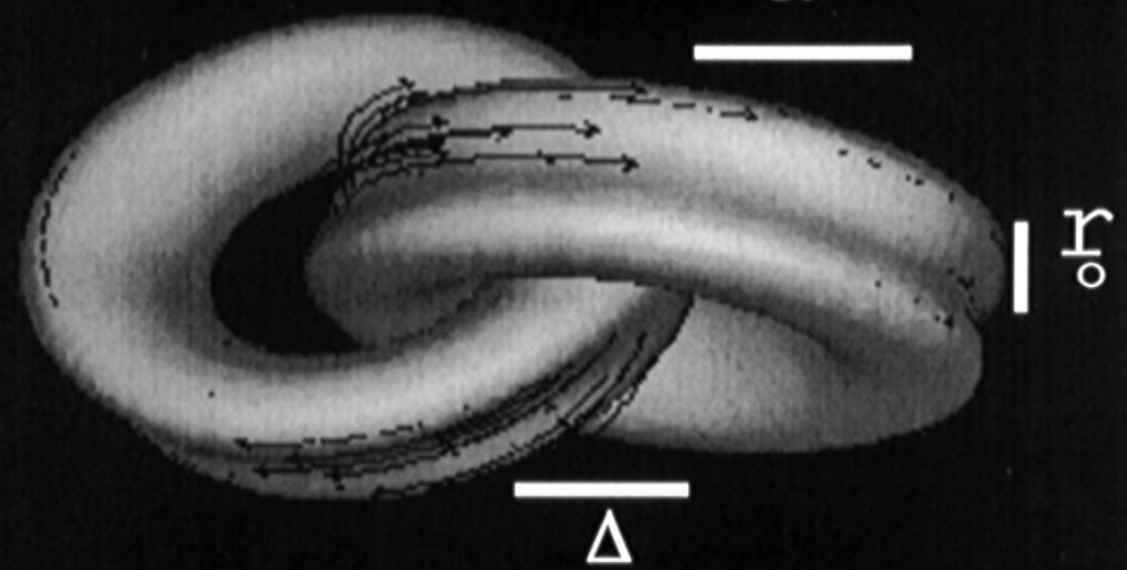
$$H = \pm 2\Phi_1\Phi_2$$

Therefore the unit is Maxwell squared

$$H_1 = \int_{L_1} \mathbf{A} \cdot d\ell \int_{S_1} \mathbf{B} \cdot d\mathbf{S}$$

$$= \int_{S_2} \nabla \times \mathbf{A} \cdot d\mathbf{S} = \Phi_2 \quad = \Phi_1$$

t=2



t=3



Conservation laws

$$\xi_M \sim \langle \mathbf{A} \cdot \mathbf{B} \rangle t^{2/3}$$

$$\text{cm} \sim (\text{cm}^3/\text{s}^2)^{1/3} \text{s}^{2/3}$$

Magnetic helicity
Anastrophy (2-D)

Hosking integral

Saffman integral

Loitsyansky integral

$$\langle \mathbf{A} \cdot \mathbf{B} \rangle$$

$$\text{cm}^3 \text{s}^{-2}$$

$$\xi_M(t) \propto \langle \mathbf{A} \cdot \mathbf{B} \rangle^{1/3} t^{2/3}$$

$$\langle A_z^2 \rangle$$

$$\text{cm}^4 \text{s}^{-2}$$

$$\xi_M(t) \propto \langle A_z^2 \rangle^{1/4} t^{1/2}$$

$$I_H$$

$$\text{cm}^9 \text{s}^{-4}$$

$$\xi_M(t) \propto I_H^{1/9} t^{4/9}$$

$$I_S$$

$$\text{cm}^5 \text{s}^{-2}$$

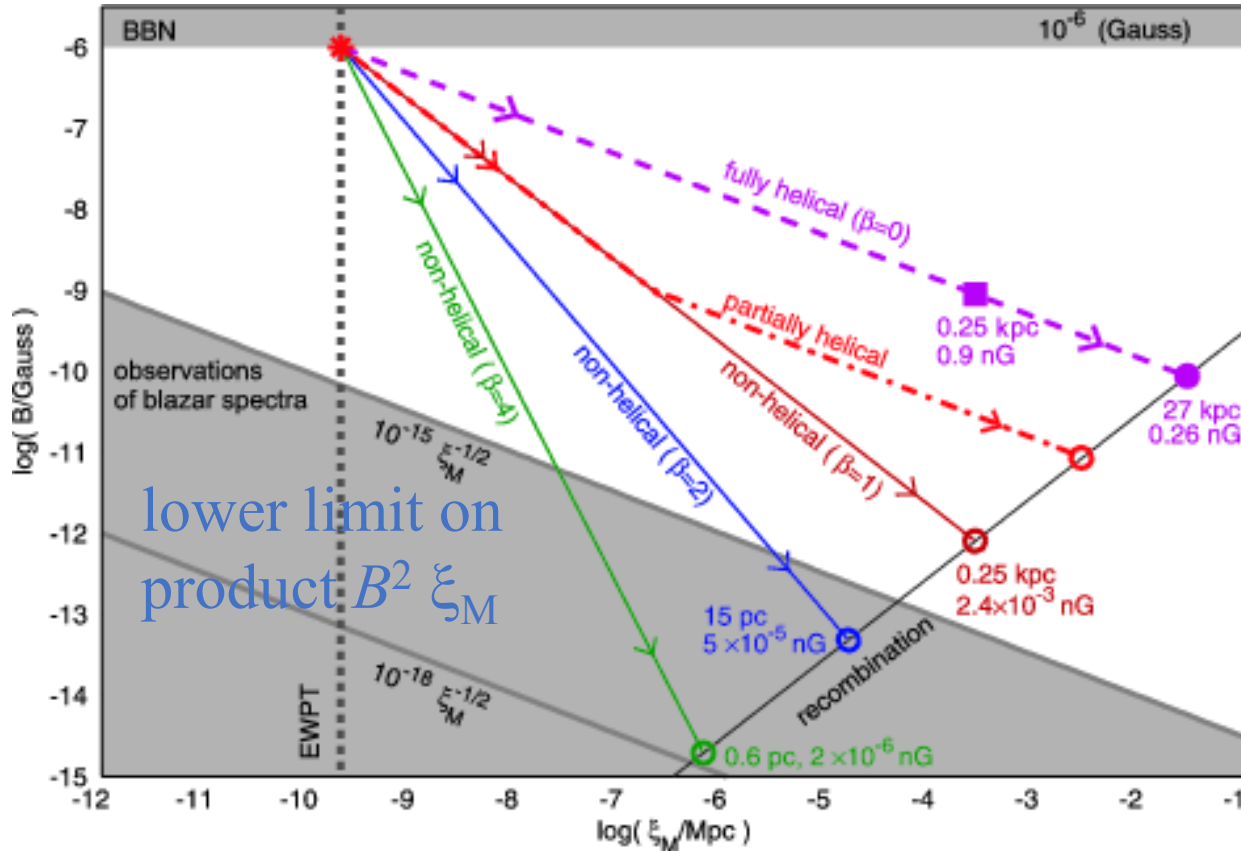
$$\xi_M(t) \propto I_S^{1/5} t^{2/5}$$

$$I_L$$

$$\text{cm}^7 \text{s}^{-2}$$

$$\xi_M(t) \propto I_S^{1/7} t^{2/7}$$

AB, Kahnashvili, ..., Vachaspati (2017)



Magnetic energy dependence
Parametric representation

magnetic energy

$$\kappa = p/2q$$

$$\mathcal{E}_M(t) \propto \langle \mathbf{A} \cdot \mathbf{B} \rangle^{2/3} t^{-2/3}$$

$$\propto \xi_M^{-1/2}$$

$$\mathcal{E}_M(t) \propto \langle A_z^2 \rangle^{1/2} t^{-1}$$

$$\propto \xi_M^{-1}$$

$$\mathcal{E}_M(t) \propto I_H^{2/9} t^{-10/9}$$

$$\propto \xi_M^{-5/4}$$

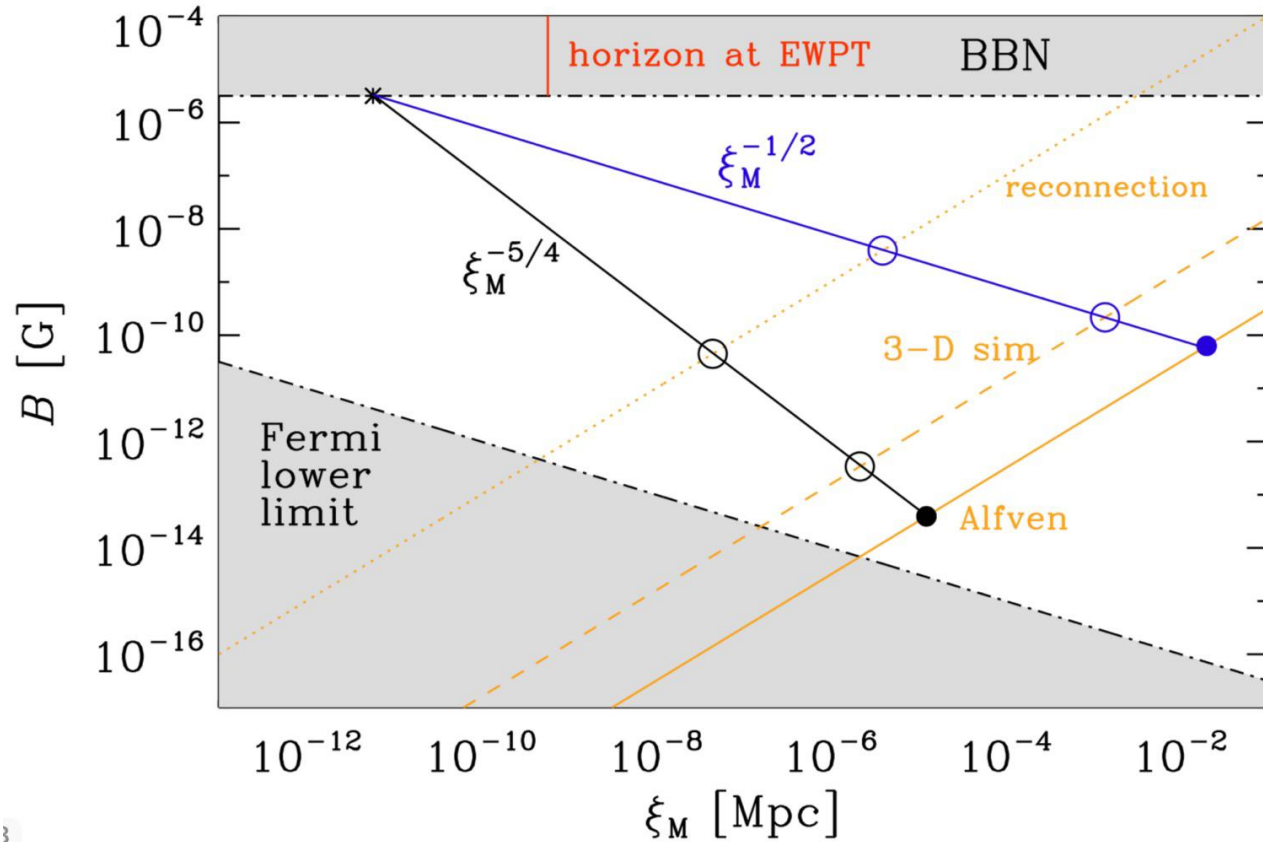
$$\mathcal{E}_M(t) \propto I_S^{2/5} t^{-6/5}$$

$$\propto \xi_M^{-3/2}$$

$$\mathcal{E}_M(t) \propto I_S^{2/7} t^{-10/7}$$

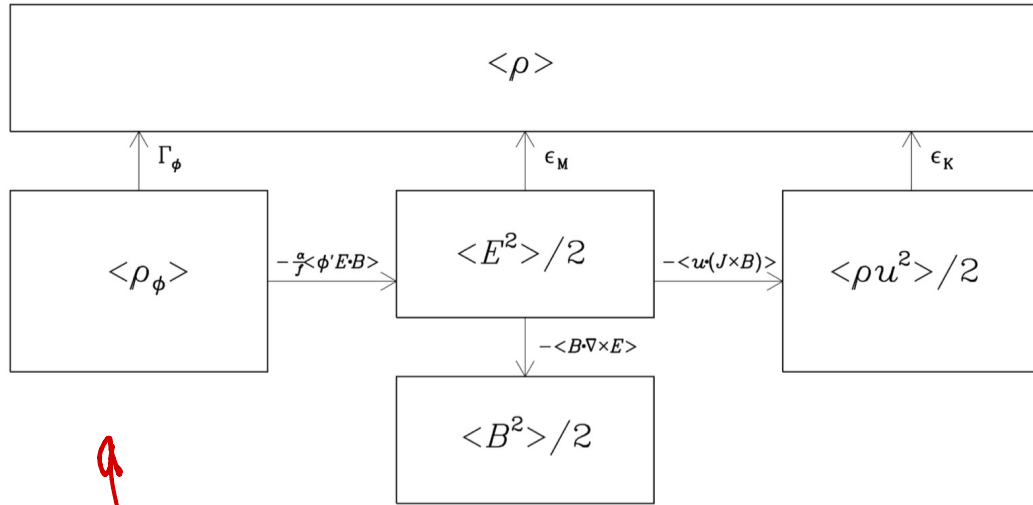
$$\propto \xi_M^{-5/2}$$

Resistive slow-down of turbulent decay



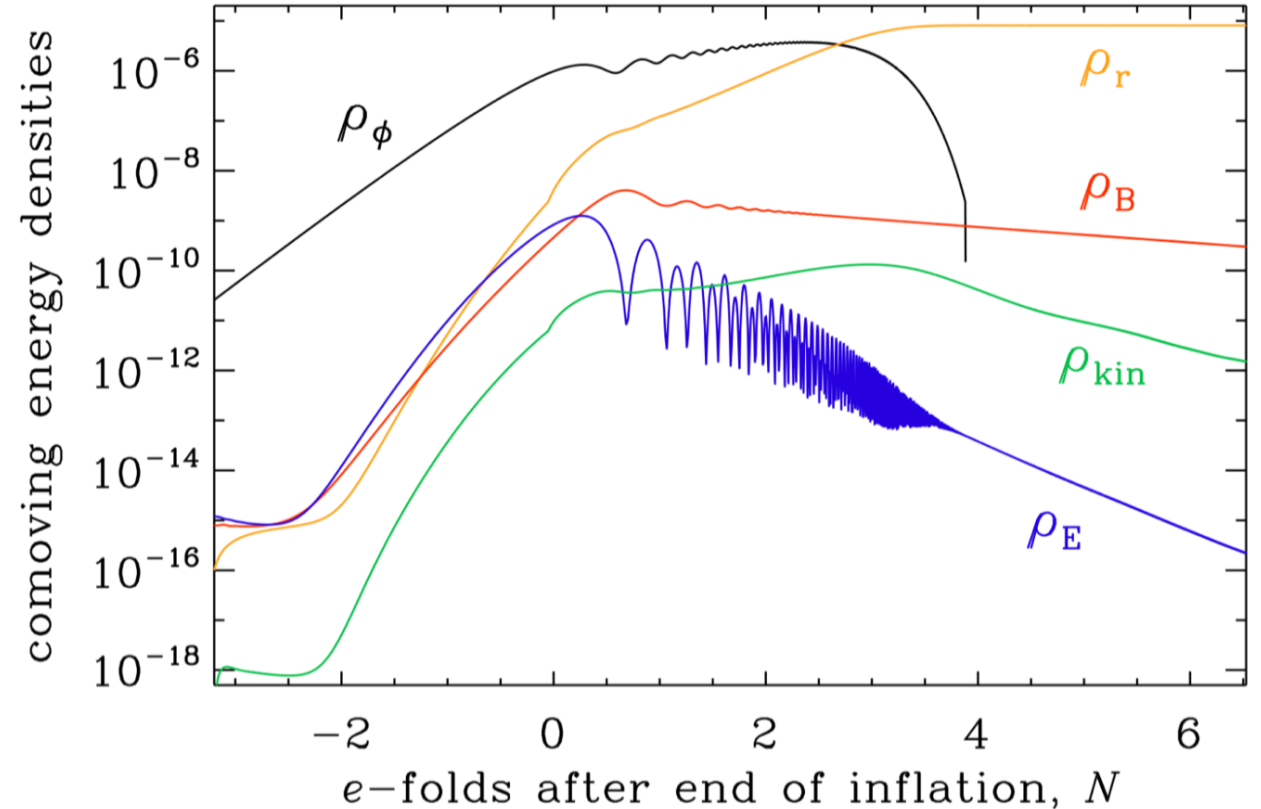
- Endpoints under assumption that decay time = Alfven time
- Use: decay time = recombination time
- Possibility: decay time \gg Alfven time
- \rightarrow Premature endpoint of evolution

Connecting magnetic decay with axion inflation



a

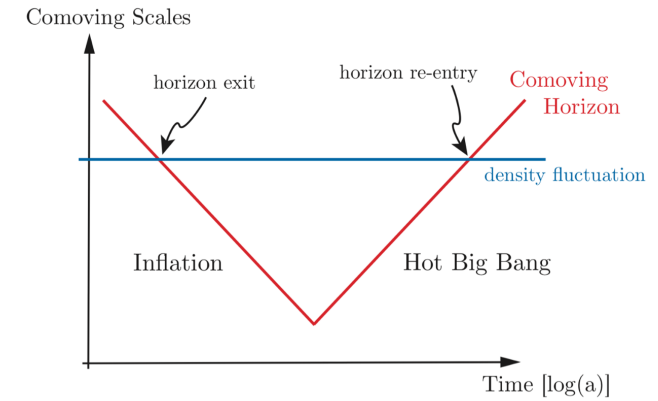
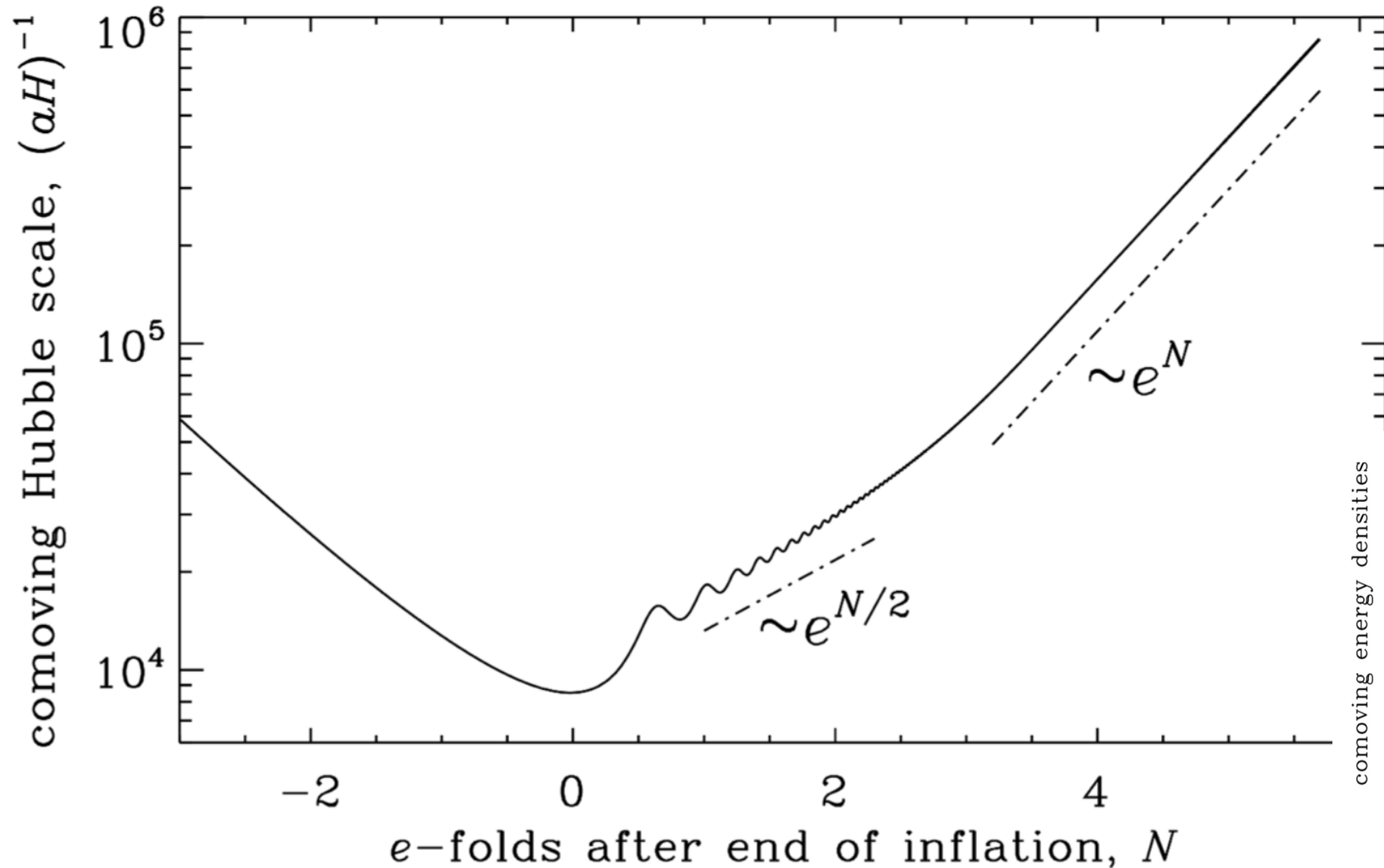
$$\begin{aligned} \partial_\tau^2 \phi + 2\mathcal{H}\partial_\tau \phi - \nabla^2 \phi + a^2 \frac{dV}{d\phi} &= \frac{\alpha}{a^2 f} \mathbf{E} \cdot \mathbf{B}, \\ \partial_\tau \mathbf{E} - \text{rot } \mathbf{B} + \frac{\alpha}{f} (\partial_\tau \phi \mathbf{B} + \nabla \phi \times \mathbf{E}) + \mathbf{J} &= 0, \\ \nabla \cdot \mathbf{E} = -\frac{\alpha}{f} \nabla \phi \cdot \mathbf{B}, \quad \nabla \cdot \mathbf{B} &= 0, \\ \partial_\tau \mathbf{B} + \text{rot } \mathbf{E} &= 0, \\ \mathcal{H}^2 = \frac{8\pi}{3m_{\text{Pl}}^2} a^2 (\rho_\phi + \rho_E + \rho_B + \rho_\chi), \end{aligned}$$



NORDITA-2025-030

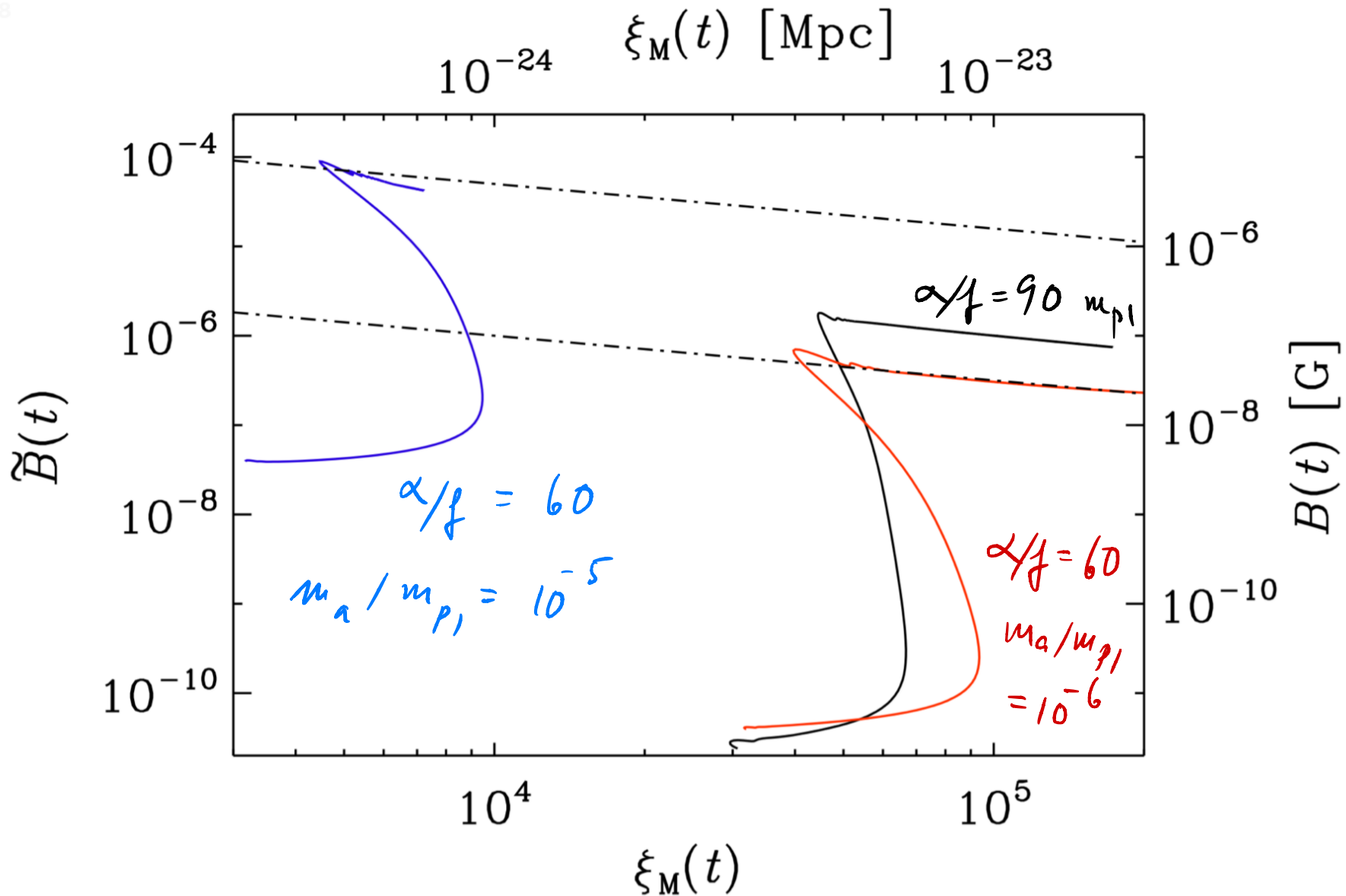
Schwinger effect in axion inflation on a lattice

Building a radiation-dominated universe



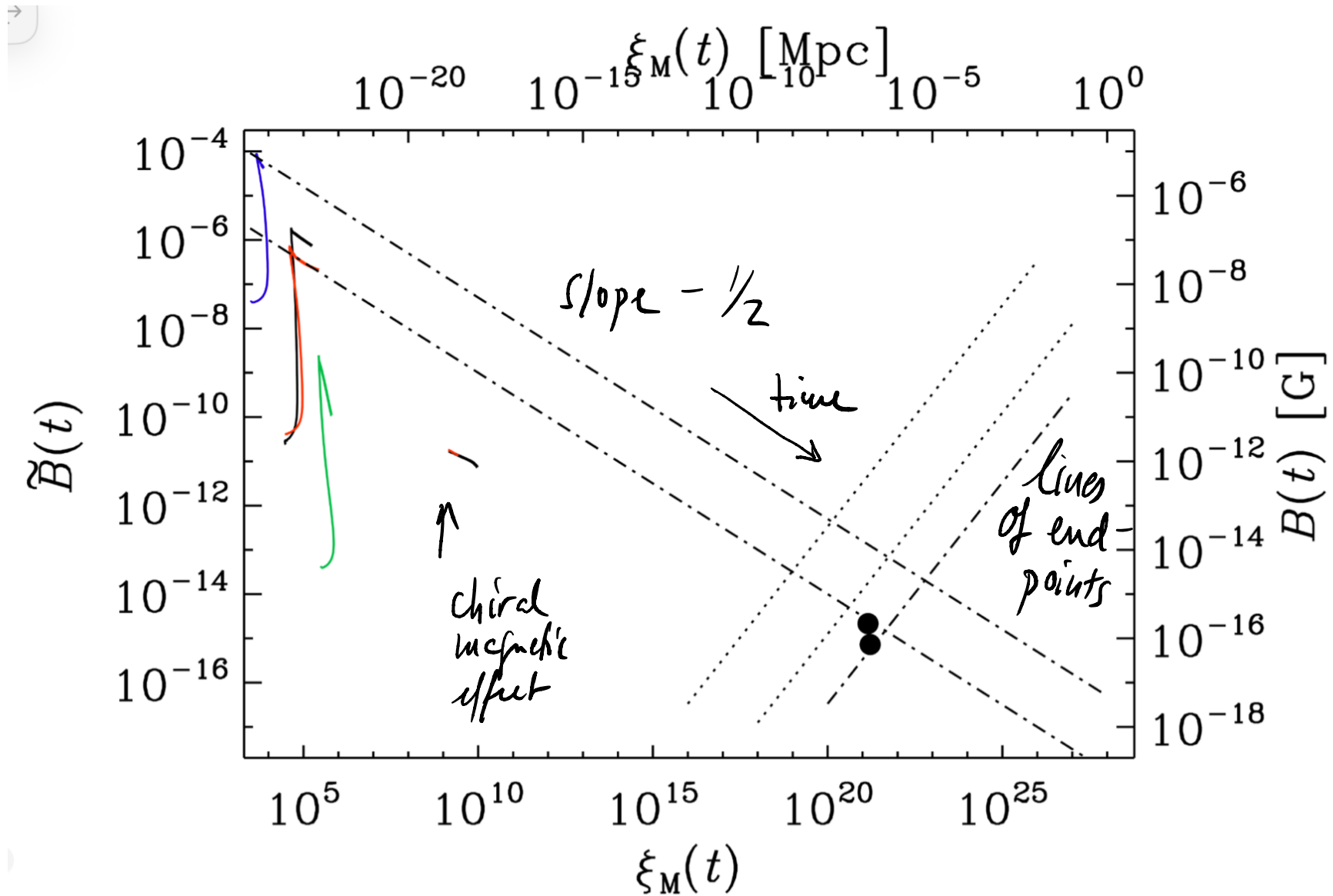
Diagnostic diagram

18



Complicated evolution
before entering inverse
cascade

.... until today



Conclusions

tools are available to explore “full” intergalactic / cosmological magnetic fields parameter space, from the moment of creation to recombination and throughout structure formation up to $z = 0$

Interplay:

Gamma-ray, radio, CRs

GWs, CMB

simulations

